

Scale setting on the 2+1+1 HISQ ensembles: progress report

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[Fermilab Lattice and MILC Collaborations]

Lattice 2024
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Outline

- FLAG: gradient flow scales
- The gradient flow: definitions, corrections, integration
- Relative scale and the integrated autocorrelation time
- Absolute scale: a few examples — w_0/r_1 , $w_0 f_\pi$, $w_0 M_\Omega$
- Conclusion

FLAG 2023 update: gradient flow scales

Collaboration	Ref.	N_f	publication status	chiral extrapolation	continuum extrapolation	finite volume	physical scale	$\sqrt{t_0}$ [fm]	w_0 [fm]
ETM 21	[43]	2+1+1	A	★	★	★	f_π	0.14436(61)	0.17383(63)
CalLat 20A	[115]	2+1+1	A	★	★	★	m_Ω	0.1422(14)	0.1709(11)
BMW 20	[119]	1+1+1+1	A	★	★	★	m_Ω		0.17236(29)(63)[70]
ETM 20	[1057]	2+1+1	C	★	★	★	f_π		0.1706(18)
MILC 15	[116]	2+1+1	A	★	★	★	$F_{p4s}(f_\pi)^\#$	0.1416(+8/-5)	0.1714(+15/-12)
HPQCD 13A	[40]	2+1+1	A	★	○	★	f_π	0.1420(8)	0.1715(9)
RQCD 22	[1058]	2+1	P	★	★	★	m_Ξ	0.1449(+7/-9)	
CLS 21	[1059]	2+1	C	★	★	★	f_π, f_K	0.1443(7)(13)	
CLS 16	[117]	2+1	A	○	★	★	f_π, f_K	0.1467(14)(7)	
QCDSF/UKQCD 15B	[718]	2+1	P	○	○	○	$m_P^{SU(3)}$	0.1511(22)(6)(5)(3)	0.1808(23)(5)(6)(4)
RBC/UKQCD 14B	[10]	2+1	A	★	★	★	m_Ω	0.14389(81)	0.17250(91)
HotQCD 14	[120]	2+1	A	★	★	★	$r_1(f_\pi)^\#$		0.1749(14)
BMW 12A	[118]	2+1	A	★	★	★	m_Ω	0.1465(21)(13)	0.1755(18)(4)

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MILC HISQ ensembles

$\approx a$ (fm)	Key	β	am'_l	am'_s	am'_c	$(L/a)^3 \times (T/a)$	L (fm)	M_π (MeV)	$M_\pi L$	N_{conf}
0.15	$m_s/5$	5.80	0.013	0.065	0.838	$16^3 \times 48$	2.45	305	3.8	1020
0.15	$m_s/10$	5.80	0.0064	0.064	0.828	$24^3 \times 48$	3.67	214	4.0	1000
0.15	physical	5.80	0.00235	0.0647	0.831	$32^3 \times 48$	4.89	131	3.3	1000
0.12	$m_s/5$	6.00	0.0102	0.0509	0.635	$24^3 \times 64$	2.93	305	4.5	1040
0.12	unphysA	6.00	0.0102	0.03054 [†]	0.635	$24^3 \times 64$	2.93	304	4.5	1020
0.12	small	6.00	0.00507	0.0507	0.628	$24^3 \times 64$	2.93	218	3.2	1020
0.12	$m_s/10$	6.00	0.00507	0.0507	0.628	$32^3 \times 64$	3.91	217	4.3	1000
0.12	large	6.00	0.00507	0.0507	0.628	$40^3 \times 64$	4.89	216	5.4	1028
0.12	unphysB	6.00	0.01275	0.01275 [†]	0.640	$24^3 \times 64$	2.93	337	5.0	1020
0.12	unphysC	6.00	0.00507	0.0304 [†]	0.628	$32^3 \times 64$	3.91	215	4.3	1020
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0.12	unphysG	6.00	0.0088725	0.022815 [†]	0.628	$32^3 \times 64$	3.91	282	5.6	1020
0.12	physical	6.00	0.00184	0.0507	0.628	$48^3 \times 64$	5.87	132	3.9	999
0.09	$m_s/5$	6.30	0.0074	0.037	0.440	$32^3 \times 96$	2.81	316	4.5	1005
0.09	$m_s/10$	6.30	0.00363	0.0363	0.430	$48^3 \times 96$	4.22	221	4.7	999
0.09	physical	6.30	0.0012	0.0363	0.432	$64^3 \times 96$	5.62	129	3.7	484

Table from Fermilab-MILC, 1712.09262

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'2013



Table from Fermilab-MILC, 1712.09262

MILC HISQ ensembles

0.06	$m_s/5$	6.72	0.0048	0.024	0.286	$48^3 \times 144$	2.72	329	4.5	1016
0.06	$m_s/10$	6.72	0.0024	0.024	0.286	$64^3 \times 144$	3.62	234	4.3	572
0.06	physical	6.72	0.0008	0.022	0.260	$96^3 \times 192$	5.44	135	3.7	842
0.042	$m_s/5$	7.00	0.00316	0.0158	0.188	$64^3 \times 192$	2.73	315	4.3	1167
0.042	physical	7.00	0.000569	0.01555	0.1827	$144^3 \times 288$	6.13	134	4.2	420
0.03	$m_s/5$	7.28	0.00223	0.01115	0.1316	$96^3 \times 288$	3.09	309	4.8	724

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'2015

'2020



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'2015
'2020
↓

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- Additionally:
 - Retuned physical m_l/m_s at $a = 0.15, 0.12$ and 0.09 (CalLat) fm with larger statistics. CalLat, 2011.12166
 - Larger volume $128^3 \times 96$ at physical m_l/m_s $a = 0.09$ fm.
 - 6 ensembles at $a = 0.06$ and 0.09 fm with lighter-than-physical strange quark mass.

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The gradient flow

- Smoothing of the original gauge field $U_{x,\mu}$ towards stationary points of the action S^f :

Lüscher, 1006.4518

$$\frac{dV_{x,\mu}}{dt} = - \left\{ \partial_{x,\mu} S^f(t) \right\} V_{x,\mu}, \quad V_{x,\mu}(t=0) = U_{x,\mu},$$

where the flow action $S^f = S_{Wilson}$ or $S_{Symanzik}$.

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- (We have not experimented with the Zeuthen flow.) Ramos, Sint, 1508.05552
- Scale setting:

$$t^2 \langle S^o(t) \rangle \Big|_{t=t_0} = \text{Const} \quad \text{or} \quad \left[t \frac{d}{dt} t^2 \langle S^o(t) \rangle \right]_{t=w_0^2} = \text{Const},$$

where the observable $S^o = S_{clover}$ or S_{Wilson} or $S_{Symanzik}$.

- In practice $\text{Const} = 0.3$.

