

Neutrino Physics

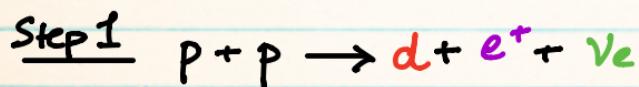
Neutrino Oscillations in matter

Jessica Turner

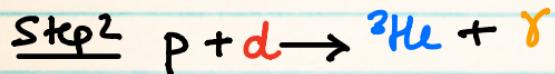
Solar neutrinos - PP chain

- The Sun shines by making hydrogen into helium. There are two main ways of doing this: pp chain or CNO cycle

pp-chain I



Sun most loses its energy (98.5 %) through pp-chain. This is pp-I chain and occurs around 85 % of the time



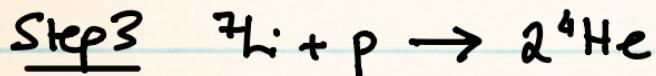
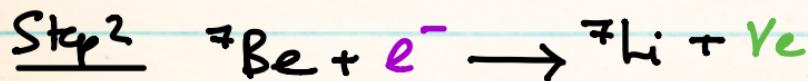
Step 3 repeat steps 1 & 2



Solar neutrinos - PP chain

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pp-chain II



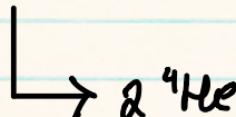
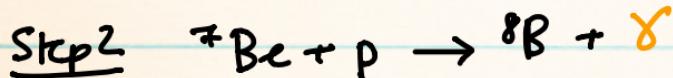
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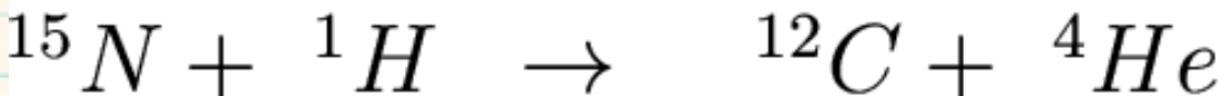
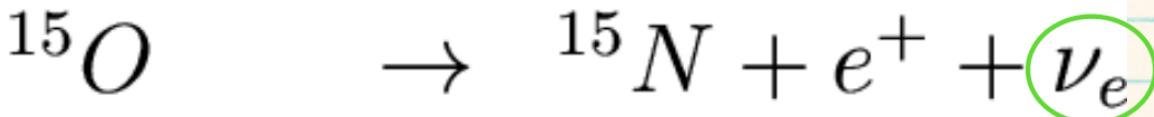
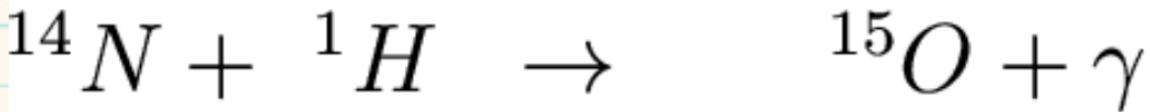
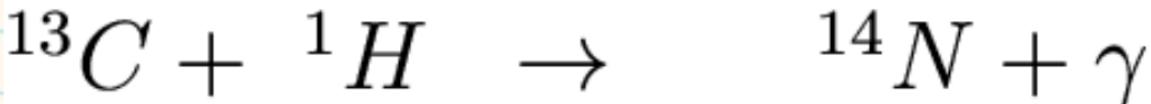
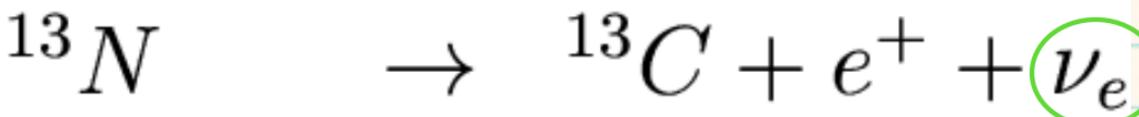
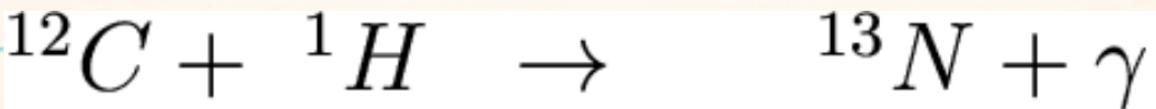
pp-chain III



