Photoproduction

MC4EIC 2024

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Outline

• MC modelling of photoproduction and some selected data comparisons

Outline

- 1. Introduction
- 2. Structure of (quasi-)real photons
- 3. Jet production
- 4. Single-particle observables
- 5. Multiparticle correlations
- 6. Summary & Outlook



[figure by P. Skands]

Introduction

Photoproduction

- High-energy collisions with real photons as beam particles
- In the $Q^2 \rightarrow 0$ limit can factorize photon flux, f_{γ} , from the hard interaction

 $\mathrm{d}\sigma^{AB} = f^{A}_{\gamma}(x,Q^2) \,\otimes\, \mathrm{d}\sigma^{\gamma B}$

- Can be studied in colliders with charged beams
 - e^+e^- : LEP, γ - γ , also γ^* - γ
 - e-p: HERA, γ -p, EIC: γ -p and γ -A
 - p-p: LHC, γ -p, γ - γ
 - p-A: LHC, γ -p, γ - γ
 - A-A: LHC, γ -A, γ - γ



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Equivalent photon approximation

Photon fluxes for different beams

• In case of a point-like lepton we have

$$f_{\gamma}^{I}(x, Q^{2}) = rac{lpha_{\mathsf{em}}}{2\pi} rac{(1+(1-x)^{2})}{x} rac{1}{Q^{2}}$$

• For protons need to account the form factor

$$f_{\gamma}^{p}(x,Q^{2}) = \frac{\alpha_{\rm em}}{2\pi} \frac{(1+(1-x)^{2})}{x} \frac{1}{Q^{2}} \frac{1}{(1+Q^{2}/Q_{0}^{2})}$$



where $Q_0^2 = 0.71 \,\text{GeV}^2$ (Drees-Zeppenfeld) \Rightarrow Large Q^2 heavily suppressed

• With heavy nuclei use *b*-integrated point-like-charge flux

$$f_{\gamma}^{A}(x) = \frac{2\alpha_{\rm EM}Z^{2}}{x \pi} \left[\xi \, K_{1}(\xi) K_{0}(\xi) - \frac{\xi^{2}}{2} \left(K_{1}^{2}(\xi) - K_{0}^{2}(\xi) \right) \right]$$

where $\xi = b_{\rm min}\, {\rm x}\, m$ where $b_{\rm min}$ reject nuclear overlap $\Rightarrow Q^2 \ll 1\,{\rm GeV^2}$

Structure of real photons

Photon structure at $Q^2 \approx 0 \text{ GeV}^2$



Partonic structure of resolved (anom. + VMD) photon encoded in photon PDFs

$$f_i^{\gamma}(\mathbf{x}_{\gamma},\mu^2) = f_i^{\gamma,\text{dir}}(\mathbf{x}_{\gamma},\mu^2) + f_i^{\gamma,\text{anom}}(\mathbf{x}_{\gamma},\mu^2) + f_i^{\gamma,\text{VMD}}(\mathbf{x}_{\gamma},\mu^2)$$

- $f_i^{\gamma,\text{dir}}(\mathbf{x}_{\gamma},\mu^2) = \delta_{i\gamma}\delta(1-\mathbf{x}_{\gamma})$
- $f_i^{\gamma,\text{anom}}(x_{\gamma},\mu^2)$: Perturbatively calculable
- $f_i^{\gamma, \text{VMD}}(x_{\gamma}, \mu^2)$: Non-perturbative, fitted or vector-meson dominance (VMD)

Factorized cross section

$$\mathrm{d}\sigma^{\gamma \mathsf{A} \to k \mathsf{I} + \mathsf{X}} = f_i^{\gamma}(\mathsf{x}_{\gamma}, \mu^2) \otimes f_j^{\mathsf{A}}(\mathsf{x}_{\mathsf{p}}, \mu^2) \otimes \mathrm{d}\sigma^{ij \to k \mathsf{I}}$$