The gradient flow

- For a given combination of the dynamical action, flow action and the observable the leading discretization effects can be canceled at tree level:

$$t^2 S(t) \rightarrow t^2 S_{corr}(t) = \frac{t^2 S(t)}{1 + \sum_{m=1}^4 C_m(a^{2m}/t^m)}$$

Fodor et al, 1406.0827

- Expansion in a^2/t

$$\langle t^2 S(t) \rangle_a = \frac{3(N^2 - 1)g_0^2}{128\pi^2} (C(a^2/t) + O(g_0^2))$$

The gradient flow

	<i>SWS</i>	<i>WWC</i>	<i>SSS</i>	<i>SWW</i>	<i>WSW</i>	<i>WSC</i>
C_2	$1/72$	$-1/24$	$-1/24$	$-1/24$	$5/72$	$-7/72$
C_4	$7/320$	$-1/512$	$1/32$	$1/32$	$23/1280$	$19/2560$
C_6	$-8539/1935360$	$-1/5120$	$-283/27648$	$-283/27648$	$2077/483840$	$-2237/1935360$
C_8	$76819/18579456$	$-1/65536$	$3229/442368$	$3229/442368$	$16049/9289728$	$14419/74317824$
	<i>SSW</i>	<i>WWW</i>	<i>WSS</i>	<i>WWS</i>	<i>SWC</i>	<i>SSC</i>
C_2	$-7/72$	$1/8$	$1/8$	$13/72$	$-5/24$	$-19/72$
C_4	$35/768$	$3/128$	$3/128$	$13/384$	$167/2560$	$145/1536$
C_6	$-5131/276480$	$13/2048$	$13/2048$	$277/30720$	$-58033/1935360$	$-12871/276480$
C_8	$10957/884736$	$77/32768$	$77/32768$	$323/98304$	$457033/24772608$	$52967/1769472$

The gradient flow

	<i>SWS</i>	<i>WWC</i>	<i>SSS</i>	<i>SWW</i>	<i>WSW</i>	<i>WSC</i>
C_2	$1/72$	$-1/24$	$-1/24$	$-1/24$	$5/72$	$-7/72$
C_4	$7/320$	$-1/512$	$1/32$	$1/32$	$23/1280$	$19/2560$
C_6	$-8539/1935360$	$-1/5120$	$-283/27648$	$-283/27648$	$2077/483840$	$-2237/1935360$
C_8	$76819/18579456$	$-1/65536$	$3229/442368$	$3229/442368$	$16049/9289728$	$14419/74317824$
	<i>SSW</i>	<i>WWW</i>	<i>WSS</i>	<i>WWS</i>	<i>SWC</i>	<i>SSC</i>
C_2	$-7/72$	$1/8$	$1/8$	$13/72$	$-5/24$	$-19/72$
C_4	$35/768$	$3/128$	$3/128$	$13/384$	$167/2560$	$145/1536$
C_6	$-5131/276480$	$13/2048$	$13/2048$	$277/30720$	$-58033/1935360$	$-12871/276480$
C_8	$10957/884736$	$77/32768$	$77/32768$	$323/98304$	$457033/24772608$	$52967/1769472$

- Corrections for the relevant gauge-flow-observable combinations that we measure.

Integration of the flow

- We use (6,4) 2N-storage Runge-Kutta method.

Bazavov, 2007.04225

Bazavov, Chuna, 2101.05320

Berland, Bogey, Bailly, Computers and Fluids (2006)

- For all ensembles we integrate the flow at two step sizes $\Delta t = 1/20$, $1/40$ to fully control the global integration error.

Integration error

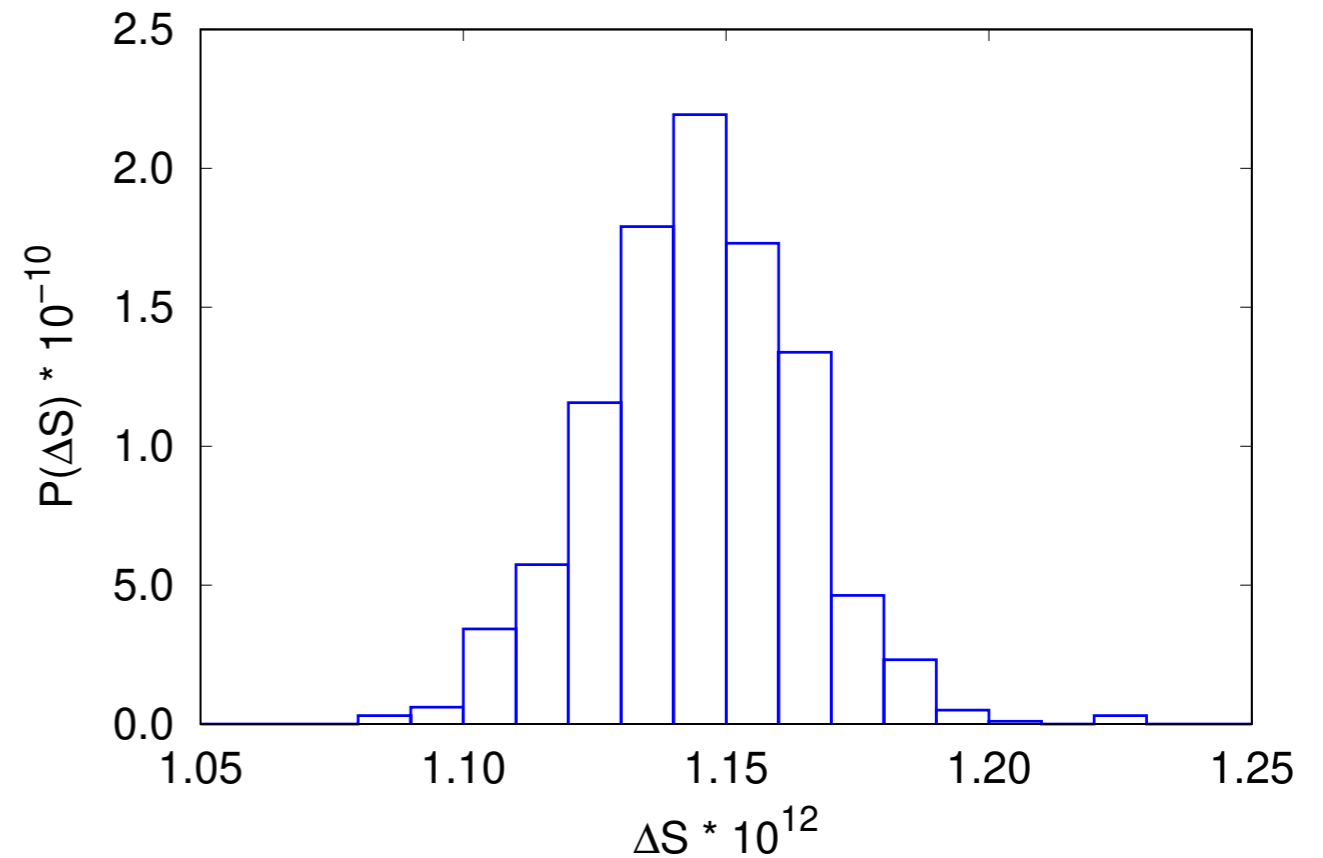
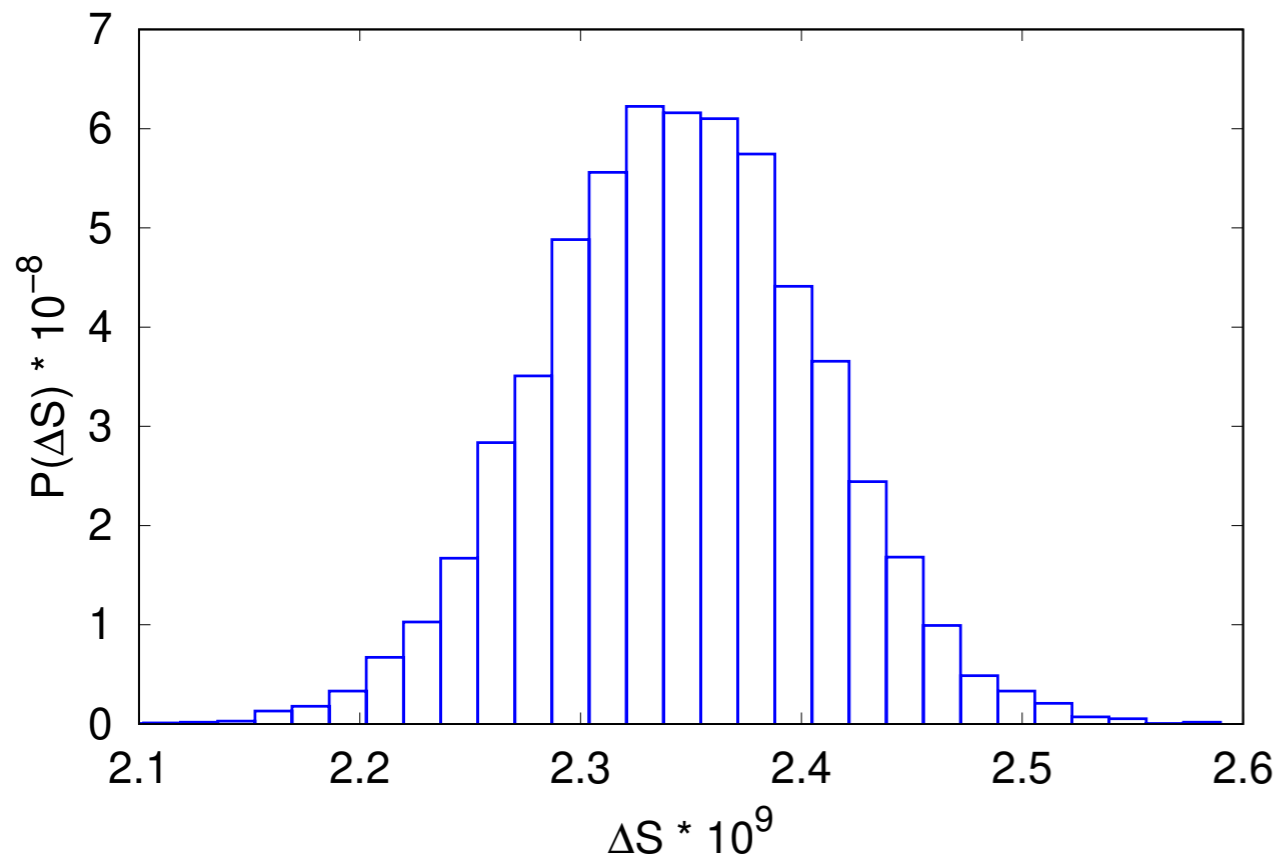
- Define the integration error as