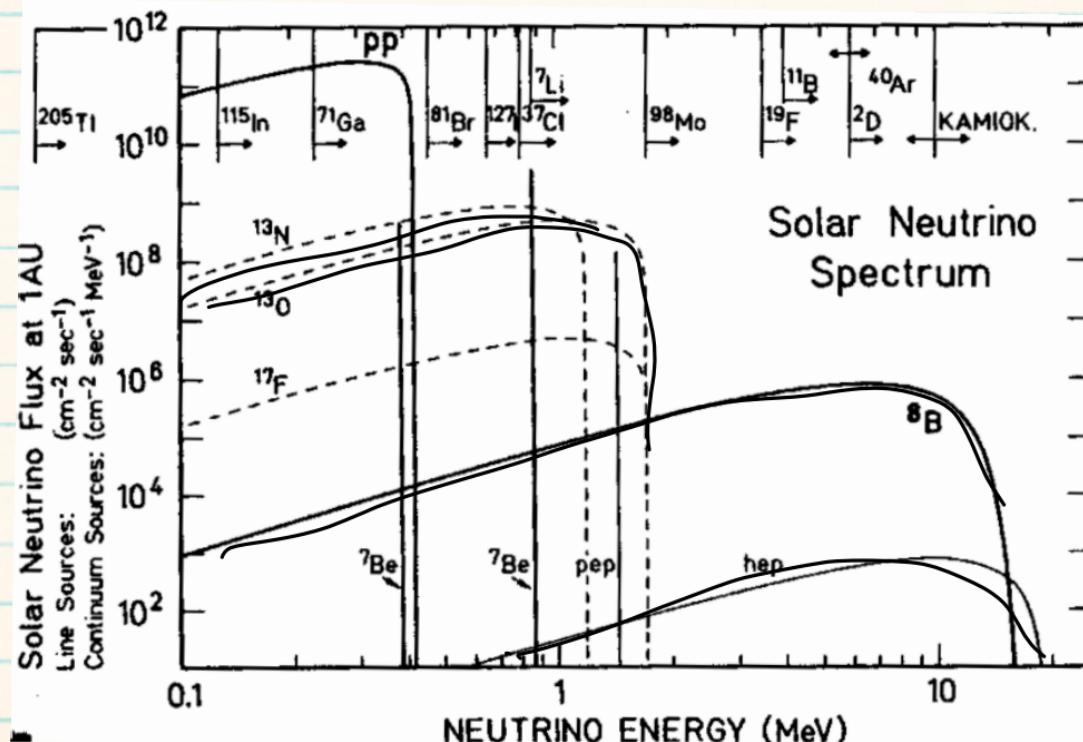
pp-III chain occurs around 0.3 % of the time



Solar neutrinos - CNO cycle



- Many nuclear processes (pp chain and CNO cycle) produce electron neutrinos
- Energies of the neutrinos will differ, depending on the reaction.



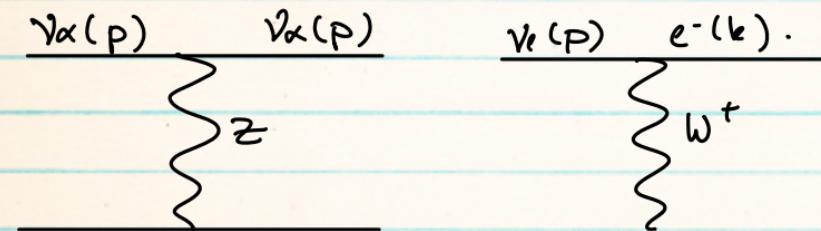
Solar neutrinos

- 1964 Homestake experiment (headed by Davis & Bahcall) detected solar neutrinos but there were approximately 2/3 less than expected from Bahcall's Standard Solar Model prediction.
- It was initially proposed that the solar models were wrong.
- Or that two experiments Homestake and GALLEX were wrong!
- As you can guess, the resolution to this problem is neither!

