

# How Primordial Black Holes constrain leptogenesis



Università degli Studi di Napoli Federico II

Based on works

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in collaboration with

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And work in preparation in collaboration with

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Primordial Black Holes constrain (resonant) leptogenesis

 $M_N < 10^3 \,\mathrm{GeV}$ 

### Primordial Black Holes constrain (resonant) leptogenesis $M_N < 10^3 \, {\rm GeV}$





#### Leptogenesis - overview

$$\mathcal{L}_{\text{seesaw}} = \frac{1}{2} \bar{N}_i^c \hat{M}_{ij} N_j - Y_{\ell i} \bar{L}_\ell \tilde{\phi} N_i + \text{h.c.},$$

$$10^{-2}$$

Low scale



 $10^{16}$ 



Low scale



 $10^{16}$ 



Occurs at very high temperatures

Theoretically appealing



Far beyond the reach of direct detection

Low scale

 $M_N(\text{GeV})$ 



Occurs at low temperatures

Requires some degree

of (unexplained) degeneracy

Within the reach of experiments

Occurs at very high temperatures

Theoretically appealing

 $10^{16}$ 

Far beyond the reach of direct detection

Low scale





Leptogenesis is difficult to constrain at all scales

Occurs at low temperatures

Occurs at very high temperatures

Requires some degree of (unexplained) degeneracy

Within the reach of experiments

Theoretically appealing

 $10^{16}$ 

Far beyond the reach of direct detection

Low scale

 $M_N(\text{GeV})$ 

$$\begin{pmatrix} M_{N_1} & 0 & 0 \\ 0 & M_{N_2} & 0 \\ 0 & 0 & M_{N_3} \end{pmatrix}$$
$$M_{N_i} \approx M_{N_j}, M_{N_k}$$
$$M_{N_i} - M_{N_j} = \Delta M_{ij}$$

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$$\begin{pmatrix} M_{N_1} & 0 & 0 \\ 0 & M_{N_2} & 0 \\ 0 & 0 & M_{N_3} \end{pmatrix}^{10^{-4}}$$
 Total active-sterile  $_{10^{-5}}$  Standard  $U^2 = U^2(x, y, m_h, M)$   $U^2 = U^2(x, y, m_h, M)$   $M_{N_i} \approx M_{N_j}, M_{N_k}$   $M_{N_i} - M_{N_j} = \Delta M_{ij}$   $U^{10^{-10}}$   $U^{1$ 



#### How do PBHs affect leptogenesis?

**Production of RHNs** 

Cosmological evolution

Entropy injection

### $\frac{\rm Production~of~RHNs}{\rm dt} \stackrel{\rm dM_{\rm PBH}}{\propto} M_{\rm PBH}^{-2}$



Boosted RHNs produced by PBH

## $\frac{\rm Production \ of \ RHNs}{\rm dt} \ \frac{\rm dM_{\rm PBH}}{\rm dt} \propto M_{\rm PBH}^{-2}$

 $M_{N_i} \leq T_{\text{PBH}}$ 

Boosted RHNs decay seeding asymmetry

Boosted RHNs produced by PBH



 $\frac{\mathrm{dY}_{\mathrm{B}}}{\mathrm{dt}} = \epsilon \, \Gamma_{N_i}^{\mathrm{PBH}} \mathcal{N}_{N_i}$ 



#### **Cosmological evolution**

#### PBHs redshift like matter



#### **Cosmological evolution**

PBHs redshift like matter Large populations dominate the energy budget



#### **Cosmological evolution**

PBHs redshift like matter Large populations dominate the energy budget

Increased Hubble rate close to evaporation



#### Sphaleron freeze-out



#### Sphaleron freeze-out











#### Mutual exclusion limits



#### Mutual exclusion limits



#### Mutual exclusion limits



#### Hope on the Gravitational Wave horizon?



Bhaumik, Jain, Lewicki
















## **PBH Hot-spots formation**

Hawking radiation deposits energy at particular r



## **PBH Hot-spots formation**



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## **PBH Hot-spot profiles**



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#### **PBH Hot-spot profiles**



 $\lambda(T(r)) = \frac{1}{\Gamma_X(T(r))} \longrightarrow \Gamma_X(r) = \Gamma_S(r) + \Gamma_D(r)$ 

 $\frac{\mathbf{L}}{\chi(T(r))} \longrightarrow \Gamma_X(r) = \Gamma_S(r) + \Gamma_D(r)$  $\lambda(T(r))$ Insufficient

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Insufficient

$$P(r) = e^{-\int_0^r \Gamma_X(r') \mathrm{d}r'}$$

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$$P(r) = e^{-\int_0^r \Gamma_X(r') \mathrm{d}r'}$$

$$P\left(\lambda\right) \equiv \frac{1}{2}$$





$$\mathcal{N}_{X}^{\text{escape}} = \int_{\alpha_{\text{form}}}^{\alpha_{\text{evap}}} \frac{\log(10)}{H} P(r_{\text{HS}}) \Gamma_{\text{PBH} \to X} d\alpha$$









# Outlook

Where are PBHs incompatible with alternative mechanisms of baryogenesis (type-2 seesaw, GUT baryogenesis)

Does the standard treatment of entropy injection need modification in the presence of hot-spots?

How do hot-spots cool after PBH evaporation?

Is Hawking radiation ever able to equilibriate across the entire universe?