$$\Delta S \equiv \langle S^o(t, \Delta t = 1/40) \rangle \Big|_{t=w_0^2} - \langle S^o(t, \Delta t = 1/20) \rangle \Big|_{t=w_0^2}$$

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- The integration error on the physical mass ensembles at $a = 0.12$ fm (left) and $a = 0.042$ fm (right).

Autocorrelations

- Define the autocorrelation function for an observable \mathcal{O} :

$$C(n) \equiv \langle \mathcal{O}_0 \mathcal{O}_n \rangle - \langle \mathcal{O} \rangle^2$$

- The integrated autocorrelation time

$$\tau_{int} = 1 + 2 \sum_{n=1}^{N-1} \left(1 - \frac{n}{N} \right) \frac{C(n)}{C(0)}, \quad \sigma^2(\bar{\mathcal{O}}) = \frac{\sigma^2(\mathcal{O})}{N} \tau_{int}$$

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- The integrated autocorrelation time

$$\tau_{int} = 1 + 2 \sum_{n=1}^{N-1} \left(1 - \frac{n}{N}\right) \frac{C(n)}{C(0)}, \quad \sigma^2(\bar{\mathcal{O}}) = \frac{\sigma^2(\mathcal{O})}{N} \tau_{int}$$

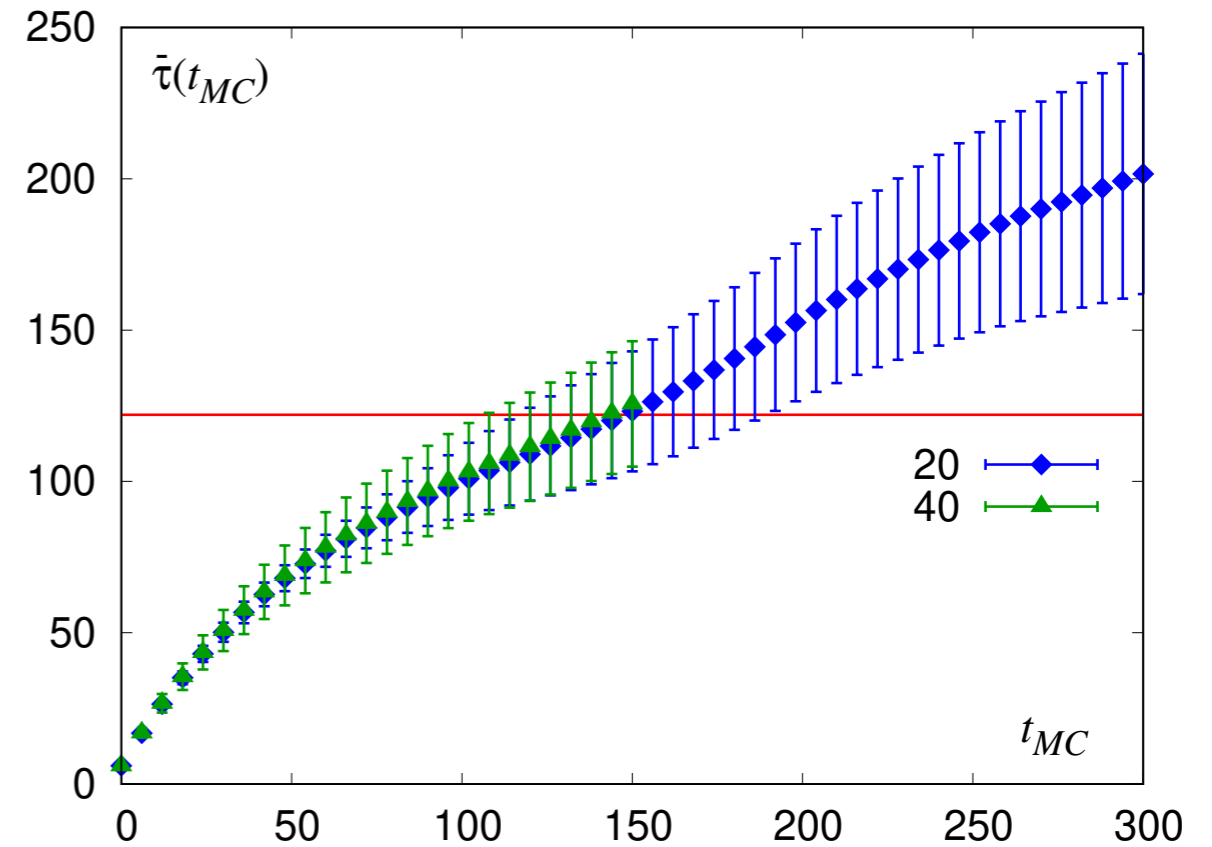
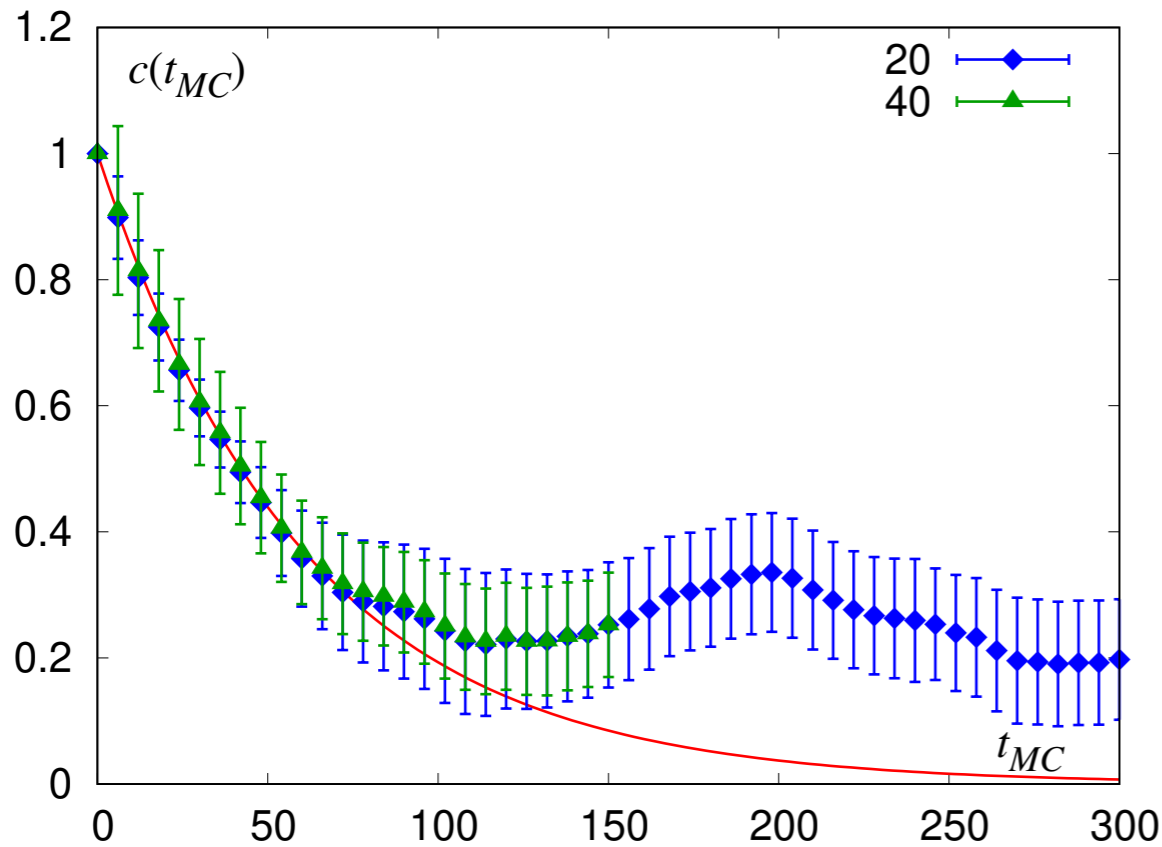
- Window method to estimate the integrated autocorrelation time

