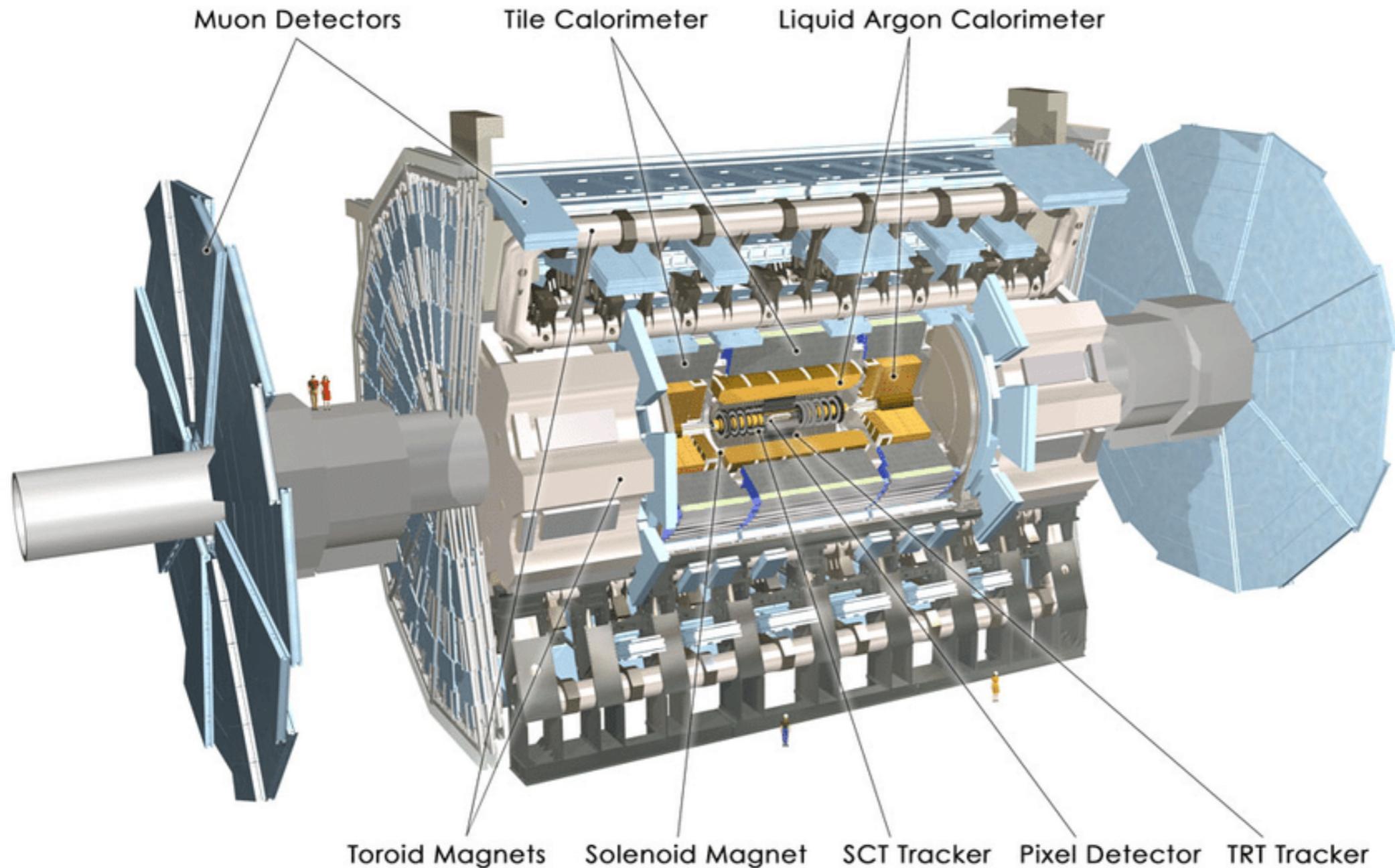


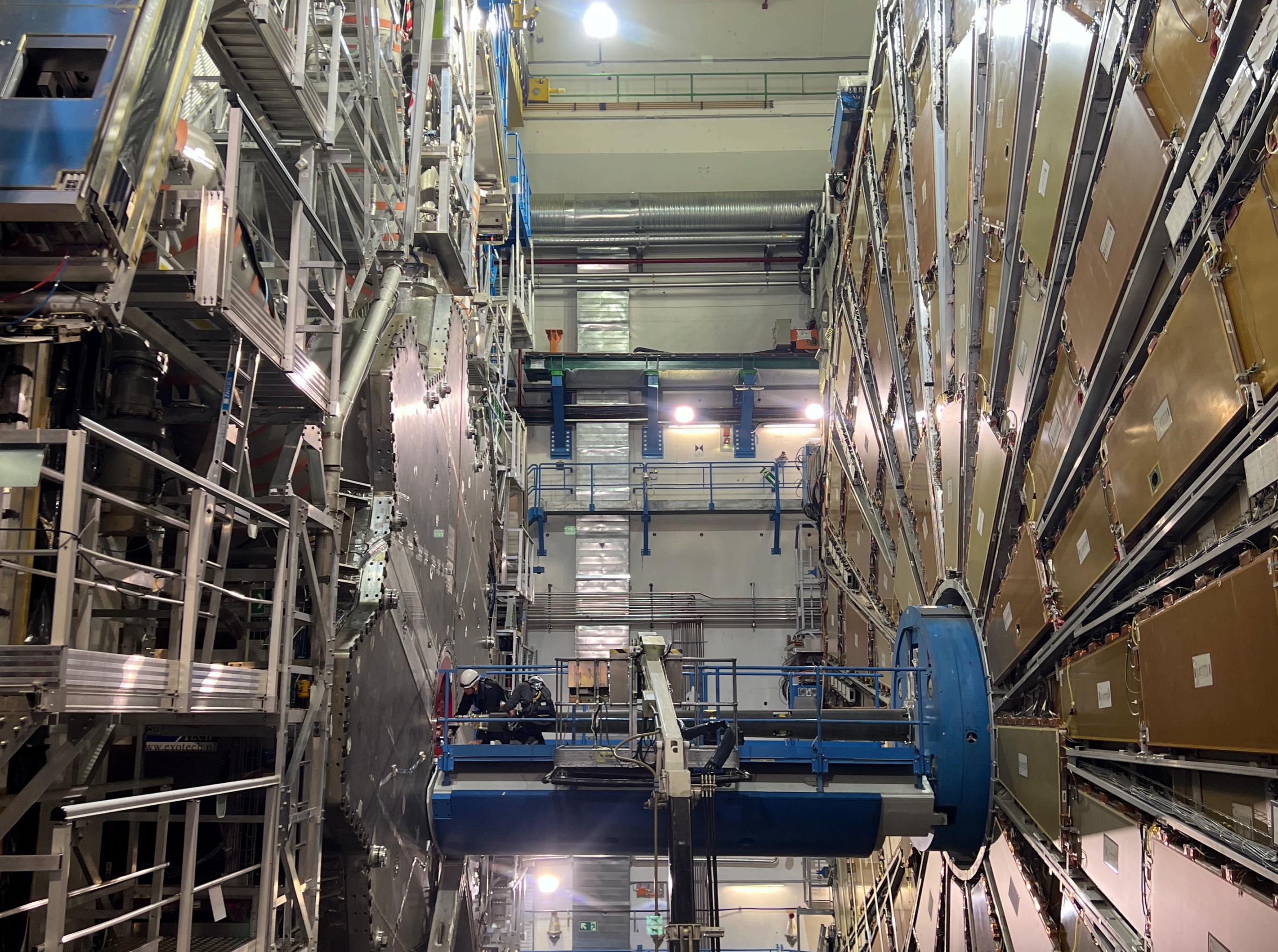
# Observation of Entanglement \_\_\_\_\_ in \_\_\_\_\_ Top Quark Pairs at ATLAS

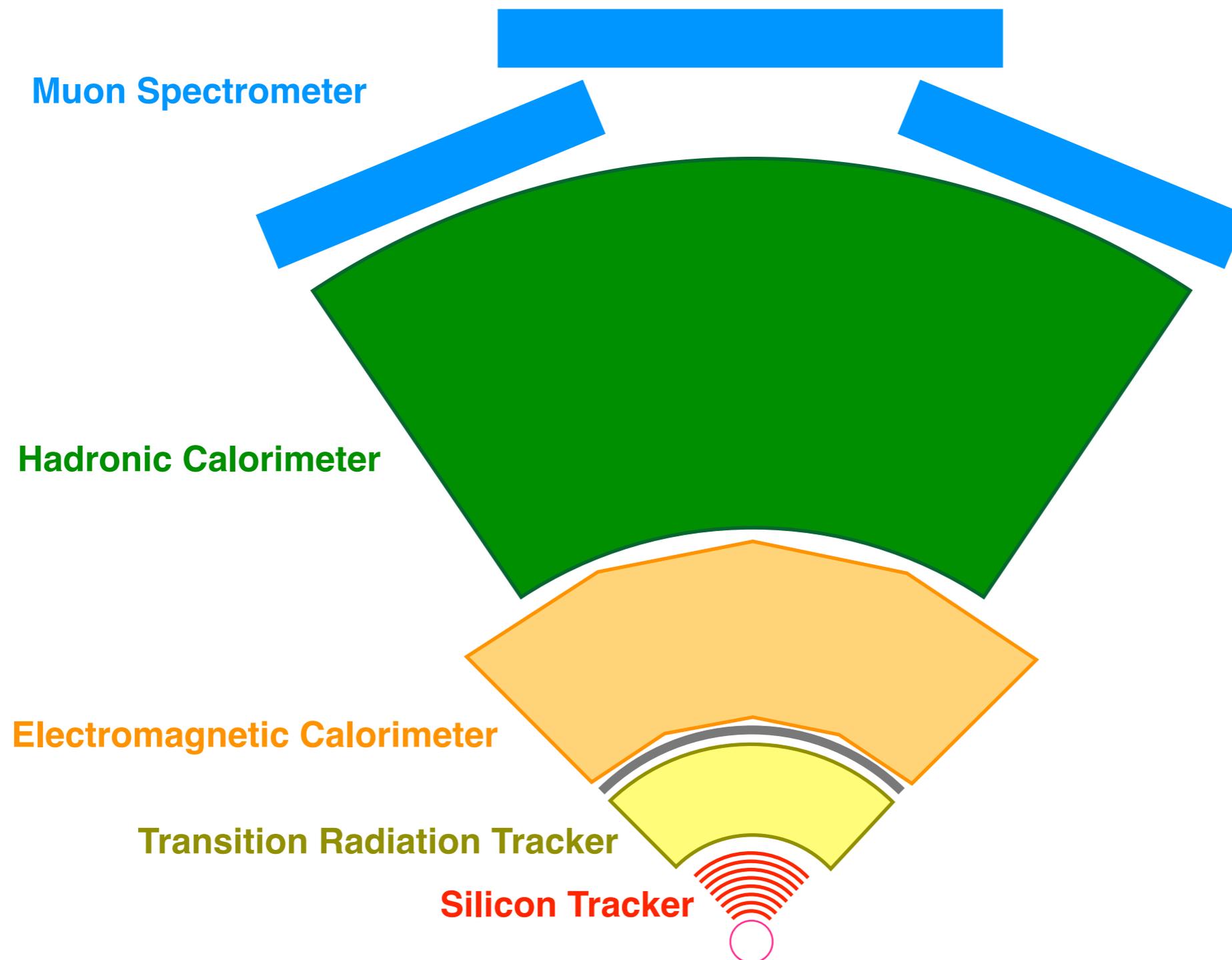
Oxford, 01/10/2024

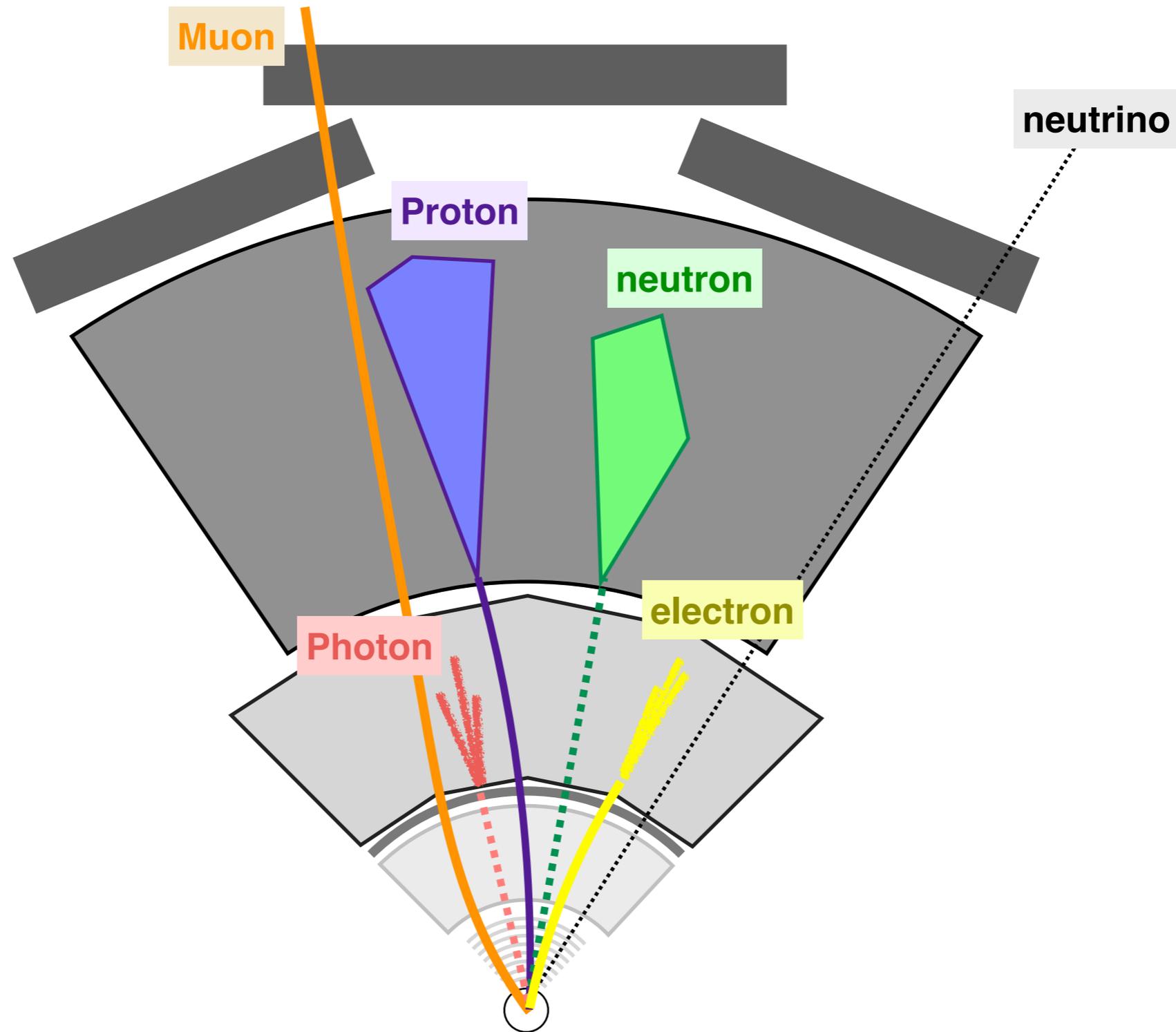
- **Top Physics and properties measurements at hadron colliders**
- **Quantum information in high energy particle physics**
- **The recent ATLAS result**
- **Implications and expected future results**

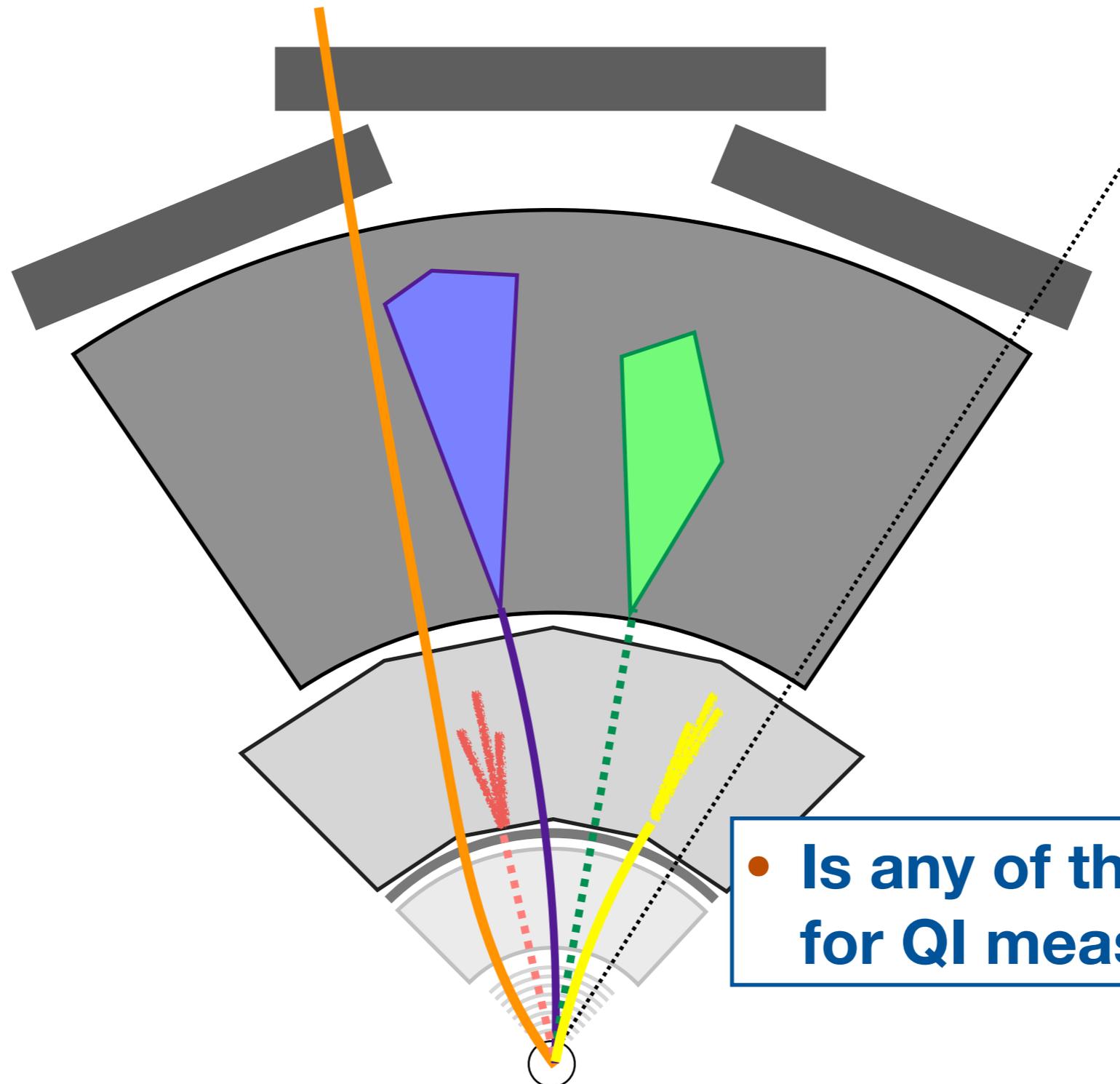


- **A Toroidal Lhc ApparatuS:**  
*“the best experiment with the worst acronym.”*

