# Quantum tests in collider physiscs

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# TIME vs. REALITY:

A novel quantum effect in K° - K° Entaglement



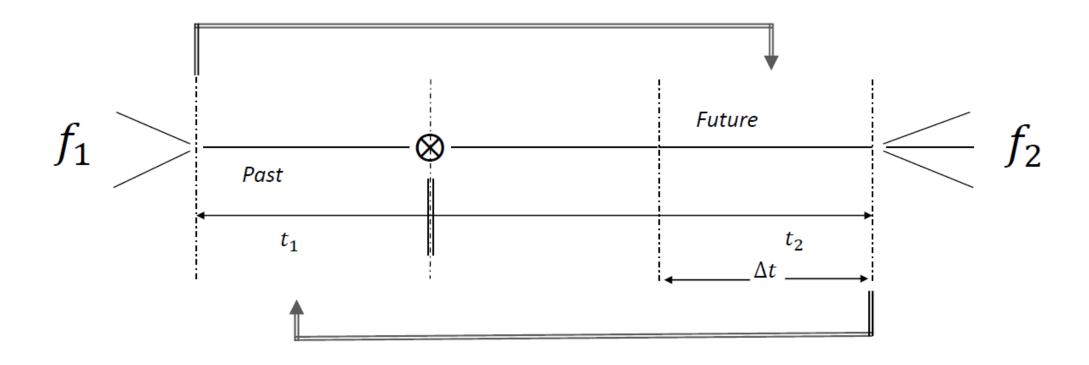
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# What is THE NOVELTY beyond Entanglement in Quantum Optics?

- $ightharpoonup \Delta$  F = 2 Mixing  $(K^0 \overline{K^0}, B^0 \overline{B^0}, ...)$
- CP Violation Mixing Decay InterferenceDecay
- Non-Trivial Time Evolution: Anton Zeilinger
  Production ⇒ Entangled ⇒ Interference ⇒ Decoherence
  with rich distinct information from one or double decay on the three regimes
- > States with definite Mass and Lifetime  $\lambda = M i\Gamma/2$ ,  $\Delta M \neq 0$ ,  $\Delta \Gamma \neq 0$  are those with **definite Time Evolution**.
- $\triangleright$  Existence of B-Factory and  $\Phi$ -Factory Facilities

- I. TIME HISTORY of Entangled System: from Production to its fate
- TIME REVERSAL in Δt for unstable particles



- II. POST-TAG of Past-decayed state: Entanglement times t<sub>1</sub>
- K<sub>s</sub> TAG

#### **NOVEL EFFECTS**

- (1) As a **Tool** for the BYPASS of (otherwise) NO-GO THEOREMS
  - 1.1 The Conundrum of Time Reversal and CPT for Unstable Particles
  - 1.2 What is a K<sub>s</sub> experimentally?
- (2) The discovery of new quantum phenomena:

#### **SURVIVING CORRELATION - IN - TIME FROM FUTURE TO PAST**

It comes definite from measurement in the future  $t_2$ , when the system is no-longer entangled, to the state –depending on  $t_2$  (!?) - of the partner in the past  $t_1$ , before its decay when it was entangled and "unspeakable".

It is asymmetric compared to the correlation from past to future.

If EPR  $\rightarrow$  Spooky Action at a Distance  $\rightarrow$  Bell Theorem  $\rightarrow$  end of Hidden Variables and proof of "Lack of Local Realism"  $\rightarrow$  Quantum Information,

then  $\rightarrow$  What about the novel correlation - in - time?  $\rightarrow$  Spooky Action to the Past  $\rightarrow$ ???

