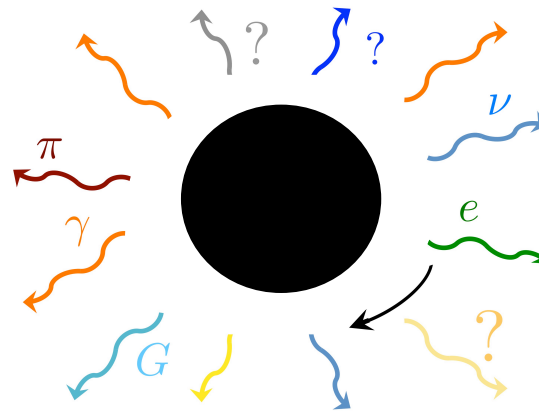


Evaporating Kerr black holes as probes of new physics (in the string axiverse)



João G. Rosa

University of Coimbra, Portugal
arXiv:2312.09261 (with Marco Calzà)

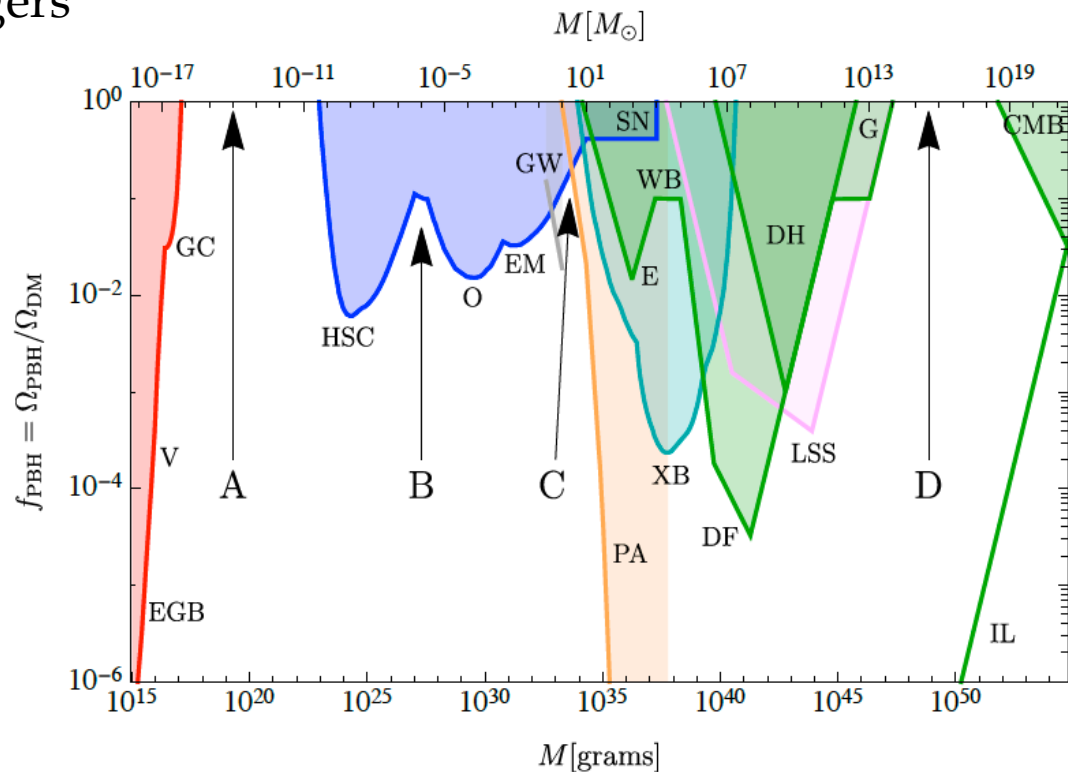
New Horizons in Primordial Black Holes, Edinburgh
18 June 2024

Primordial black holes

Several cosmological scenarios predict formation of primordial BHs (non-standard inflation, phase transitions, etc) [c.f. Escrivá & Kuhnel (2022)]

Typical scenarios predict broad PBH mass spectrum that could explain:

- Dark matter
- OGLE ultrashort microlensing events
- LIGO-Virgo-Kagra BBH mergers

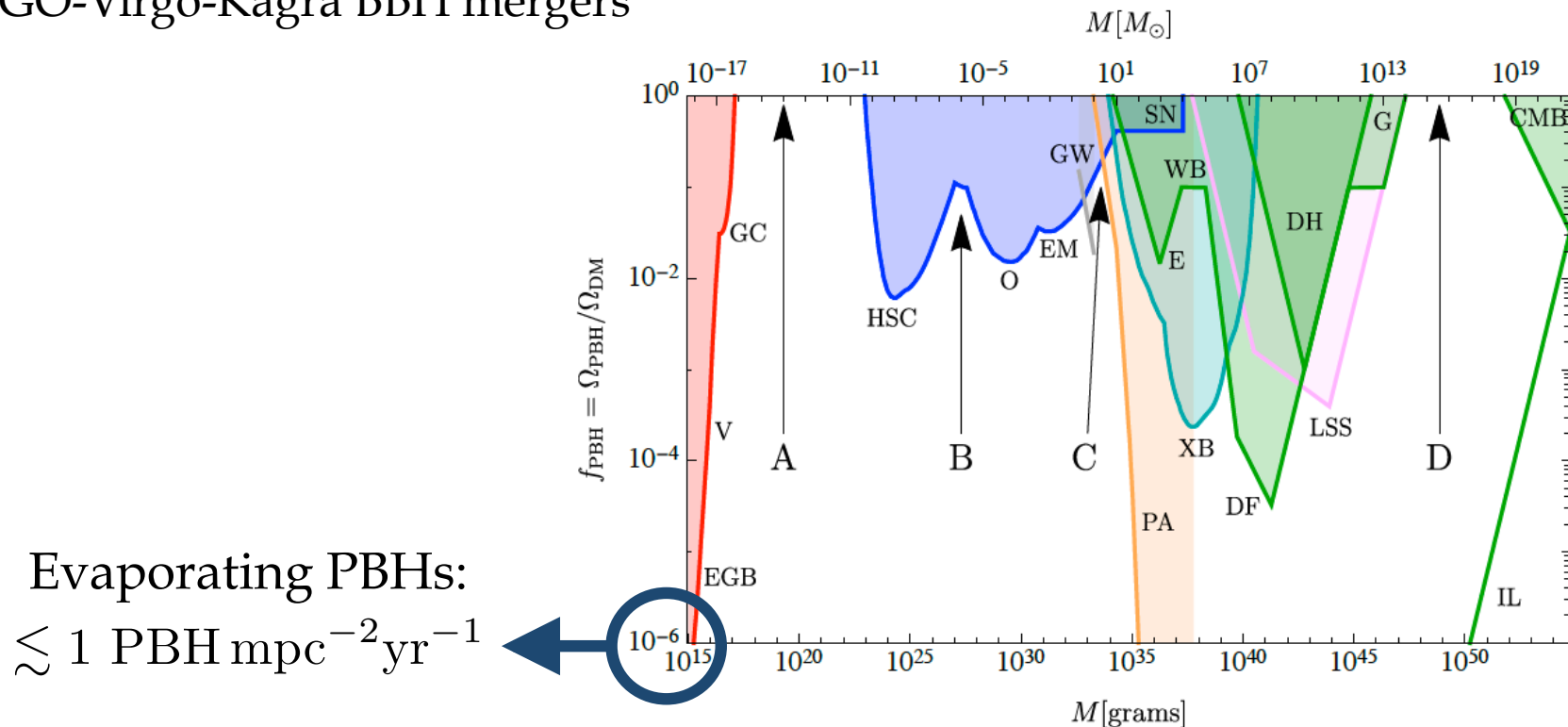


Primordial black holes

Several cosmological scenarios predict formation of primordial BHs (non-standard inflation, phase transitions, etc) [c.f. Escrivá & Kuhnel (2022)]