PDFs for resolved photons

DGLAP equation for photons

• Additional term due to $\gamma \rightarrow q\overline{q}$ splittings

$$\frac{\partial f_i^{\gamma}(x,Q^2)}{\partial \log(Q^2)} = \frac{\alpha_{\text{em}}}{2\pi} e_i^2 \mathsf{P}_{i\gamma}(x) + \frac{\alpha_{\text{s}}(Q^2)}{2\pi} \sum_j \int_x^1 \frac{\mathrm{d}z}{z} \, \mathsf{P}_{ij}(z) \, f_j(x/z,Q^2)$$

where $P_{i\gamma}(x) = 3(x^2 + (1 - x)^2)$ for quarks, 0 for gluons (LO)



Evolution equation and ISR for resolved photons

ISR probability based on DGLAP evolution

• Add a term corresponding to $\gamma \rightarrow q\overline{q}$ to (conditional) ISR probability

$$\mathrm{d}\mathcal{P}_{a\leftarrow b} = \frac{\mathrm{d}Q^2}{Q^2} \frac{\alpha_{\rm s}}{2\pi} \frac{x' f_a^{\gamma}(x',Q^2)}{x f_b^{\gamma}(x,Q^2)} P_{a\rightarrow bc}(z) \,\mathrm{d}z + \frac{\mathrm{d}Q^2}{Q^2} \frac{\alpha_{\rm em}}{2\pi} \frac{e_b^2 \, P_{\gamma\rightarrow bc}(x)}{f_b^{\gamma}(x,Q^2)}$$

- Corresponds to ending up to the beam photon during evolution
 - \Rightarrow Parton originated from the point-like (anomalous) part of the PDFs
 - No further ISR or MPIs below the scale of the splitting
 - Implemented for the default Simple Shower in Pythia 8



Jet photoproduction

Jet photoproduction

Factorize the cross section

• Direct processes: Convolute photon flux f_{γ} with proton PDFs f_i^p and partonic coefficient functions $d\hat{\sigma}$

 $\mathrm{d}\sigma^{\mathrm{ep} o kl + X} = f^{\mathrm{e}}_{\gamma}(x, \mathrm{Q}^2) \, \otimes \, f^{\mathrm{p}}_{j}(x_{\mathrm{p}}, \mu^2) \, \otimes \, \mathrm{d}\hat{\sigma}^{\gamma j o kl}$

Resolved processes: Convolute also with photon PDFs

 $\mathrm{d}\sigma^{\mathrm{ep}\to kl+X} = f_{\gamma}^{\mathrm{e}}(\mathbf{x}, \mathbf{Q}^2) \otimes f_i^{\gamma}(\mathbf{x}_{\gamma}, \mu^2) \otimes f_j^{\mathrm{p}}(\mathbf{x}_{\mathrm{p}}, \mu^2) \otimes \mathrm{d}\sigma^{ij \to kl}$

- In case of $\gamma\text{-}\gamma$ convolute with two photon fluxes

Generate parton showers and MPIs for resolved events

- Sample x and Q^2 , setup γ -p sub-system with $W_{\gamma p}$
- Evolve γ -p (γ - γ) as any hadronic collision (with MPIs)
- Add beam remants for resolved photons



Jet photoproduction

A comparison study between Herwig, Sherpa and Pythia for jet photoproduction

[I. Helenius, P. Meinzinger, S. Plätzer, P. Richardson: in progress]



- Start with hard-process partons in LEP-like setup
- Pythia and Sherpa in agreement when identical $\alpha_{\rm s}$, PDFs, photon flux, sub-processes with massive partons...

Jet photoproduction in LEP [I. Helenius, P. Meinzinger, S. Plätzer, P. Richardson: in progress]



- Large LO \rightarrow NLO correction in Sherpa, uncertainty from scale variations
- Herwig LO close to Sherpa LO, NLO with matchbox underway
- Pythia result with highest cross section, still within the uncertainties

• Summary of the modelling differences between the generators

Property	Pythia	Sherpa	Herwig
Flux	LL	NLL	LL
$\alpha_{s}(M_{Z}^{2})$	0.130, 1-loop running	0.118, 3-loop running	
PDFs	CJKL	SAS2M	SAS2M
Remnants	forced splittings/PS rejection	PS rejection	forced splitting
$\gamma ightarrow {\it q} ar q$ Splitting	yes	no	no
MPI tuning	preliminary γ -p/ γ - γ tune	untuned	untuned

