Introduction to QED & QCD Tutorial Questions

1 Tuesday Tutorial

1. Suppose we have a plane-wave solution to the Klein-Gordon equation of the form

$$
\phi(\boldsymbol{x},t) = A e^{-i(\omega t - \boldsymbol{k}\cdot\boldsymbol{x})}.
$$

Use the Klein-Gordon equation to find the *dispersion relation*, i.e. find ω in terms of **k**. How do you interpret the two solutions?

Show that these solutions are eigenstates of the energy operator, $i\partial_t$, and the 3-momentum operator, −*i*∇.

2. Show that the Dirac γ -matrices defined in the lectures:

$$
\gamma^0 = \beta, \qquad \gamma^k = \beta \, \alpha^k,
$$

obey the hermiticity relation

$$
(\gamma^\mu)^\dagger = \gamma^0 \gamma^\mu \gamma^0.
$$

3. When evaluating cross sections, you will frequently need to manipulate Dirac matrices. Using the anti-commutation relations for the γ -matrices, show that in 4 dimensions:

(i)
$$
\gamma^{\mu}\gamma_{\mu} = 4
$$
,
\n(ii) $\gamma^{\mu}\gamma^{\nu}\gamma_{\mu} = -2\gamma^{\nu}$,
\n(iii) $\gamma^{\mu}\gamma^{\nu}\gamma^{\lambda}\gamma_{\mu} = 4g^{\nu\lambda}$,
\n(iv) $\gamma^{\mu}\gamma^{\nu}\gamma^{\lambda}\gamma^{\rho}\gamma_{\mu} = -2\gamma^{\rho}\gamma^{\lambda}\gamma^{\nu}$.

How do these change in arbitrary dimensions where $g^{\mu\nu}g_{\mu\nu} = \delta^{\mu}_{\mu} = d$?

4. Verify the orthonormality and completeness relations for the solutions of the Dirac equation:

$$
\overline{u}_r(\boldsymbol{p})u_s(\boldsymbol{p})=-\overline{v}_r(\boldsymbol{p})v_s(\boldsymbol{p})=2m\,\delta^{rs},\quad \overline{u}_r(\boldsymbol{p})v_s(\boldsymbol{p})=\overline{v}_r(\boldsymbol{p})u_s(\boldsymbol{p})=0,
$$

and

$$
\sum_{r=1}^2 u_r(\mathbf{p}) \overline{u}_r(\mathbf{p}) = (\not\!{p} + m), \qquad \sum_{r=1}^2 v_r(\mathbf{p}) \overline{v}_r(\mathbf{p}) = (\not\!{p} - m).
$$

5. Show that the Dirac hamiltonian, $H = \alpha \cdot p + \beta m$, commutes with the total angular momentum operator

$$
[\boldsymbol{L} + \boldsymbol{S}, H] = 0\,,
$$

where $\mathbf{L} = \mathbf{x} \times \mathbf{p}$ is the orbital angular momentum and \mathbf{S} is the spin operator

$$
\boldsymbol{S} = \frac{1}{2} \left(\begin{array}{cc} \boldsymbol{\sigma} & 0 \\ 0 & \boldsymbol{\sigma} \end{array} \right) .
$$

6. Using the plane-wave solutions of the Dirac equation given in the lectures, show that for $p = (0, 0, p_z)$

$$
S_z u_1 = \frac{1}{2} u_1
$$
, $S_z u_2 = -\frac{1}{2} u_2$, $S_z v_1 = \frac{1}{2} v_1$ and $S_z v_2 = -\frac{1}{2} v_2$,

where S_z is the *z*-component of the spin operator.

2 Wednesday Tutorial

Note: if pressed for time ignore the "arbitrary gauge" part in Q7. Q8 is partly a repetition of what was already presented in the lectures.