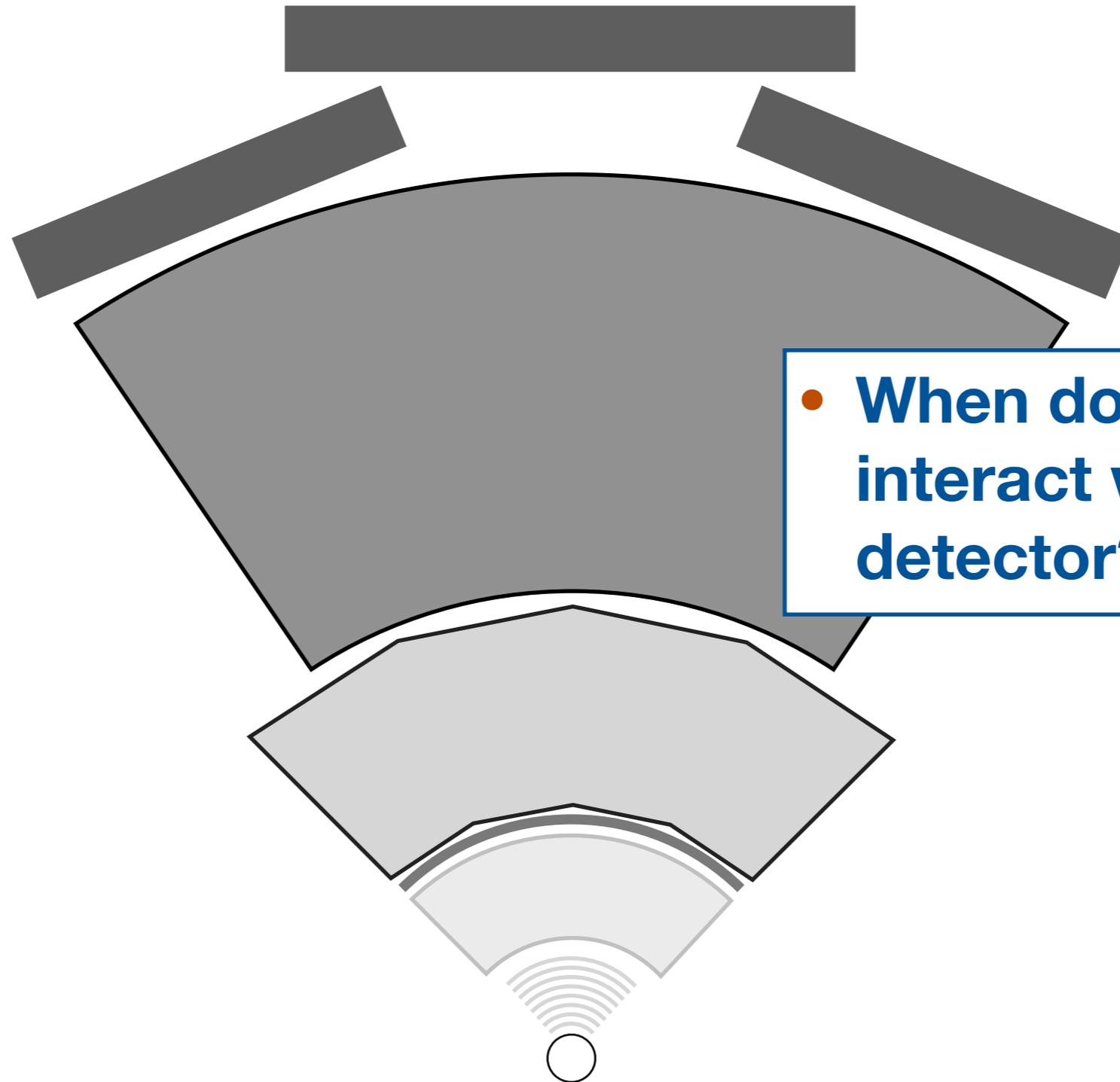




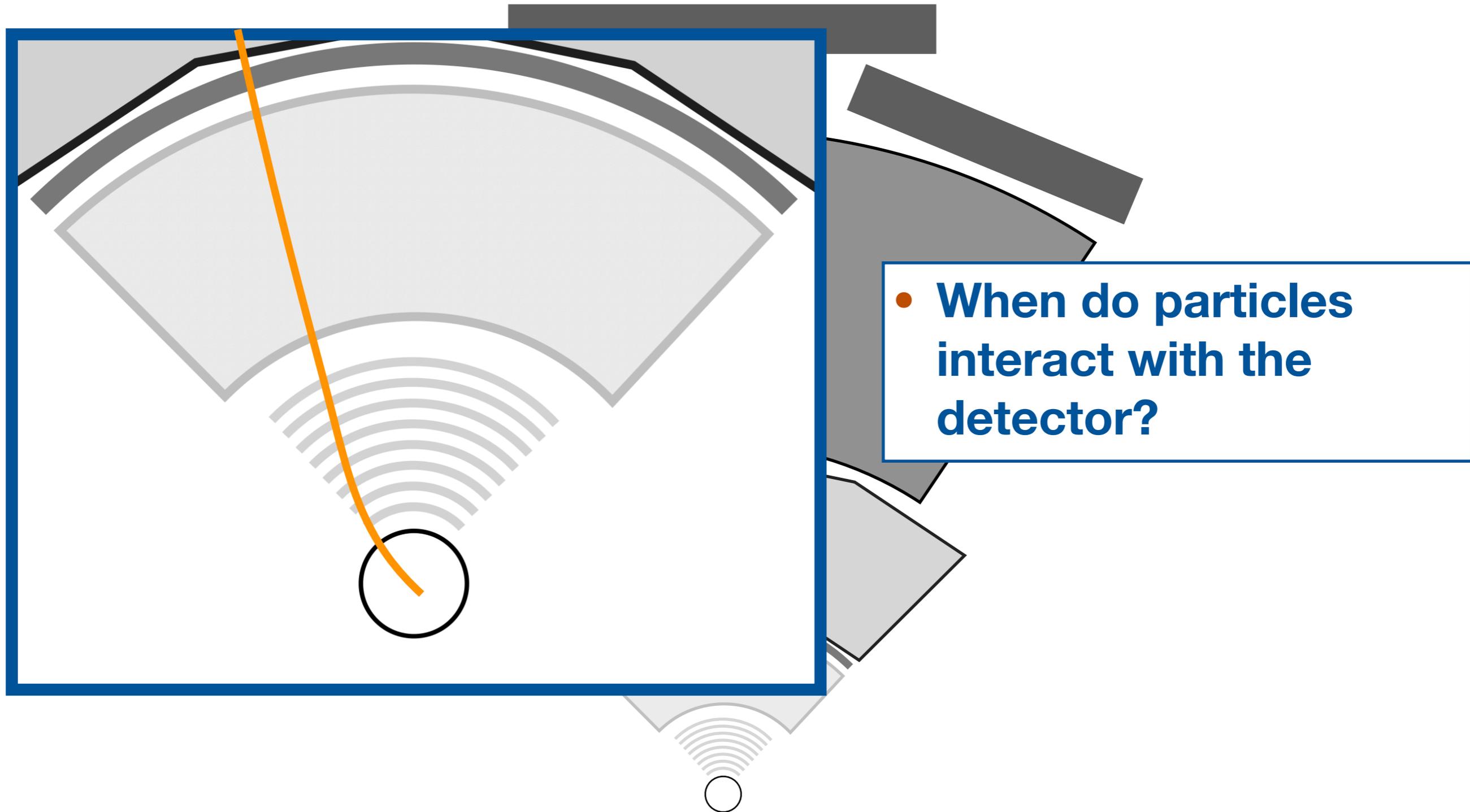


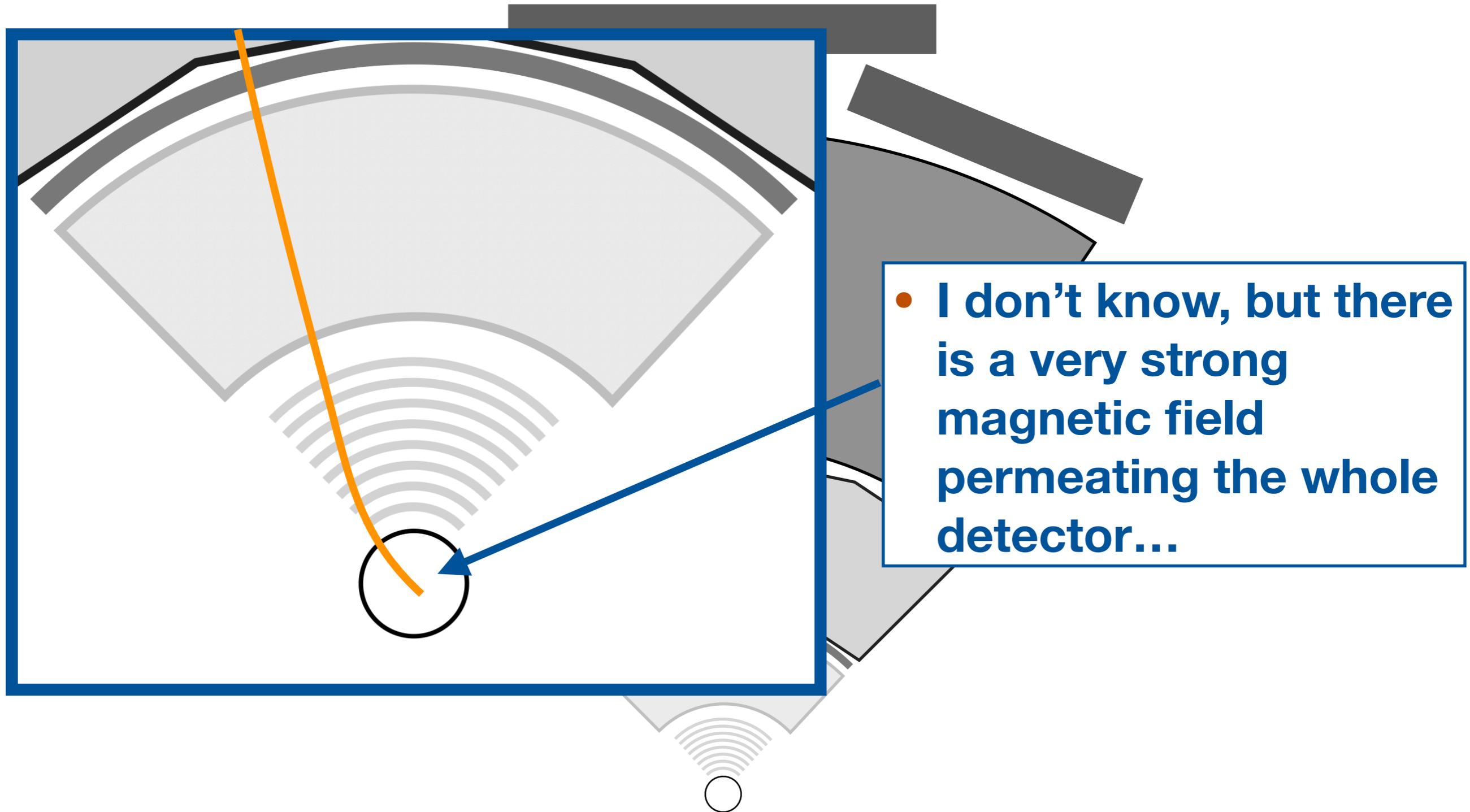


- Is any of this relevant for QI measurements?

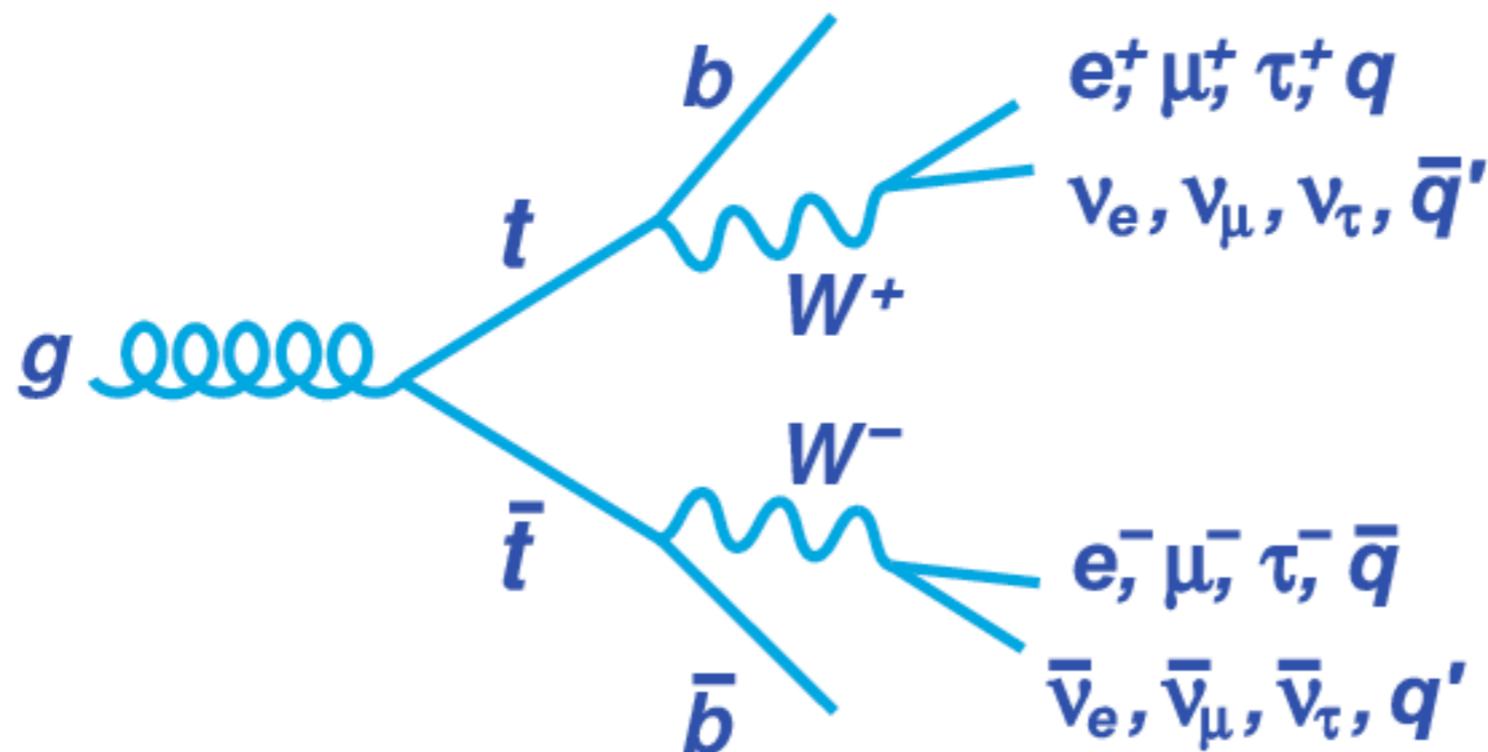


- **When do particles interact with the detector?**

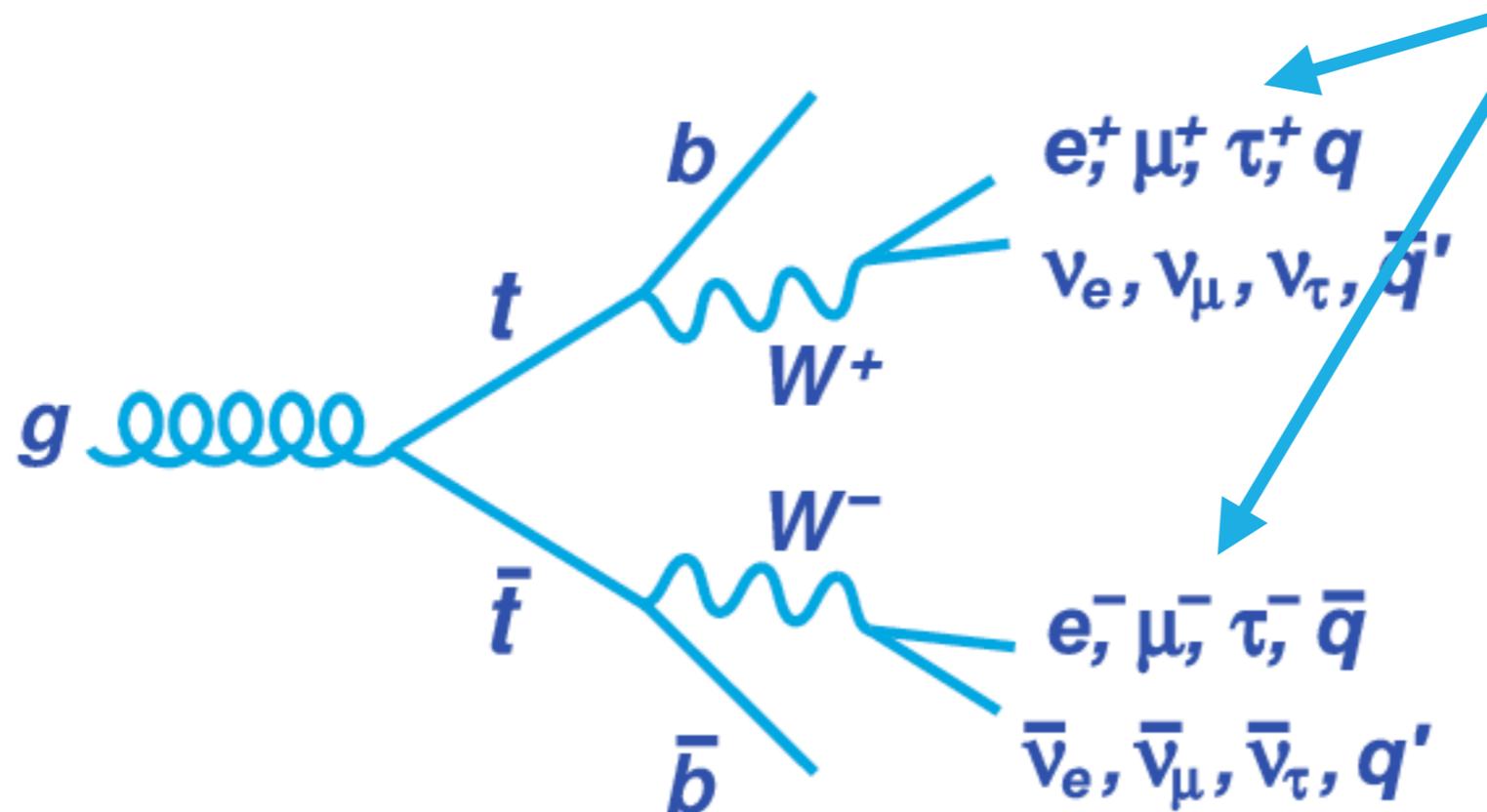




- This measurement uses dileptonic  $t\bar{t}$  events, which are difficult to reconstruct.

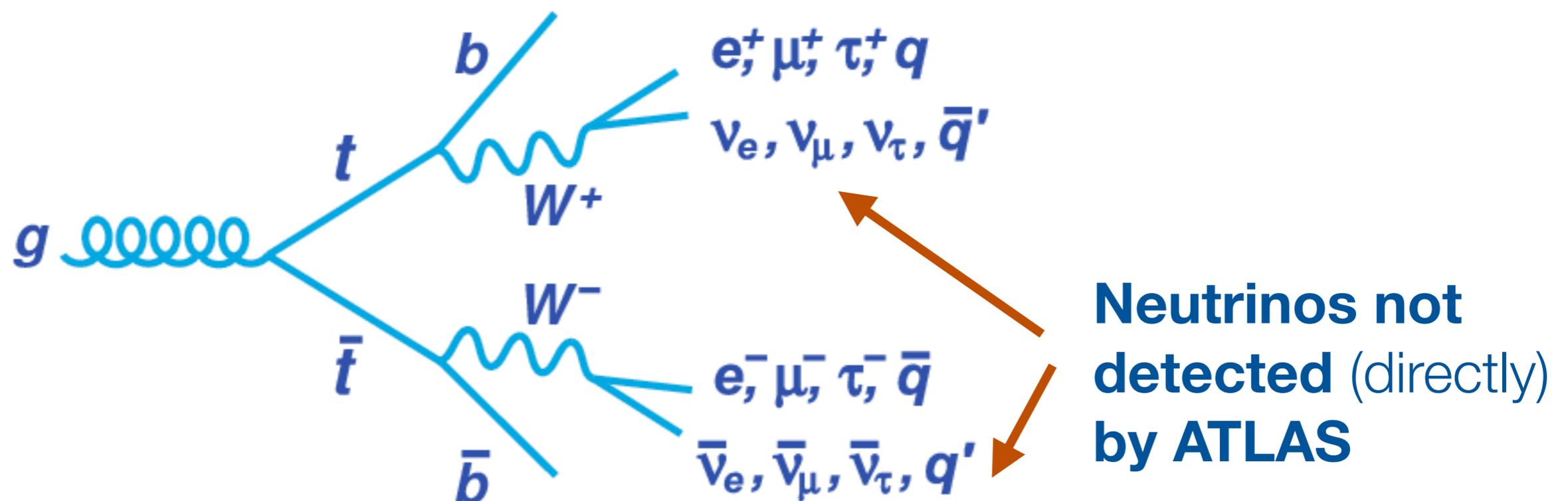


- Why is it hard to reconstruct top quarks?

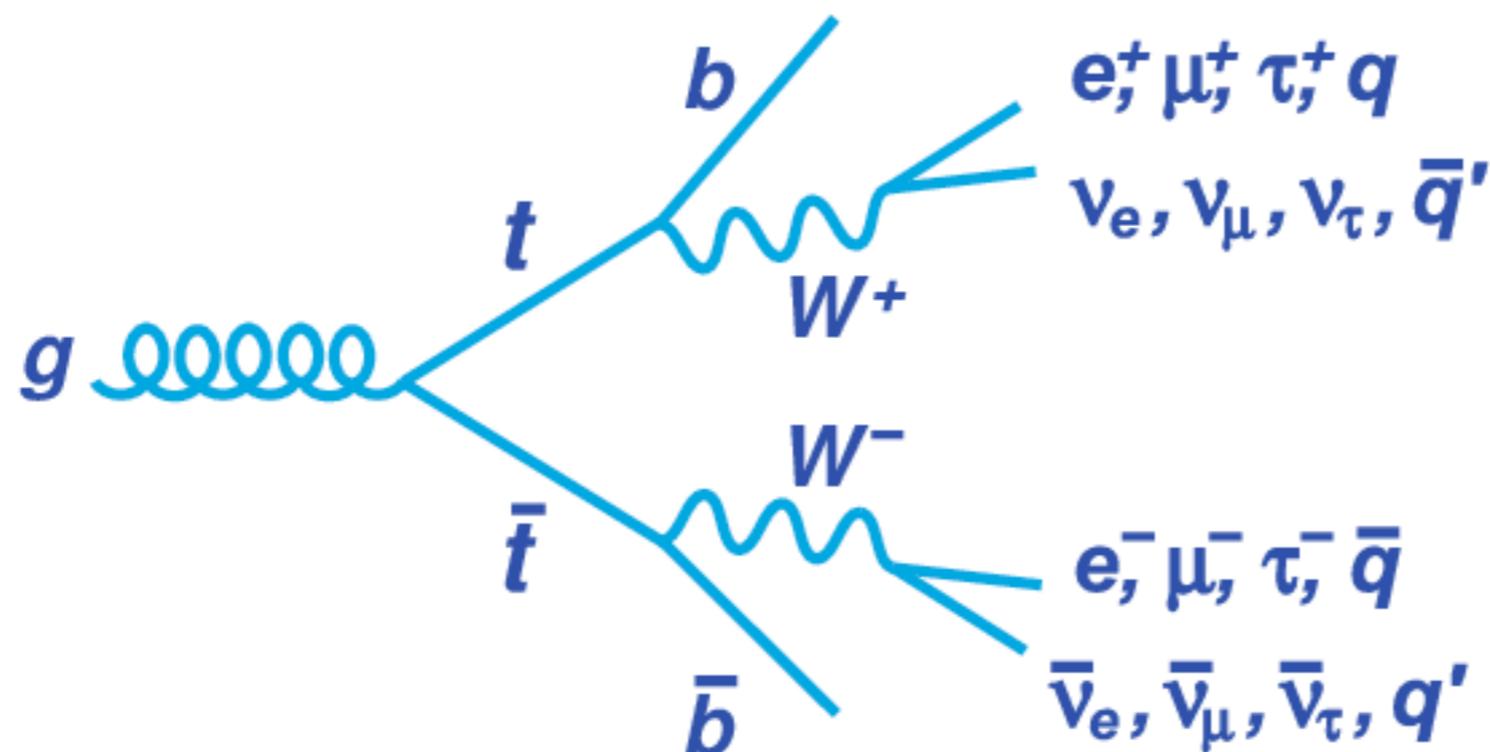


**Charged leptons are the perfect spin analyser!**

- Why is it hard to reconstruct top quarks?

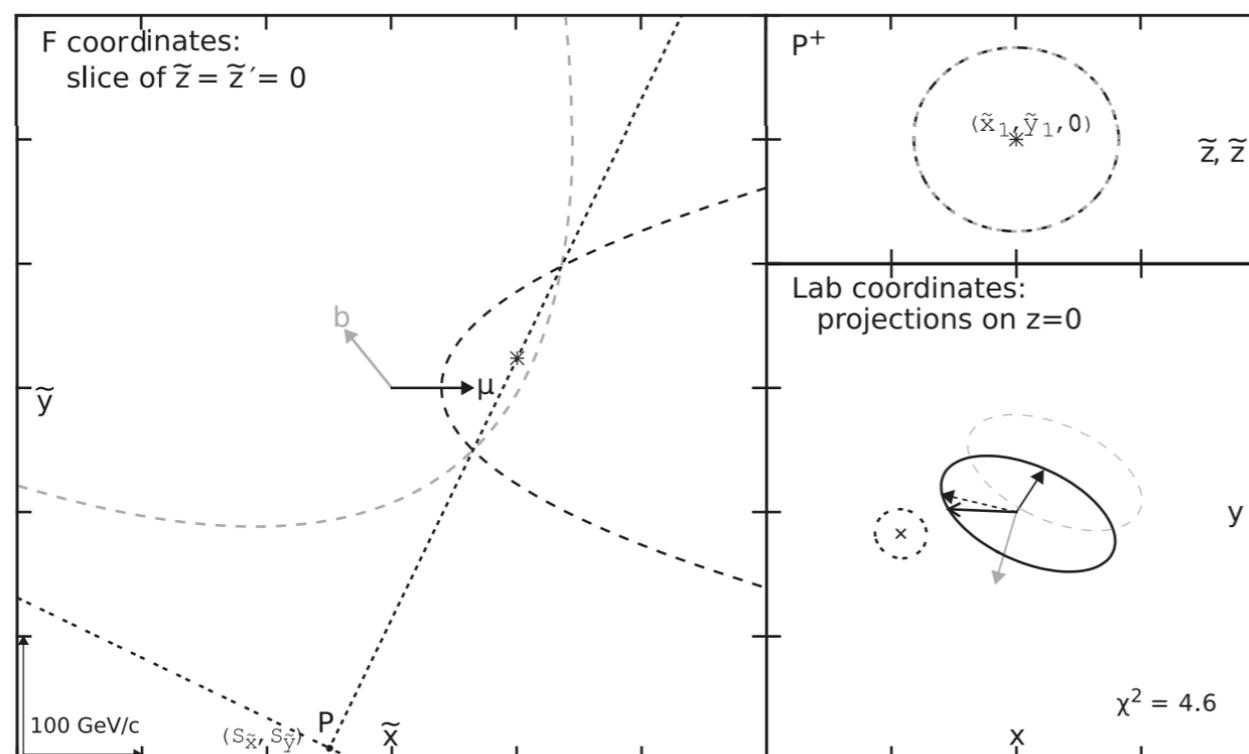


- Why is it hard to reconstruct top quarks?



- **ATLAS selects events with two charged leptons in the final state** (+ 1 or more b-tagged jets).

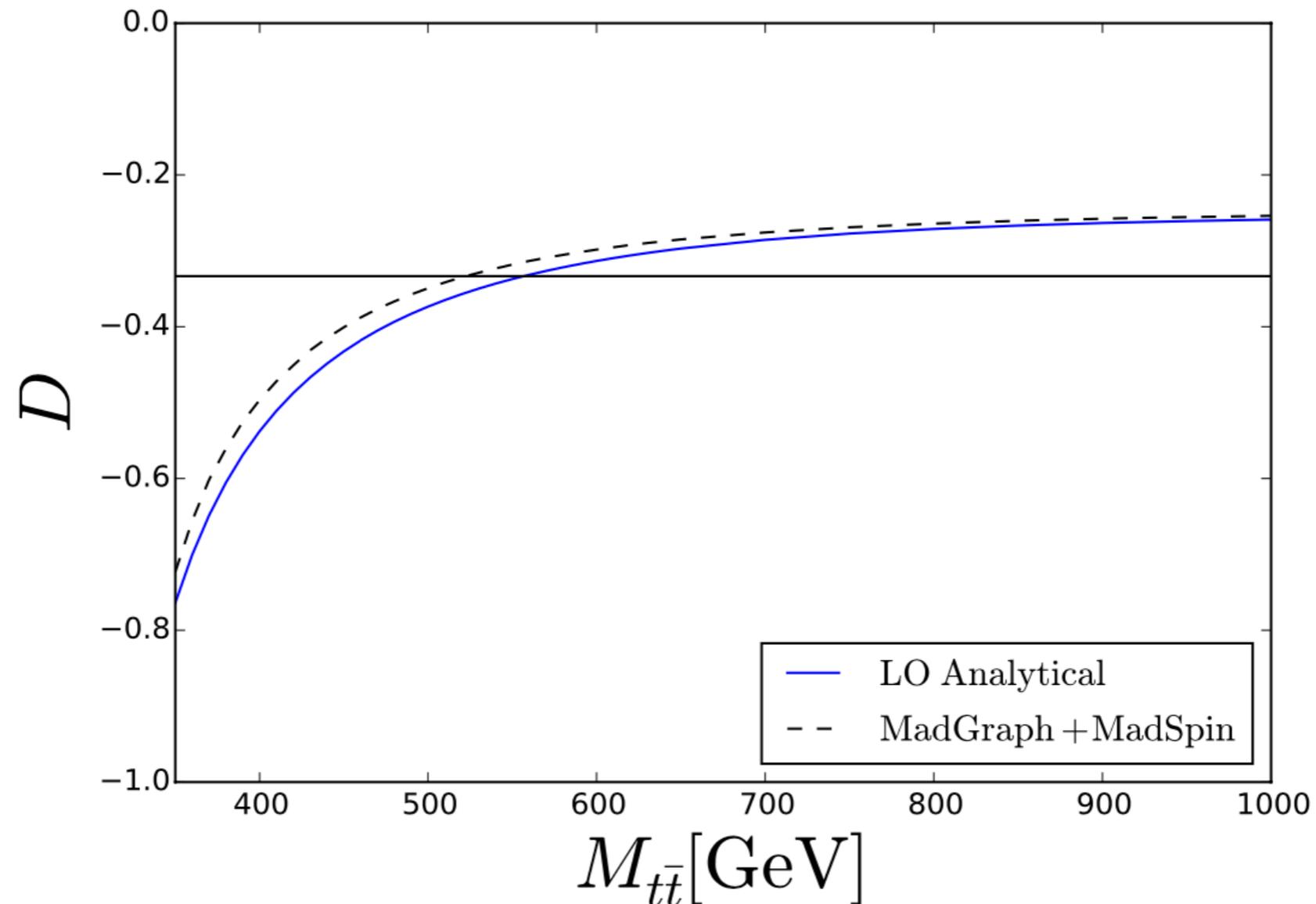
- In order to measure  $D$ , we need to fully reconstruct both **tops** (we need measure  $\cos(\Phi)$  in parent top rest frames).  
➔ This means somehow dealing with two neutrinos
- There are a number of methods to achieve this, but this measurements relies heavily on the “Ellipse method”.



- Employs a geometry approach to analytically solve the system using linear algebra.
- Some other numerical methods used in small number of events.