$$\tau_{int}(n) = 1 + 2 \sum_{n'=1}^n \frac{C(n')}{C(0)}$$

- If the autocorrelation function is a single exponential

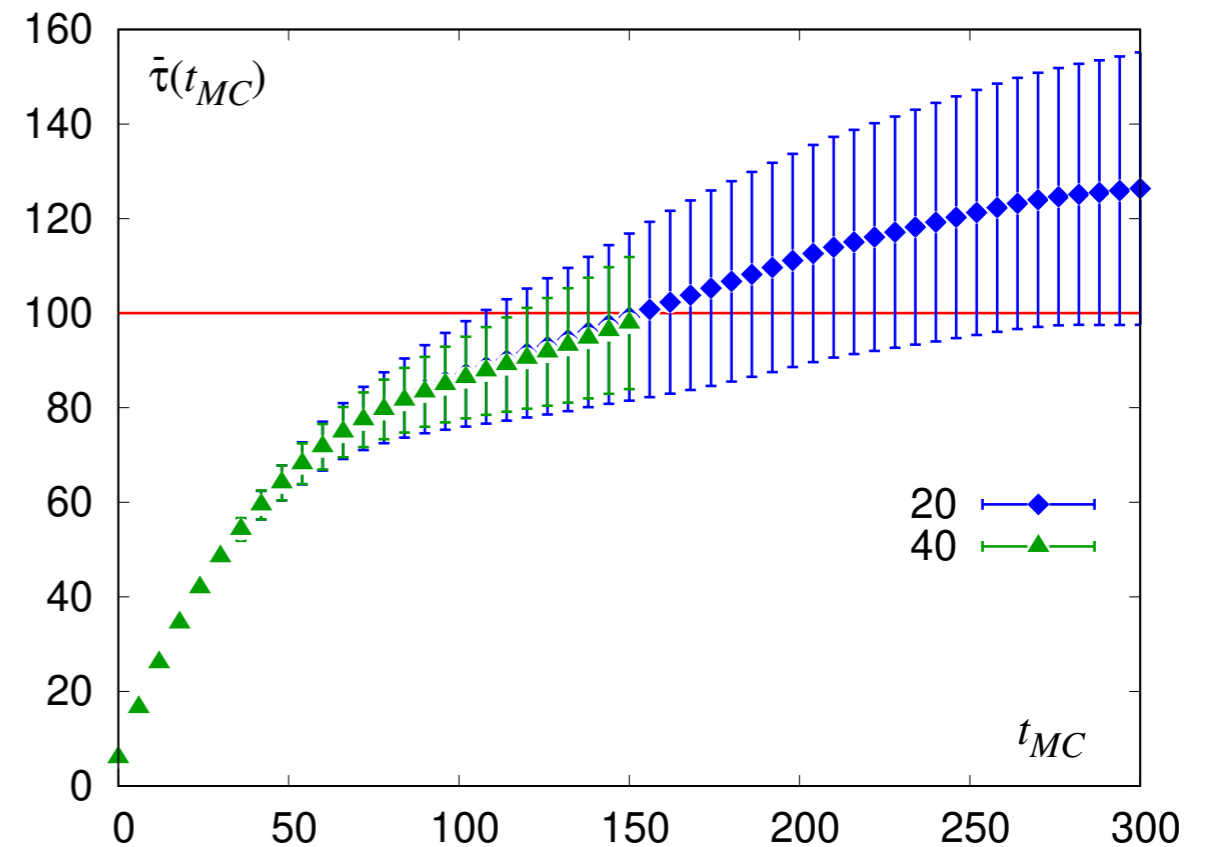
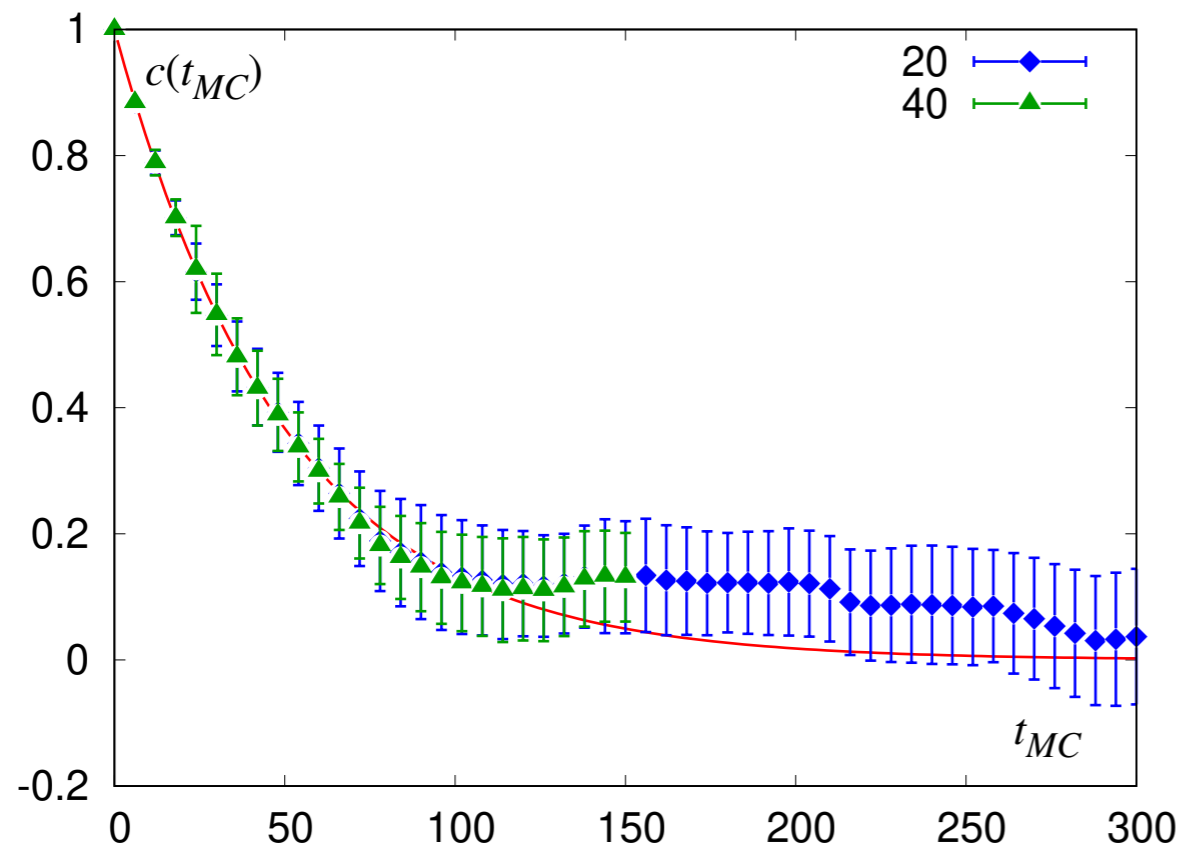
$$C(n) = C(0) \exp(-an) \quad \text{then} \quad \tau_{int}^1 = \frac{e^a + 1}{e^a - 1}$$