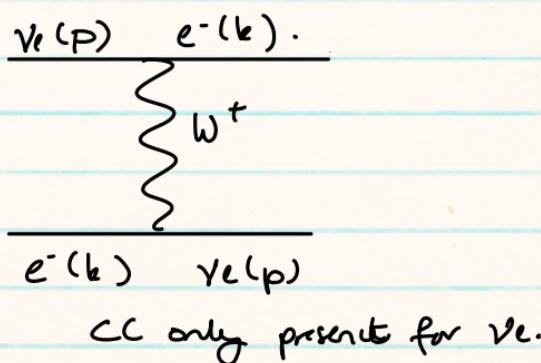
Neutrino Oscillations in matter

ν oscillation in matter is affected by the background.

For simplified two neutrino case ; ν_e & ν_μ behave differently in matter.



$\alpha = e, \mu, \tau$
NC same for all flavours



- ν oscillations are affected by **COHERENT SCATTERING** with background (no momentum exchange between ν & background).
- coherent scattering is elastic scattering at $\theta=0^\circ$

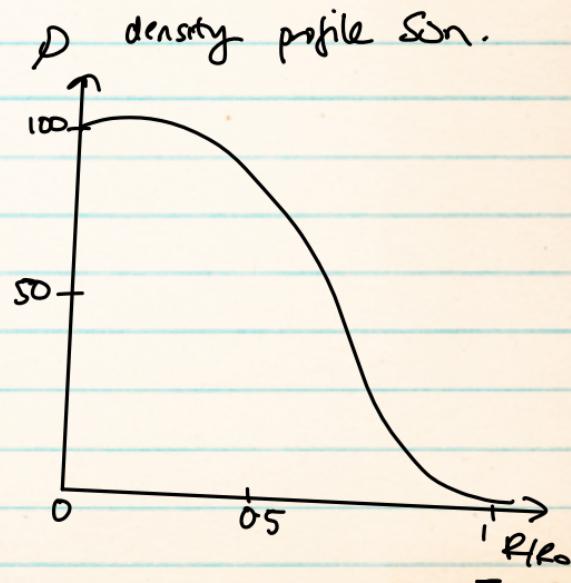
Neutrino Oscillations in matter

The background creates a **MATTER POTENTIAL** for the ν .

$$V_{NC} = \frac{GF N_n}{\sqrt{2}} \quad N_n \equiv \# \text{ density of nucleons}$$

$$V_{CC} = \sqrt{2} GF N_e. \quad N_e = \text{number density of electrons.}$$

Medium	density (g/cm^3)	$V_{CC} (\text{eV})$
solar	~ 100	$\sim 10^{-12}$
Earth's core	~ 10	$\sim 10^{-13}$
supernovae	$\sim 10^{14}$	~ 1



Neutrino Oscillations in matter

$$i \frac{\partial}{\partial t} |14\rangle = H |14\rangle$$

two neutrino case
 $|14\rangle = \begin{pmatrix} \nu_e(t) \\ \nu_\mu(t) \end{pmatrix}$

$$|\nu_e\rangle = \cos |\nu_1\rangle + \sin |\nu_2\rangle \quad \theta = \theta_{\text{vac.}}$$

$$|\nu_\mu\rangle = -\sin |\nu_1\rangle + \cos |\nu_2\rangle$$

$$\text{ie } \nu_\alpha = \sum_i U_{\alpha i} \nu_i$$

Hamiltonian for ν-osc in vacuum:

$$\begin{aligned} \langle \nu_\alpha | H_v | \nu_\beta \rangle &= \left\langle \sum_i U_{\alpha i} \nu_i | H_v | \sum_j U_{\beta j} \nu_j \right\rangle \\ &= \sum_i \hat{U}_{\alpha i} \langle \nu_i | H_v | \nu_i \rangle U_{\beta i} = \sum_i U_{\alpha i}^* E_i U_{\beta i} \\ \alpha &= e, \mu. \end{aligned}$$