#### **OUTLINE**

- > Entangled two-body C=- neutral meson system
- > Time Evolution and "Survival" probalility: the Total Width
- $\succ$  The state  $|K_{\rightarrow f}\rangle$  not decaying to f. **The K<sub>L</sub> tag**
- ➤ The Conundrum of **Time Reversal** —and CPT- **for Unstable Particles**: NO-GO and its Bypass (in 1999): **The Conceptual Basis**

- From the observation of second decay  $f_2$  at  $t_2$  to the partner state before its decay at  $t_1$ . SURPRISE of the "initial" state depending on  $t_2$ .
- > The K<sub>s</sub> tag
- > Conclusion: An epistemological open question

## **ENTANGLED** C = - neutral meson system

 $\blacktriangleright$  Actually existing at DA $\phi$ NE with  $\Phi \to K^0 \, \overline{K}^0$ ,

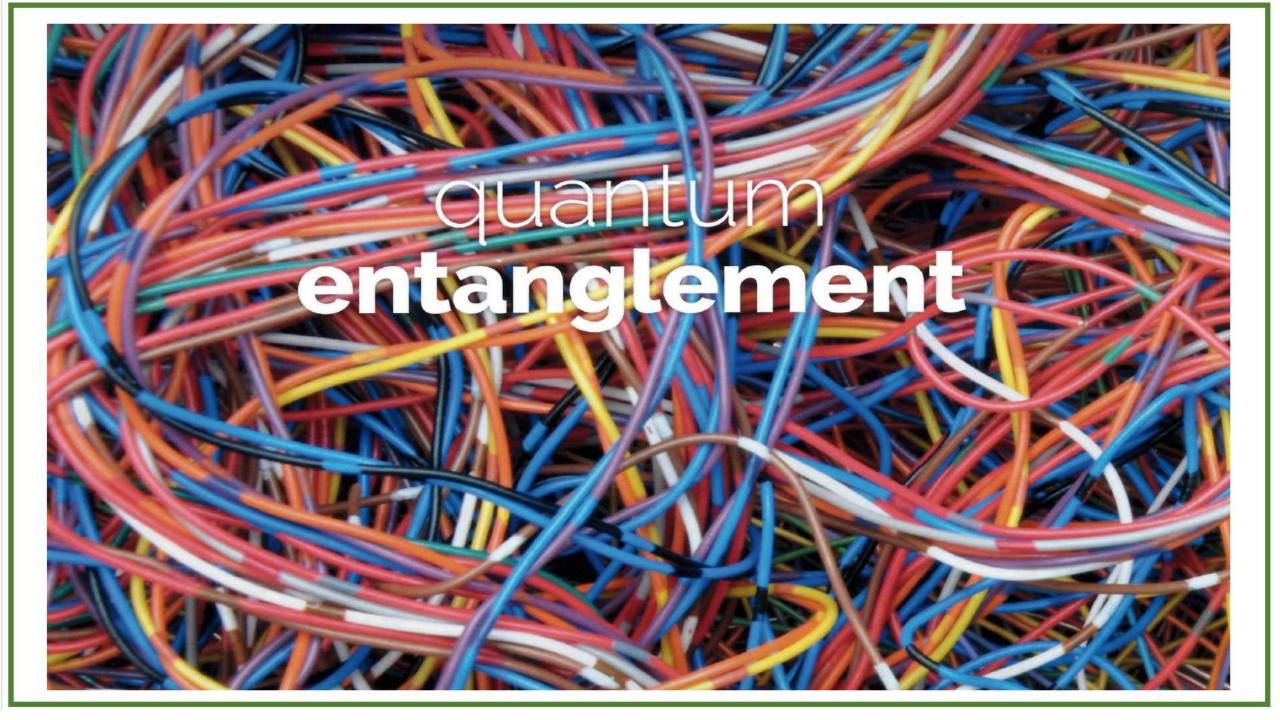
at BABAR and BELLE with  $\Upsilon(4S) \longrightarrow B^0 \, \bar{B}^0$ 

$$C\mathcal{P}=+ \Rightarrow |i(t=0)\rangle = \frac{1}{\sqrt{2}}\{|K^0\rangle|\overline{K}^0\rangle - |\overline{K}^0\rangle|K^0\rangle\}$$

with particle 1 decaying at  $t_1$ , particle 2 decaying at  $t_2 > t_1$ 

- Fiven With Mixing,  $|i(t)\rangle$  does not generate any  $K^0$   $K^0$ , nor  $\overline{K}^0$   $\overline{K}^0$ , due to antisymetry (not valid for symmetric C=+!)
- ➤ Time Evolution → definite in terms of non-orthogonal eigenstates of the non-normal Hamiltonian