Autocorrelations: $a = 0.06$ fm, 300 MeV pion



- MC time series: $\sim 6,000$ time units
- Observable: Clover action density at $t \sim w_0^2$
- Normalized autocorrelation function (left) and integrated autocorrelation time $\tau_{int}(t_{MC})$ (right)
- Single-exponential fit: $\tau_{int}^1 = 122 \pm 31$

Autocorrelations: $a = 0.042$ fm, physical pion

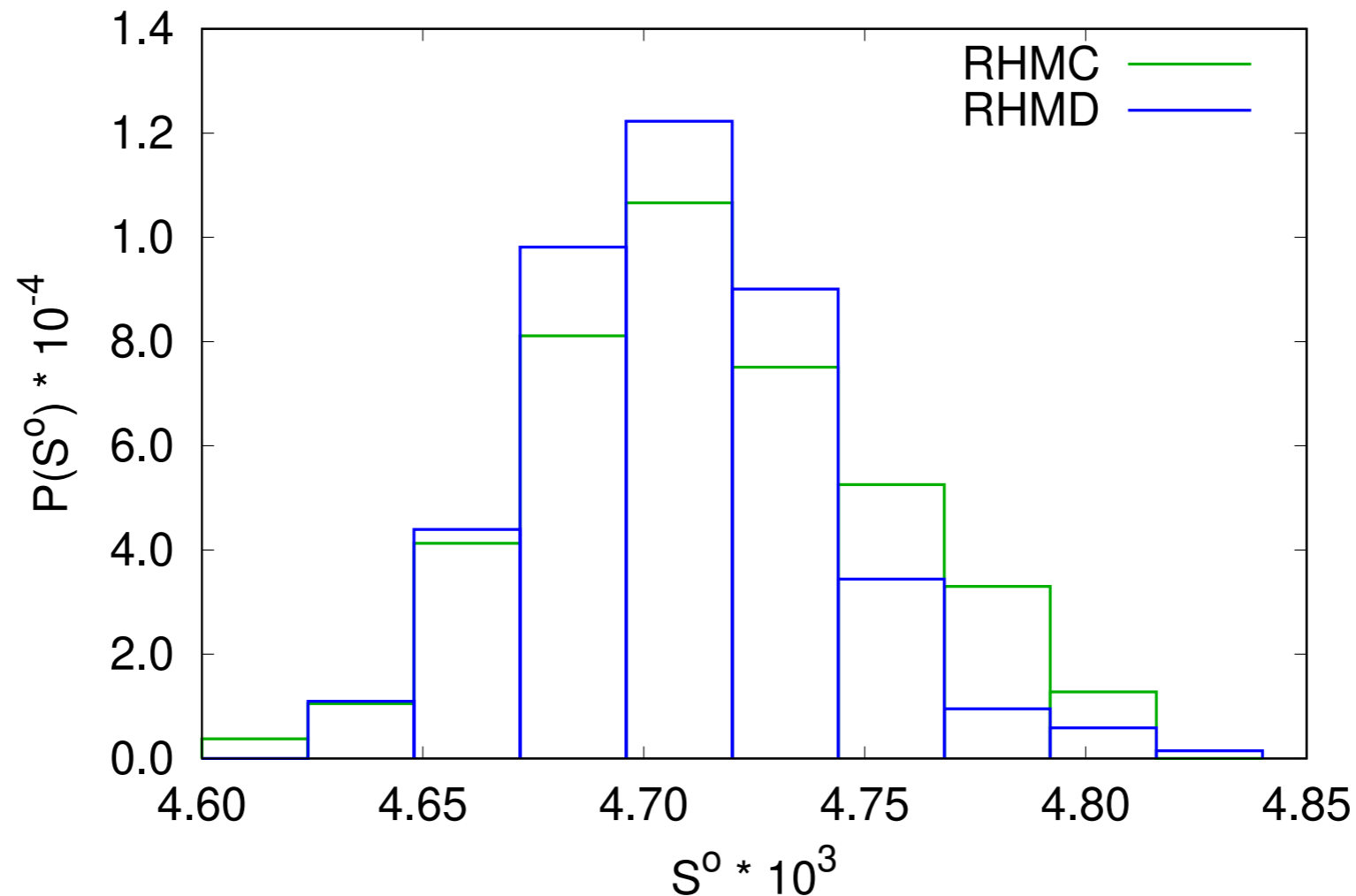


- MC time series: $\sim 6,000$ time units
- Observable: Clover action density at $t \sim w_0^2$
- Normalized autocorrelation function (left) and integrated autocorrelation time $\tau_{int}(t_{MC})$ (right)
- Single-exponential fit: $\tau_{int}^1 = 100 \pm 12$

Relative scale

- Statistical uncertainty:
 - Propagated with jackknife on binned data.
 - Bin size is extrapolated to infinity.