Neutrino Oscillations in matter

$$\langle \nu_\alpha | H_\nu | \nu_\beta \rangle = \begin{pmatrix} c_\theta & -s_\theta \\ s_\theta & c_\theta \end{pmatrix} \begin{pmatrix} E_1 & 0 \\ 0 & E_2 \end{pmatrix} \begin{pmatrix} c_\theta & s_\theta \\ -s_\theta & c_\theta \end{pmatrix}$$

ex.

$$= \begin{pmatrix} H_{\alpha\alpha} & H_{\alpha\beta} \\ H_{\beta\alpha} & H_{\beta\beta} \end{pmatrix}$$

let

$$|\vec{p}_1| = |\vec{p}_2| = |\vec{p}|.$$

$$\begin{aligned} H_{\alpha\alpha} &= E_1 c_\theta^2 + E_2 s_\theta^2 = \sqrt{|\vec{p}|^2 + m_1^2} c_\theta^2 + \sqrt{|\vec{p}|^2 + m_2^2} s_\theta^2 \\ &= |\vec{p}| \left(\left(1 + \frac{m_1^2}{|\vec{p}|^2} \right)^{\frac{1}{2}} c_\theta^2 + \left(1 + \frac{m_2^2}{|\vec{p}|^2} \right)^{\frac{1}{2}} s_\theta^2 \right) \\ &= |\vec{p}| \left(c_\theta^2 + \frac{m_1^2}{2|\vec{p}|^2} c_\theta^2 + s_\theta^2 + \frac{m_2^2}{2|\vec{p}|^2} s_\theta^2 \right) \end{aligned}$$

Neutrino Oscillations in matter

$$= |\vec{p}| \left(1 + \frac{m_1^2 c\theta^2}{2|\vec{p}|^2} + \frac{m_2^2 s\theta^2}{2|\vec{p}|^2} \right)$$

$$= |\vec{p}| + \frac{m_1^2}{2|\vec{p}|} \left(\frac{1 + \cos 2\theta}{2} \right) + \frac{m_2^2}{2|\vec{p}|} \left(\frac{1 - \cos 2\theta}{2} \right)$$

$$= |\vec{p}| + \frac{m_1^2 + m_2^2}{4|\vec{p}|} + \frac{m_1^2 - m_2^2}{4|\vec{p}|} \cos 2\theta$$

$$H_{\text{eff}} = |\vec{p}| + \frac{m_1^2 + m_2^2}{4|\vec{p}|} - \frac{\cos 2\theta \Delta m^2}{4|\vec{p}|} \quad \Delta m^2 = m_2^2 - m_1^2$$

Neutrino Oscillations in matter

Ex

$$H_{\text{eff}} = H_{\text{p}\bar{\nu}} = -\frac{\Delta m^2}{4|\vec{p}|} \sin 2\theta .$$

$$\Rightarrow H_V = \frac{\Delta m^2}{4|\vec{p}|} \begin{pmatrix} -\cos(2\theta) & -\sin(2\theta) \\ -\sin(2\theta) & \cos(2\theta) \end{pmatrix} + \left(|\vec{p}| + \frac{m_1^2 + m_2^2}{4|\vec{p}|} \right) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

consider only the term as relative phases
matter.