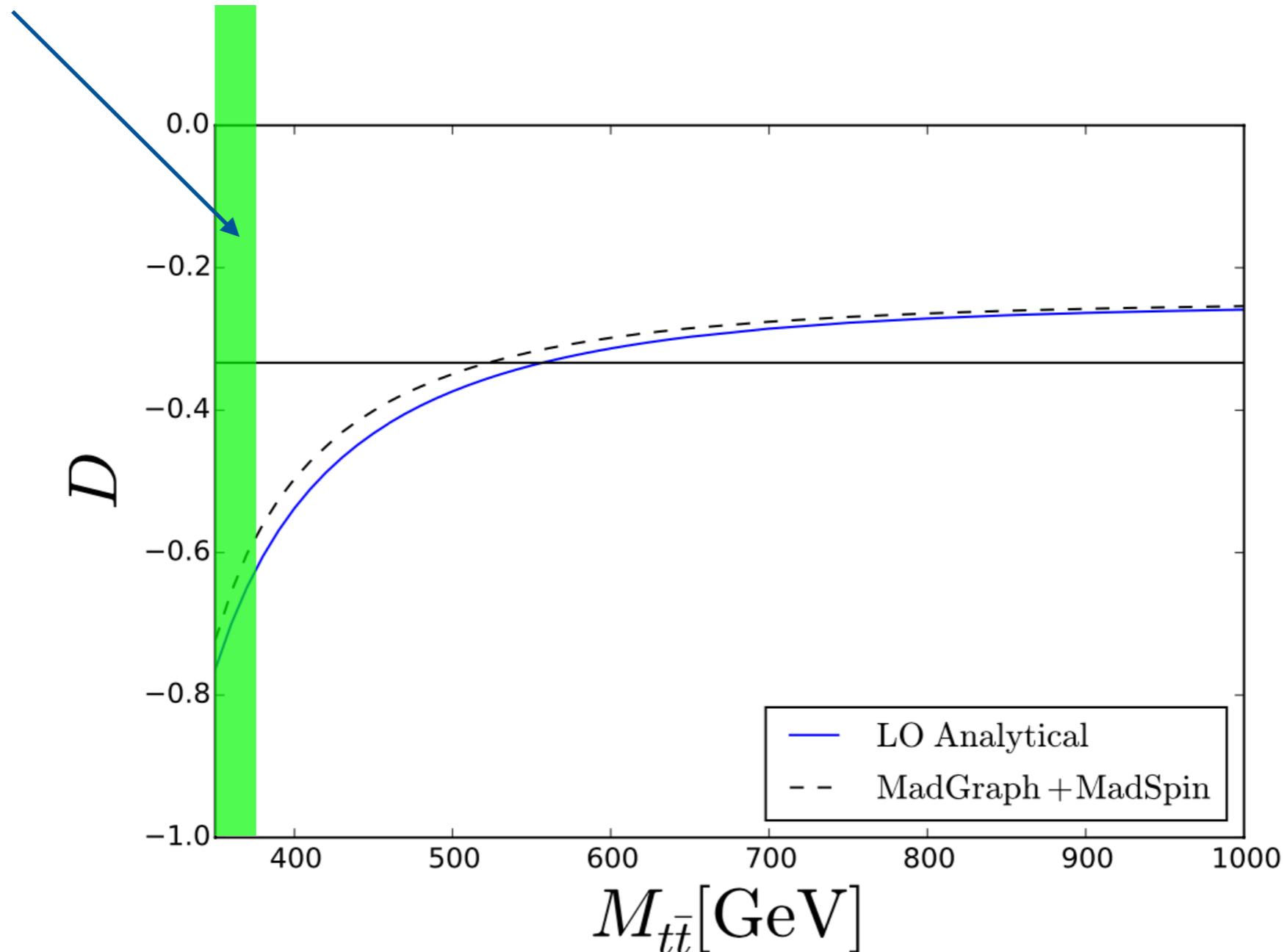
Nucl.Instrum.Meth.A 736 (2014) 169-178

- We split our measurement based on  $m_{t\bar{t}}$ :



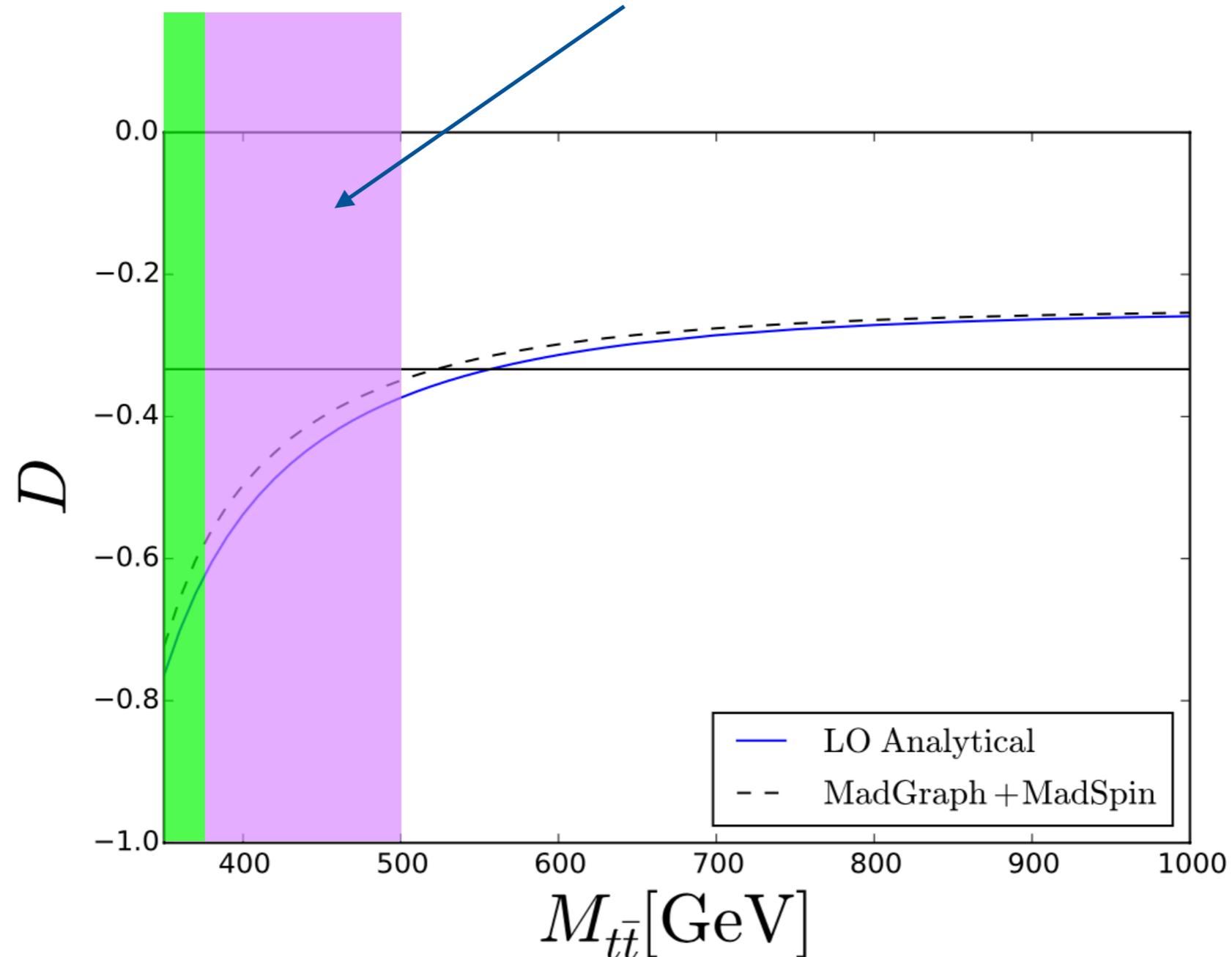
- We split our measurement based on  $m_{t\bar{t}}$ :

SR: 340 - 380 GeV



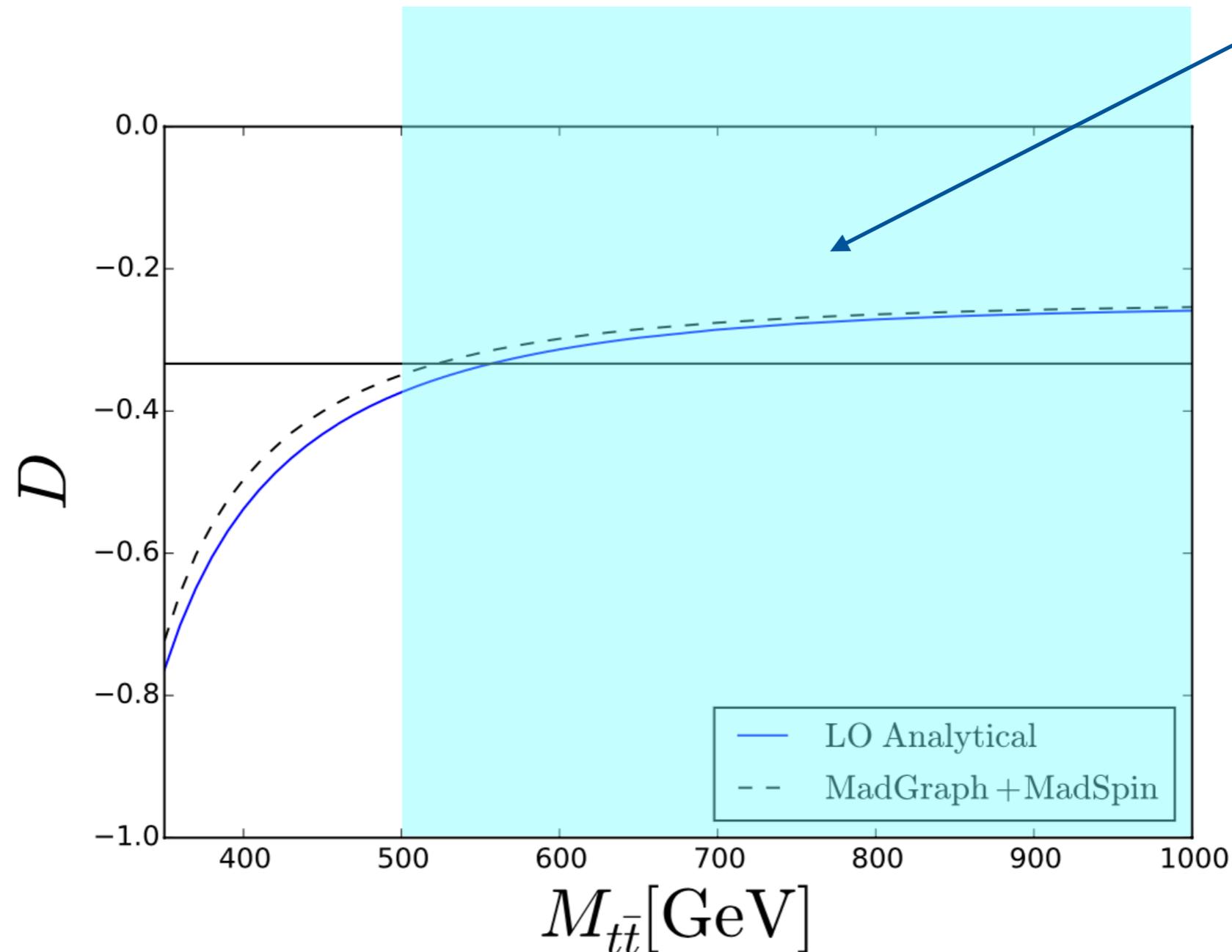
- We split our measurement based on  $m_{t\bar{t}}$ :

**VR1: 380 - 500 GeV**



- We split our measurement based on  $m_{t\bar{t}}$ :

**VR2:**  $> 500$  GeV

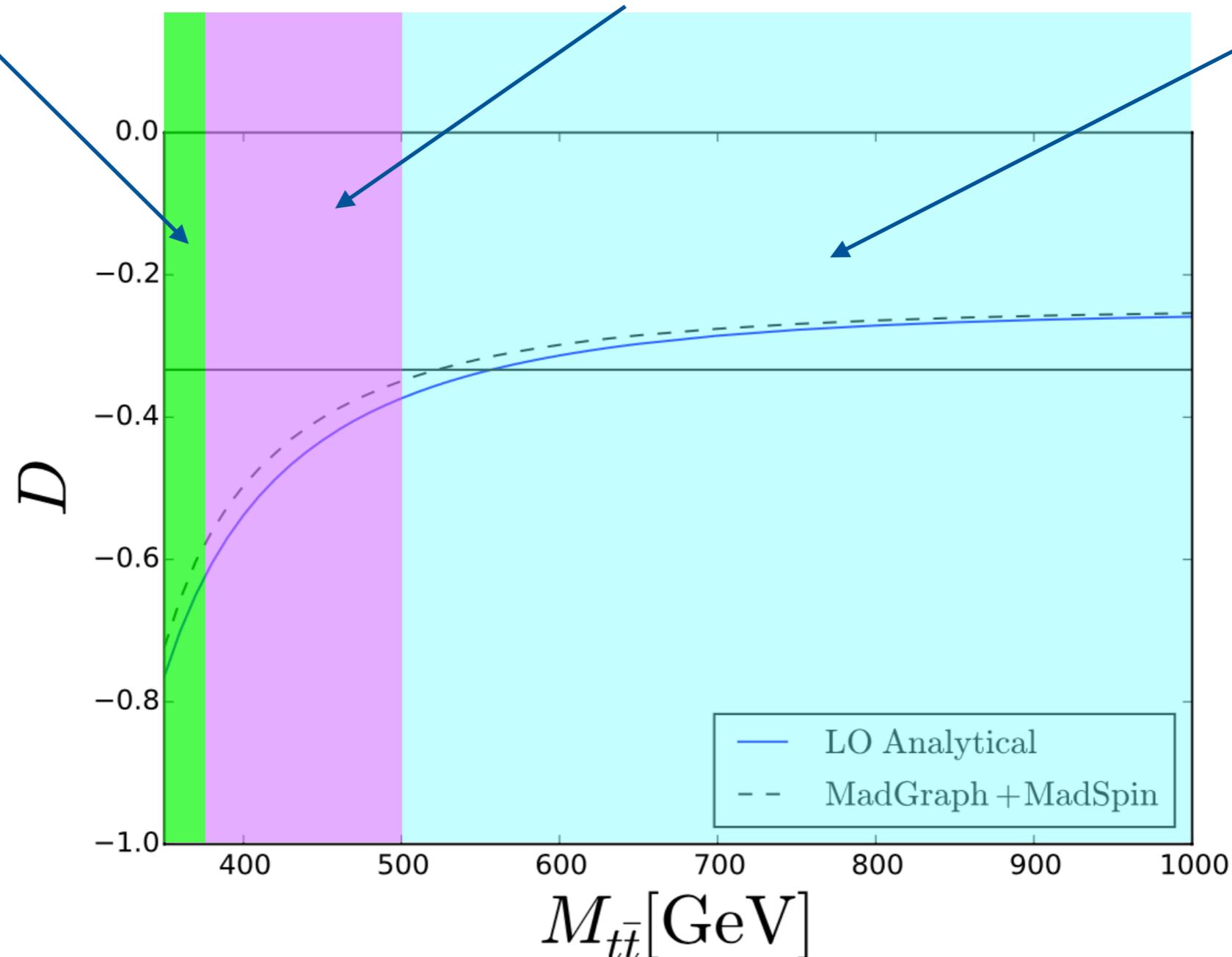


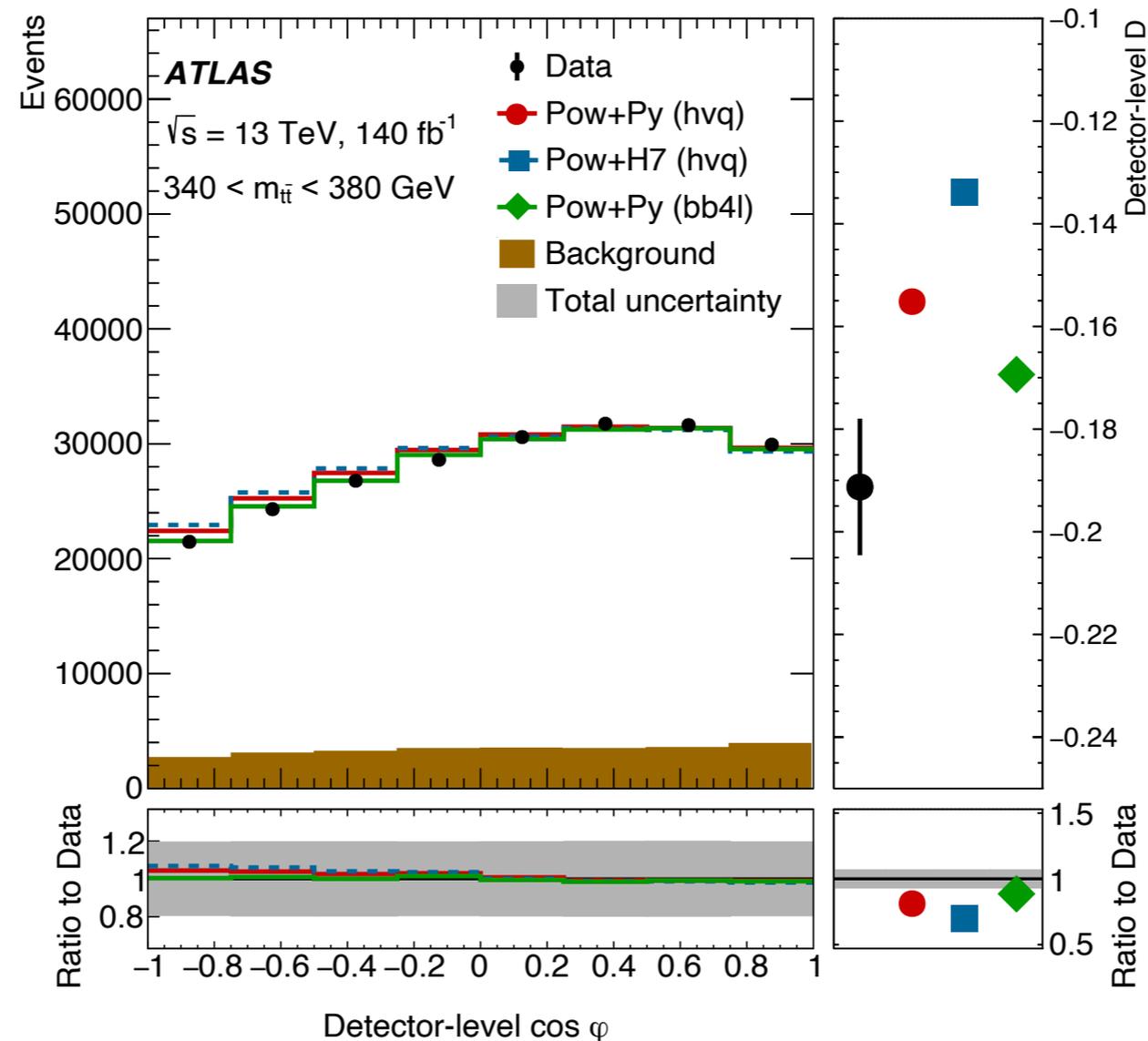
- We split our measurement based on  $m_{t\bar{t}}$ :

**SR:** 340 - 380 GeV

**VR1:** 380 - 500 GeV

**VR2:** > 500 GeV





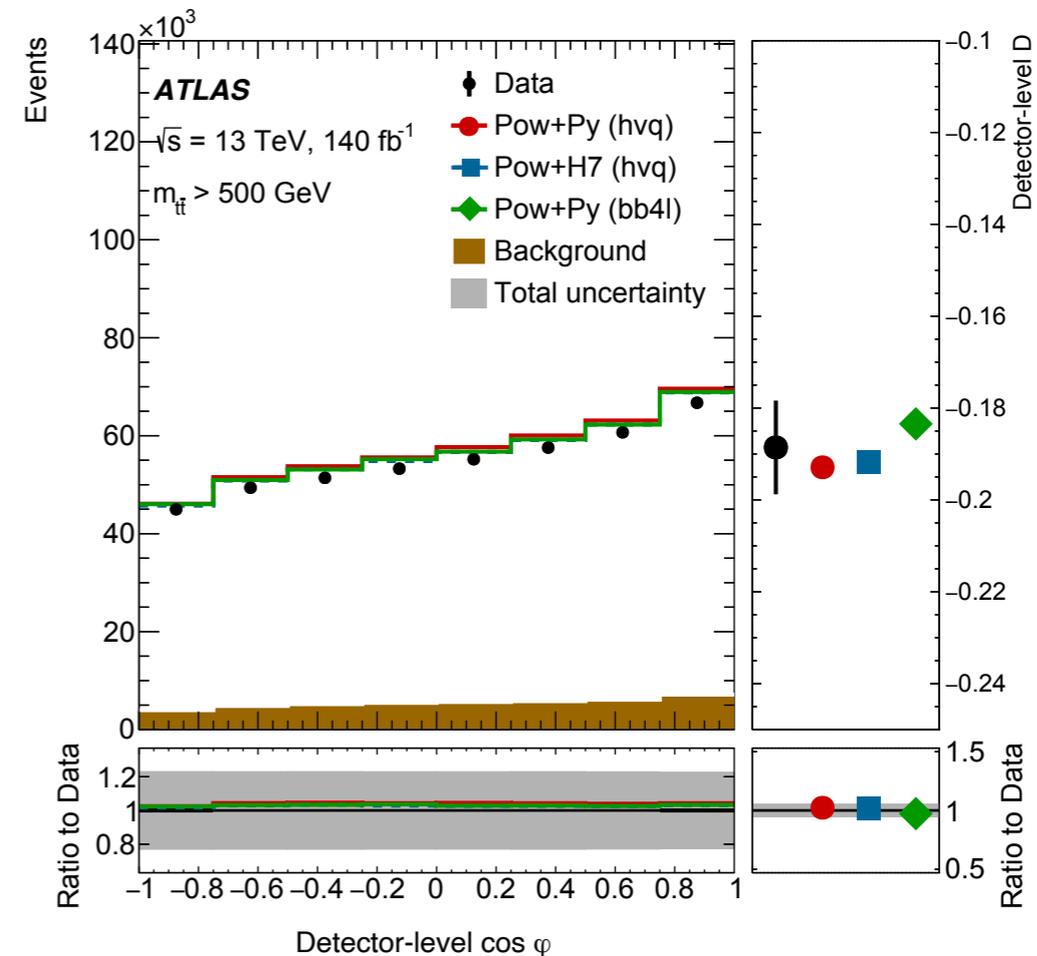
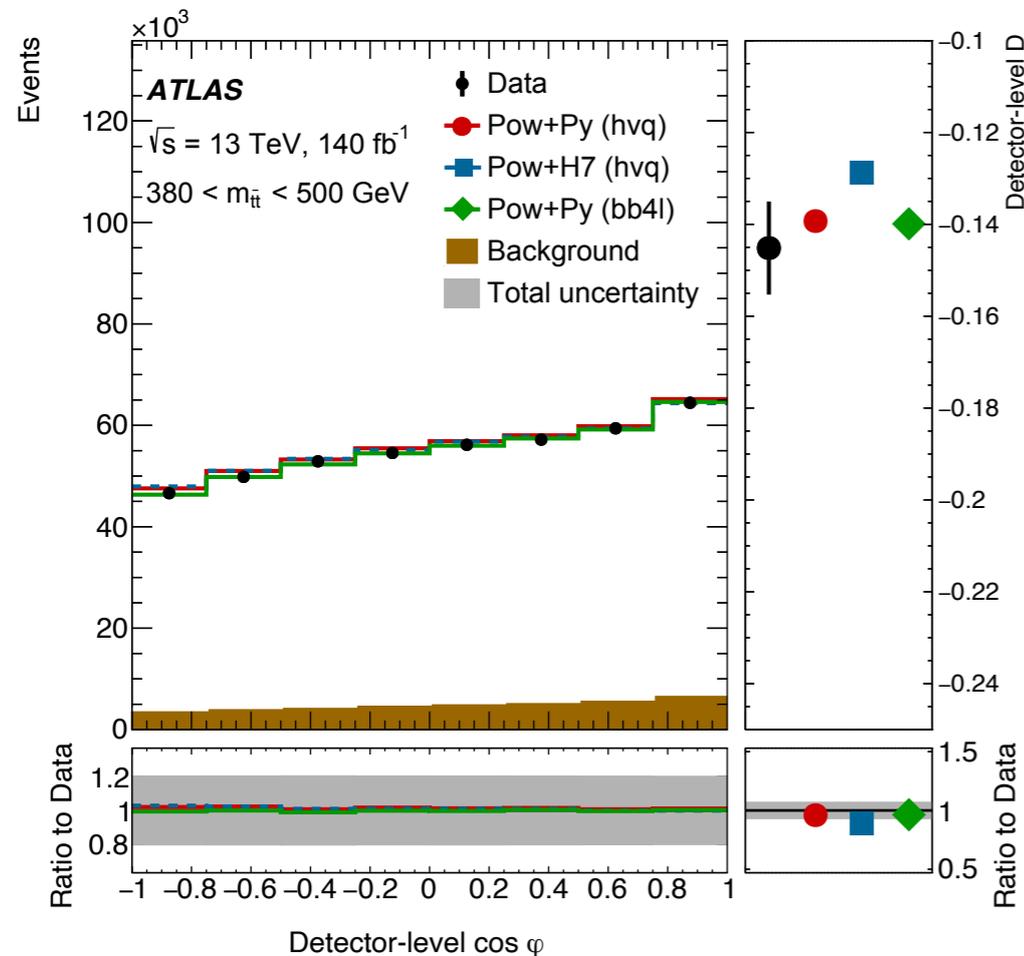
- **Events are selected with exactly 1 electron and 1 muon.**
- **Require 1 or more b-tagged jets (85% W.P):**
  - ➔ loose working point to ensure high stats in signal region.