RHMC vs RHMD

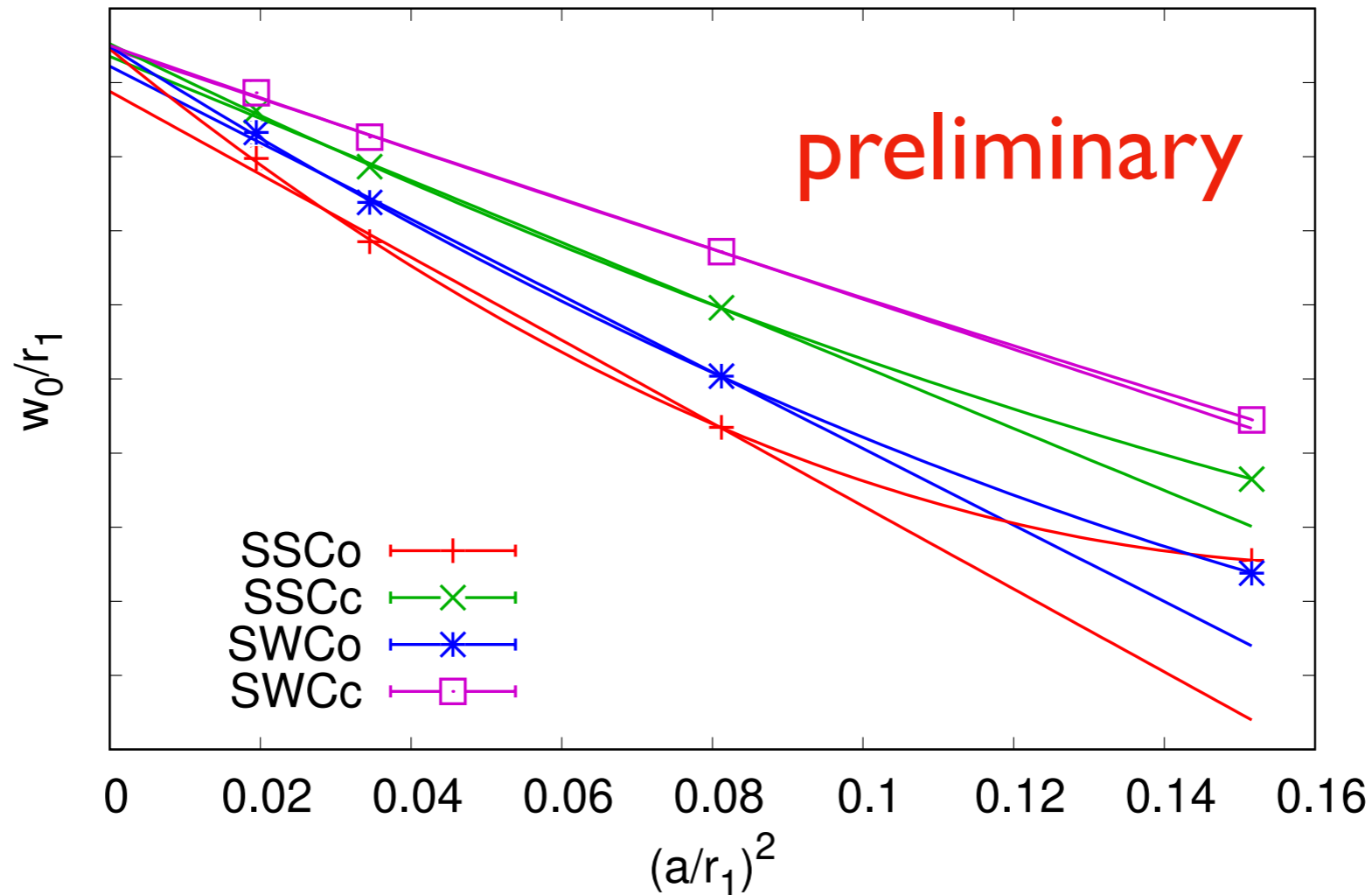


- Histogram of the clover observable at $t \simeq w_0^2$ on the $m_\pi = 200$ MeV $a = 0.06$ fm ensemble
- w_0/a in SSCc: 2.9557(34) RHMC vs 2.9520(47) RHMD

Absolute scale

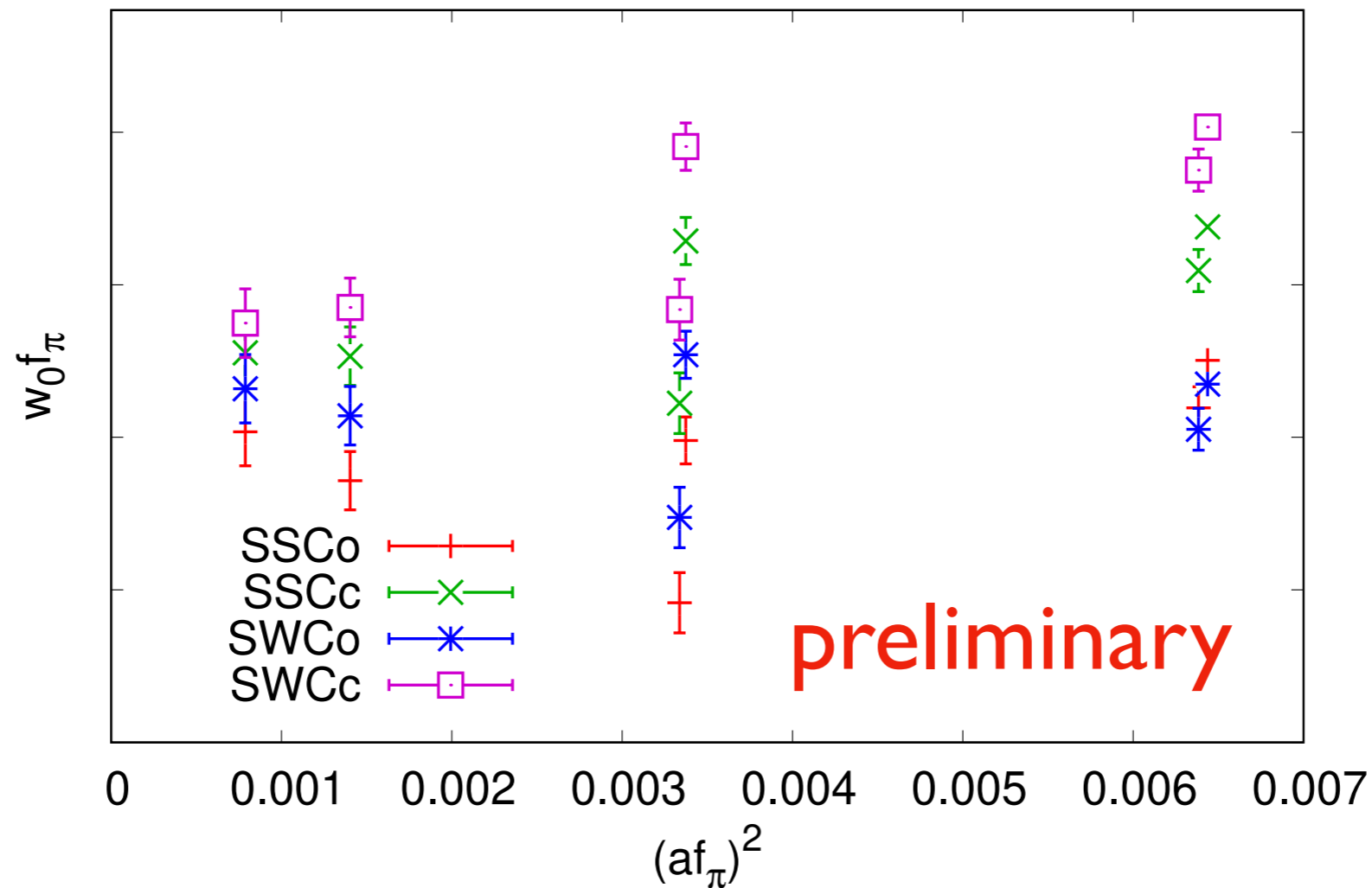
- Our plan:
 - $w_0 f_\pi$ on all ensembles (also as a crosscheck of our earlier work, MILC, PRD 93 (2016), 1503.02769).
 - $w_0 M_\Omega$ on the physical mass ensembles.