1×1 .

$$\text{Since } E = |\vec{p}| \Rightarrow$$

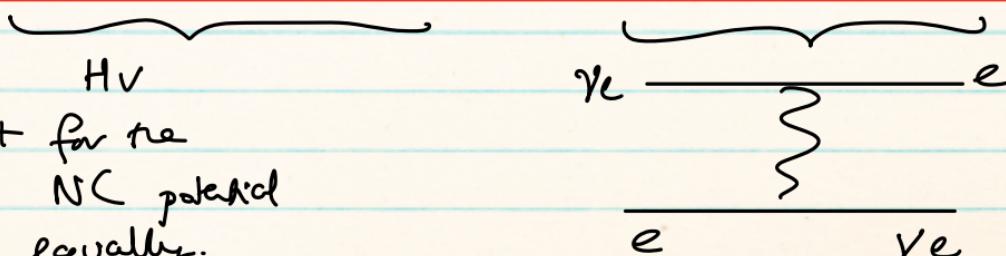
$$H_V = \frac{\Delta m^2}{4E} \begin{pmatrix} -\cos(2\theta) & -\sin(2\theta) \\ -\sin(2\theta) & \cos(2\theta) \end{pmatrix}$$

Neutrino Oscillations in matter

$$i \frac{\partial}{\partial t} \begin{pmatrix} \nu_e(t) \\ \nu_{\mu}(t) \end{pmatrix} = H_V \begin{pmatrix} \nu_e(t) \\ \nu_{\mu}(t) \end{pmatrix} = \frac{\Delta m^2}{4E} \begin{pmatrix} -\cos(2\alpha) & -\sin(2\alpha) \\ -\sin(2\alpha) & \cos(2\alpha) \end{pmatrix} \begin{pmatrix} \nu_e(t) \\ \nu_{\mu}(t) \end{pmatrix}$$

In matter :

$$H_M = \frac{\Delta m^2}{4E} \begin{pmatrix} -\cos(2\alpha) & -\sin(2\alpha) \\ -\sin(2\alpha) & \cos(2\alpha) \end{pmatrix} + G_F N e \sqrt{2} \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$



NB we only account for the CC potential as the NC potential affects ν_e & ν_{μ} equally.

Neutrino Oscillations in matter

We can diagonalise H_m to find energy eigenstates in matter

$$U_m^+ E_{\text{DIAG}} H_m U_m = H_m$$
$$\Rightarrow E_{\text{DIAG}} = U_m H_m U_m^+$$

θ_m ≡ mixing angle in matter.

$$\begin{pmatrix} E_m & 0 \\ 0 & E_{2m} \end{pmatrix} = \begin{pmatrix} \cos \theta_m & \sin \theta_m \\ -\sin \theta_m & \cos \theta_m \end{pmatrix} \begin{pmatrix} \frac{\Delta m^2}{4E} \cos 2\theta + \sqrt{2} G_F N_e & -\frac{\Delta m^2}{4E} \sin 2\theta \\ -\frac{\Delta m^2}{4E} \sin 2\theta & \frac{\Delta m^2}{4E} \cos 2\theta \end{pmatrix}$$
$$\begin{pmatrix} \cos \theta_m & -\sin \theta_m \\ \sin \theta_m & \cos \theta_m \end{pmatrix}$$

⇒ the (1,2) & (2,1) entry must be zero.

Neutrino Oscillations in matter

(1,2) entry:

$$-\frac{\Delta m^2}{4E} \cos^2 \theta_m \sin 2\theta - \frac{\sqrt{2} G_F N_e}{2} \underbrace{\cos \theta_m \sin \theta_m}_{\frac{1}{2} \sin 2\theta_m} + \frac{\Delta m^2}{2E} \cos 2\theta \cos \theta_m \sin \theta_m \\ + \frac{\Delta m^2}{4E} \sin^2 \theta_m \sin \theta_m^2 = 0$$

$$= -\frac{\Delta m^2}{4E} \sin 2\theta \cos 2\theta_m - \frac{\sqrt{2} G_F N_e}{2} \sin 2\theta_m + \frac{\Delta m^2}{4E} \cos 2\theta \sin 2\theta_m = 0$$

$$= -\frac{\Delta m^2}{4E} \sin 2\theta \cos 2\theta_m + \sin 2\theta_m \left(\frac{\Delta m^2}{4E} \cos 2\theta - \frac{1}{\sqrt{2}} G_F N_e \right) = 0$$

$$\Rightarrow \frac{\Delta m^2}{4E} \sin 2\theta \cos 2\theta_m = \sin 2\theta_m \left(\frac{\Delta m^2}{4E} \cos 2\theta - \frac{1}{\sqrt{2}} G_F N_e \right)$$