Typical scenarios predict broad PBH mass spectrum that could explain:

- Dark matter
- OGLE ultrashort microlensing events
- LIGO-Virgo-Kagra BBH mergers



Evaporating primordial black holes

PBHs born with masses $M_i \sim 10^{12}$ kg evaporate through Hawking emission completely in 14 Gyrs, so we could hope to observe their final “explosion”

$$T_H = \frac{M_P^2}{8\pi M} \simeq \left(\frac{10^7 \text{ kg}}{M} \right) \text{ TeV}$$

PBHs could emit new particle species in their final seconds, and this speeds up evaporation [Baker & Thamm (2022-23)]

Standard lore:

PBHs quickly lose their spin as they evaporate

 not in string theory!

Lots of new information about new physics from PBH spin

String axiverse

Dimensional reduction of NS & RR p-form fields leads to 4D axions:
[Arvanitaki et al. (2010)]

$$B_{\mu\nu} = \frac{1}{2\pi} \sum_i a_i(x) \omega_{\mu\nu}^i(y) \quad \int_{C_j} \omega^i = \delta_j^i$$

- # axions = # compact cycles in the 6 extra-dimensions
- perturbative shift symmetries inherited from p-form gauge fields
- exponentially small masses from non-perturbative effects (assuming one QCD axion)

Expect $O(10^2 - 10^3)$ light axions in realistic string compactifications!

Evaporation of Kerr black holes

Mass and angular momentum loss functions:

$$\mathcal{F}(\tilde{a}) = - \left(\frac{M^3}{M_P^4} \right) \frac{\dot{M}}{M}, \quad \mathcal{G}(\tilde{a}) = - \left(\frac{M^3}{M_P^4} \right) \frac{\dot{J}}{J}$$

$$\begin{pmatrix} \mathcal{F} \\ \mathcal{G} \end{pmatrix} = \sum_{i,l,m} \frac{1}{2\pi} \int_0^\infty d\bar{\omega} \frac{\Gamma_{lm}^s}{e^{(\omega - m\Omega_H)/T_H} \pm 1} \begin{pmatrix} \bar{\omega} \\ m\tilde{a}^{-1} \end{pmatrix}$$

Regge trajectory:

$$\frac{d \log \tilde{a}}{d \log M} = \frac{\mathcal{G}}{\mathcal{F}} - 2 \equiv \mathcal{H}(\tilde{a}) \quad J = \tilde{a} \frac{M^2}{M_P^2}$$

Approximation:

only include each particle for $T_H > m$ and neglect their mass

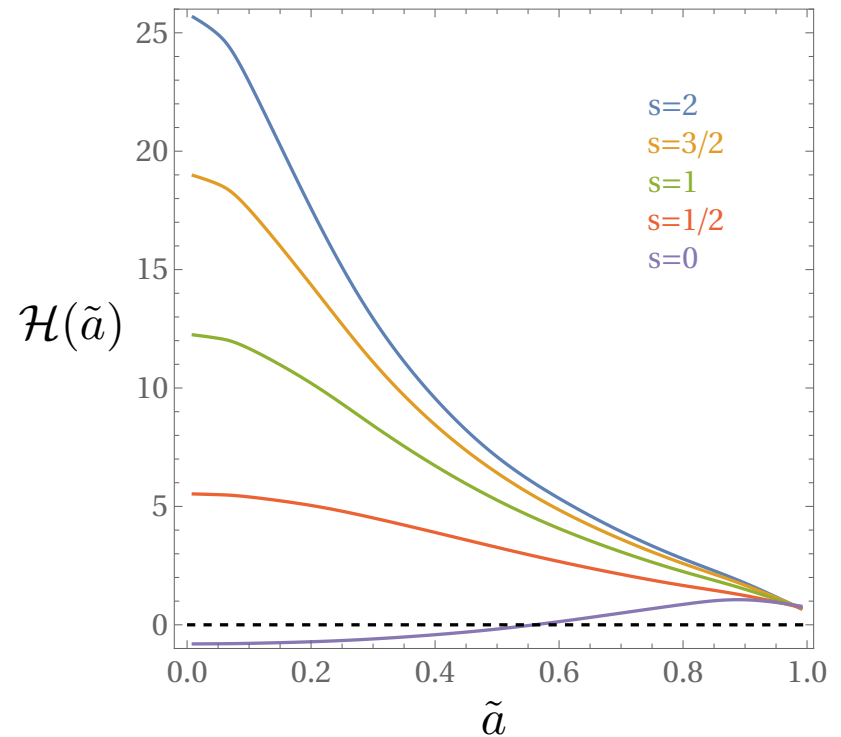
Evaporation of Kerr black holes

Scalar emission can spin up a PBH for $\tilde{a} < 0.555$, since $l=m=0$ mode is dominant.
[Chambers, Hiscock & Taylor (1997)]

In the SM the only scalars are pions and Higgs doublet...

... but in string theory there should $O(10^2-10^3)$ light axions!
[Arvanitaki et al. (2010)]