- Top decay products have “spin analysing power”:

	$b$ -quark	$W^+$	$l^+$	$\bar{d}$ -quark or $\bar{s}$ -quark	$u$ -quark or $c$ -quark
$\alpha_i$ (LO)	-0.410	0.410	1.000	1.000	-0.310
$\alpha_i$ (NLO)	-0.390	0.390	0.998	0.930	-0.310

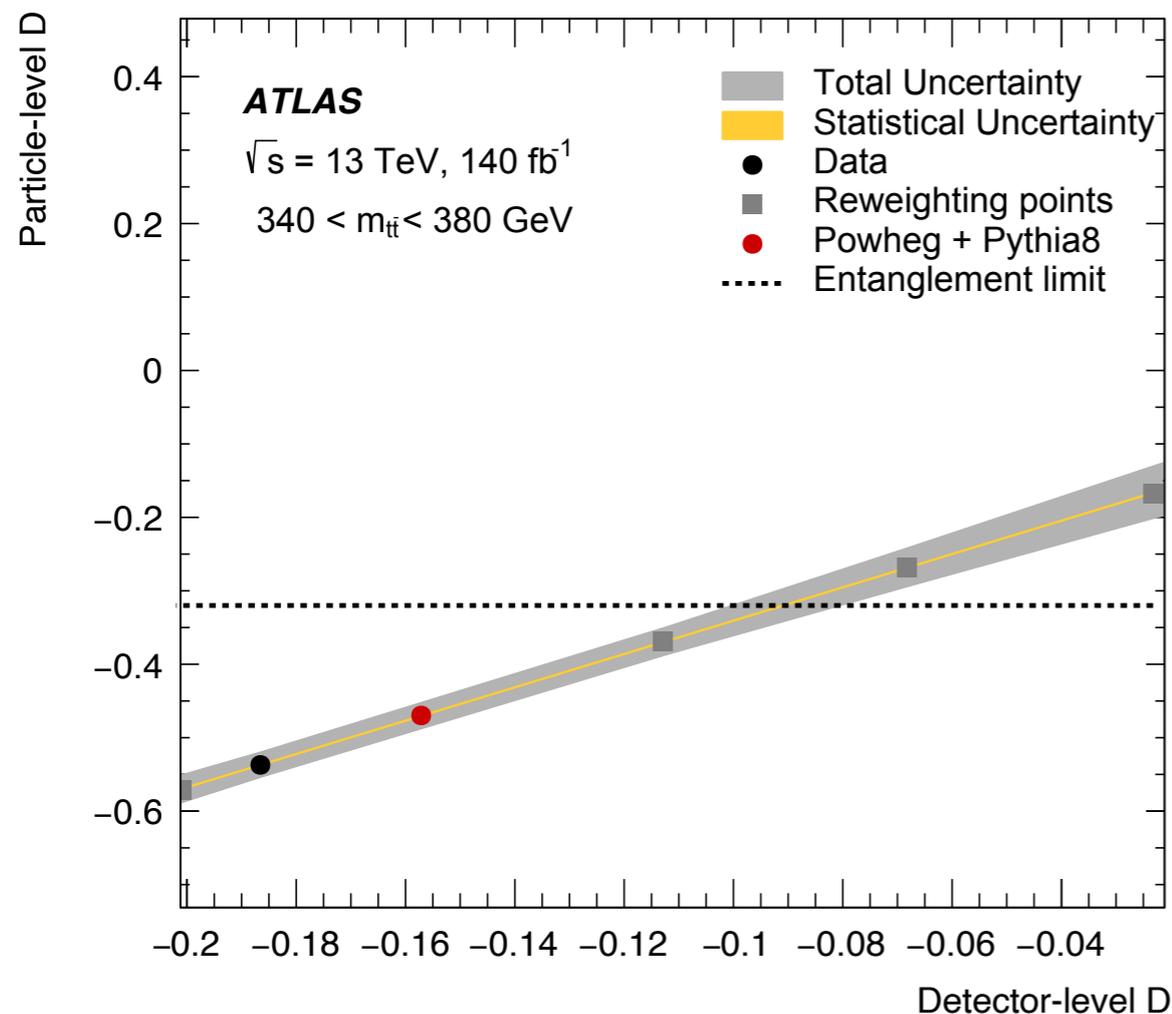
- $\alpha_i = 1$  means a particle carries the full spin information.
- $\alpha_i = 0$  means a particle carries none of the spin info.
- Almost all published spin measurements in top physics use the leptonic decay mode:
  - ➔ easiest to identify experimentally.
- In ~~Run3~~ **now!** we will start to see results using down-type jets.
- Interesting question about implications of these not being exactly 1.

- **This selection is a very robust one** (similar selection used in dozens of analyses).



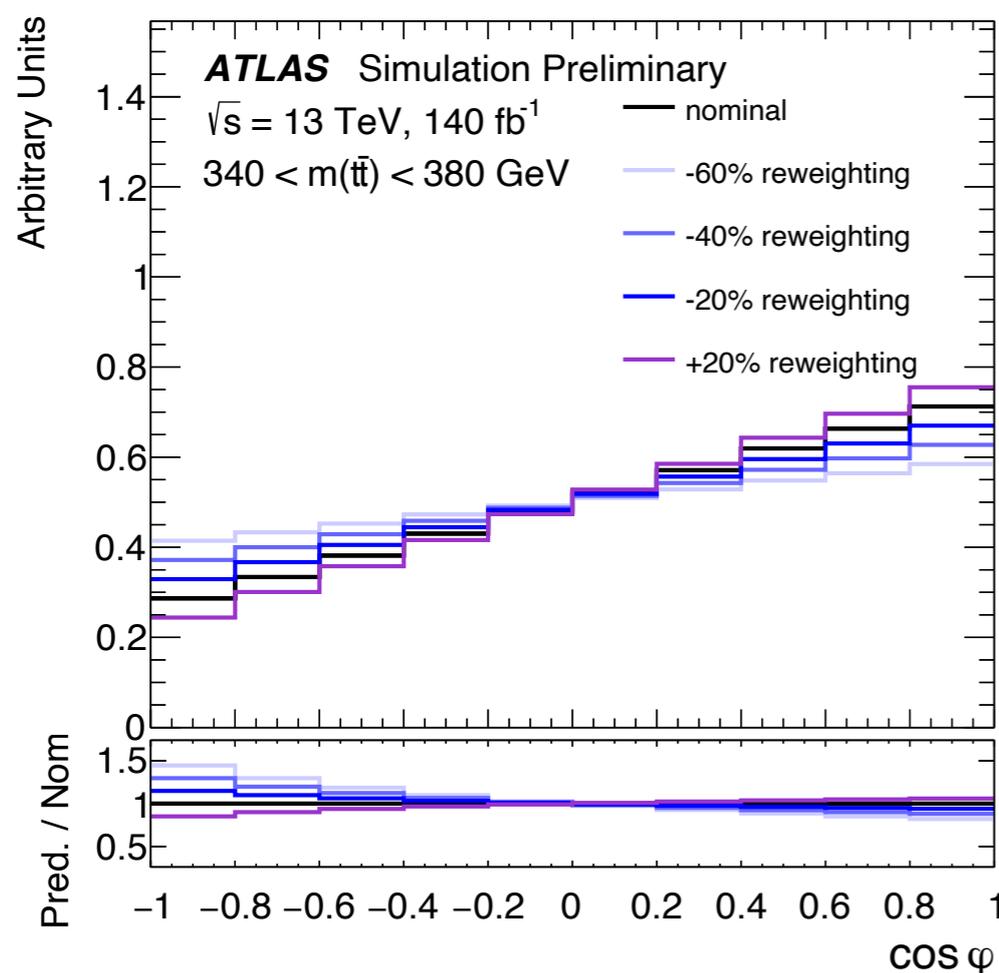
- **Very good overall agreement between the number of signal+background events and the observed number of events in data.**

- We somehow need to correct our observed D for detector effects to some ‘truth’ level (particle in this case):  
➔ We achieve this with a calibration curve.



- To construct this curve we need to change the amount of entanglement in our MC.
- We create 5 hypothesis points corresponding to the SM and 4 different reweighting points:  
(+20%, -20%, -40%, -60%)

- How these alternative hypothesis points are constructed is one of the key points of the measurement.
- We cannot dial entanglement up or down in the MC, so we reweight the  $\cos(\Phi)$  distribution as a function of  $m(t\bar{t})$ .



- If this is not done correctly, the relation:

$$D = \frac{\text{tr}[C]}{3} = -3 \cdot \langle \cos(\phi) \rangle$$

does not hold.

- The method we have used ensures that this relationship remains correct.

- **The relative size of the systematics is not fixed and changes at each hypothesis point:**

Source of uncertainty	$\Delta D_{\text{observed}}(D = -0.537)$	$\Delta D$ [%]	$\Delta D_{\text{expected}}(D = -0.470)$	$\Delta D$ [%]
Signal modeling	0.017	3.2	0.015	3.2
Electrons	0.002	0.4	0.002	0.4
Muons	0.001	0.2	0.001	0.1
Jets	0.004	0.7	0.004	0.8
<i>b</i> -tagging	0.002	0.4	0.002	0.4
Pile-up	< 0.001	< 0.1	< 0.001	< 0.1
$E_{\text{T}}^{\text{miss}}$	0.002	0.4	0.002	0.4
Backgrounds	0.005	0.9	0.005	1.1
Total statistical uncertainty	0.002	0.3	0.002	0.4
Total systematic uncertainty	0.019	3.5	0.017	3.6
Total uncertainty	0.019	3.5	0.017	3.6

- **As with most top measurements, we are limited by signal modelling** (also note that the relative uncertainty depends on *D*).

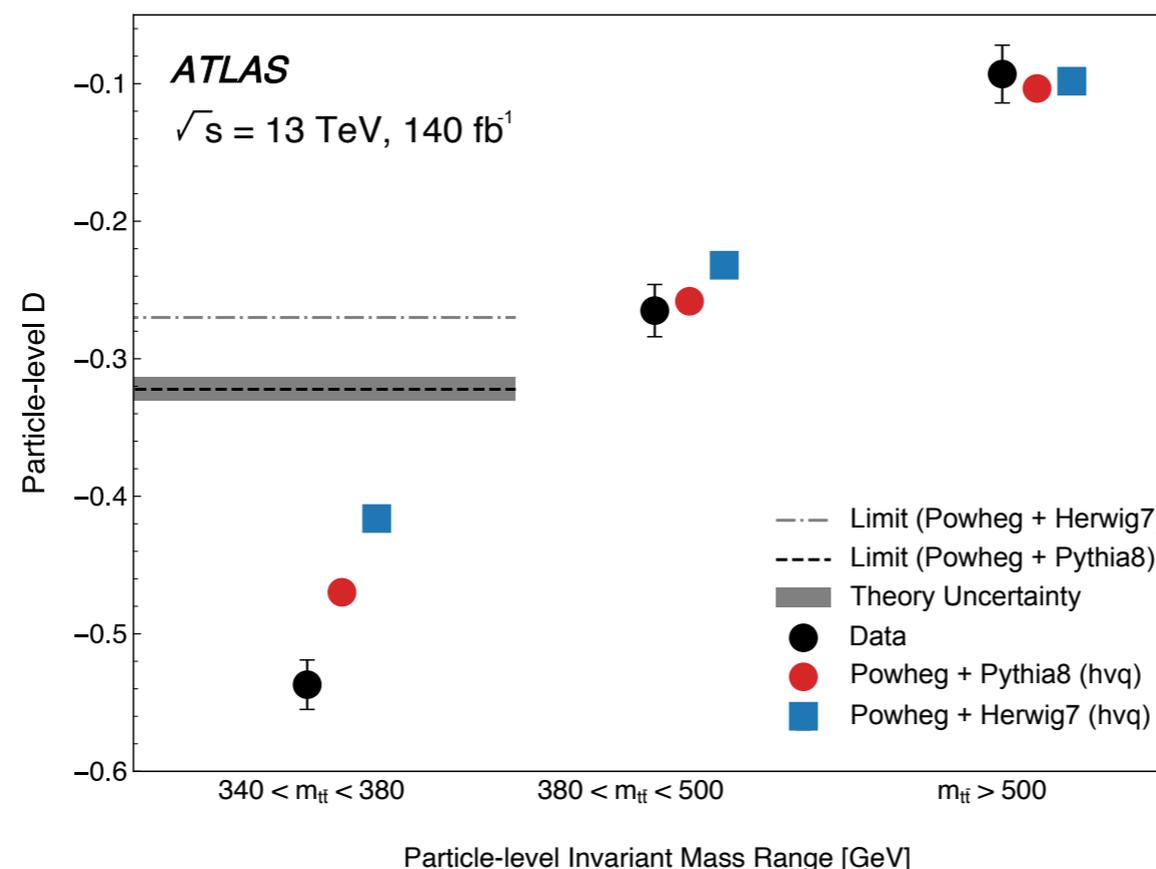
- **We have a large suite of MC modelling related systematic uncertainties:**

Systematic uncertainty source	Relative size (for SM $D$ value)
Top-quark decay	1.6%
Parton distribution function	1.2%
Recoil scheme	1.1%
Final-state radiation	1.1%
Scale uncertainties	1.1%
NNLO reweighting	1.1%
$p_{T\text{hard}}$ setting	0.8%
Top-quark mass	0.7%
Initial-state radiation	0.2%
Parton shower and hadronization	0.2%
$h_{\text{damp}}$ setting	0.1%

- **Colour reconnection, string vs cluster fragmentation, spin correlation in parton shower, EW shower were all tested but found to be negligible effects.**

- The observed (expected) results are:

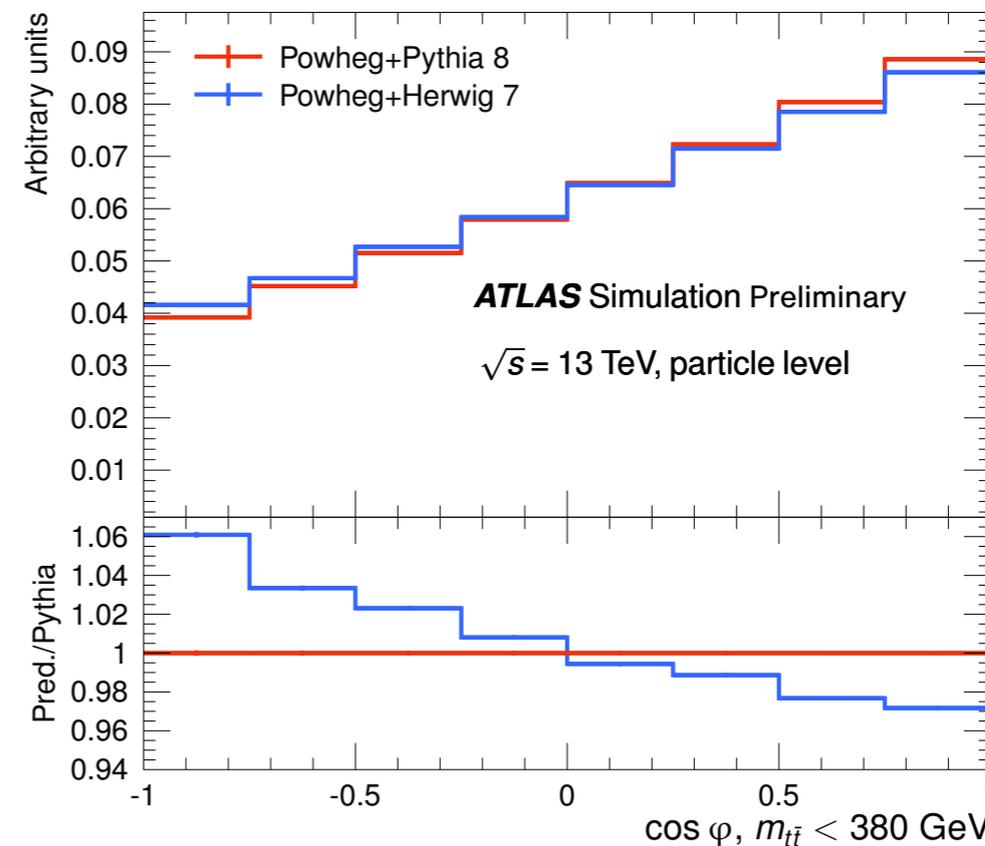
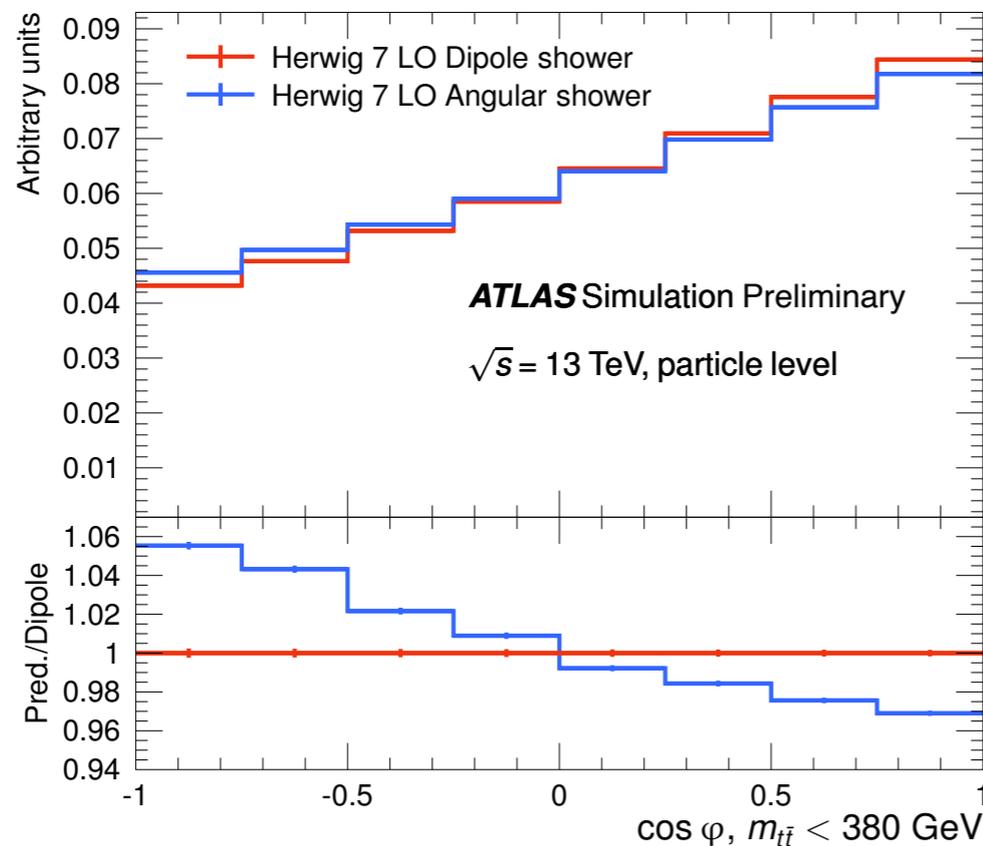
**SR**  $D = -0.537 \pm 0.002$  [stat.]  $\pm 0.018$  [syst.] ( $-0.470 \pm 0.002$  [stat.]  $\pm 0.016$  [syst.]),  
**VR1**  $D = -0.265 \pm 0.001$  [stat.]  $\pm 0.019$  [syst.] ( $-0.258 \pm 0.001$  [stat.]  $\pm 0.019$  [syst.]),  
**VR2**  $D = -0.093 \pm 0.001$  [stat.]  $\pm 0.021$  [syst.] ( $-0.103 \pm 0.001$  [stat.]  $\pm 0.021$  [syst.]),



[Nature 633 \(2024\) 542](#)

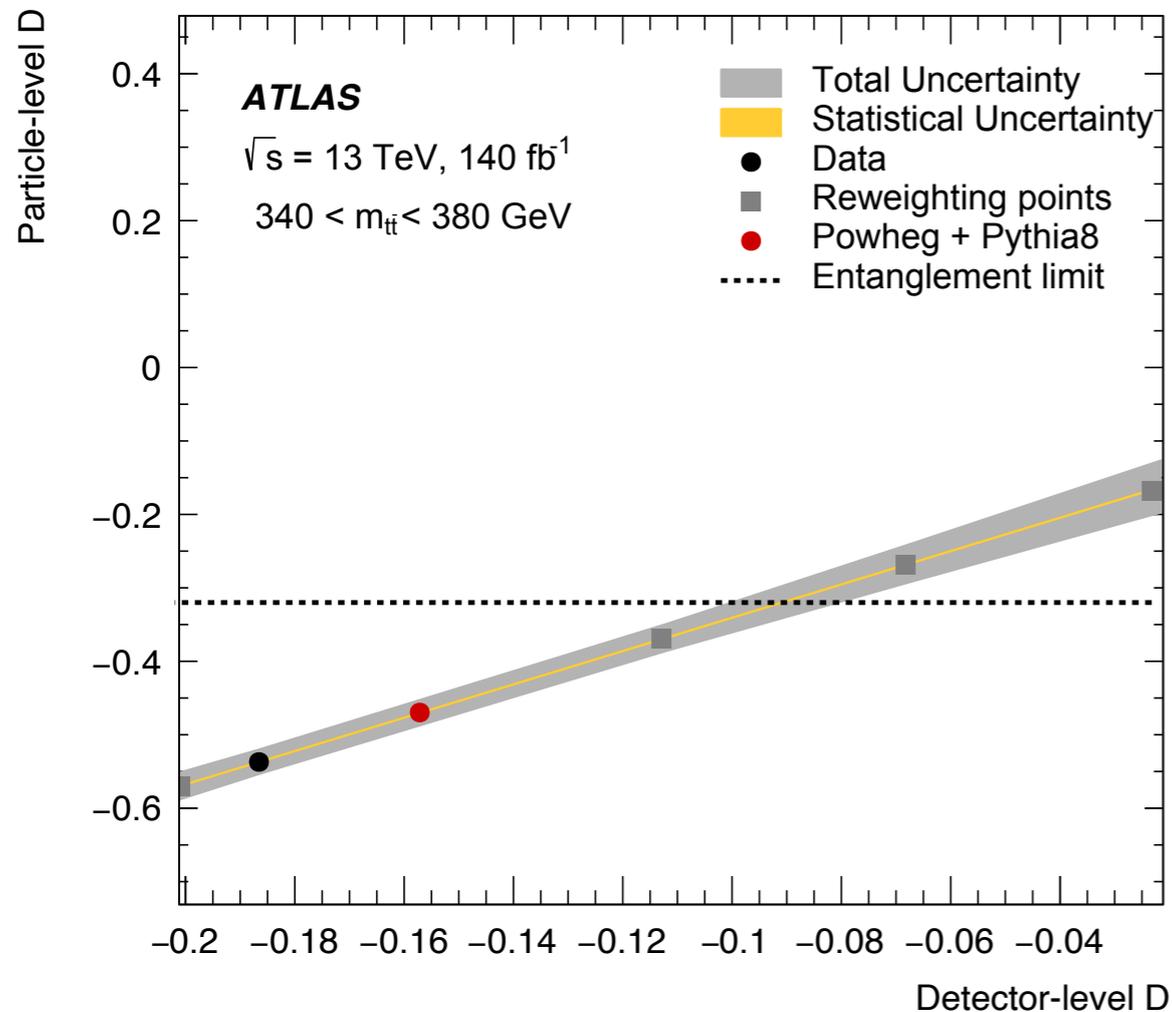
- The observed results excludes the entanglement limit at (much) more than 5 sigma significance.

- **Difference seems to come from the ordering of the shower.**



- **Angular ordered showers have a large effect compared to dipole showers.**
- **Doesn't effect detector corrections significantly.**

- This difference between parton showers IS included in the calibration curve!