- 7. Draw all the tree-level diagrams for Bhabha-scattering, $e^+(p) e^-(k) \to e^+(p') e^-(k')$ and give the expression for the scattering amplitude, i *M*, in Feynman gauge. What happens in an arbitrary gauge?
- 8. (a) Show that the process $e^+(k')e^-(k) \to \mu^+(p')\mu^-(p)$, in the limit $m_e \to 0$, has a matrixelement-squared given by

$$
\overline{|\mathcal{M}|}^2 = \frac{1}{4} \frac{e^4}{(k+k')^4} \text{Tr} [k' \gamma^{\mu} k \gamma^{\nu}] \text{Tr} [(\not p + M) \gamma_{\mu} (\not p' - M) \gamma_{\nu}],
$$

when summed and averaged over final and initial spins, where *M* is the mass of the muon.

(b) Show that

$$
\left(\frac{d\sigma}{d\Omega}\right)_{e^+e^-\to\mu^+\mu^-} = \frac{\overline{|\mathcal{M}|}^2}{64\pi^2 s} \sqrt{1 - \frac{4M^2}{s}}\,,
$$

where $s = (k + k')^2$.

(c) The traces evaluate to (check if you have time!)

$$
\overline{|\mathcal{M}|}^2 = \frac{8e^4}{s^2} \left[(pk)^2 + (pk')^2 + M^2(kk') \right] .
$$

Move to the centre-of-mass frame and let the scattering angle be θ . Show that

$$
\left(\frac{d\sigma}{d\Omega}\right)_{e^+e^-\to\mu^+\mu^-}=\frac{e^4}{64\pi^2s}\sqrt{1-\frac{4M^2}{s}}\left[1+\left(1-\frac{4M^2}{s}\right)\cos^2\theta+\frac{4M^2}{s}\right]
$$

.

(d) Find an expression for the total cross section in the high-energy limit where the mass of the muon can be neglected.

3 Thursday Tutorial

Note: if pressed for time focus on Q9, Q10, Q12 a).

- 9. Write the amplitude for Compton scattering $e(p)\gamma(k) \rightarrow e(p')\gamma(k')$ in the form $i\mathcal{M} =$ $M_{\mu\nu}\varepsilon^{*\mu}(k')\varepsilon^{\nu}(k)$. Verify that this is gauge-invariant.
- 10. In the lectures, we found the matrix element squared for unpolarised Compton scattering was

$$
\overline{\left|\mathcal{M}\right|}^2 = 2e^4\left(\frac{pk}{pk'} + \frac{pk'}{pk} + 2m^2\left(\frac{1}{pk} - \frac{1}{pk'}\right) + m^4\left(\frac{1}{pk} - \frac{1}{pk'}\right)^2\right).
$$

Working in the centre-of-mass system, in the limit where the electron mass *m* can be neglected, show that the matrix element squared is dominated by backward scattering, $\theta \simeq \pi$, where θ is the scattering angle of the photon.

11. Use the matrix-element squared for Compton scattering to obtain the matrix-element squared for the annihilation process $e^+e^- \rightarrow \gamma\gamma$. Again work in the centre-of-mass frame and show that, in the high-energy limit $E \gg m$,

$$
\overline{|\mathcal{M}|}^2 \simeq 4e^4 \frac{1 + \cos^2 \theta}{\sin^2 \theta}.
$$

Figure 1: One-loop corrections to a $q\bar{q}g$ -vertex.

12. Consider the diagrams in figure 1. Show that the colour factors are given by

(a)
$$
t^c t^a t^b \delta^{bc} = -\frac{1}{2N_c} t^a
$$
, and (b) $i f^{abc} t^b t^c = -\frac{1}{2} C_A t^a$

respectively.

- 13. Calculate the summed and averaged matrix-element squared, $|\overline{\mathcal{M}}|^2$, for the quark-scattering process $ud \rightarrow ud$.
- 14. Solve the one-loop β -functions for QCD and QED:

$$
\mu^2 \frac{d\alpha_s}{d\mu^2} = -\frac{11C_A - 2n_f}{12\pi} \alpha_s^2, \quad \text{and} \quad \mu^2 \frac{d\alpha}{d\mu^2} = \frac{1}{3\pi} \alpha^2,
$$

using as initial condition the value of the couplings at the *Z* mass. Sketch the solutions as a function of μ^2 .