$$\begin{aligned} \left| K_{S,L} \right\rangle \alpha \left[ \left( 1 + \epsilon_{S,L} \right) \left| K^0 \right\rangle \right. & \pm \left( 1 - \epsilon_{S,L} \right) \left| \overline{K}^0 \right\rangle \right], \\ \mathscr{L} \mathcal{P} &\longrightarrow \left\langle K_S \middle| K_L \right\rangle \right. & \simeq \epsilon_L + \epsilon_S^*, \\ \epsilon &= (\epsilon_S + \epsilon_L)/2 \longrightarrow \mathcal{T}, \qquad \delta = (\epsilon_S - \epsilon_L)/2 \longrightarrow \mathscr{LP} T \end{aligned}$$



#### TIME EVOLUTION $|i(t)\rangle$

- > The entangled state is **non-separable** in parts:
- (i) "which is which" is not defined;
- (ii) the two parts are not definite: any two linerally independent combinations. Only the state  $|i\rangle$  is definite: the state of each part is "unspeakable".
- > The time evolution, written as

$$|i(t=0)\rangle = N/\sqrt{2} \left\{ |K_S\rangle \ |K_L\rangle - |K_L\rangle |K_S\rangle \right\} \ , \ |N|^2 = (1 - |\langle K_S|K_L\rangle|^2)^{-1} \Rightarrow |i(t)\rangle = e^{-i(\lambda_S + \lambda_L)t} \ |i(t=0)\rangle = (1 - |\langle K_S|K_L\rangle|^2)^{-1} \Rightarrow |i(t)\rangle = e^{-i(\lambda_S + \lambda_L)t} \ |i(t=0)\rangle = (1 - |\langle K_S|K_L\rangle|^2)^{-1} \Rightarrow |i(t)\rangle = e^{-i(\lambda_S + \lambda_L)t} \ |i(t=0)\rangle = (1 - |\langle K_S|K_L\rangle|^2)^{-1} \Rightarrow |i(t)\rangle = e^{-i(\lambda_S + \lambda_L)t} \ |i(t=0)\rangle = (1 - |\langle K_S|K_L\rangle|^2)^{-1} \Rightarrow |i(t)\rangle = e^{-i(\lambda_S + \lambda_L)t} \ |i(t=0)\rangle = (1 - |\langle K_S|K_L\rangle|^2)^{-1} \Rightarrow |i(t)\rangle = e^{-i(\lambda_S + \lambda_L)t} \ |i(t=0)\rangle = (1 - |\langle K_S|K_L\rangle|^2)^{-1} \Rightarrow |i(t)\rangle = e^{-i(\lambda_S + \lambda_L)t} \ |i(t=0)\rangle = (1 - |\langle K_S|K_L\rangle|^2)^{-1} \Rightarrow |i(t)\rangle = e^{-i(\lambda_S + \lambda_L)t} \ |i(t=0)\rangle = e^{-i(\lambda_S + \lambda_L)t} \ |i(t=0)\rangle = (1 - |\langle K_S|K_L\rangle|^2)^{-1} \Rightarrow |i(t)\rangle = e^{-i(\lambda_S + \lambda_L)t} \ |i(t=0)\rangle = e^{-i(\lambda_S + \lambda$$

The Survival Probability  $P(t_1) = \|i(t = t_1)\|^2 = e^{-\Gamma t_1}$ , Total Width  $\Gamma = \Gamma_S + \Gamma_L$ 

- $|i(t)\rangle$  is unaltered, it reamains the same: NO INTEREST BEFORE THE FIRST DECAY. The considered observable has been the **Double Decay Rate Intensity** I ( $f_1$ ,  $f_2$ ;  $\Delta t$ )!
- > Careful! P(t<sub>1</sub>) iff nothing else is observed in the future
- > How to inquire in the "unspeakable" regime ?