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SM is extended by at least two gauge singlet Right Handed Neutrinos (RHNs)

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 $m_{\nu} \approx -v_{\rm EW}^2 Y M_N^{-1} Y^T$ Provides mechanism of

neutrino mass generation

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 $\mathbf{v} = \frac{1}{2} \bar{N}_{i}^{c} \hat{M}_{ij} N_{j} - Y_{\ell i} \bar{L}_{\ell} \tilde{\phi} N_{i} + \text{h.c.},$   $\sum_{\text{EWSB}} \int_{m_{D} \ll M_{N}} m_{D} \ll M_{N}$  $m_{\nu} \approx -v_{\rm EW}^2 Y M_N^{-1} Y^T$ 

> Provides mechanism of neutrino mass generation

Fits neutrino mass splitting and mixing data

 $\mathcal{L}_{\text{seesaw}}$ 

$$\Delta m_{\rm sol}^2 = m_m^2 - m_l^2$$
$$\Delta m_{\rm atm}^2 = m_h^2 - m_l^2$$

 $\mathbf{v} = \frac{1}{2} \bar{N}_{i}^{c} \hat{M}_{ij} N_{j} - \underbrace{Y_{\ell i} \bar{L}_{\ell} \tilde{\phi} N_{i} + \text{h.c.}}_{\text{evsb}},$ SM is extended by at least two gauge singlet Right Handed Neutrinos (RHNs)

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 $\mathcal{L}_{\text{seesaw}}$ 

$$\Delta m^2_{\rm sol} = m^2_m - m^2_l \qquad \qquad {\rm free \ parameter} \\ \Delta m^2_{\rm atm} = m^2_h - m^2_l \qquad \qquad {\rm free \ parameter} \\ \qquad {\rm free \ parameter} \\$$

$$\mathcal{L}_{ ext{seesaw}} = rac{1}{2} ar{N_i}^c \hat{M}_{ij} N_j - egin{matrix} Y_{\ell i} ar{L}_\ell ilde{\phi} N_i + ext{h.c.}, \ & \swarrow \ Y = rac{1}{v_0} U_{ ext{PMNS}} \cdot \sqrt{\hat{m}_
u} \cdot R^T \cdot \sqrt{\hat{M}}, \end{cases}$$

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u} \cdot R^T \cdot \sqrt{\hat{M}}, \ N_R = egin{pmatrix} 0 & \cos(z_{23}) & \sin(z_{23}) \ 0 & -\sin(z_{23}) & \cos(z_{23}) \ 1 & 0 & 0 \end{pmatrix}$$

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u} \cdot R^T \cdot \sqrt{\hat{M}}, \ X = egin{pmatrix} 0 & \cos(z_{23}) & \sin(z_{23}) \ 0 & -\sin(z_{23}) & \cos(z_{23}) \ 1 & 0 \ \end{pmatrix} \ z_{23} = x + i y$$

$$\mathcal{L}_{\text{seesaw}} = \frac{1}{2} \bar{N}_{i}^{c} \hat{M}_{ij} N_{j} - \underbrace{Y_{\ell i} \bar{L}_{\ell} \tilde{\phi} N_{i} + \text{h.c.}}_{Y = \frac{1}{v_{0}} U_{\text{PMNS}} \cdot \sqrt{\hat{m}_{\nu}} \cdot R^{T} \cdot \sqrt{\hat{M}},$$

$$\stackrel{\ell}{\underset{\ell}{\overset{\ell}{\overset{\ell}{\overset{N_{2,3}}{\overset{\ell}{\overset{N_{2,3}}{\overset{\ell}{\overset{\Gamma}{\overset{N_{2,3}}{\overset{\Gamma}{\overset{N_{2,3}}{\overset{\ell}{\overset{\Gamma}{\overset{N_{2,3}}{\overset{N_{2,3}}{\overset{\Gamma}{\overset{N_{2,3}}}{\overset{N_{2,3}}{\overset{N_{2,3}}{\overset{N_{2,3}}{\overset{N_{2,3}}{\overset{N_{2,3}}{\overset{N_{2,3}}{\overset{N_{2,3}}{\overset{N_{2,3}}{\overset{N_{2,3}}{\overset{N_{2,3}}{\overset{N_{2,3}}{\overset{N_{2,3}}{\overset{N_{2,3}}{\overset{N_{2,3}}{\overset{N_{2,3}}{\overset{N_{2,3}}}{\overset{N_{2,3}}{\overset{N_{2,3}}}{\overset{N_{2,3}}}{\overset{N_{2,3}}{\overset{N_{2,3}}}{\overset{N_{2,3}}}{\overset{N_{2,3}}}{\overset{N_{2,3}}}{\overset{N_{2,3}}}{\overset{N_{2,3}}{\overset{N_{2,3}}}{\overset{N_{2,3}}{\overset{N_{2,3}}}{\overset{N_{2,3}}}}{\overset{N_{2,3}}}{\overset{N_{2,3}}}{\overset{N_{2,3}}}{\overset{N_{2,3}}}{\overset{N_{2,3}}}}{\overset{$$





## Leptogenesis – High Scale Models

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 $M_{N_1} > 10^9 \text{GeV}$ 

Davidson-Ibarra limit

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 $M_{N_1} > 10^9 \text{GeV}$ 

Davidson-Ibarra limit

Far out of the reach of direct detection

Dynamics at very high scales

Fewer relevant degrees of freedom

Mass hierarchy can be arbitrary

## Free parameters

**Common RHN** mass  $M_{ij}$ Mass splitting -ow scale ratio  $\mathcal{X}$ Real and imaginary parts of mixing Y angle

 $m_h$  Mass of heaviest active neutrino

## Free parameters



-ow scale

 $M_{N_1}$ scale  $\mathcal{X}$ High Y Пh

#### RHN mass

Real and imaginary parts of mixing angle

Mass of heaviest active neutrino

#### Mutual exclusion limits – High scale



## Mutual exclusion limits – High scale

High scale


## Mutual exclusion limits – High scale

High scale