$$\frac{d \log \tilde{a}}{d \log M} = \mathcal{H}(\tilde{a})$$

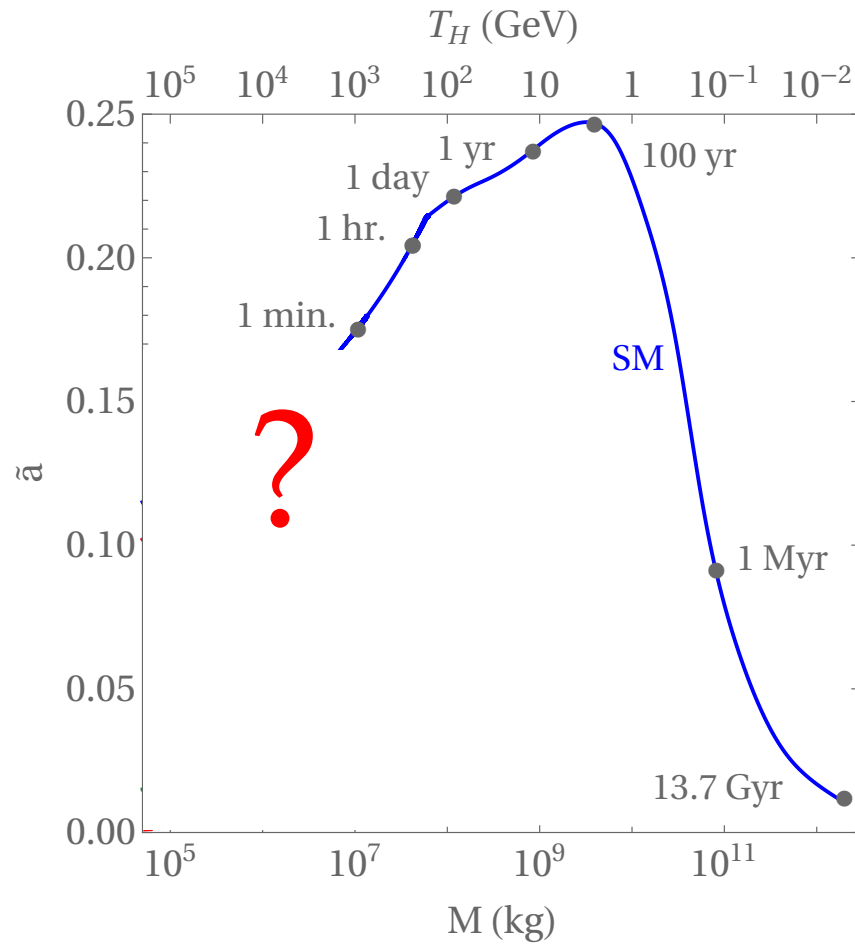


[c.f. Page (1976)]

PBHs can spin up in the string axiverse!

[Calzà, March-Russell & JGR, 2110.13602]

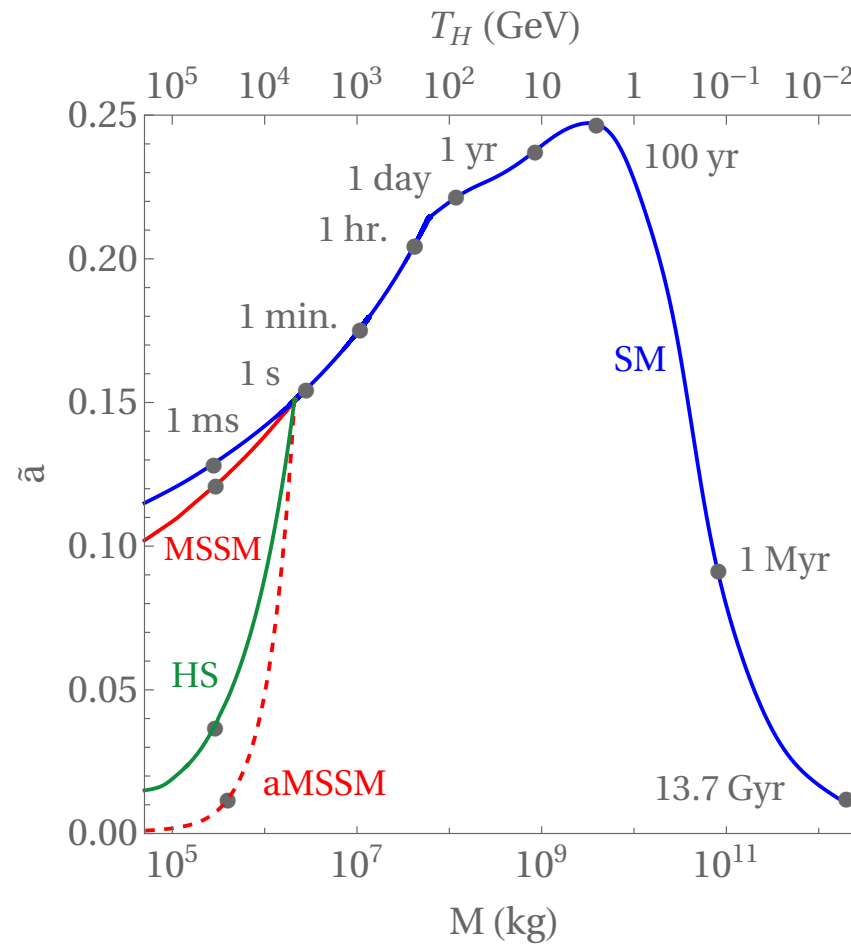
Kerr PBHs in the string axiverse



SM = Standard Model + graviton + 400 axions

$$(n_0, n_{1/2}, n_1, n_{3/2}) = (400+4, 45, 12, 0)$$

Adding new physics @ 5 TeV



MSSM: $(n_0, n_{1/2}, n_1, n_{3/2}) = (400+4+4+90, 45+12+4, 12, 1)$ (+128 d.o.f.)

aMSSM: $(n_0, n_{1/2}, n_1, n_{3/2}) = (400+4+4+90+400, 45+12+8+400, 12, 1)$ (+1336 d.o.f.)

HS: $(n_0, n_{1/2}, n_1, n_{3/2}) = (400+4+4, 45+45, 12+12, 0)$ (+118 d.o.f.)

Tracking the PBH evolution

If we detect a PBH and track $M(t)$ and $\tilde{a}(t)$, we can determine the loss rates:

$$\begin{pmatrix} \mathcal{F} \\ \mathcal{F}' \\ \mathcal{G} \\ \mathcal{G}' \end{pmatrix} = \begin{pmatrix} \mathcal{F}_0 & \mathcal{F}_{1/2} & \mathcal{F}_1 & \mathcal{F}_{3/2} \\ \mathcal{F}'_0 & \mathcal{F}'_{1/2} & \mathcal{F}'_1 & \mathcal{F}'_{3/2} \\ \mathcal{G}_0 & \mathcal{G}_{1/2} & \mathcal{G}_1 & \mathcal{G}_{3/2} \\ \mathcal{G}'_0 & \mathcal{G}'_{1/2} & \mathcal{G}'_1 & \mathcal{G}'_{3/2} \end{pmatrix} \begin{pmatrix} n_0 \\ n_{1/2} \\ n_1 \\ n_{3/2} \end{pmatrix}$$

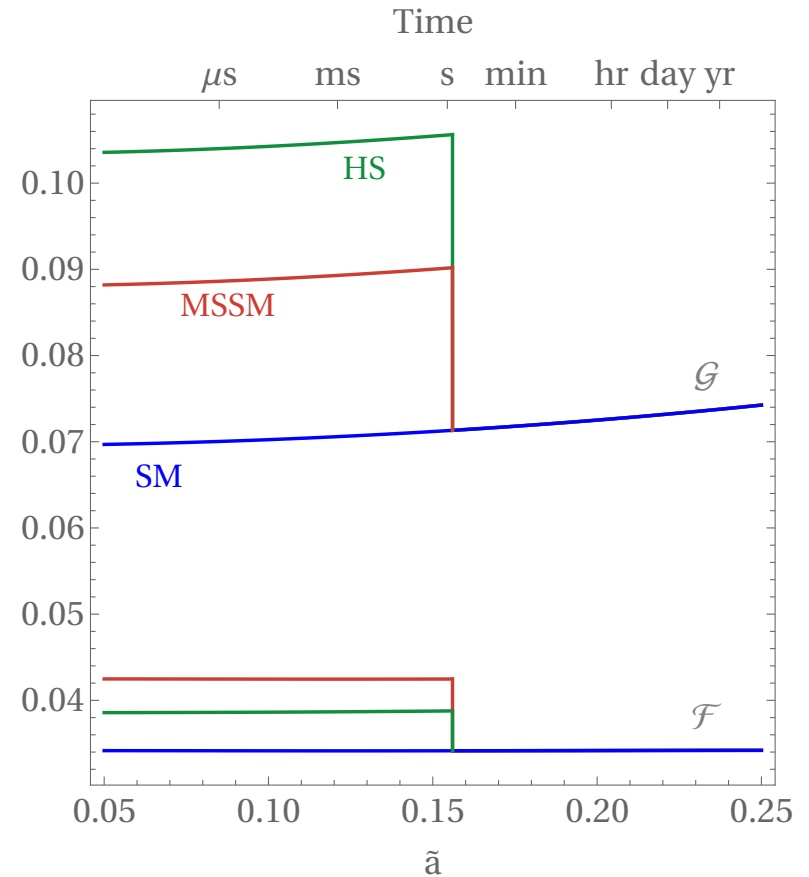
for each value of \tilde{a} .