#### FIRST DECAY $f_1 \rightarrow TAGGING$ AND FILTERING

>Any state can decay to f, but that with zero probability

$$|K_{\to f}\rangle = N_{\to f}[K_L\rangle - \eta_f |K_S\rangle]; \qquad \eta_f = \frac{\langle f| \ T \ |K_L\rangle}{\langle f| \ T \ |K_S\rangle}$$

> If you observe the first decay to  $f_1$  at  $t_1$ , proyecting  $|i(t=t_1)\rangle$  to  $f_1$ , the **living partner** (2) corresponds to the pure state

$$|K^{(2)}(t=t_1)\rangle = |K_{\leftrightarrow f_1}\rangle \iff TAG \ of \ (2)$$

This fact was always recognized for "flavour tag": First decay to  $| \cdot (| \cdot |) \rightarrow$  Partner tagged to  $\overline{K}^0(K^0)$ . It is, however, valid in general as stated!

> What for the decayed state (1)? The state before decay was undefined.

Written as a superposition of  $|K_{\rightarrow f}\rangle$  and its orthogonal  $|K_{\rightarrow f}^{\perp}\rangle$ 

Decay to 
$$f_1 \Rightarrow |K^{\perp}_{+f_1}\rangle$$
 FILTERED for (1)

**Decay Rate** given by the **decay probability** to  $f_1$  of  $|K_{+}^{\perp}f_1\rangle \equiv$  **FILTERING IDENTITY** 

#### Δt HISTORY OF THE LIVING PARTNER

- The subsequent  $\Delta t$  evolution of particle (2) and its decay to  $f_2$  are definite from the prepared tagged state.
- For  $\Delta t \le few \ \tau_s$ , one has an interference patern, because no decay channel due to CP Violation projects either  $K_s$  or  $K_L$ !
- ightharpoonup For long enough  $\Delta t$ , one has **Decoherence**  $K_L tag \Leftrightarrow |\eta_1| e^{-\Delta \Gamma \Delta t/2} \ll 1$  with a quantitative purity of the  $K_L$  -state
- $\triangleright$  The observable is the **Double Decay Rate, the Intensity** I(  $f_1$ ,  $f_2$ ;  $\Delta t$ ). Tagging of the living partner at  $t_1$  and Filtering of its state in its Decay to  $f_2$  at  $t_2$

allows to talk of  $\Delta t$  Transition Probability  $P\left(K_{\to f1} \xrightarrow{\Delta t} K_{\to f2}^{\perp}\right)$ 

"independent of the decay" and connected to I( $f_1$ ,  $f_2$ ;  $\Delta t$ ).