Neutrino Oscillations in matter

$$\frac{\Delta m^2}{4E} \sin 2\theta \cos 2\theta m = \sin 2\theta m \left(\frac{\Delta m^2}{4E} \cos 2\theta - \frac{1}{\sqrt{2}} G_F N_e \right)$$

$$\Rightarrow \tan 2\theta m = \frac{\frac{\Delta m^2}{4E} \sin 2\theta}{\left(\frac{\Delta m^2}{4E} \cos 2\theta - \frac{G_F N_e}{\sqrt{2}} \right)}$$

$$= \frac{\frac{\Delta m^2}{2E} \sin 2\theta}{\left(\frac{\Delta m^2}{2E} \cos 2\theta - \sqrt{2} G_F N_e \right)}$$

Neutrino Oscillations in matter

Show that the difference in the eigenvalues.

$$E_{1m} - E_{2m} = \sqrt{\left(\frac{\Delta m^2}{2E} \cos(2\theta) - \sqrt{2} G_F N_e\right)^2 + \left(\frac{\Delta m^2}{2E} \sin(2\theta)\right)^2}$$

To do this $\det(dI - k_m) = 0$ and rearrange bit.

Neutrino Oscillations in matter

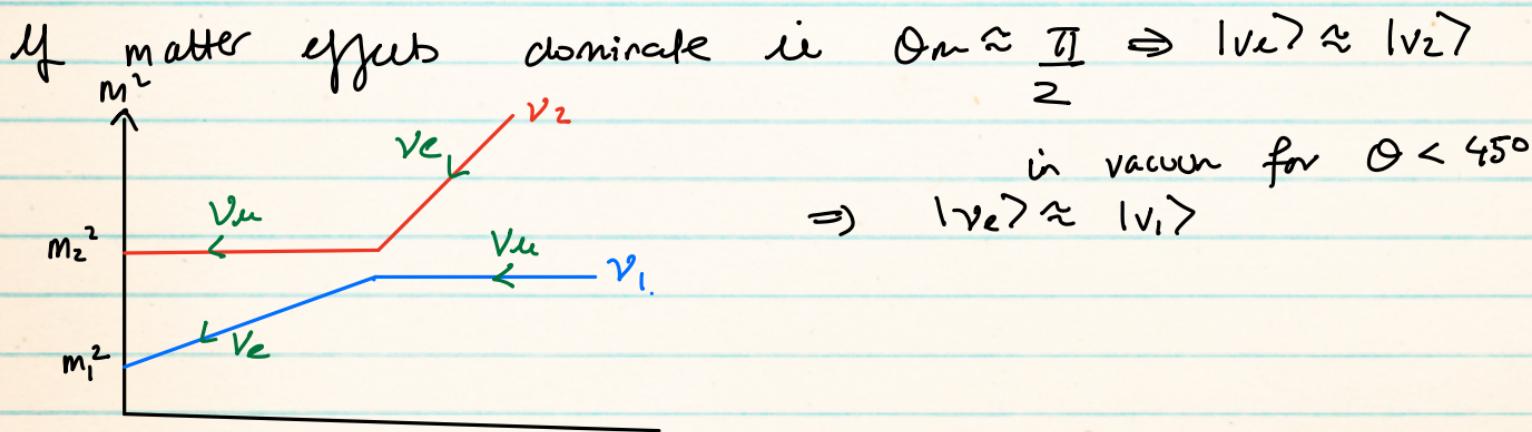
$$\tan 2\theta_m = \frac{\frac{\Delta m^2}{2E} \sin 2\theta}{\left(\frac{\Delta m}{2E} \cos 2\theta - \frac{G_F N_e}{\sqrt{2}} \right)}$$