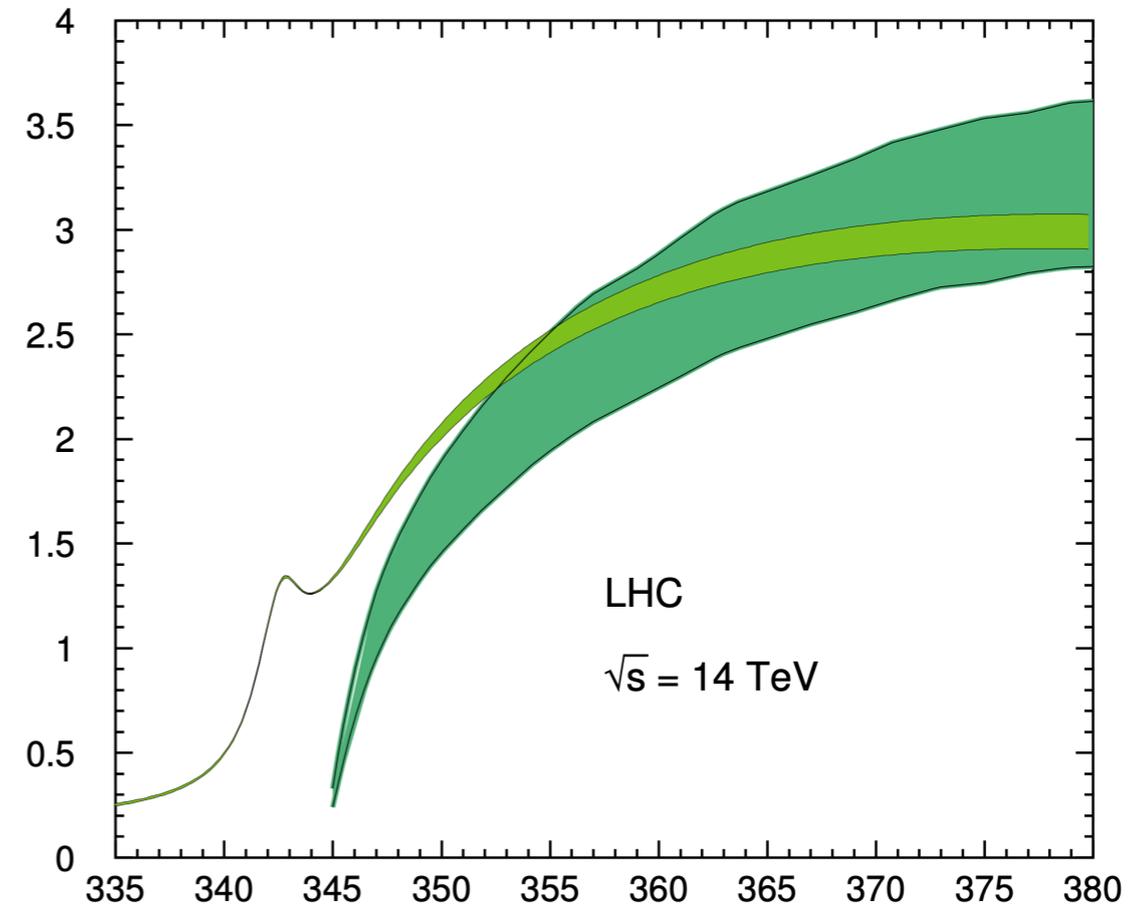
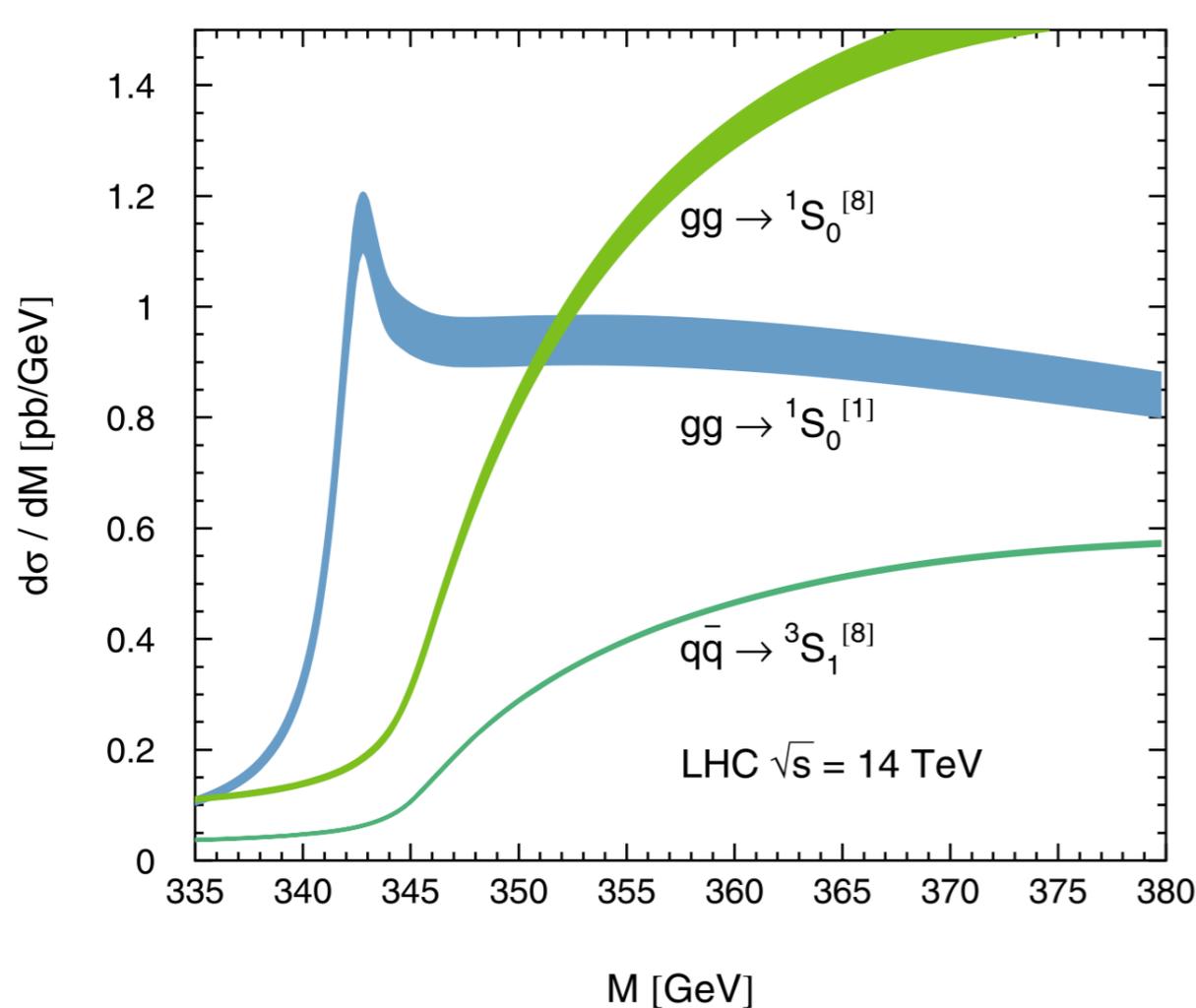
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Parton shower and hadronization	0.2%
$h_{\text{damp}}$ setting	0.1%

- Big differences in prediction don't necessarily mean large detector correction effects.

# What about 'Topponium'?

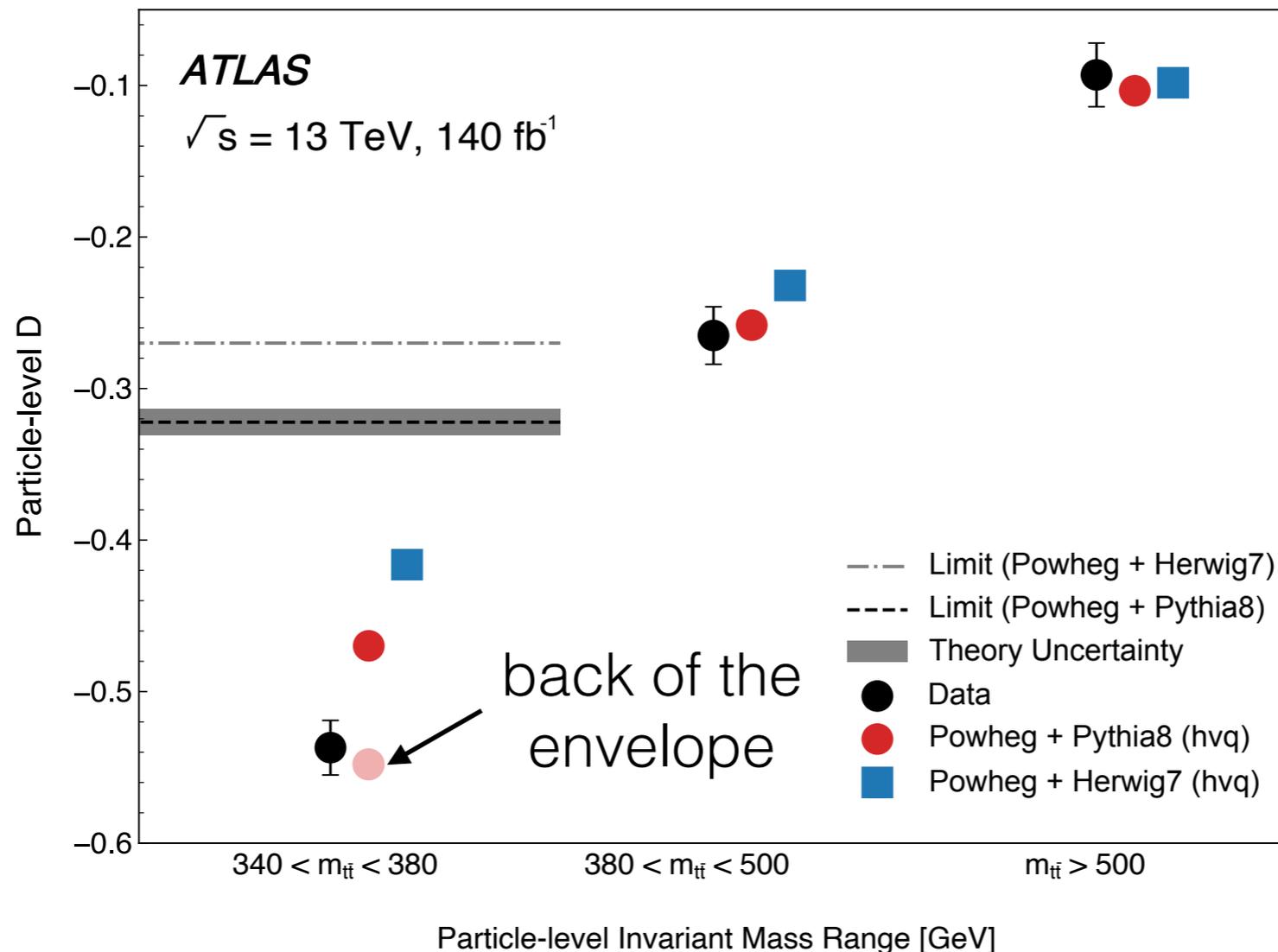
- **Bound state effects are most prevalent in the region that we care about.**

Kiyo, Kühn, Moch, Steinhauser, Uwer, 2009



- **These are not directly included in our MC simulations** (but we have attempted to introduce them as a cross-check and other uncertainties cover similar effects).

- **Effect on data correction is  $\sim 0.5\%$**  (adding it into the total uncertainty doesn't change the error within the precision we quote).



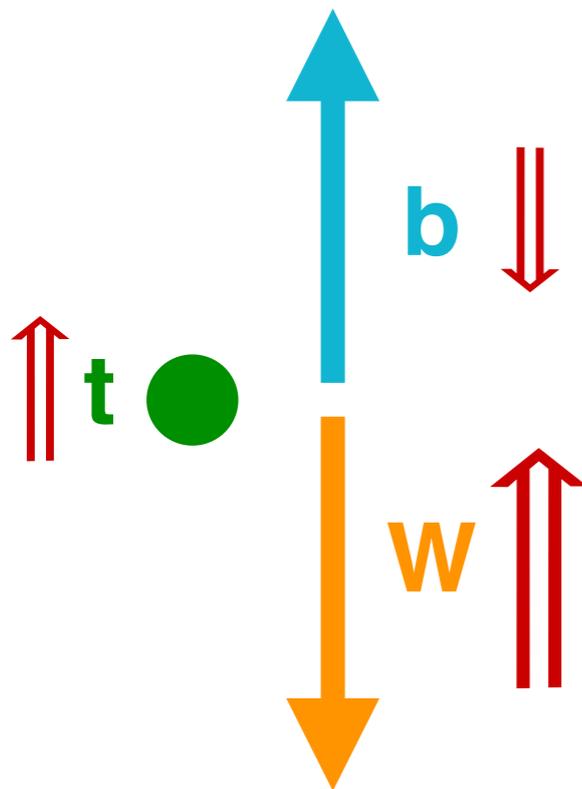
- **If added to predictions, would move them closer to data** (but not clear by how much as we cannot isolate the spin singlet).

- **ATLAS has observed quantum entanglement for the first time in a pair of fundamental quarks, at the highest lab-made energies.**
- **ATLAS has not made any claims about Bell operators or locality.**
- **This is the first step in a program to use the LHC as a tool for exploring quantum information.**
- **Important questions about how entanglement (and spin correlation) is modelled in this threshold region:**
  - ➔ **Would be a very profitable area for further study in the theory community!**

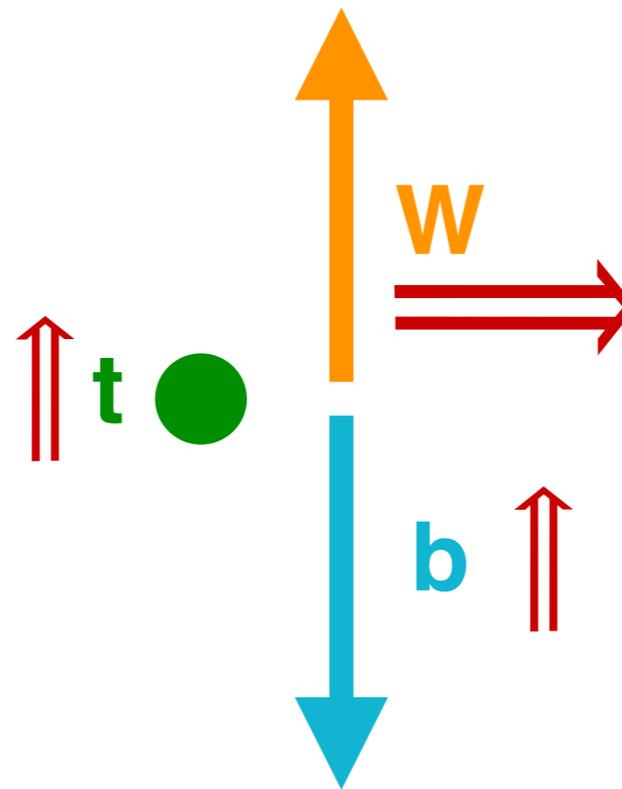
# Backup

- **W bosons act as their own polarimeters**

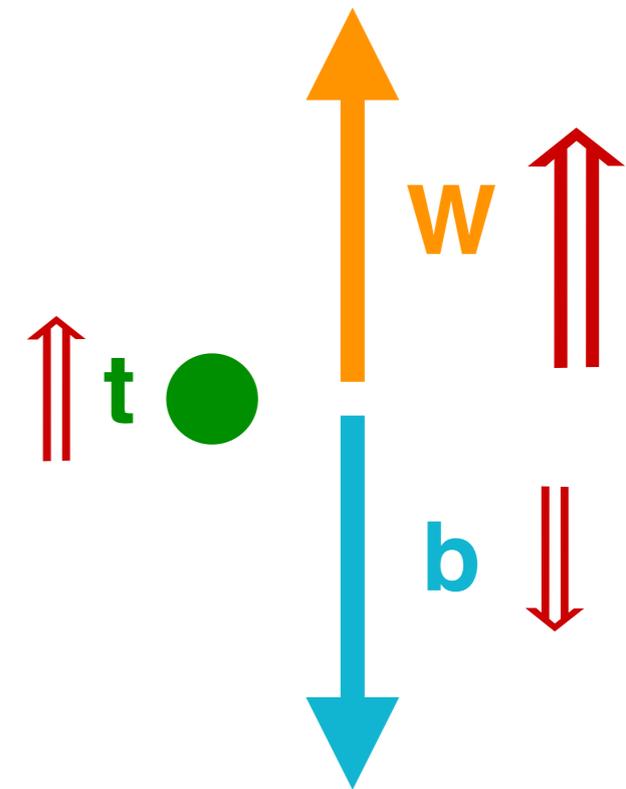
Left-handed W



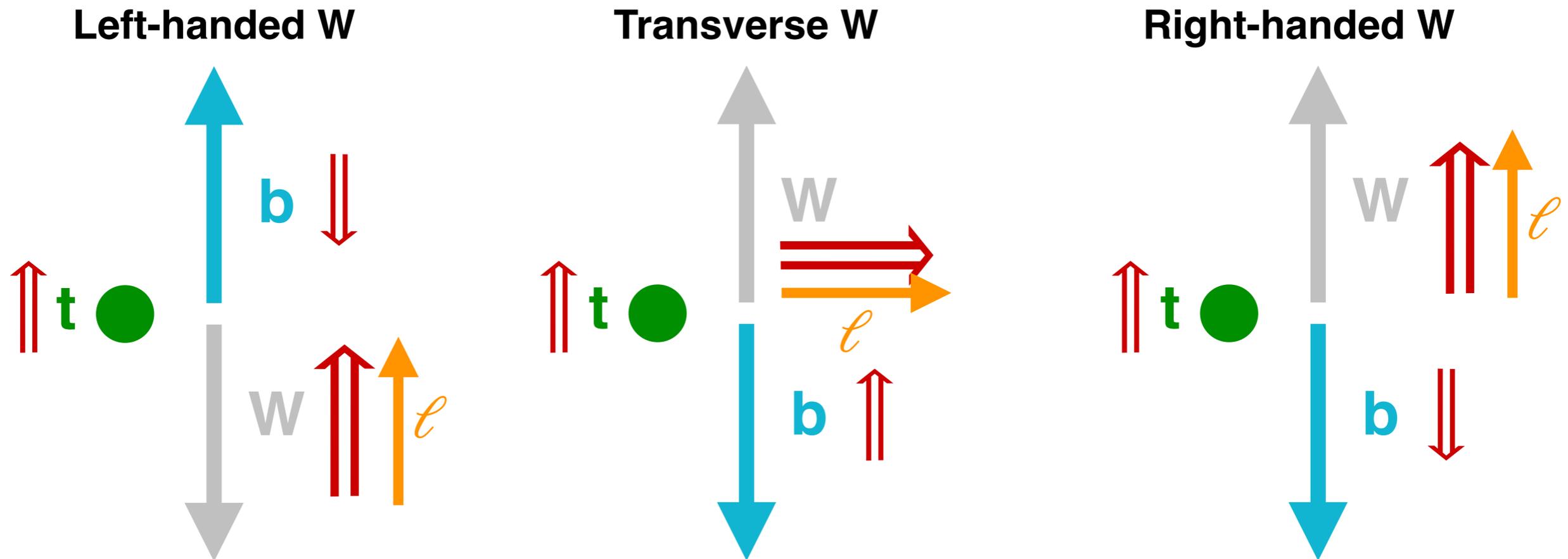
Transverse W



Right-handed W

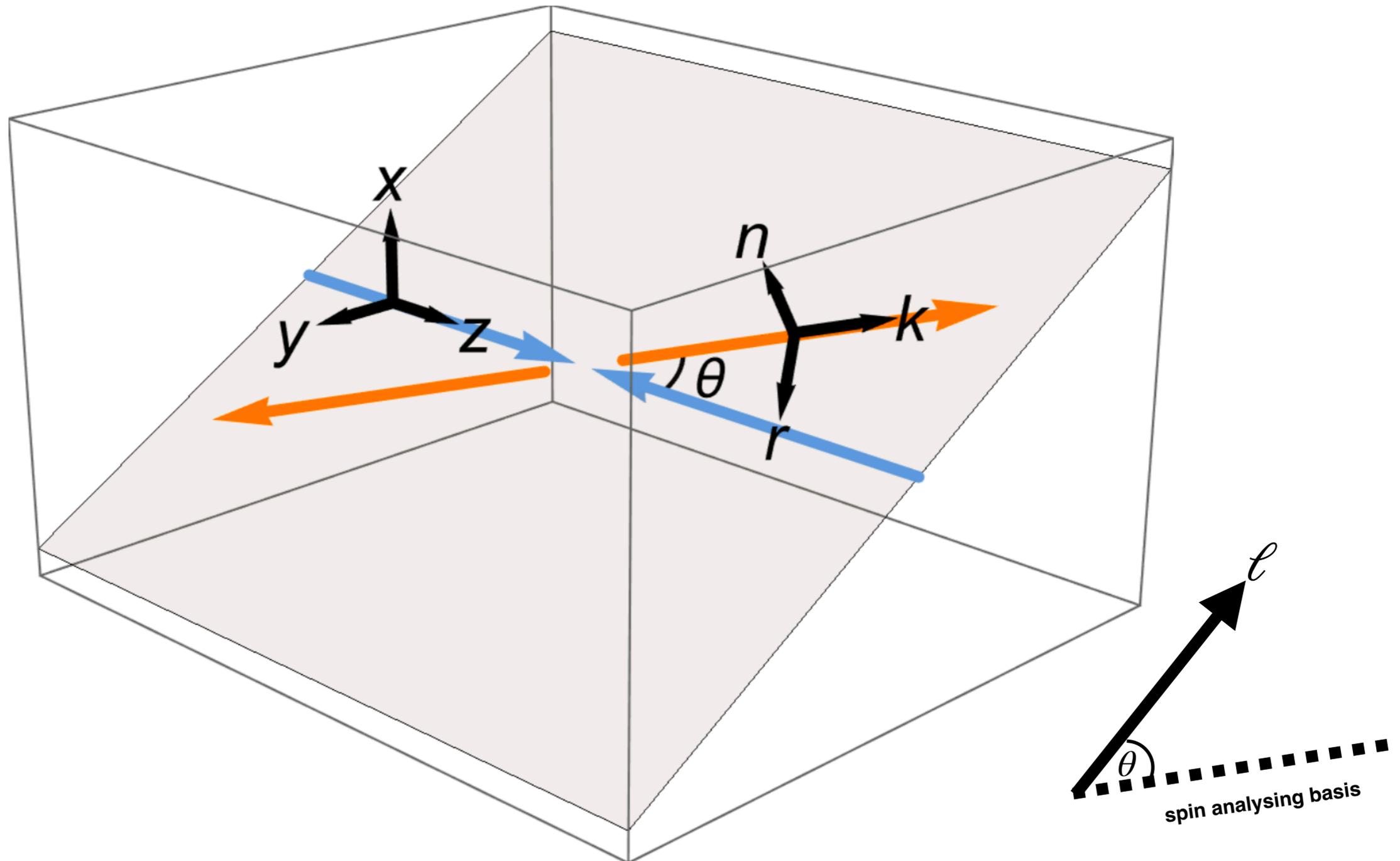


- **W bosons act as their own polarimeters**



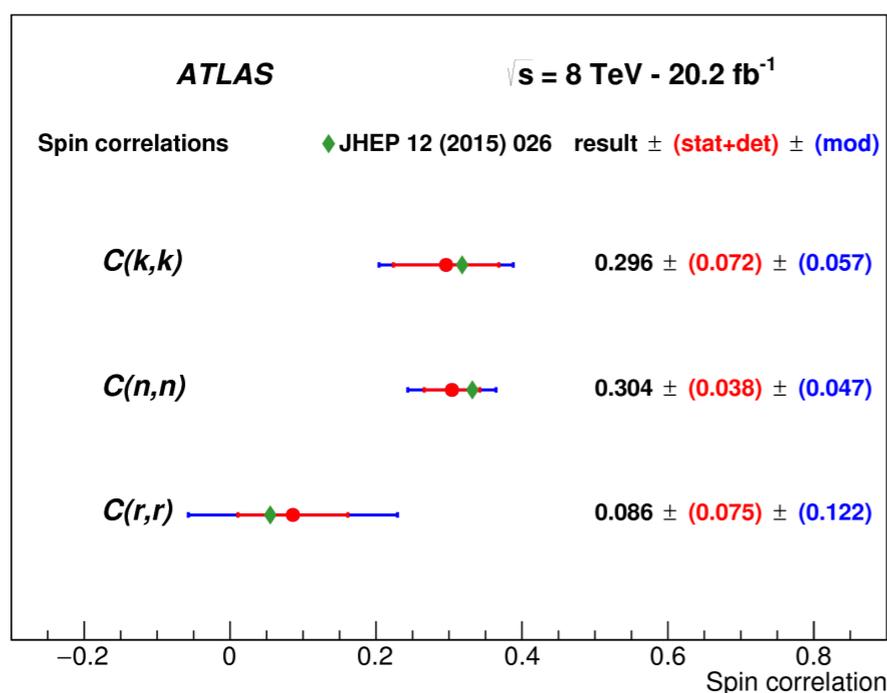
- **Their down-type decay particle momenta always points in the direction of their spin!**

- It matters how you measure these angles!