We could fully determine the particle content!

Potential problem:

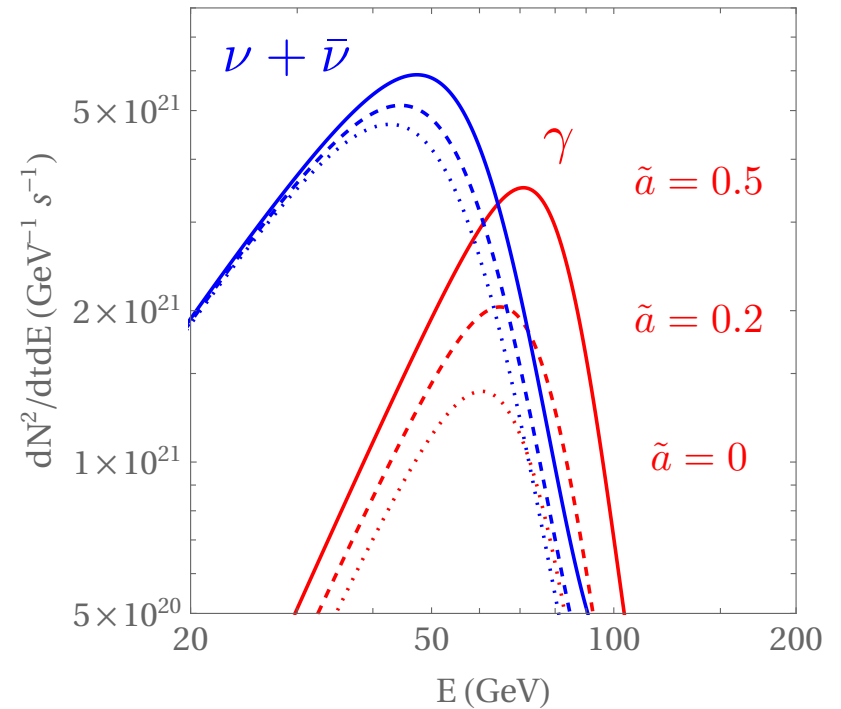
hard to measure \mathcal{F}' due to cancellation between scalar (<0) and non-scalar (>0) contributions.

Solution: measure \mathcal{G}'' !



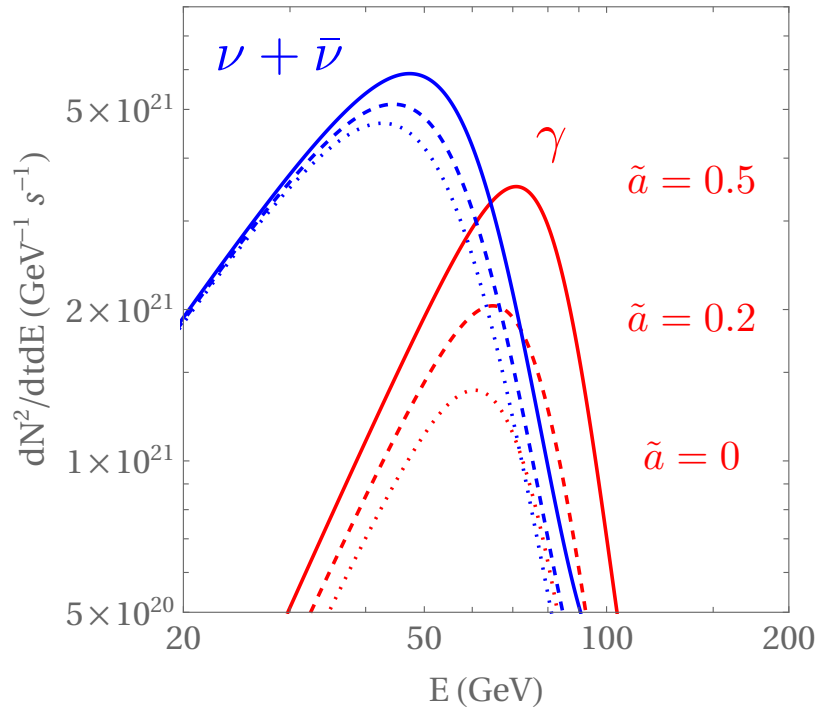
Measuring the PBH mass and spin from Hawking spectrum

- Primary photon peak (known distance)
[Calzà, March-Russell & JGR, 2110.13602]
- “Peak-valley” of primary + secondary photon spectra [Calzà & JGR, 2210.06500]
- Multipolar peaks in photon spectrum ($\tilde{a} > 0.6$) [Calzà & JGR, 2312.11930]
- Multi-messenger approach:
photon + neutrino primary emission



$$\frac{d^3 N_s}{dE dt d\Omega} = \frac{1}{4\pi} \sum_{l,m} \frac{\Gamma_{l,m}^s(\omega)}{e^{(\omega - m\Omega_H)/T_H} \pm 1} (|{}_s S_{lm}(\theta)|^2 + |{}_{-s} S_{lm}(\theta)|^2)$$

Measuring the PBH mass and spin from Hawking spectrum



$$E_{\gamma, \nu} \propto T_H \propto M^{-1}$$

$$\frac{E_{\nu}}{E_{\gamma}} = 0.705 - \frac{0.559\tilde{a}^2}{1 + 5.18\tilde{a}}$$

$$\frac{I_{\nu}}{I_{\gamma}} = 3.423 - \frac{31.05\tilde{a}^2}{1 + 7.05\tilde{a}}$$

Problem:

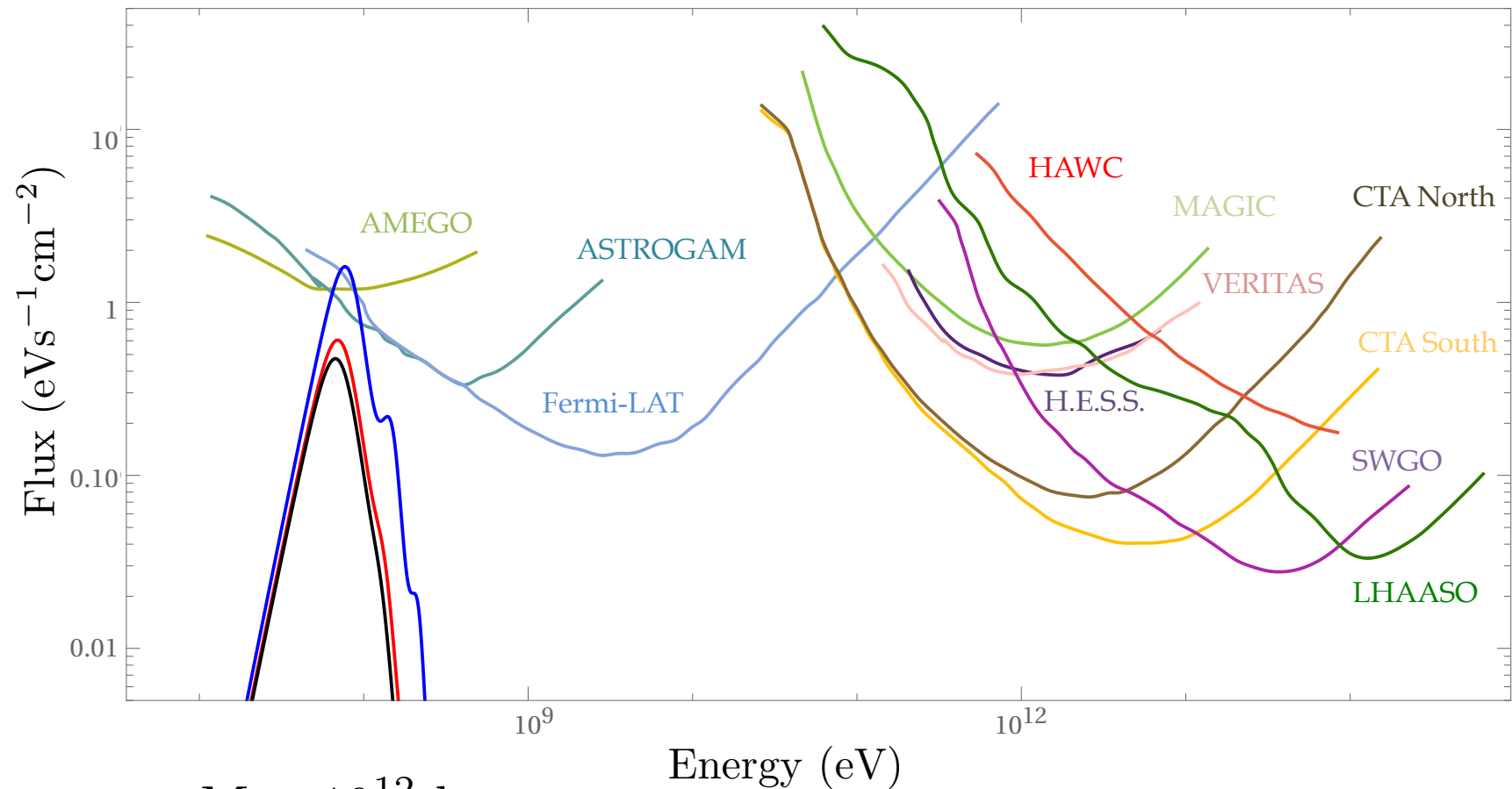
Emission is anisotropic and depends on unknown inclination of PBH axis

Solutions:

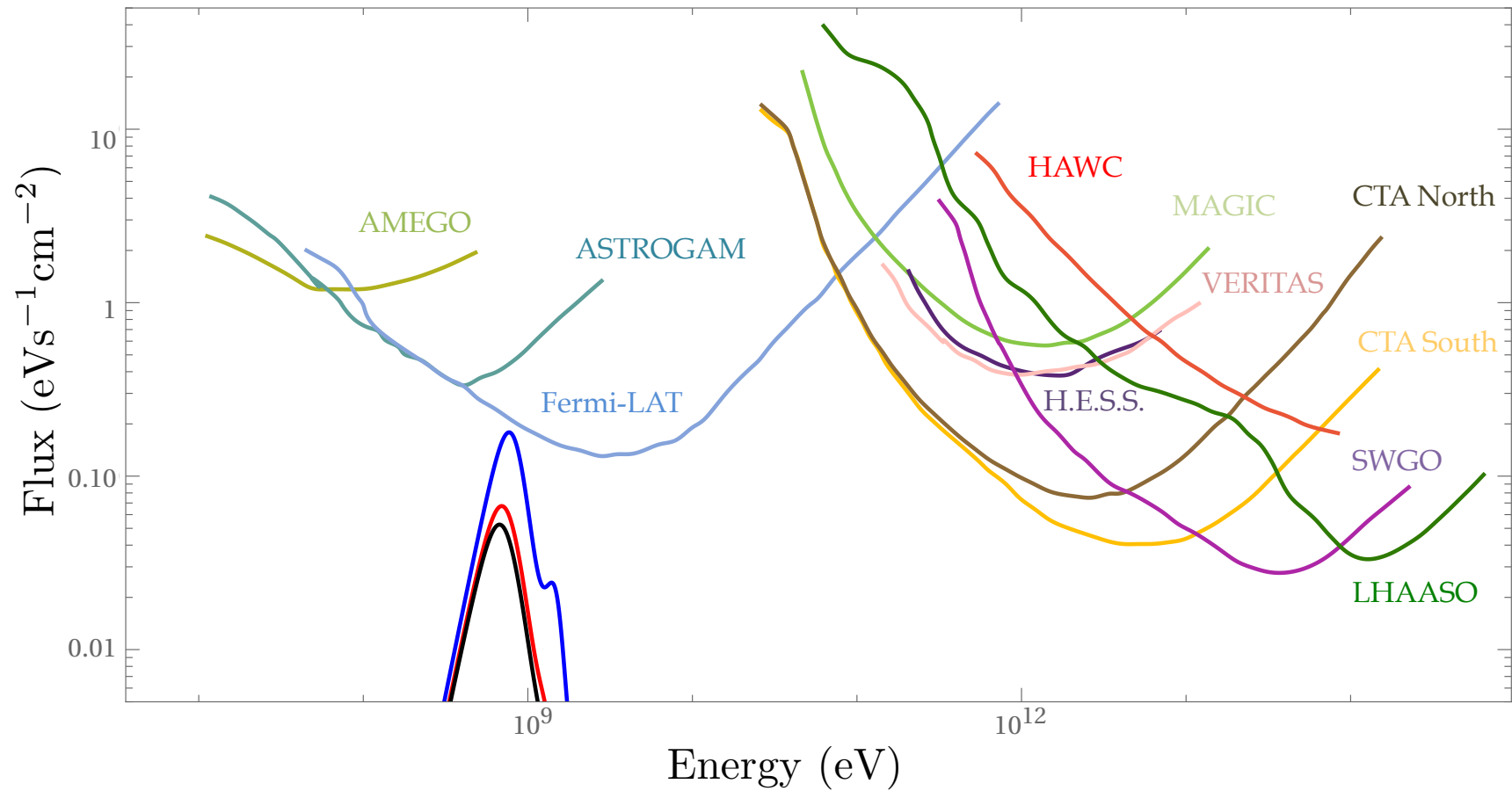
- (i) Measure $\nu - \bar{\nu}$ (& photon polarization) asymmetry [Perez-Gonzalez (2023)]
- (ii) Measure intensity modulation from PBH proper motion:

$$v_{PBH} \sim 50 \text{ AU/yr}$$

Detecting Kerr PBHs (primary γ)