$$w_0/r_1$$



- Crosscheck against the r_1 scale that has been recently determined on most of the HISQ ensembles. TUMQCD, 2206.03156
- Simple fits: linear and quadratic in a^2 .

$$w_0 f_\pi$$



- The $w_0 f_\pi$ quantity on the physical mass $a = 0.042, 0.06, 0.09$ (original and retuned) and 0.12 (original and retuned) fm ensembles.
- No corrections of the mass mistuning yet. The magnitude of the effect seems comparable to the spread of the flow-observable schemes.

Omega baryon

- We use HISQ in the valence sector for computing M_Ω .
- General challenges:
 - Signal-to-noise for baryons.
 - Excited states at early Euclidean times.
 - Staggered baryon spectroscopy.

Golterman, Smit, NPB 255 (1985)

Kilcup, Sharpe, NPB 283 (1987)

Bailey, hep-lat/0611023

Hughes, Lin, Meyer, 1912.00028

Omega baryon: selected fits

- Wall-point correlation functions.
- $N + N$ state constrained fits with $N = 1, 2, 3$.
- Scan t_{min} with fixed t_{max} .

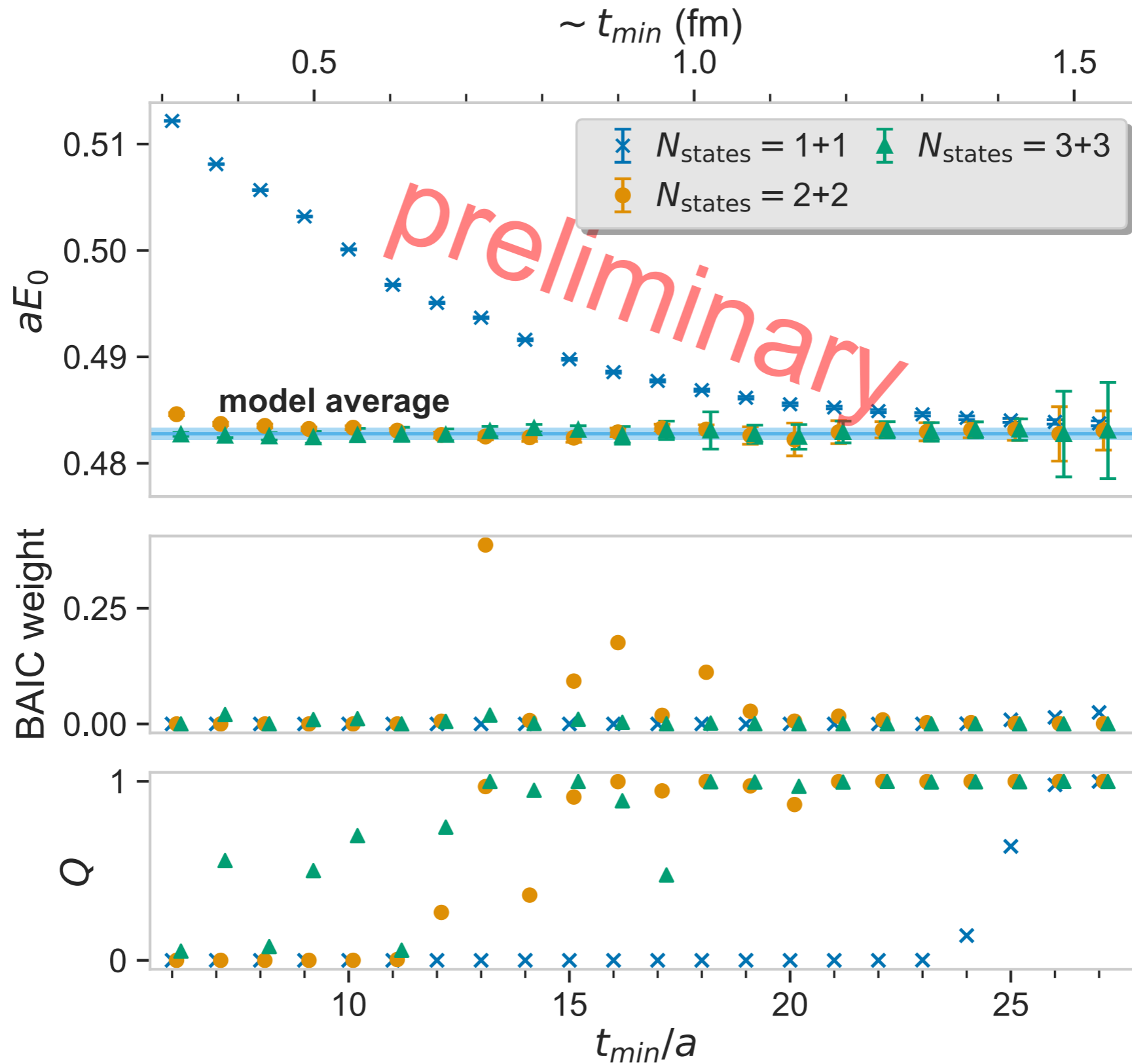
Omega baryon: selected fits

- Wall-point correlation functions.
- $N + N$ state constrained fits with $N = 1, 2, 3$.
- Scan t_{min} with fixed t_{max} .
- To avoid small eigenvalues in the correlation matrix for the $a = 0.042$ fm ensemble every third time slice is included in the fit, for the 0.06 and 0.09 fm ensembles every third time slice is dropped from the fit.