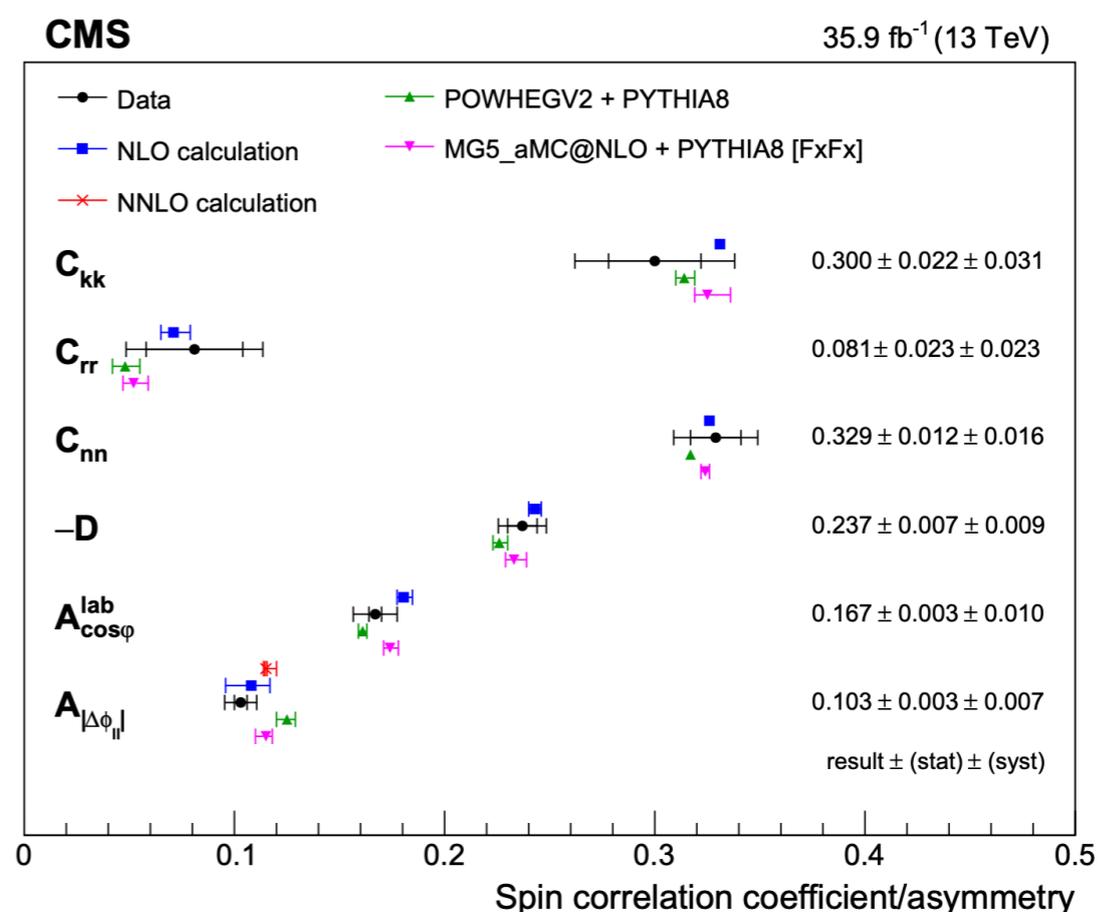


- You can build spin sensitive observables and measure them in data:

JHEP 03 (2017) 113



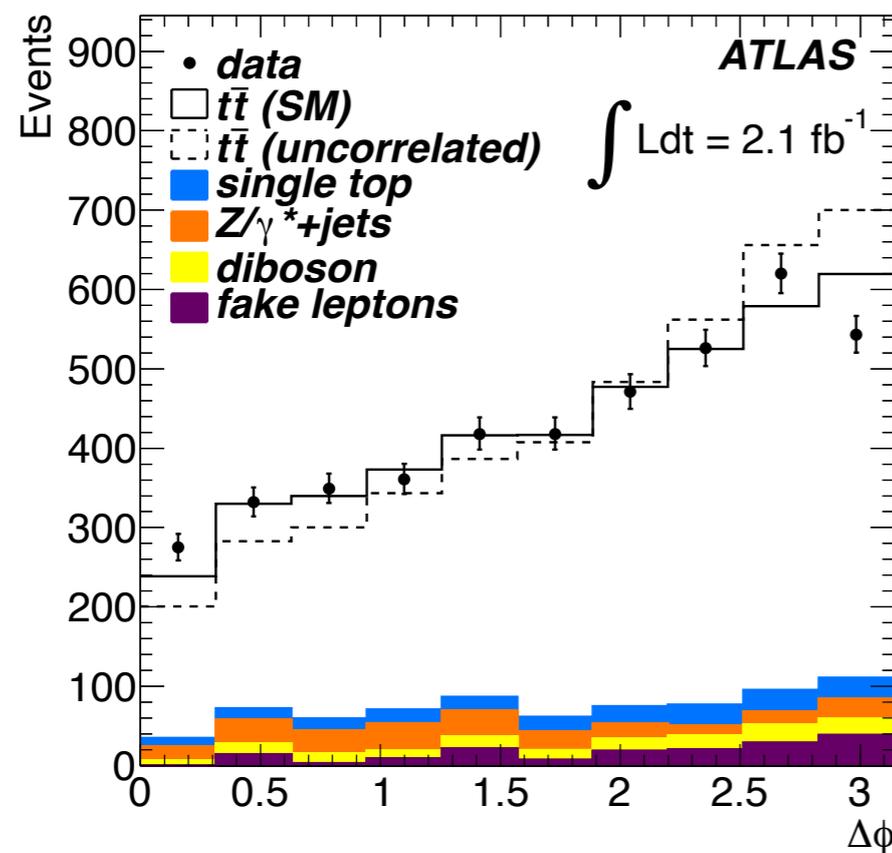
Phys. Rev. D 100, 072002 (2019)



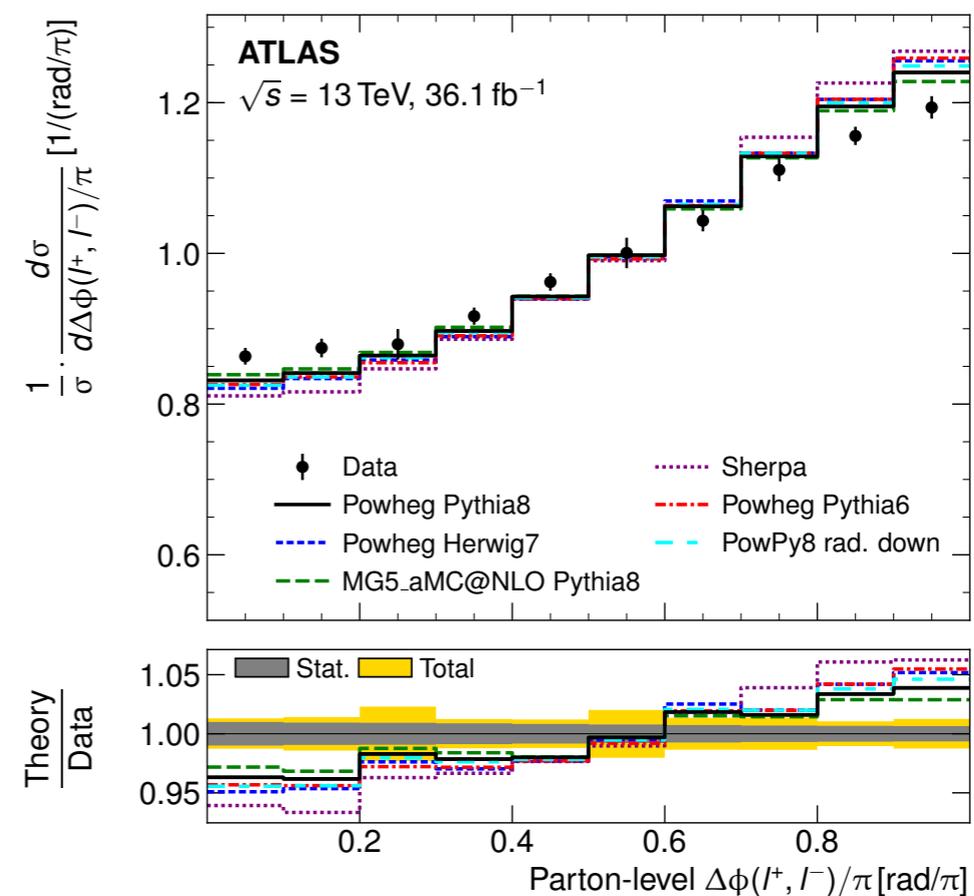
- Many more observables with interesting symmetry structures and BSM potential (ask me if you're interested).

- An easy lab-frame observable that you can build is the  $\Delta\phi$  between two leptons in dilepton events:

Phys. Rev. Lett. 108 (2012) 212001



Eur. Phys. J. C 80 (2020) 754



- Was used to discover spin correlation in tops, to exclude light stops, and currently has a  $3\sigma$  tension with the SM.

- **It's important to note that:**

**Entanglement  $\not\Rightarrow$  Bell Inequality Violation**

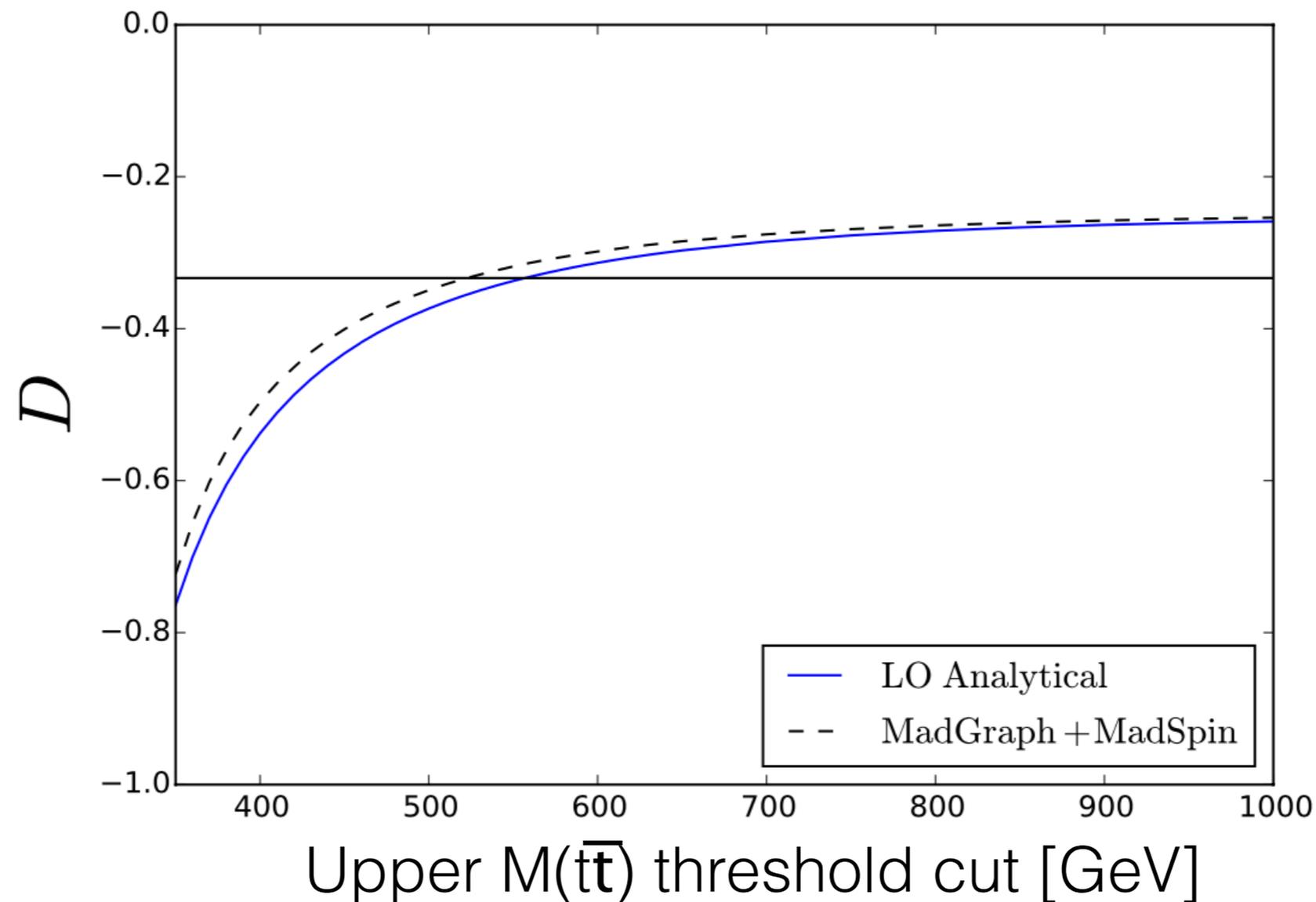
- **Something can be quantum entangled but not strongly enough measure but not to violate a Bell inequality.**

- The goal of the ATLAS measurement is to measure:

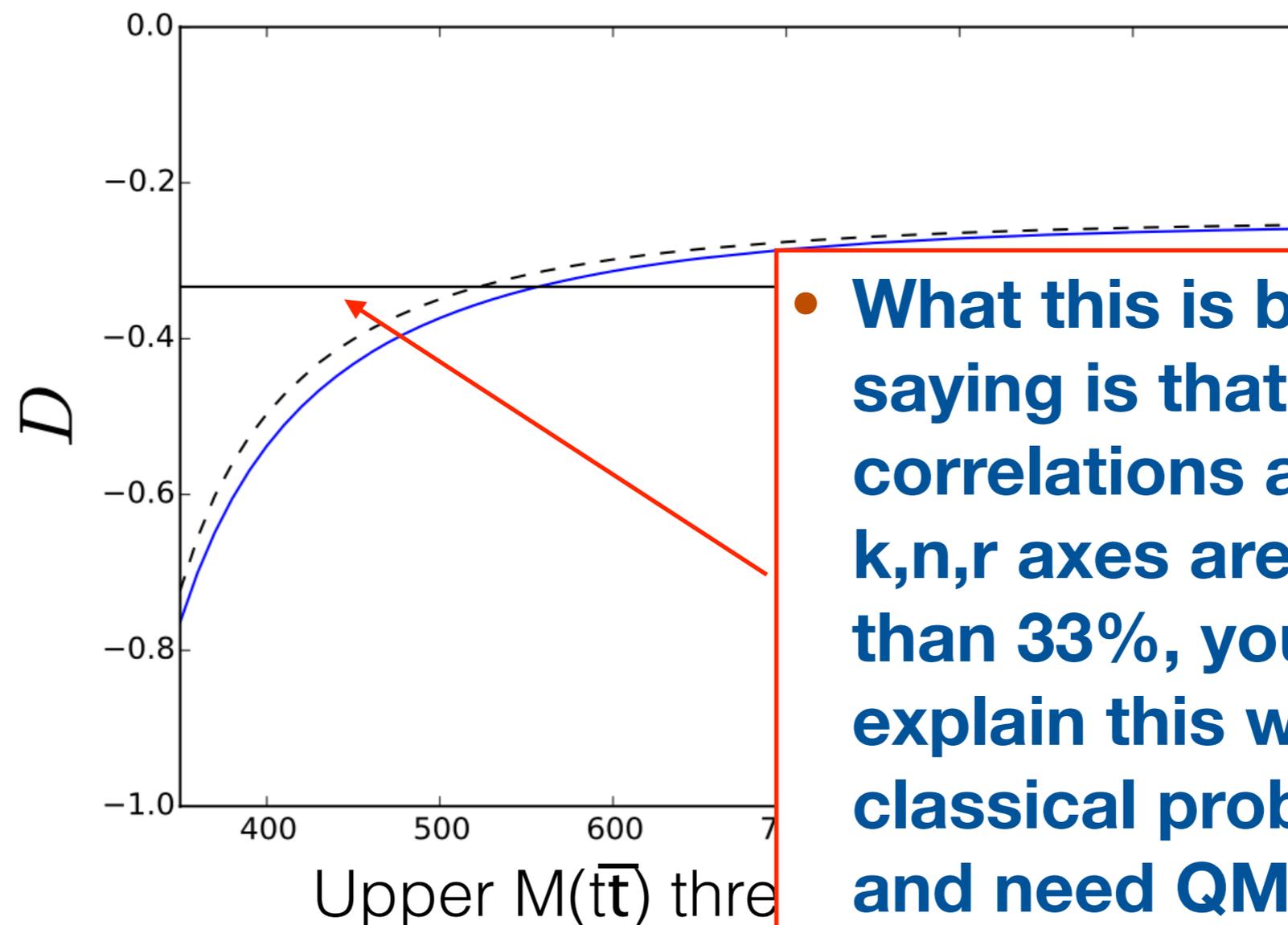
$$D = \frac{\text{tr}[C(k, n, r)]}{3} = -3 \cdot \langle \cos(\phi) \rangle$$

- Where  $\cos(\phi)$  is the dot product of the top spin analysers in their parent top rest frames.
- An observation of  $D < -1/3$  is a sufficient condition to **claim entanglement in  $t\bar{t}$  pairs** (equivalently, that their density matrices are not factorable).
- ATLAS has measured this  $D$  in  $t\bar{t}$  events using  $140 \text{ fb}^{-1}$  of 13 TeV data.

- The primary experimental challenges in this result are to reconstruct the tops with sufficient sensitivity to isolate the threshold region where tops are entangled.

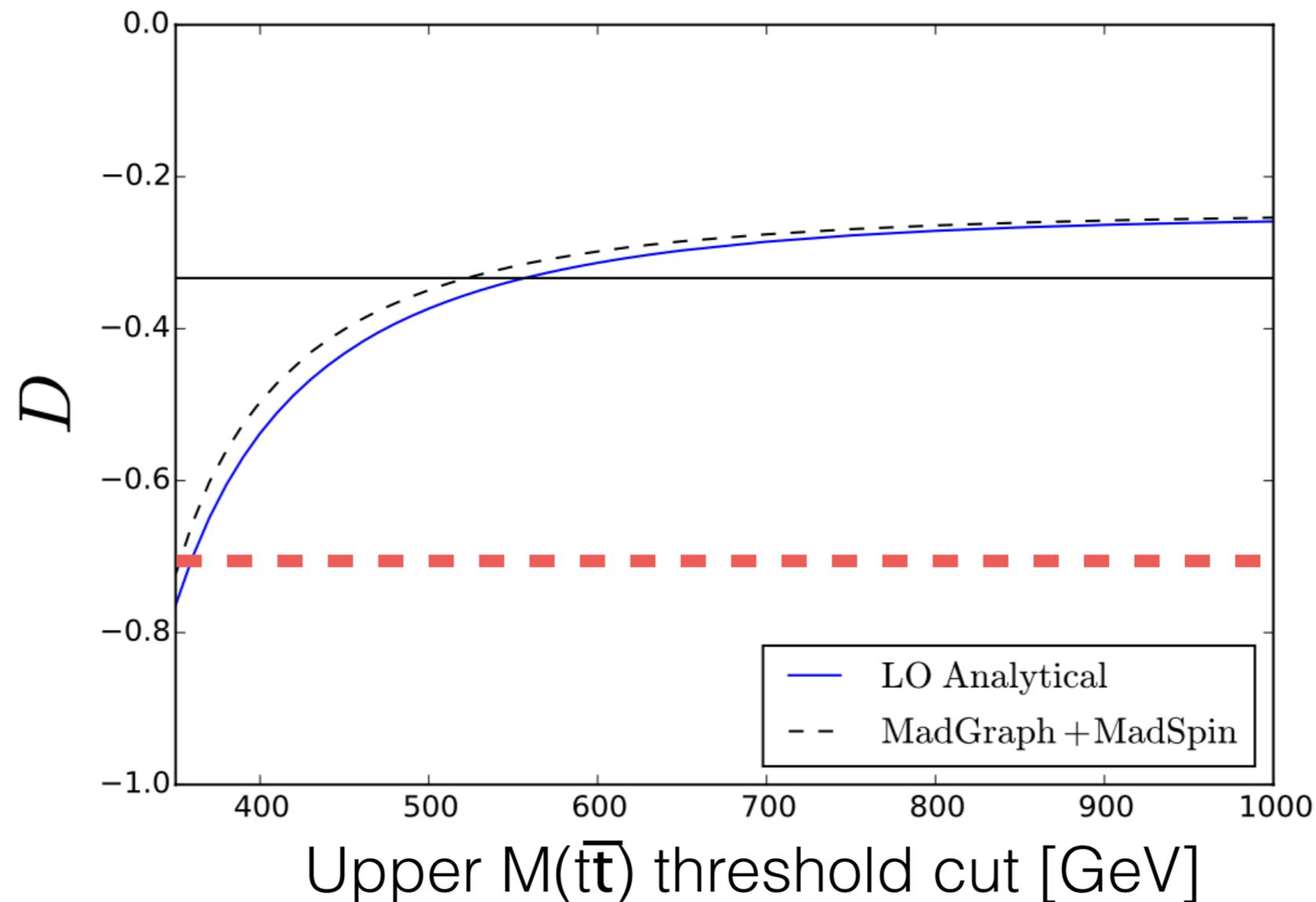


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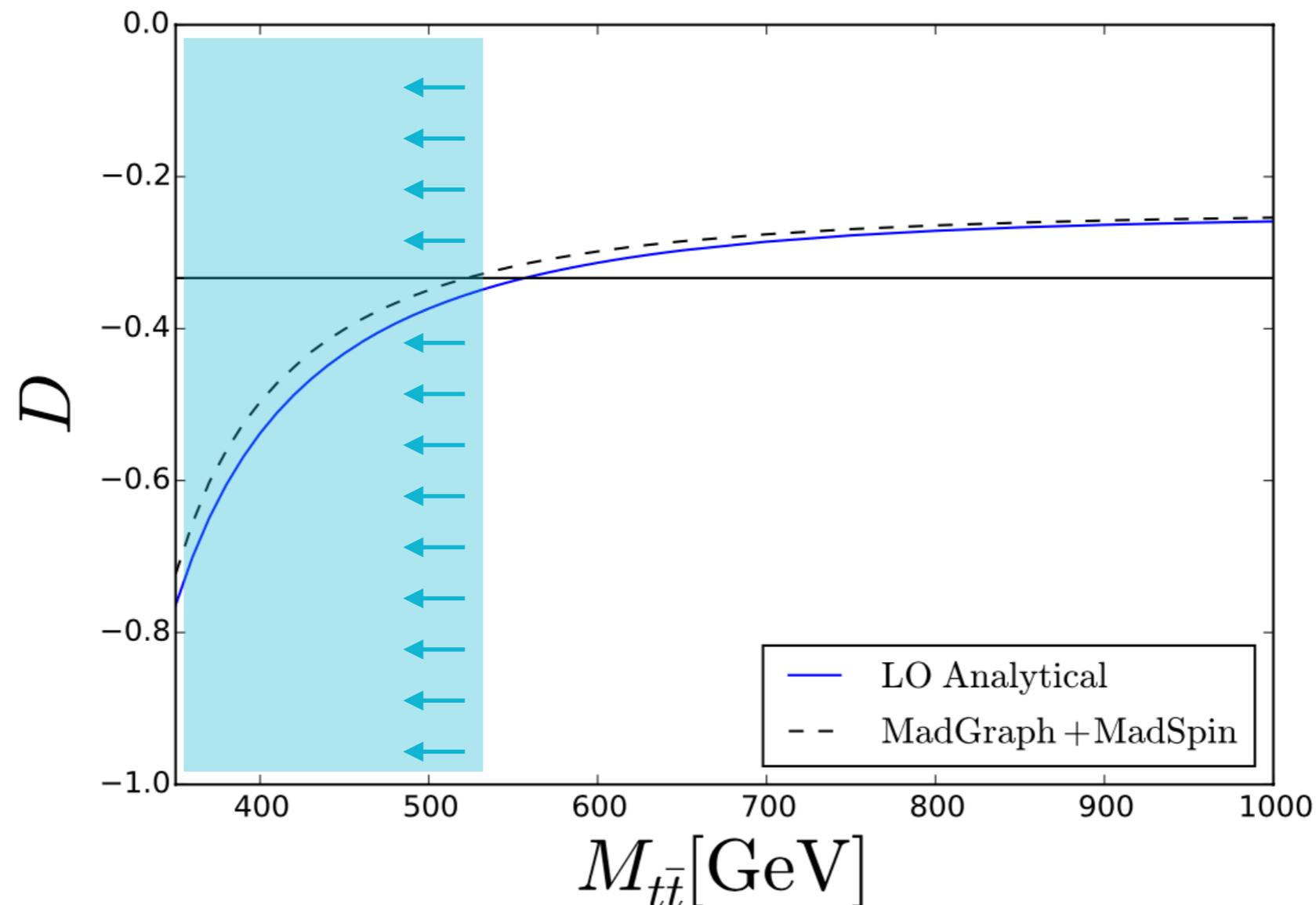


- What this is basically saying is that if the spin correlations across the  $k,n,r$  axes are stronger than 33%, you can't explain this with purely classical probabilities and need QM.