- $G_F N_e = 0 \Rightarrow \theta = \theta_m$. i.e vacuum case is recovered.
- $\frac{G_F N_e}{\sqrt{2}} = \frac{\Delta m}{2E} \cos 2\theta \Rightarrow \theta_m = \frac{\pi}{4}$ (maximal mixing independent of vacuum mixing!)
- $\frac{G_F N_e}{\sqrt{2}} \gg \frac{\Delta m}{2E} \cos 2\theta \Rightarrow \theta_m = \frac{\pi}{2}$.

Neutrino Oscillations in matter

$$|\nu_e\rangle = (\cos \theta |\nu_1\rangle + \sin \theta |\nu_2\rangle) \quad \text{in vacuum.}$$
$$|\nu_\mu\rangle = -\sin \theta |\nu_1\rangle + \cos \theta |\nu_2\rangle$$

$$|\nu_e\rangle = (\cos_m |\nu_1\rangle + \sin_m |\nu_2\rangle) \quad \text{in matter.}$$
$$|\nu_\mu\rangle = -\sin_m |\nu_1\rangle + \cos_m |\nu_2\rangle$$



Neutrino Oscillations in matter

The matter potential creates a mixing angle in matter (θ_m)
It also creates a mass squared splitting in matter. To derive
this recall the Hamiltonian in vacuum:

$$H_v = \frac{\Delta m^2}{4E} \begin{pmatrix} -\cos(2\alpha) & -\sin(2\alpha) \\ -\sin(2\alpha) & \cos(2\alpha) \end{pmatrix}$$

Δm^2 = mass squared
splitting in vacuum.

We can equally well parameterize the Hamiltonian in matter as:

$$\textcircled{1} - H_m = \frac{\Delta m_m^2}{4E} \begin{pmatrix} -\cos(2\alpha_m) & -\sin(2\alpha_m) \\ -\sin(2\alpha_m) & \cos(2\alpha_m) \end{pmatrix}$$

Δm_m^2 = mass squared
splitting in matter.

But from before we had:

$$\textcircled{2} \quad H_m = \frac{\Delta m^2}{4E} \begin{pmatrix} -\cos(2\alpha) & -\sin(2\alpha) \\ -\sin(2\alpha) & \cos(2\alpha) \end{pmatrix} + G_F N e \sqrt{2} \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix}$$

Neutrino Oscillations in matter

① and ② are equal. ∴

difference in eigenvalues of ① = difference in eigenvalues of ②

$$\text{difference in eigenvalues of } ② = \frac{\Delta m_m^2}{2E}$$

$$\text{difference in eigenvalues of } ① = \frac{\Delta m^2}{2E} \sqrt{\sin^2 2\theta + (\cos 2\theta - 2\sqrt{2} g_f N_e)^2}$$

$$\Rightarrow \Delta m_m^2 = \Delta m^2 \sqrt{\sin^2 2\theta + (\cos 2\theta - 2\sqrt{2} g_f N_e)^2}$$

$$\Rightarrow m_{1m}^2 = -\frac{1}{2} \Delta m^2 \sqrt{\sin^2 2\theta + (\cos 2\theta - 2\sqrt{2} g_f N_e)^2}$$

$$m_{2m}^2 = +\frac{1}{2} \Delta m^2 \sqrt{\sin^2 2\theta + (\cos 2\theta - 2\sqrt{2} g_f N_e)^2}$$

Solar neutrinos

- Confirmation of neutrino oscillations came in 2001 by the Sudbury Neutrino Observatory. They measured not only electron neutrino flux but all flavour neutrinos via NC interactions