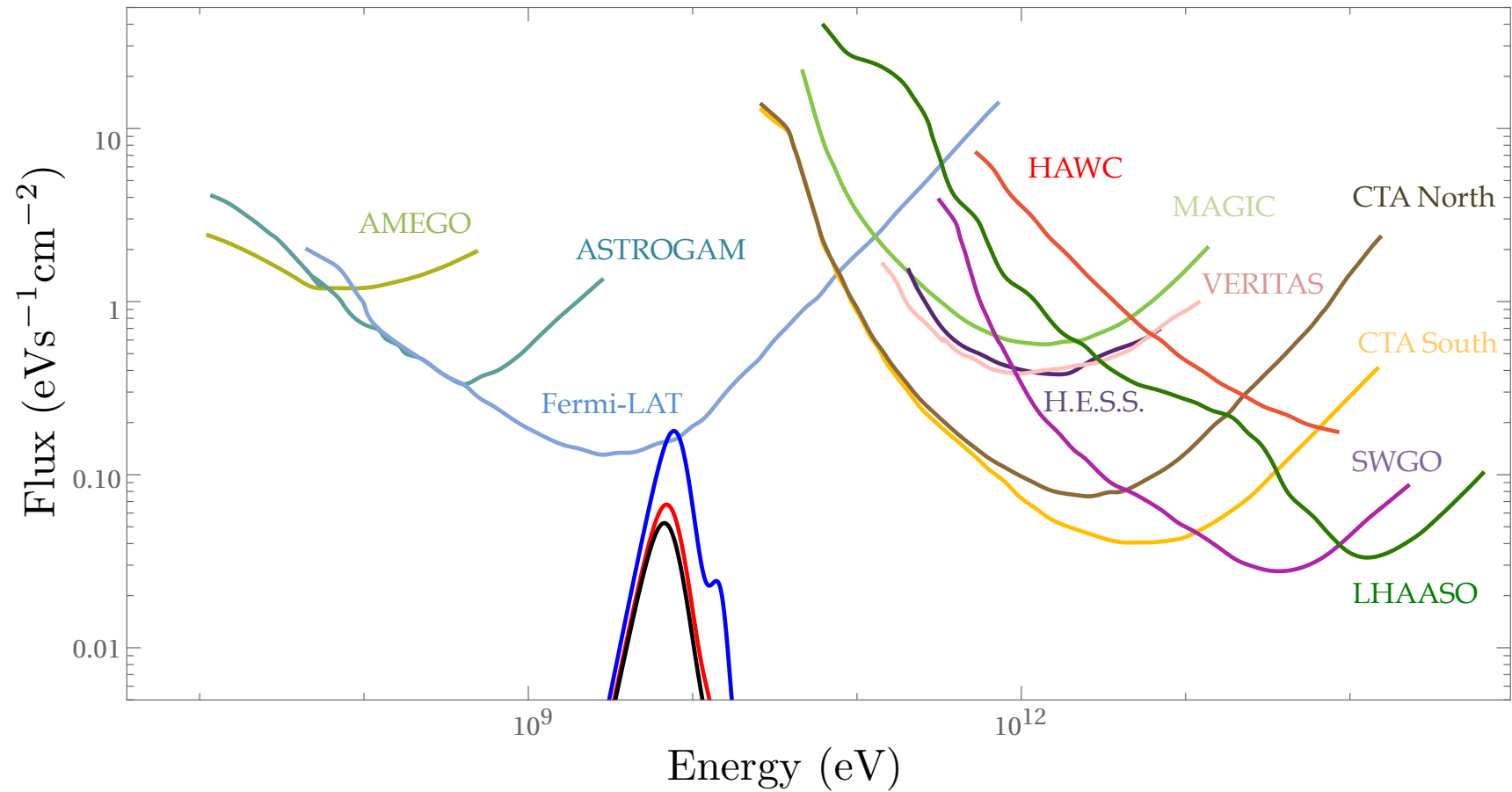
Detecting Kerr PBHs (primary γ)



$$M = 10^{11} \text{ kg}$$

$$d = 3 \times 10^{-4} \text{ pc}$$

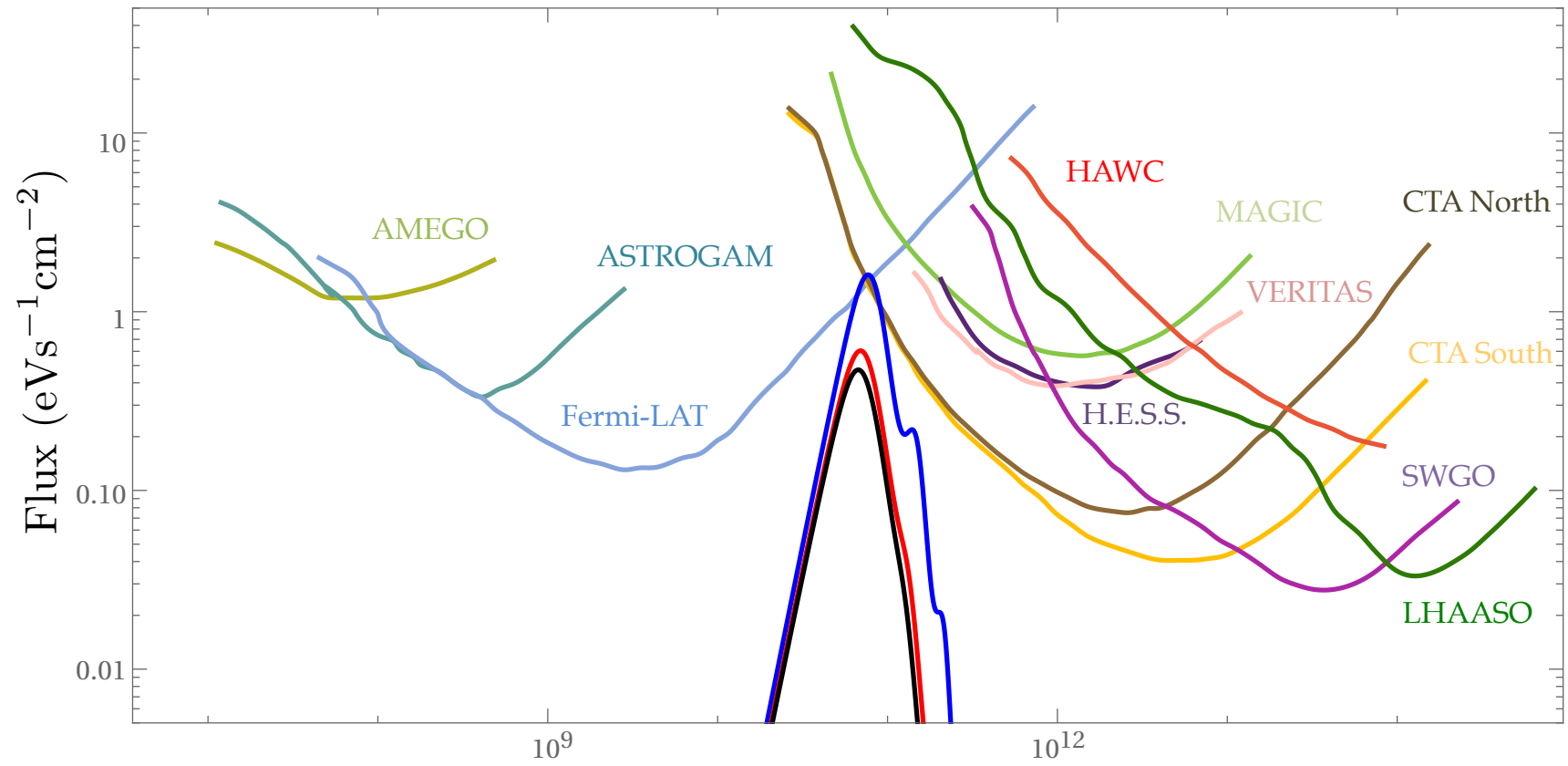
Detecting Kerr PBHs (primary γ)