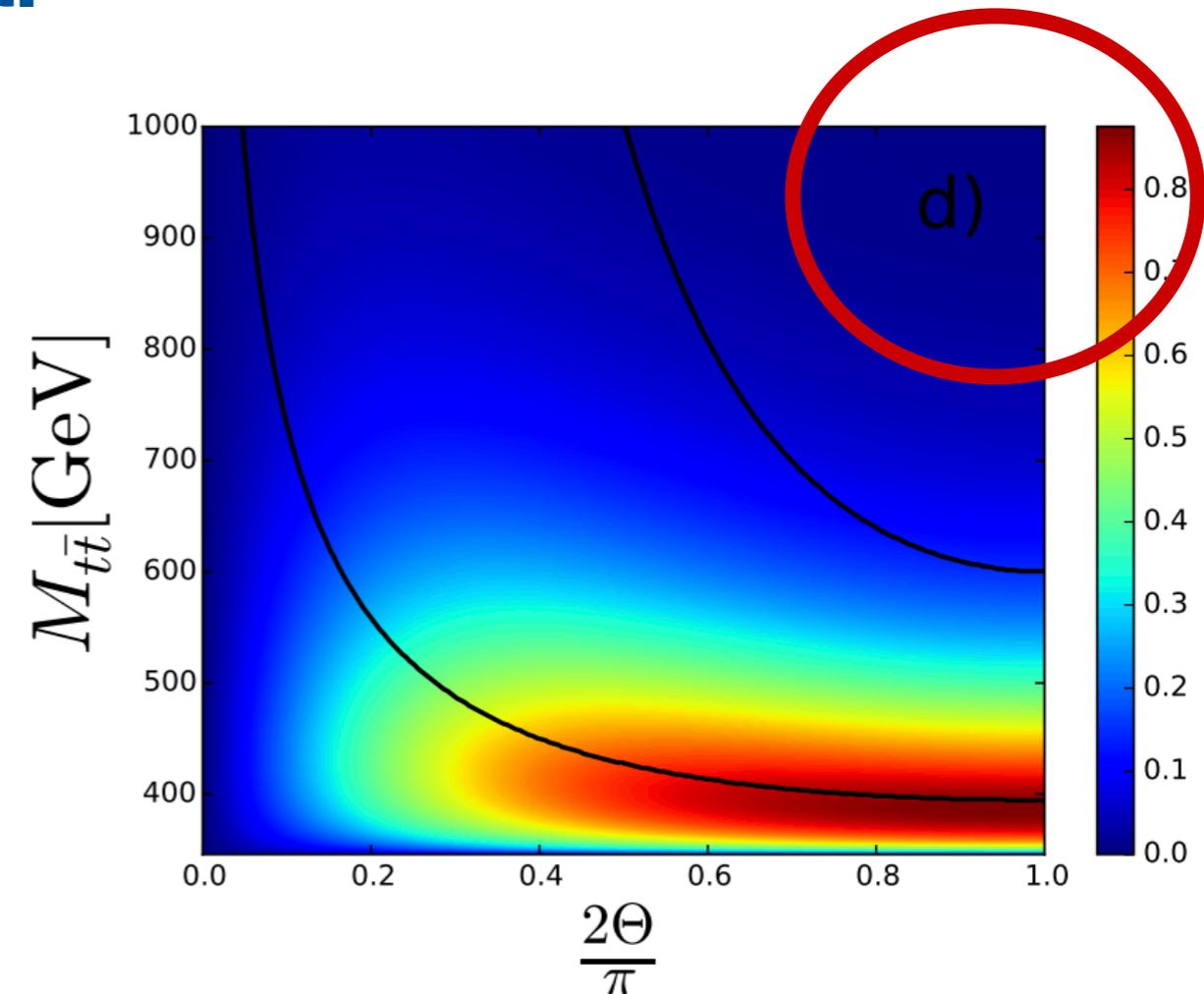
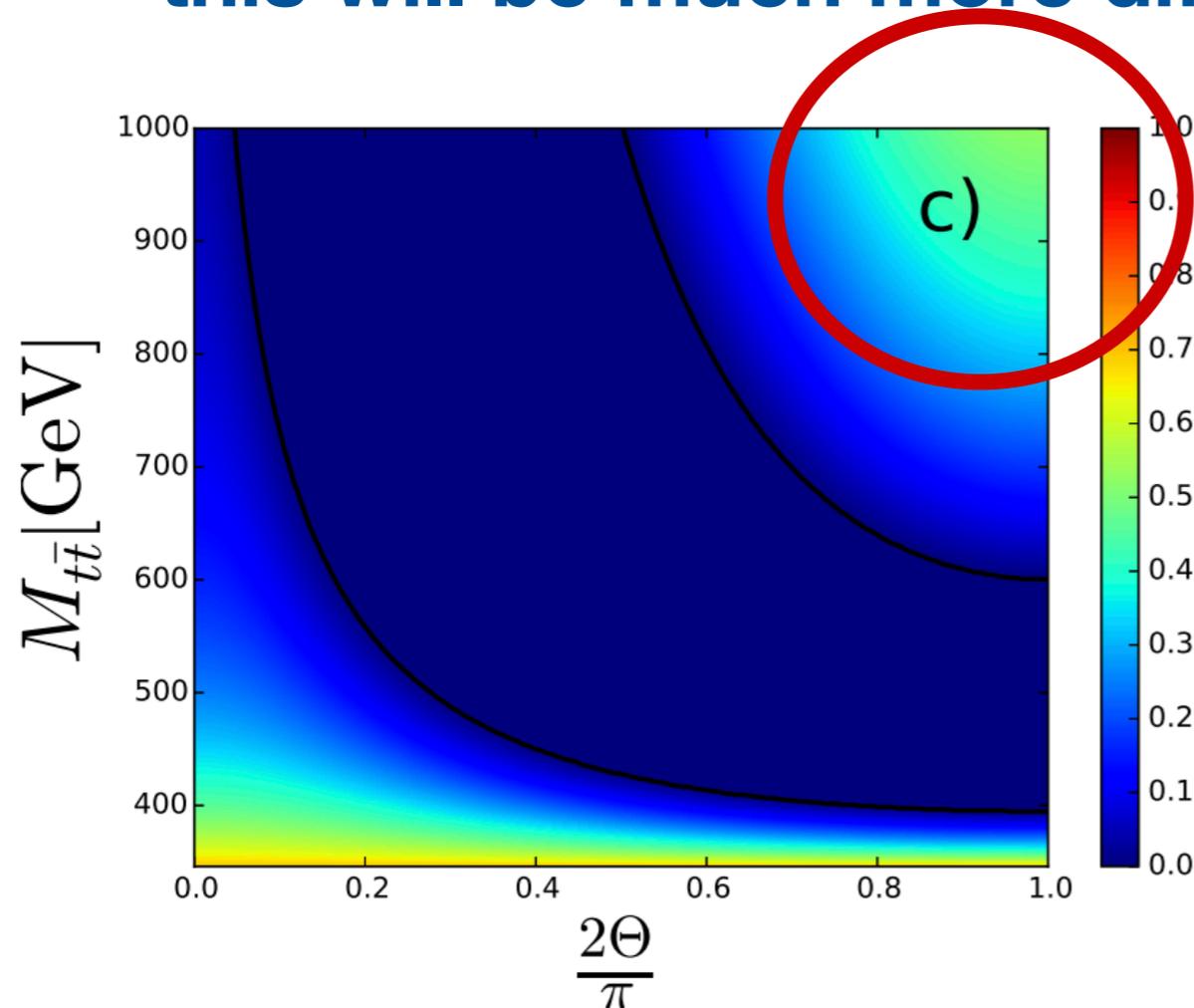
- You can actually use  $D$  to do a Bell test if you wanted to (though it isn't an optimal way).



- The primary experimental challenge is to reconstruct the tops with sufficient sensitivity to isolate the threshold region where tops are entangled.



- The next target will be to test Bell inequalities (CHSH), but this will be much more difficult.



- We have a lot of experience in looking at high-scale boosted top events, but not for spin correlation measurements.

- **What is allowed without violating relativity** (i.e. non-signalling)
  - ➔ **CHSH  $\leq 2$** : Purely classical correlations.
  - ➔ **CHSH  $\leq 2\sqrt{2}$** : Maximum allowed by QM correlations.
  - ➔ **CHSH  $\leq 4$** : Maximum allowed by non-signalling.
- **Particle physics measurements aim to minimise the dependence of detector corrections to the POI** (e.g. CHSH).
- **Easy to be sensitive to exotic values for QI observables in principle** (not trivial to test in practice).

- **Why stop at top quarks? The SM offers many more ways to explore QI, with more exotic spin states:**
- **Higgs decays:**
  - ➔ **HWW** (semileptonic and dileptonic)
  - ➔ **HZZ** (4 lepton)
- **Diboson events:**
  - ➔ **Vector boson processes** (ZZ, WW, WZ etc)
- **Other top decay modes:**
  - ➔ **Boosted semi-leptonic**
  - ➔ **Single top** [brand new idea]
- **Not just stamp collecting, each of these offers unique spin structures.**

- Why stop at top quarks? The SM offers many more ways to explore QI, with more exotic spin states:

- Higgs decays:

→ **HWW** (semileptonic and  
→ **HZZ** (4 lepton)

- Diboson events:

→ Vector boson process

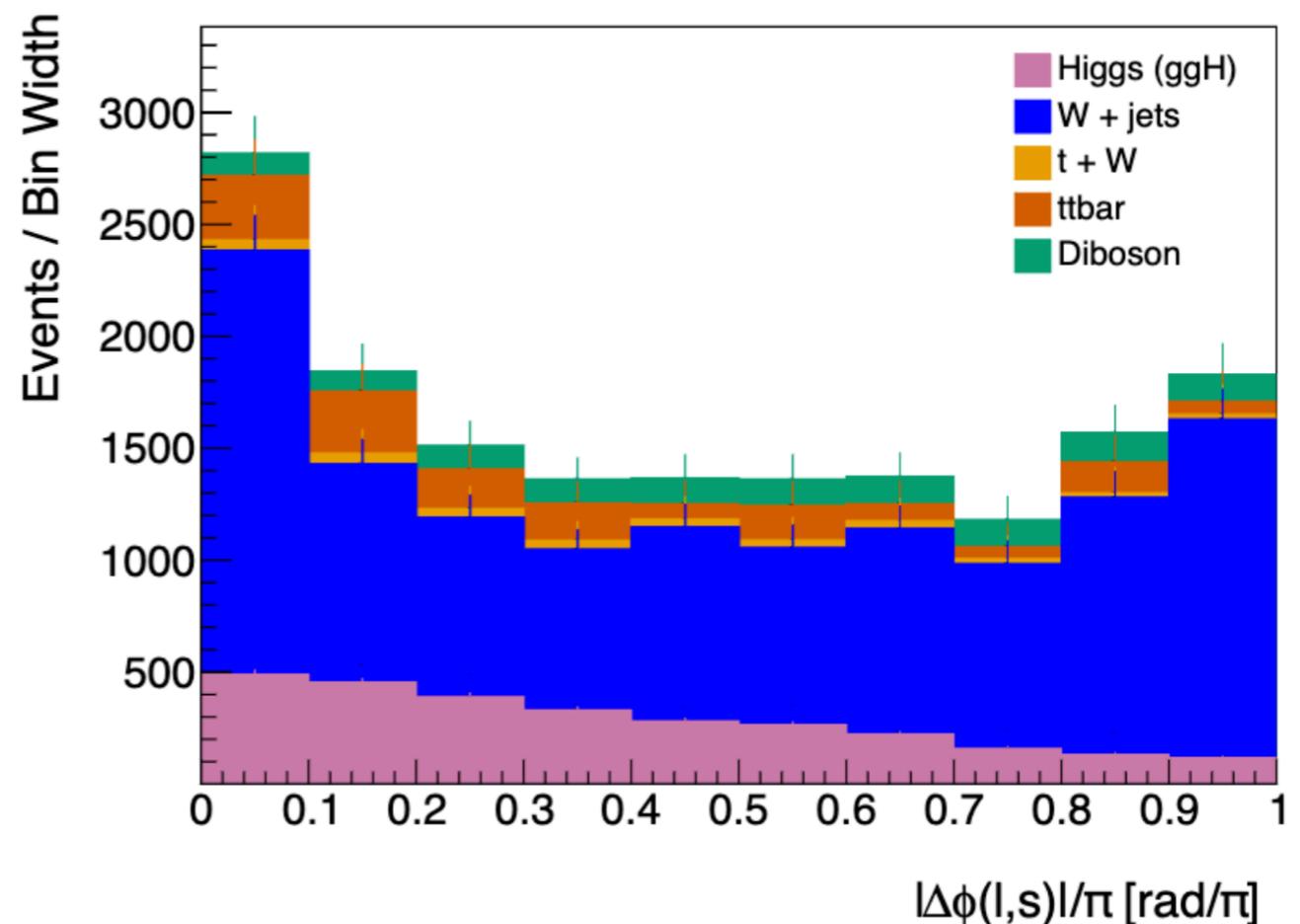
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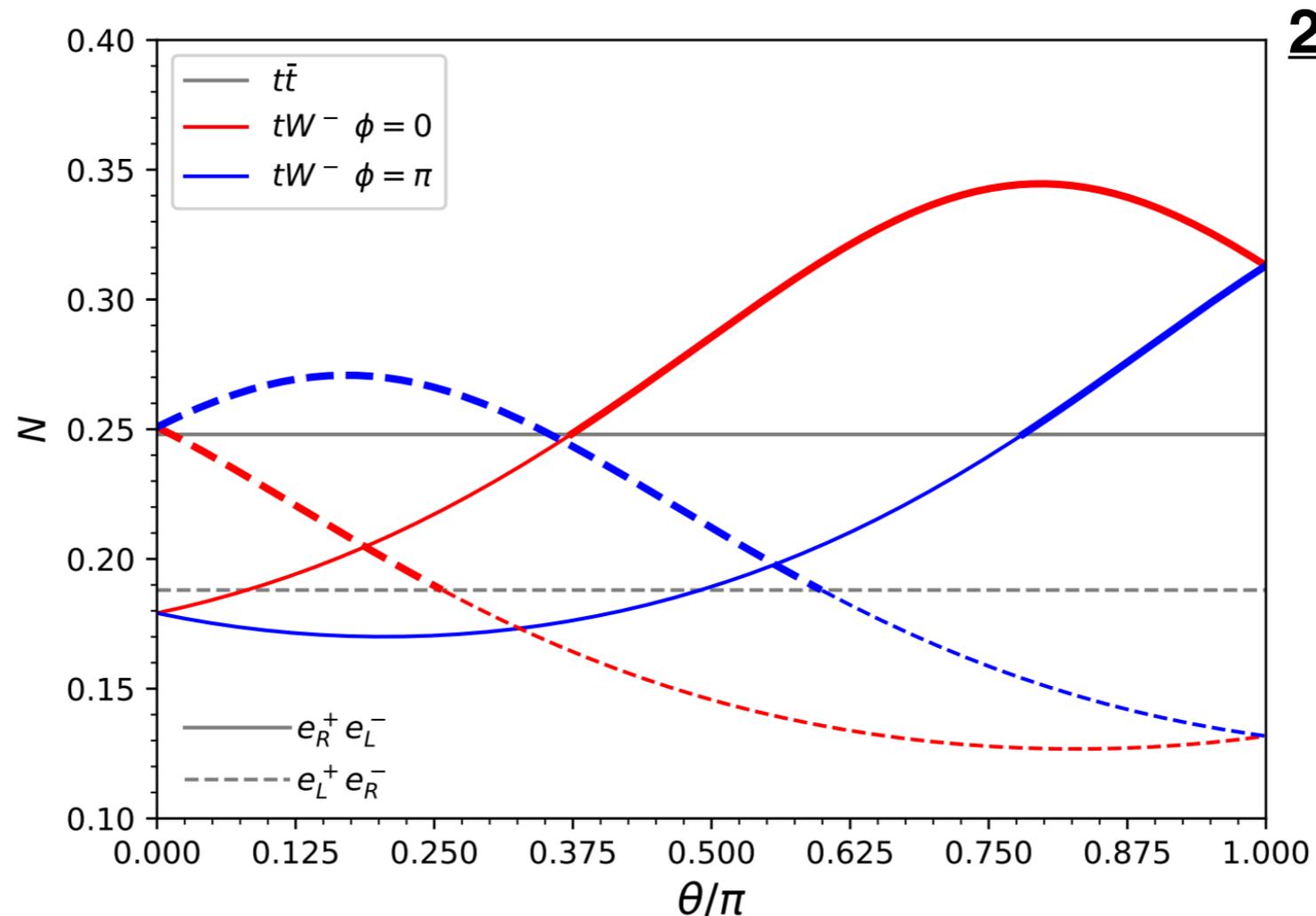
→ Single top [brand new i

- Not just stamp collecting, each of these offers unique spin structures.

**2307.13783**



- Are there things we can do that no one else can?
- “*Autodistillation*” is the idea that as particle systems decay, their entanglement gets stronger.

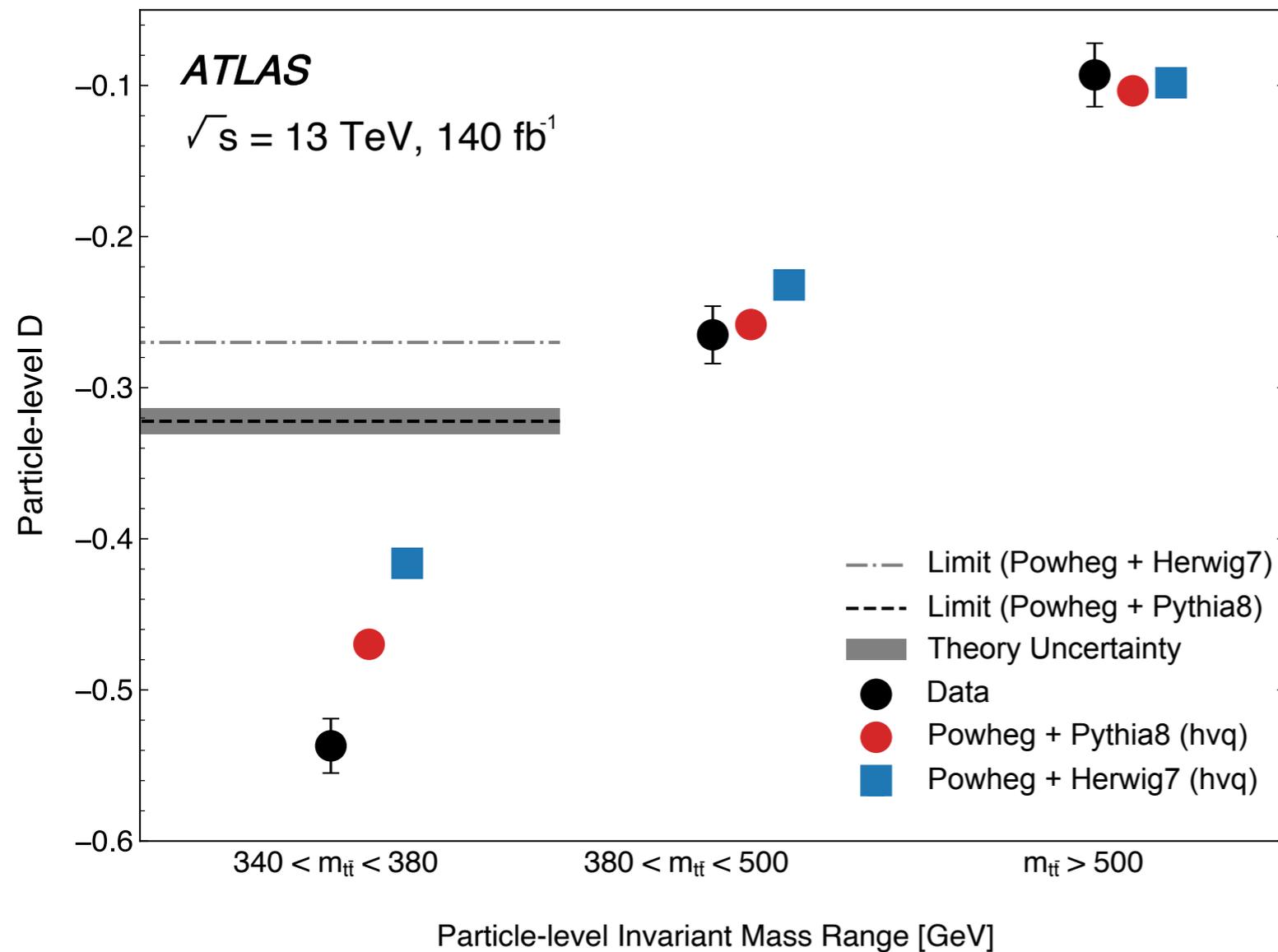


- This leads to the concept of “spin analysing power”:

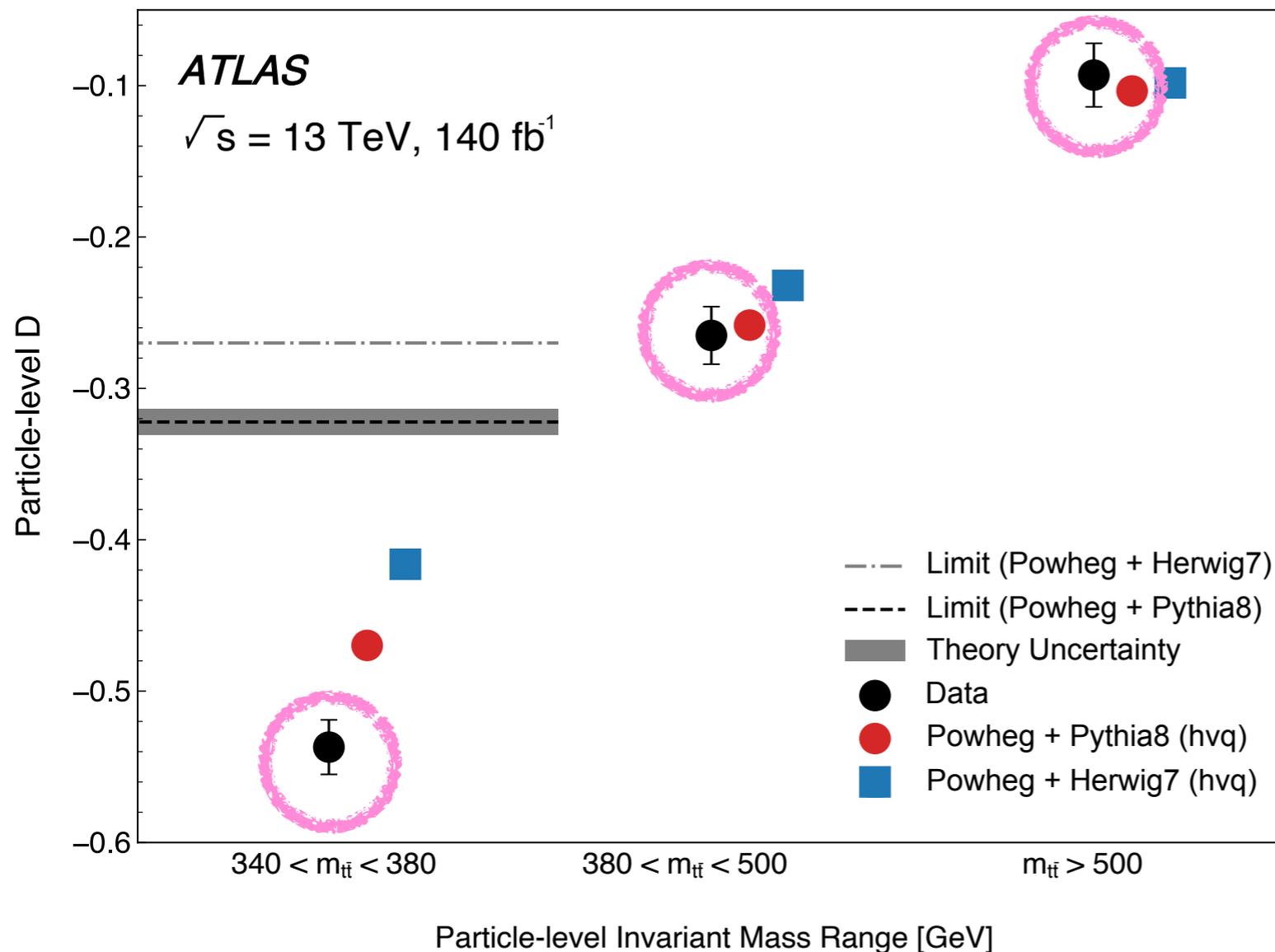
	$b$ -quark	$W^+$	$l^+$	$\bar{t}$ -quark or $\bar{s}$ -quark	$u$ -quark or $c$ -quark
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- Almost all published spin measurements in top physics use the leptonic decay mode:
  - ➔ easiest to identify experimentally.
- In Run3 we will start to see results using down-type jets.

- How reliable are the elements of this result?

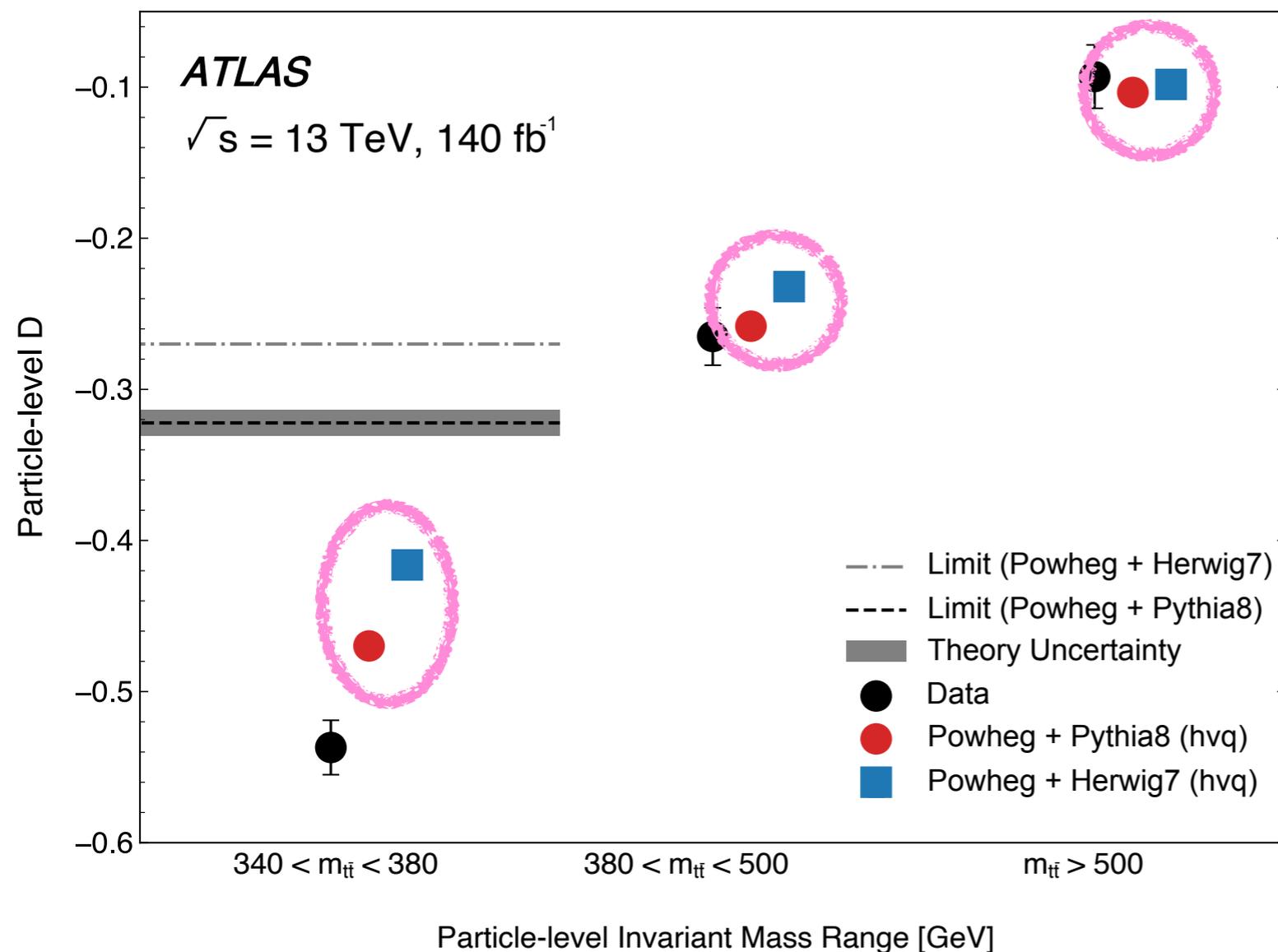


- How reliable are the elements of this result?