Jet photoproduction in HERA [I. Helenius, P. Meinzinger, S. Plätzer, P. Richardson: in progress]



- Similar hierarchy as for LEP comparisons
- Herwig results on their way

Jet photoproduction in EIC [I. Helenius, P. Meinzinger, S. Plätzer, P. Richardson: in progress]



- We set up a Rivet analysis with highest-energy EIC-kinematics with similar jet reconstruction as in HERA
- Similar observations in jet E_T spectra as with HERA kinematics
- Large differences in multiplicity distribution (lack of tuning)

Jet photoproduction in HERA



- Possible to run Rivet analyses for Pythia 6 provided in the EIC MCEG GitLab
- Enable systematic comparisons with the modern generators

Single-particle observables

Single-inclusive hadrons

- Allows to study events at lower p_T
- More sensitive to hadronization and MPI effects than inclusive-/di-jet data
- May require soft QCD modelling

H1 analysis

- Photon flux integrated out, data for γ -p at $\langle W_{\gamma p} \rangle \approx 200 \, \text{GeV}$
- Rivet routine added in 3.1.10 release



Single-inclusive hadrons

- Double-resolved photons dominate at low p_T, only these have MPIs
- Single-resolved contribute $\sim 20\%$
- Direct processes take over above $p_{\rm T}\sim 5~{\rm GeV/c}$

OPAL data

- Again potential to constrain MPI parameters
- Rivet routine added in 3.1.10 release



[OPAL: PLB 651 (2007) 92-101]

Charged-particle production in LEP



[OPAL: PLB 651 (2007) 92-101]

A handle for energy dependence (LEP/LOG-tune in Pythia)

Multi-particle correlations

Multi-particle correlations

Correlations can arise from

- Jets, particle decays, rescattering
- Initial state effects, eg. CGC
- Final-state effects
 - hydrodynamic evolution of quark-gluon plasma,
 - Collectivity in hadronization, eg. string interactions

ZEUS analysis

- High-multiplicity events ($N_{ch} > 20$)
- Reasonable agreement with MPIs in Pythia
- Rivet routine in progress



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Collectivity in $\gamma\text{-}\mathrm{p}$ at the LHC

CMS analysis for p+Pb

- Pb provides photon flux, γ-p at energies similar to HERA and beyond
- Fourier fit to $dN/d\Delta\phi$ to obtain v_2
- Finite v_2 for γ -p, in line with Pythia
- No explicit collectivity included in the model
- \Rightarrow No collective behaviour observed



Two-particle correlations in γ -A by ATLAS

- ATLAS apply template-fitting method to extract v_n from two-particle correlations
 - Perform a Fourier fit to obtain c_n's for low-multiplicity events

$$\mathsf{Y}^{\mathsf{LM}}(\Delta\phi) = c_0 + 2 \cdot \sum_{n=1}^4 c_n \cos(n\Delta\phi)$$

• Fit high multiplicity $v_{n,n}$'s on top

$$Y^{\mathsf{HM}}(\Delta\phi) = F \cdot Y^{\mathsf{LM}}(\Delta\phi) + G\left[1 + 2 \cdot \sum_{n=2}^{4} v_{n,n} \cos(n\Delta\phi)\right]$$

Free parameters c_n , $v_{n,n}$, F, G

Should subtract the "non-flow" contributions



[ATLAS: PRC 104, 014903 (2021)]

ATLAS data for v_n in γ +Pb



- Non-zero flow coefficients for γ +Pb, even after the template fit
- Expected baseline from MC simulations? Hydro or hadronization effect?

Summary

- Photoproduction provides access to rich structure of real photons
- Modern event generators can handle different contributions
- Not only HERA data available, also LEP and LHC should be included
- Number of Rivet available analyses has increased in recent years
 - \Rightarrow Enables systematic and global tuning of MC generator parameters

Open questions

- Obviously still important Rivet analysis missing, which ones?
- Interplay with photon PDF fitting and MPIs?
- UPCs at the LHC provides first collider-energy γ -A
 - \Rightarrow Origin of the collective behaviour? Relevant for EIC?

Backup slides

Jet photoproduction in LEP



- Also a more recent jet analysis in γ - γ from LEP available
- Only p_{T} with varying η range