$$M = 10^{10} \text{ kg}$$

$$d = 3 \times 10^{-3} \text{ pc}$$

Detecting Kerr PBHs (primary γ)

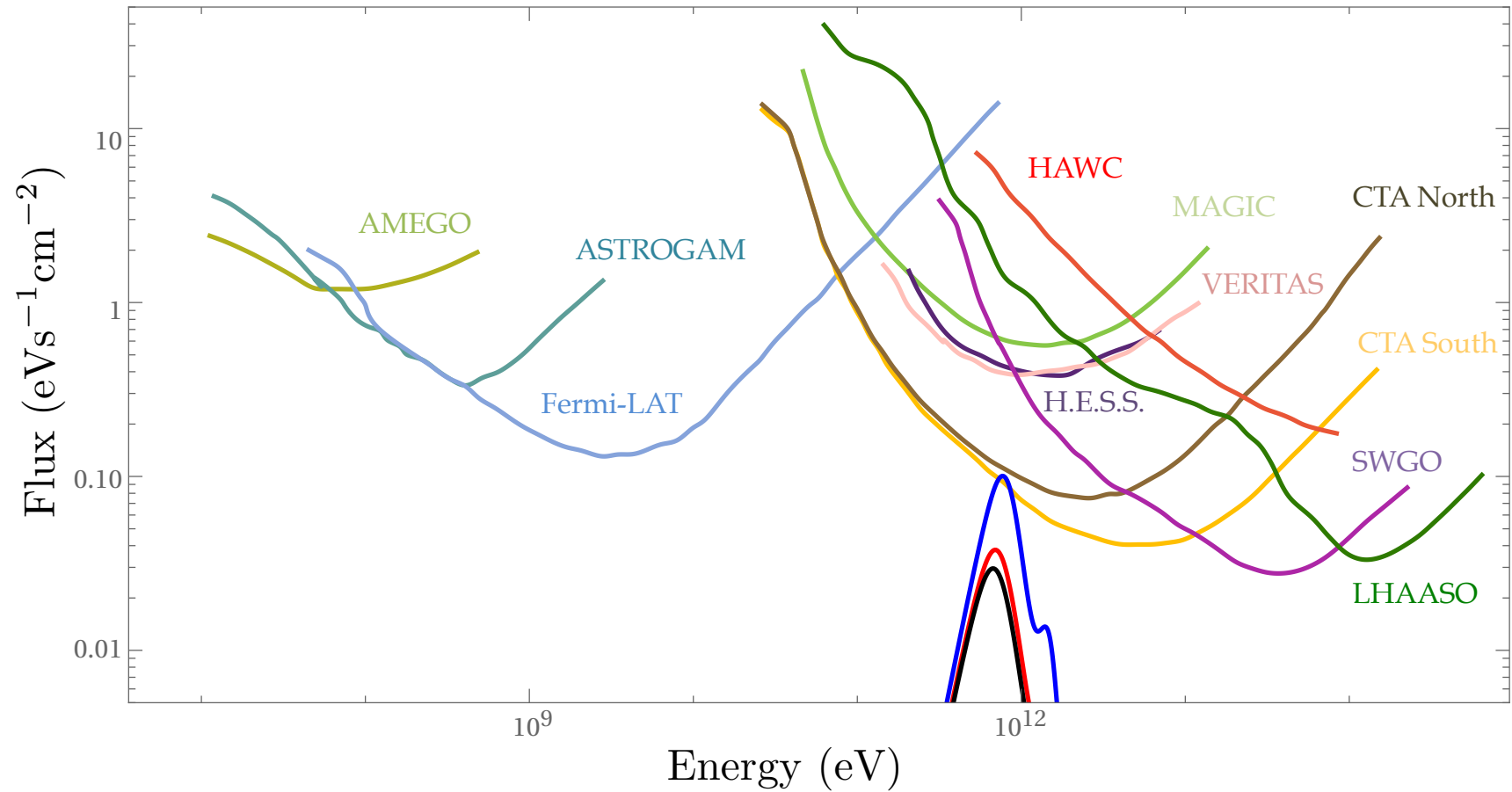


Energy (eV)