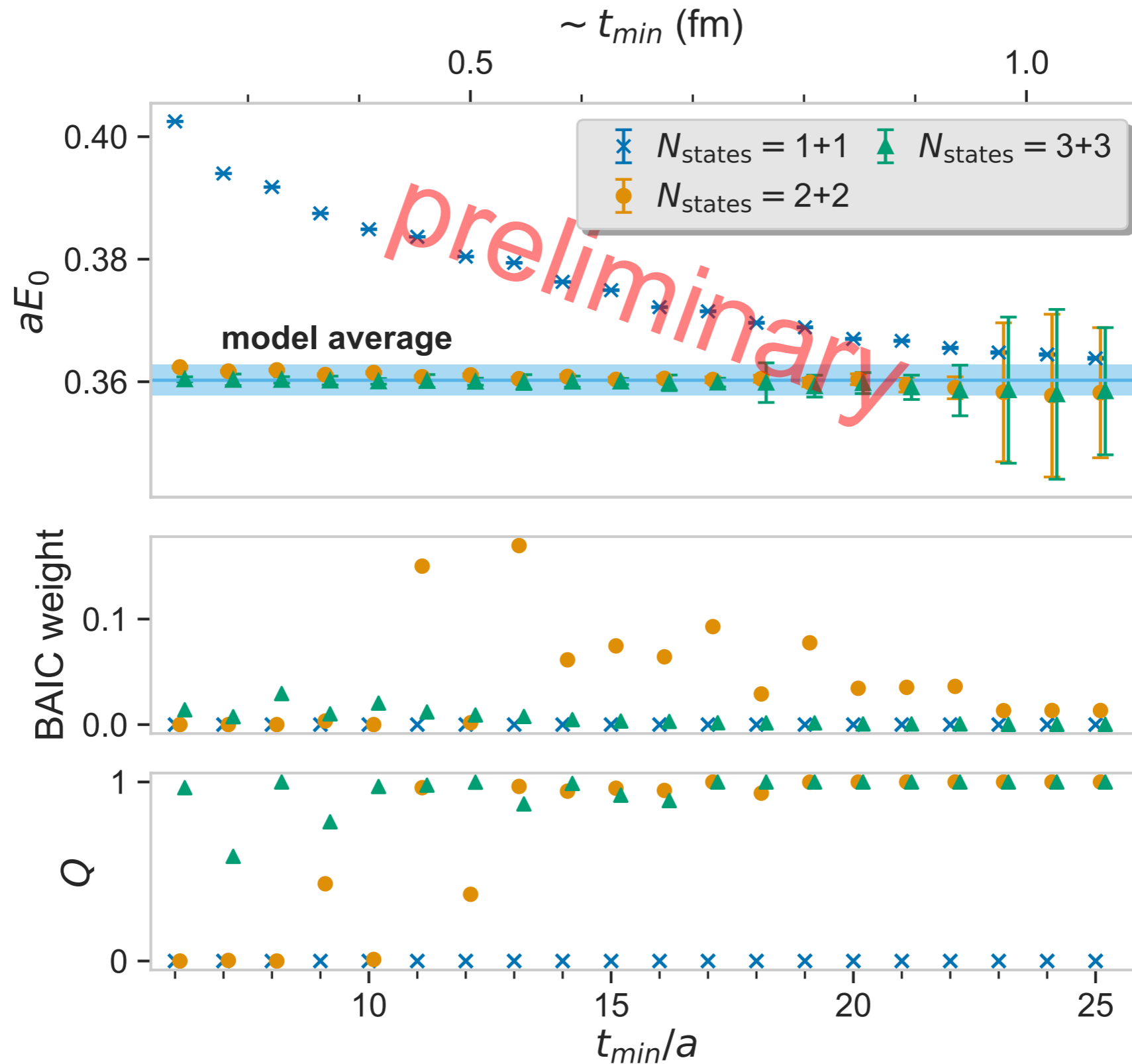
Omega baryon: selected fits

- Wall-point correlation functions.
- $N + N$ state constrained fits with $N = 1, 2, 3$.
- Scan t_{min} with fixed t_{max} .
- To avoid small eigenvalues in the correlation matrix for the $a = 0.042$ fm ensemble every third time slice is included in the fit, for the 0.06 and 0.09 fm ensembles every third time slice is dropped from the fit.
- Fit results from all t_{min} are averaged with the Bayesian Model Averaging (BMA) procedure. Jay, Neil, 2008.01069
Neil, Sitison, 2208.14983

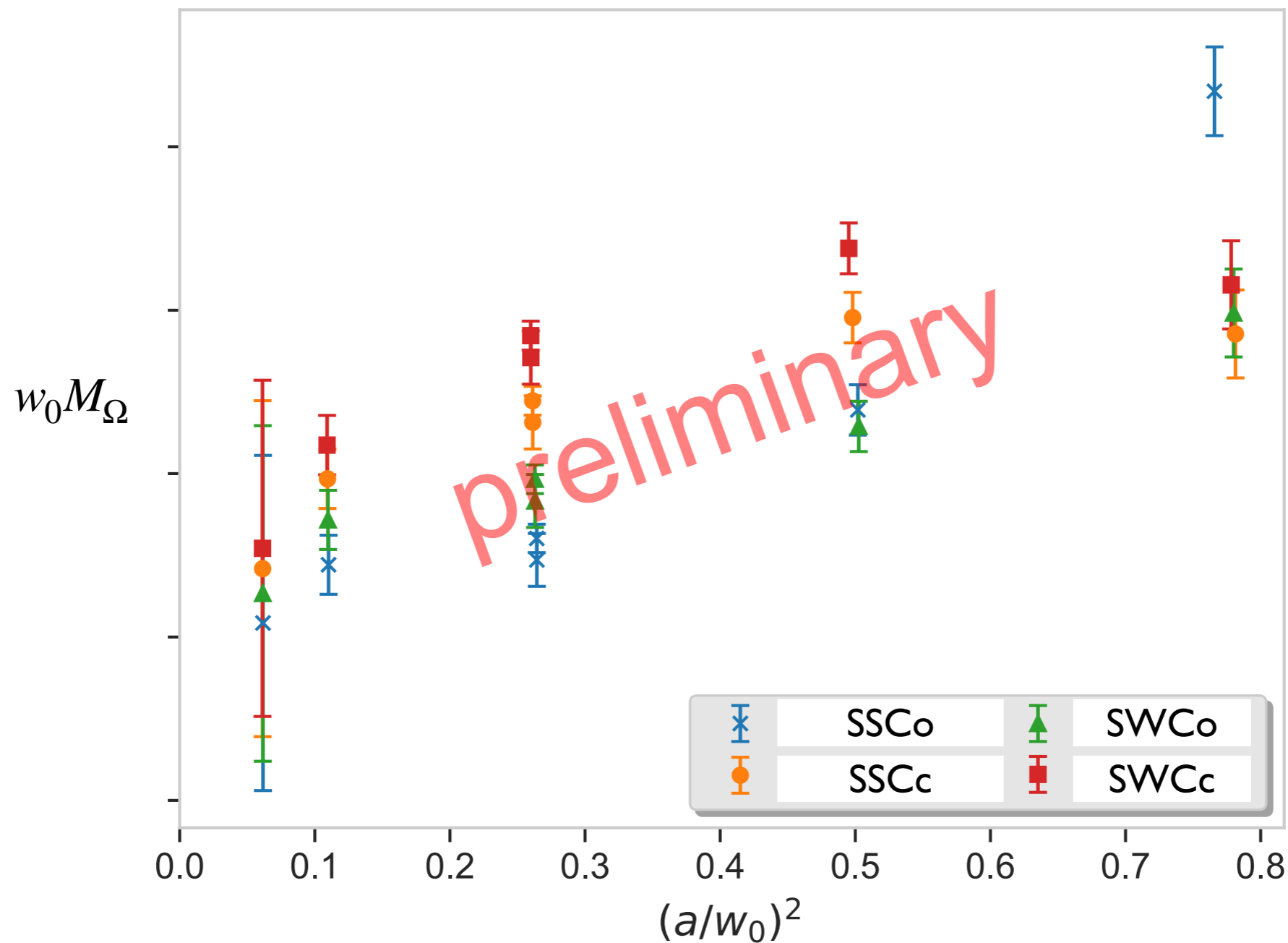
M_Ω on the $a = 0.06$ fm ensemble



M_Ω on the $a = 0.042$ fm ensemble



$w_0 M_\Omega$ on all physical mass ensembles



- Not corrected for the strange quark mass mistuning.

Conclusion

- Ongoing program of the gradient flow scales $\sqrt{t_0}/a$ and w_0/a computations for all MILC HISQ ensembles with two flow and three observable combinations.
- Ongoing computation of aM_Ω with HISQ on the physical-mass ensembles.
- Next steps:
 - Full chiral-continuum analysis of $w_0 f_\pi$.
 - Continuum extrapolation of $w_0 M_\Omega$.
 - Adding electromagnetic effects for M_Ω .