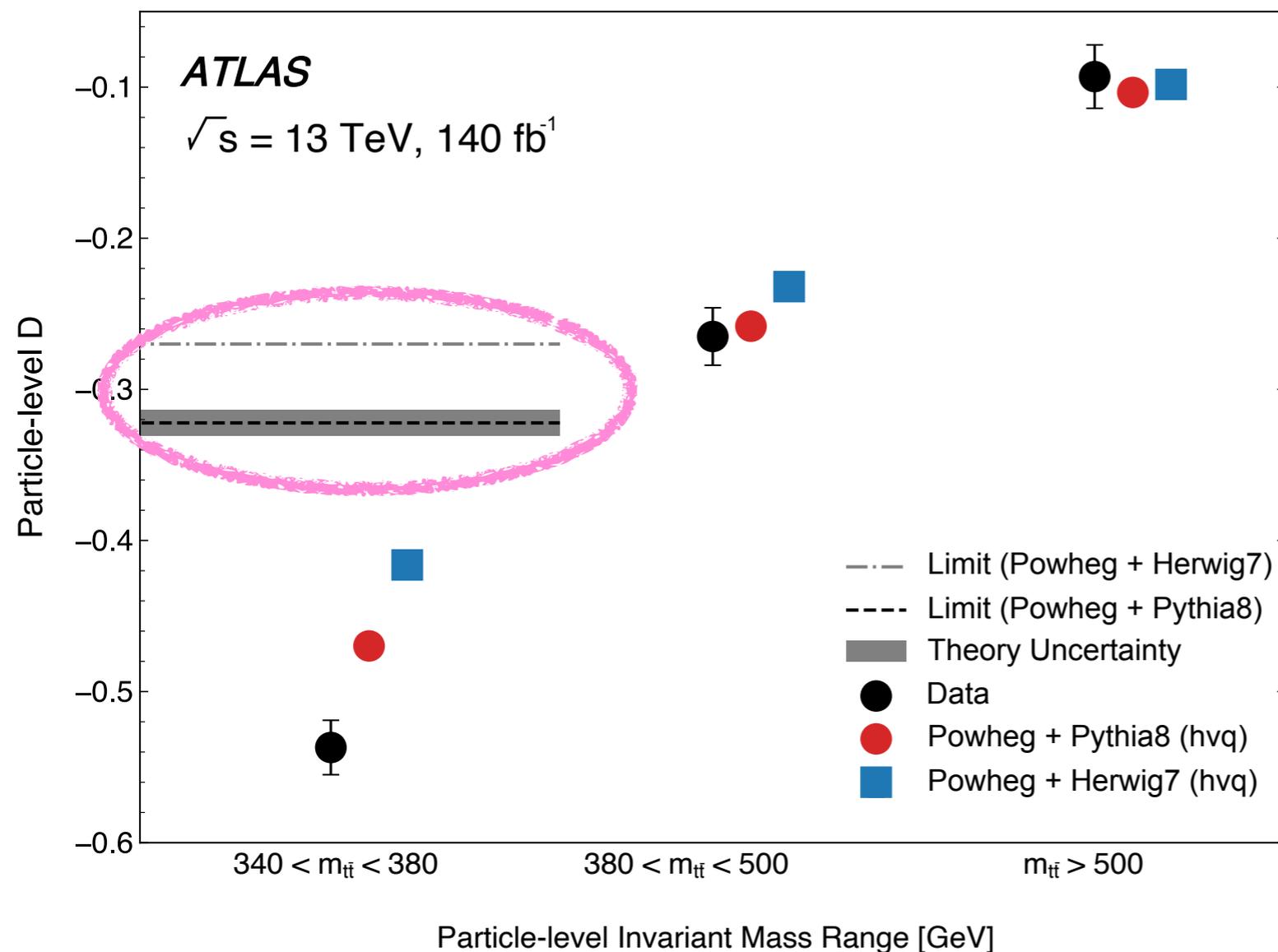
- Corrections to the data: **very reliable**
- A comprehensive and conservative (even by ATLAS's standards) **list of systematic uncertainties has been considered on all aspects of the analysis.**

- How reliable are the elements of this result?



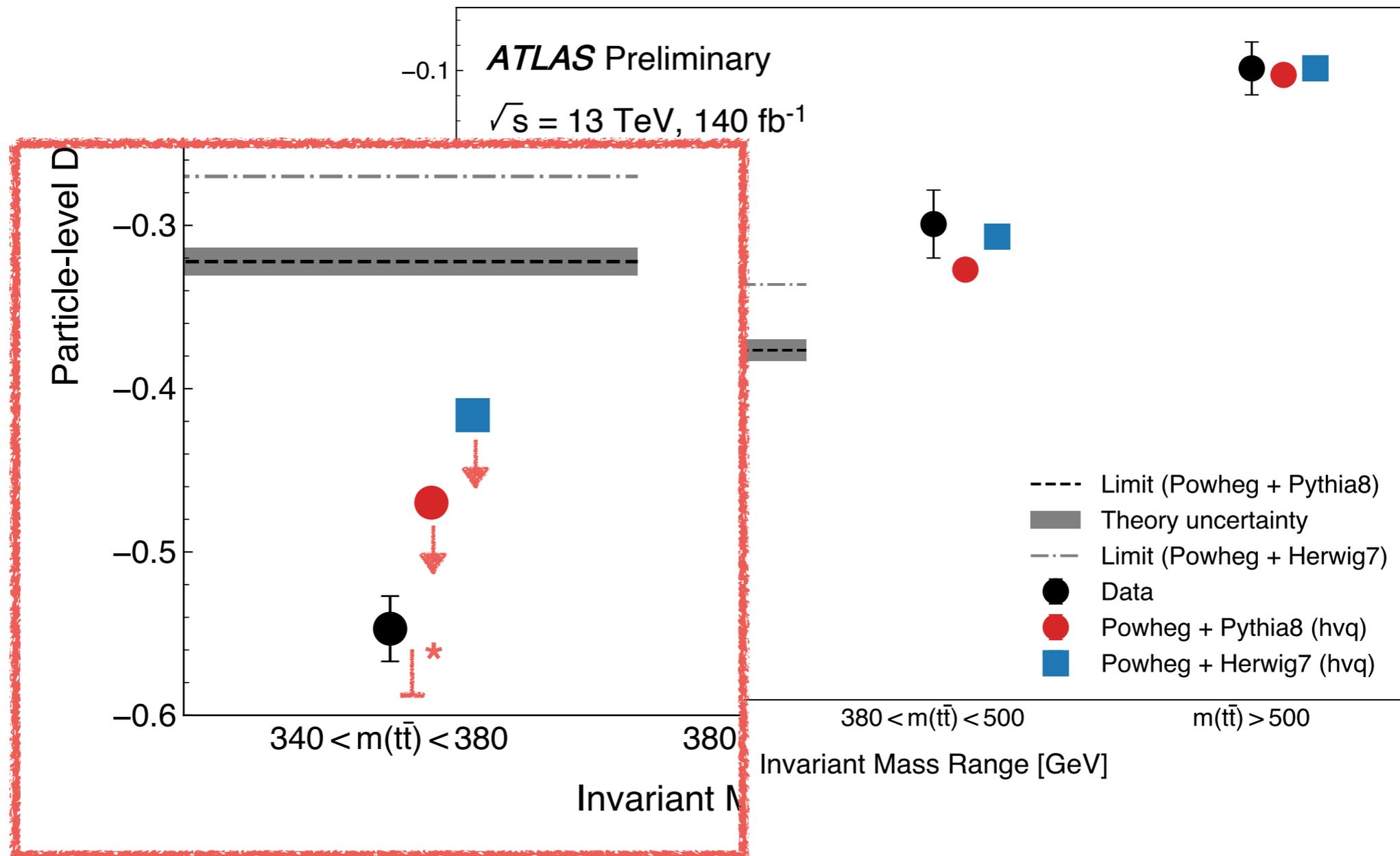
- Predictions of the SM: **Reliable but limited.**
- These predictions come from general purpose MC generators:
  - ➔ We understand them very well, but they are not designed to model threshold perfectly.

- How reliable are the elements of this result?



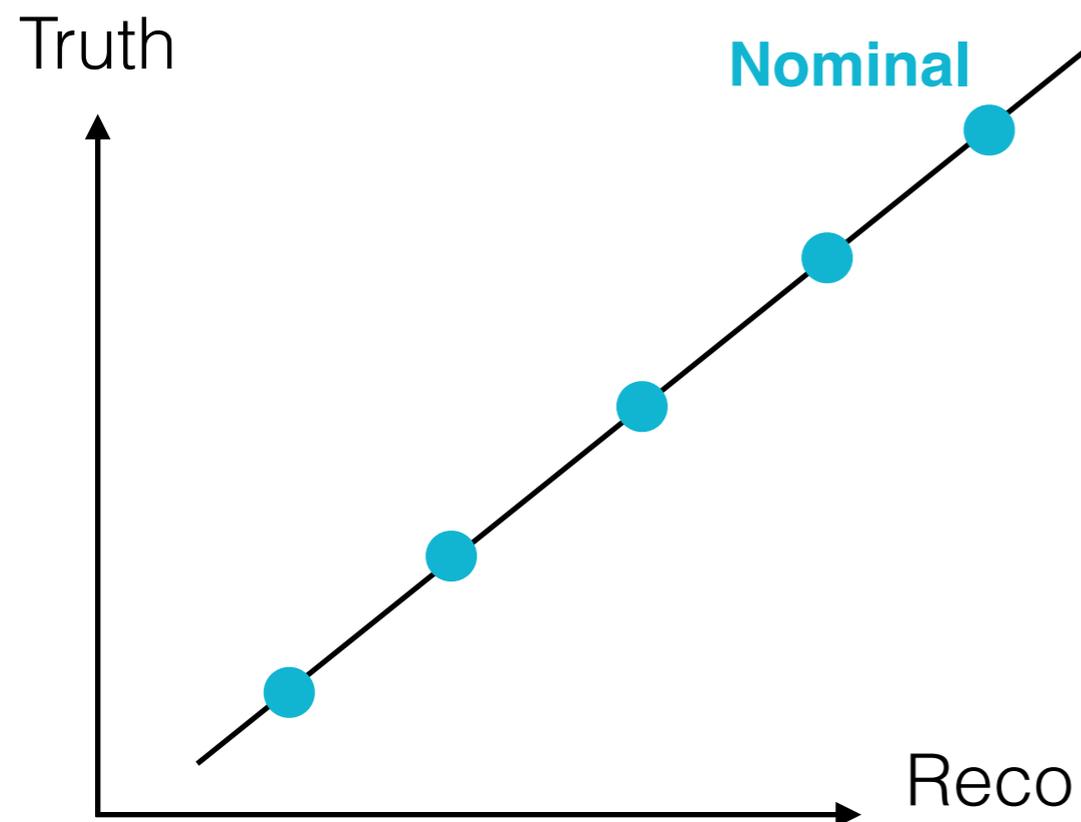
- Entanglement limits: **Reliable but limited.**
- Same limitations as predictions.
- Two models give different limits, but source is understood and we've taken the most conservative of the two.

- **Bound state effects should be increasing entanglement:**  
 ➔ Including them only makes result more significant, not less.

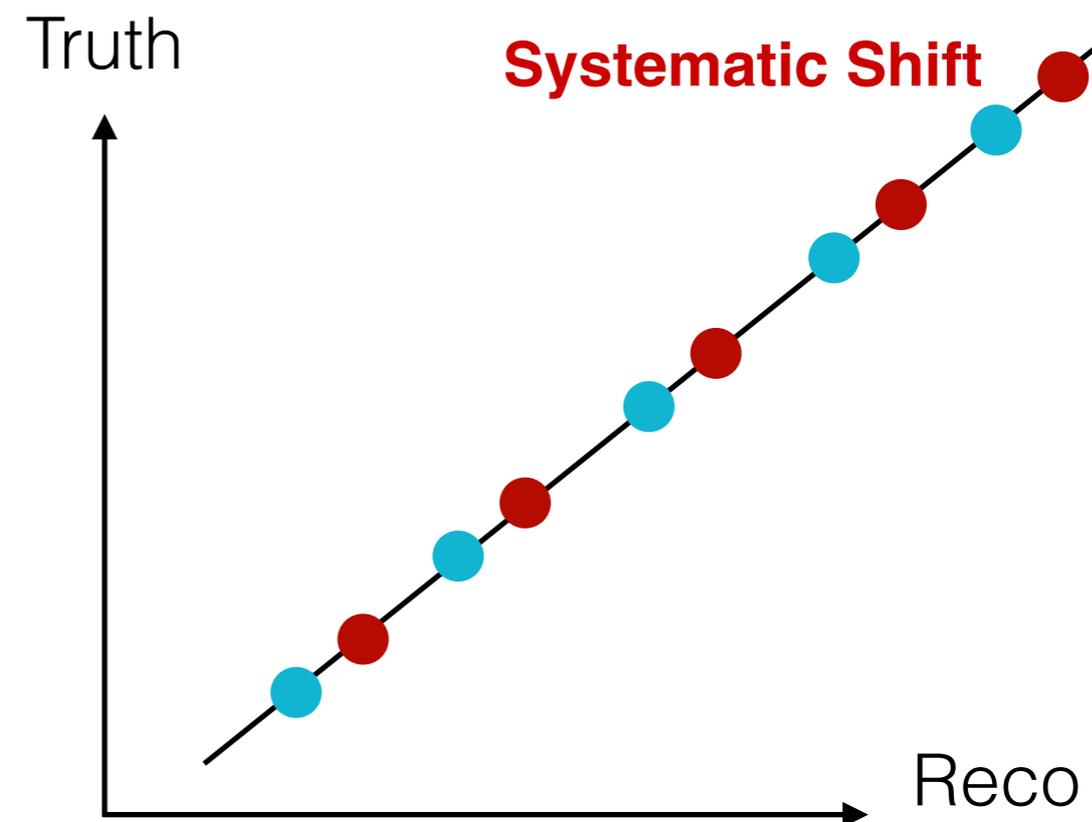
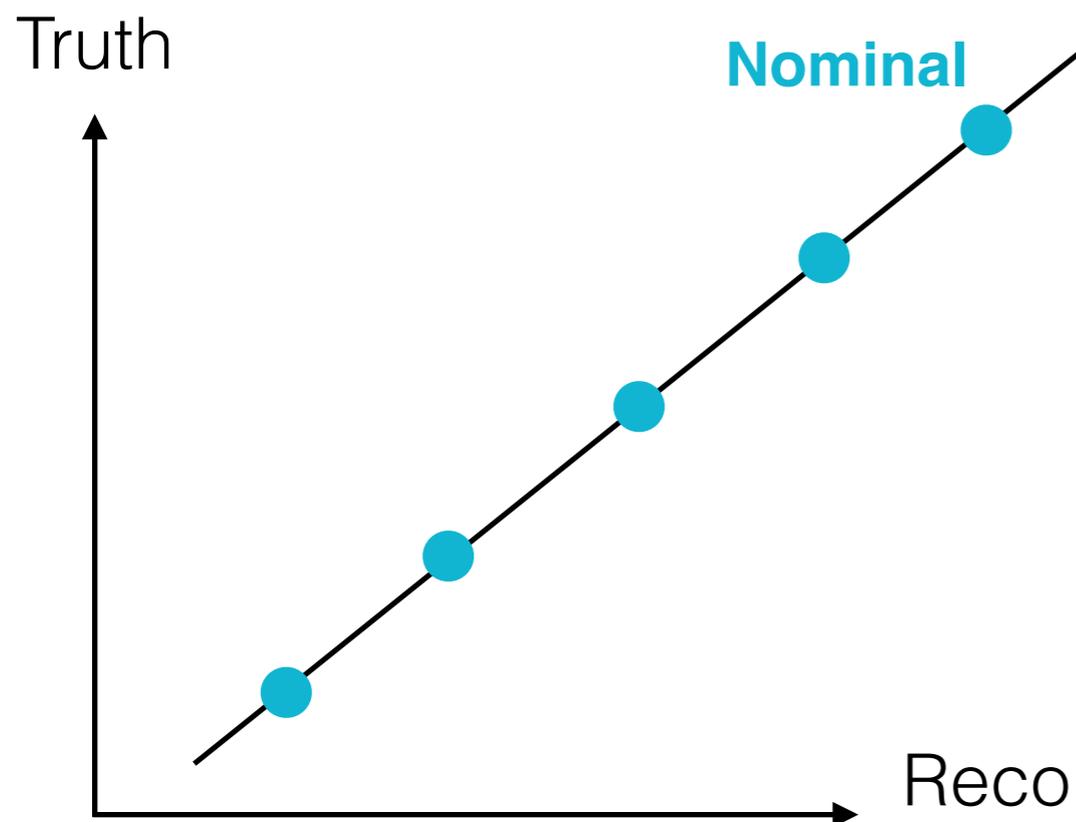


\*exaggerated, the effect on the error bars would be too small to see.

- The relative size of the systematics is not fixed and changes at each hypothesis point:

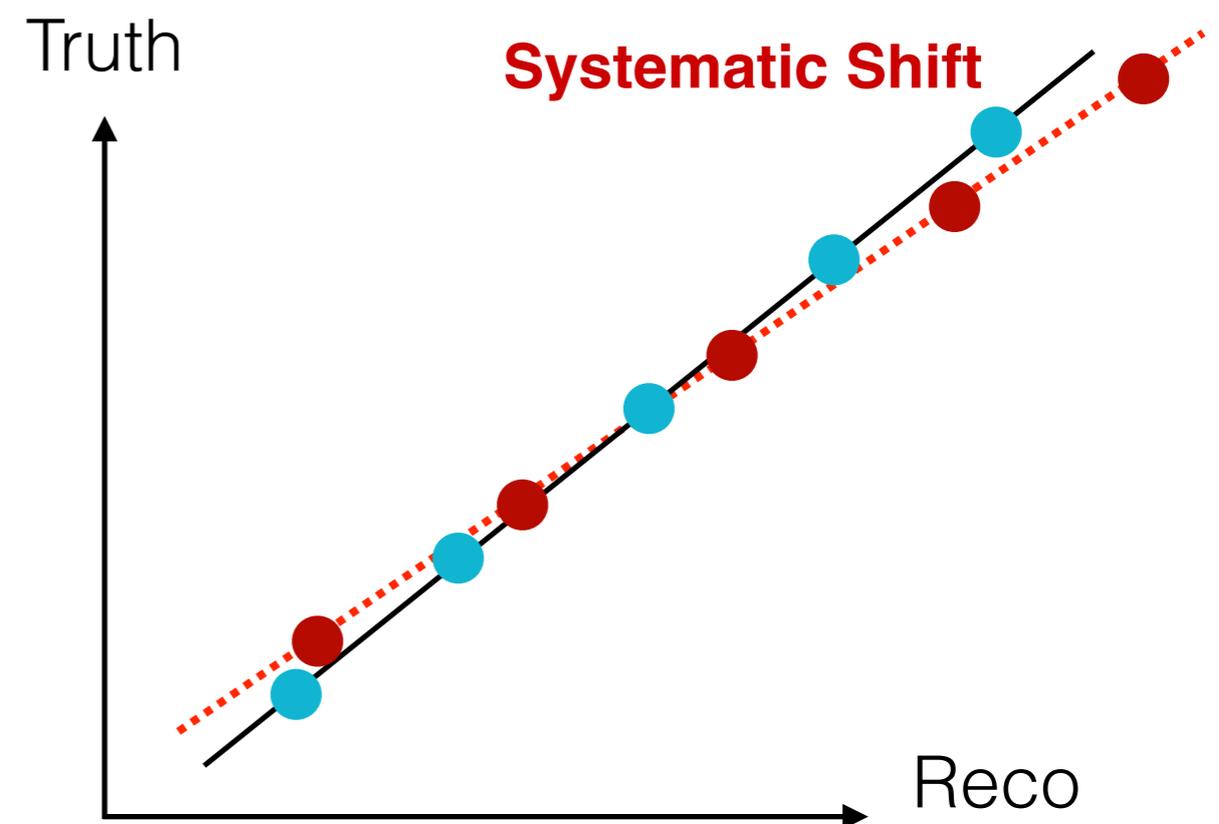
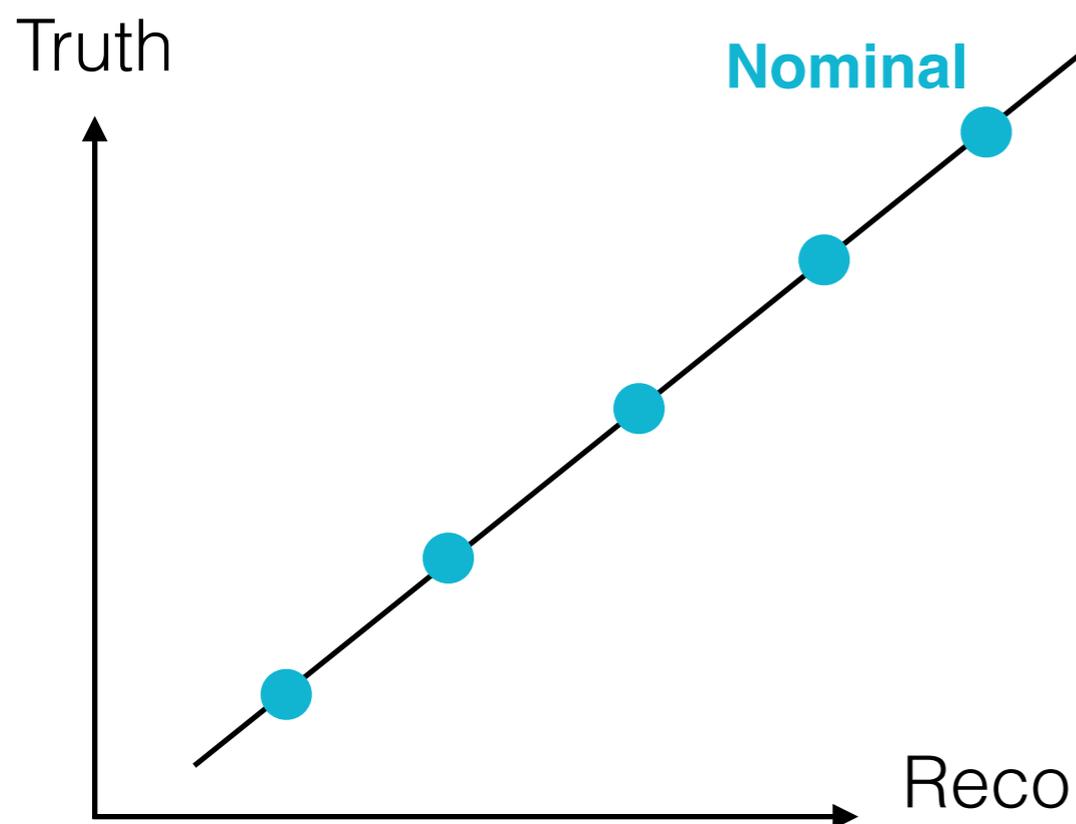


- The relative size of the systematics is not fixed and changes at each hypothesis point:



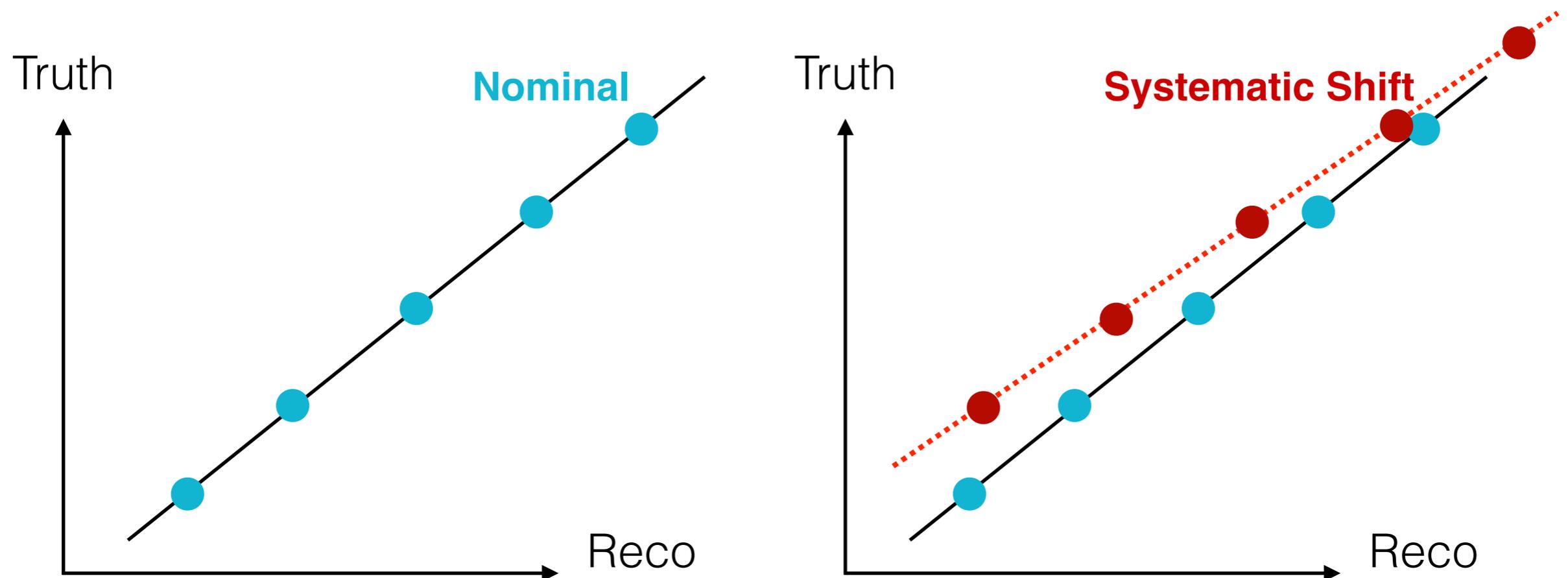
- Ideally, truth and reco shift in a correlated way, and there is no resultant uncertainty.

- The relative size of the systematics is not fixed and changes at each hypothesis point:



- In practice, most uncertainties shift reco but not truth and therefore change the slope (all detector uncertainties do this).

- The relative size of the systematics is not fixed and changes at each hypothesis point:



- In the worst case, systematics shift slope and offset and have a large effect (our dominant uncertainties behave this way).