$$M = 10^9 \text{ kg}$$

$$d = 10^{-2} \text{ pc}$$

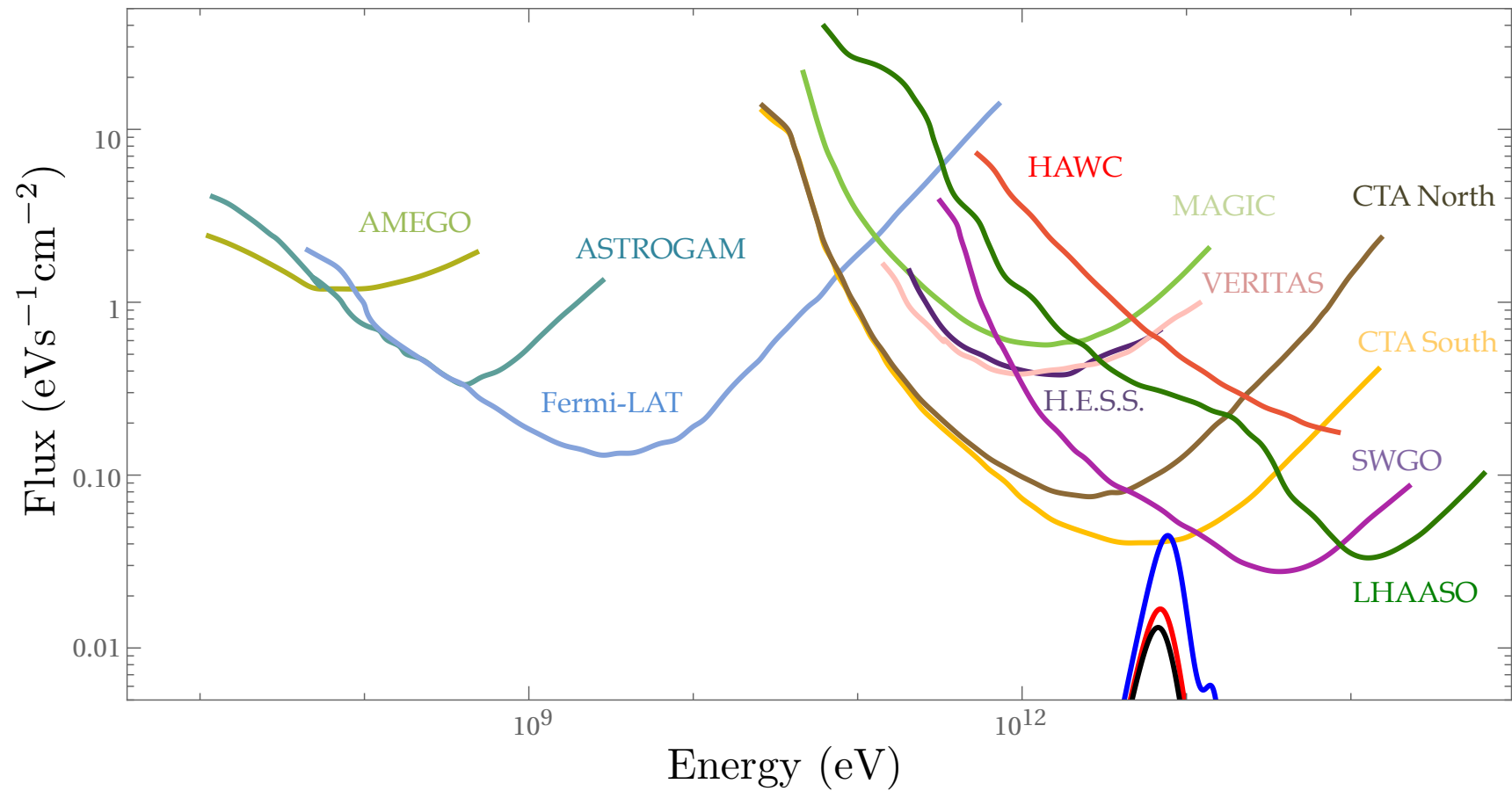
Detecting Kerr PBHs (primary γ)



$$M = 10^8 \text{ kg}$$

$$d = 0.4 \text{ pc}$$

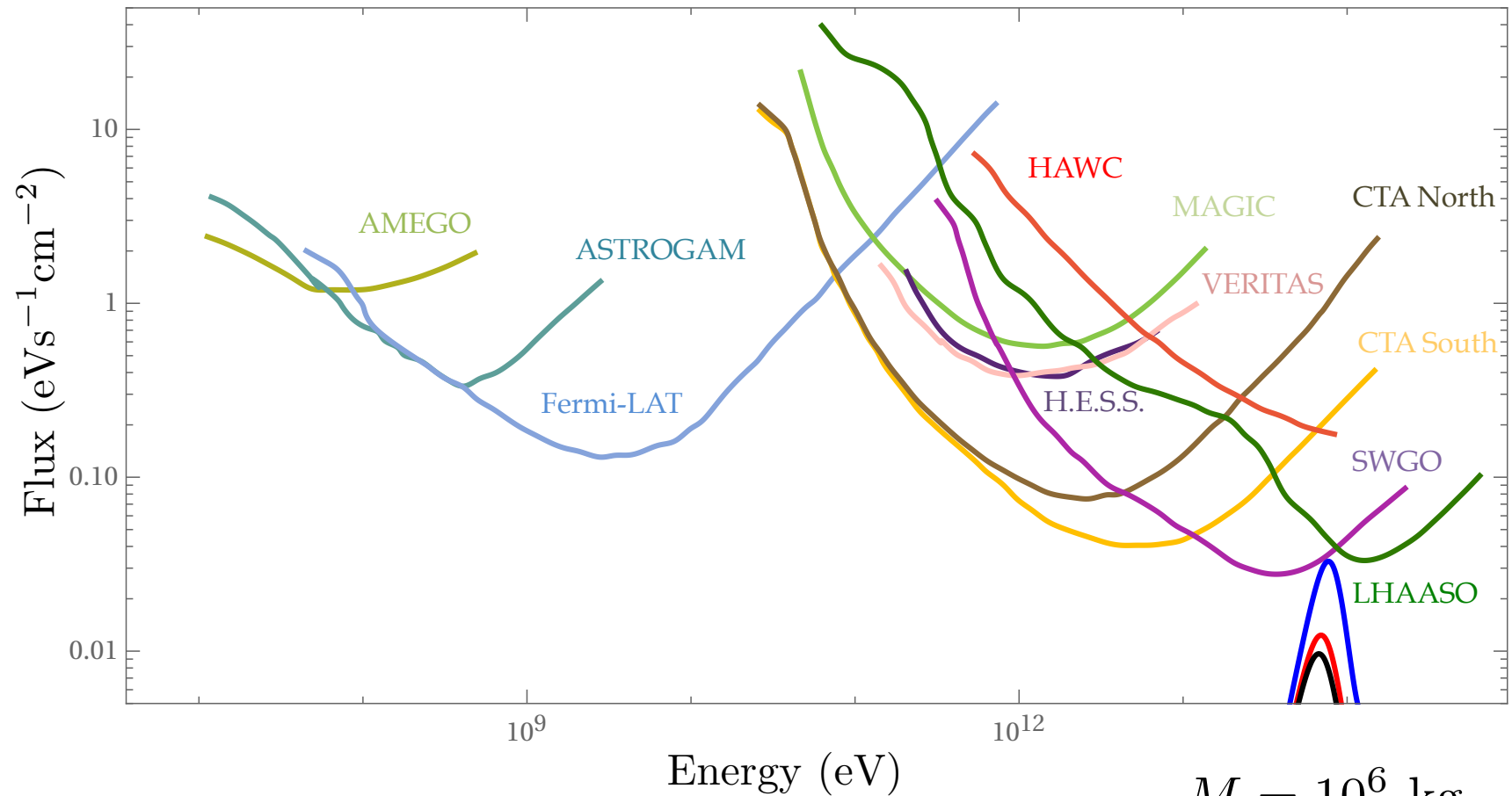
Detecting Kerr PBHs (primary γ)



$$M = 10^7 \text{ kg}$$

$$d = 6 \text{ pc}$$

Detecting Kerr PBHs (primary γ)



$$M = 10^6 \text{ kg}$$
$$d = 70 \text{ pc}$$

Take home messages

- Evaporating PBHs are fantastic particle physics laboratories
- Finding a single PBH with non-negligible spin in its last stages would be evidence for a [string axiverse](#)
- Tracking the PBH mass and spin in its last second can reveal new physics and [distinguish SM extensions \(>TeV scale\)](#)
- [Multi-messenger astronomy](#) to detect and track PBH Hawking emission:

γ -rays: H.E.S.S., Milagro, VERITAS, HAWC, Fermi-LAT, LHAASO
CTA, SWGO, AMEGO, ASTROGAM

neutrinos: IceCube,
KM3Net, P-ONE, Trident, Baikal-GVD

Further details: [arXiv:2312.09261 \[hep-ph\]](#)