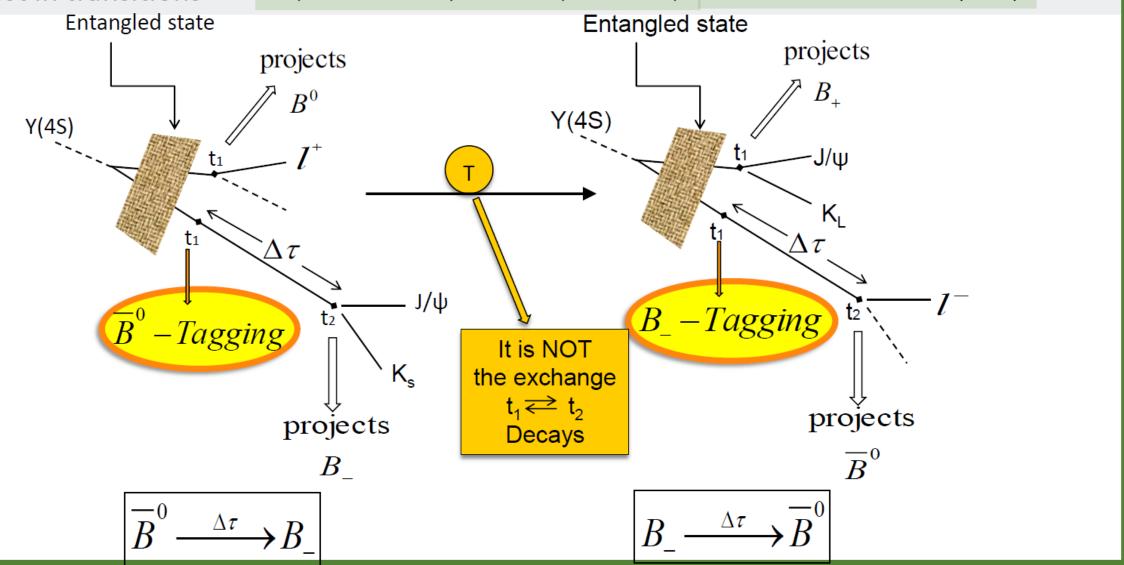
#### TR-ASYMMETRY: CONCEPTUAL BASIS FOR BYPASSING NO-GO

- ➤ Neutral Mesons  $K^0 \overline{K}^0$ ,  $B^0 \overline{B}^0$  are UNSTABLE and the Decay is irreversible.
- T and CPT, ANTIUNITARITY!, need however the exchange of initial and final states →NO-GO.
   L. Wolfenstein, PRL 1999: "The T-reverse of a decaying state is not a physical state".
- ➤ **BYPASS** M. C. Banuls, J. B., PLB 1999, NPB 2000 → Do not include the Decay Products in your Asymmetry, write it in terms of Meson States and the Decay should not be an essential ingredient for getting a non-vanishing value:
- Use the Decay as a Quantum Filtering Measurement of the Meson State ONLY: Orthogonal to Non-Decay State.
- 2) Quantum ENTANGLEMENT: Quantum Information from the First Decay to the (still alive) Partner for the Preparation of the initial Meson State: Non-Decay State if Antisymmetric entangled system.
- The test of Symmetries is made in the Time Evolution of the Partner from the first to the second decay.
  - L. Wolfenstein, IJMP E 1999: "It appears to be a true TRV Effect"

#### WHAT IS T-TRANSFORMATION EXPERIMENTALLY?

The problem is in the preparation and filtering of the appropriate initial and final meson states for a T-test in transitions

J.B., Martinez Vidal, Villanueva, JHEP 2012, COVER PAGE RMP vol. 87 (2015)



#### POST-TAG TO THE PAST DECAYED STATE

- $\gt$  In the entangled  $|i(t)\rangle$  state, there is no privilege of one of the decay times  $\gt$  Study the implications of observing the second decay to  $f_2$  at time  $t_2$
- The partner  $K^{(1)}(t=t_2)$  is tagged

$$\left|K^{(1)}(t=t_2)\right\rangle = N_1\{\eta_2|K_s\rangle - |K_L\rangle\}$$

which has not been observed! But it decayed at time t<sub>1</sub><t<sub>2</sub>

Fixing the observation  $(\eta_2, t_2)$  and evolving  $t_1$  from  $t_1$ =0 to  $t_1$ = $t_2$ , its past state

had to be

$$|K^{(1)}(t=0)\rangle = N[\eta_2 e^{-i\lambda_L t_2}|K_S\rangle - e^{-i\lambda_S t_2}|K_L\rangle]$$

DOUBLE SURPRISE! Not only there is a post-tag of the initial state, it depends on when the second decay will be observed.

## **OBSERVABILITY OF "BACK FROM THE FUTURE" EFFECT**

- $\triangleright$  Entaglement times  $t_1 < t_2$
- $\blacktriangleright$  Decay  $t_1$ -time distribution to  $f_1=f_2=f$  , at different fixed  $t_2$

$$| \langle f | K^{(1)}(t) \rangle |_{t_{2}}^{2} = K \{ e^{-\Gamma_{s}t_{1}} + |g|^{2}e^{-\Gamma_{s}t_{1}}$$

$$- 2 (\Re g) e^{-rt_{1}/2} \cos(\Delta m) t_{1}$$

$$- 2 (\Im g) e^{-rt_{1}/2} \sin(\Delta m) t_{1} \} ;$$

$$\Gamma = \Gamma_{s} + \Gamma_{t} ;$$

$$- 2M \text{ post-diction } g(t_{2}) = e^{-i(\lambda_{s} - \lambda_{s})t_{2}}$$

➤ Extract the relative PROBABILITY AMPLITUDE

#### THE K<sub>s</sub>-TAG

Decoherence is reached for large Δt before the observation of the second decay

$$e^{-\Delta\Gamma\Delta t/2}/|\eta_2| \ll 1$$

#### leading to a pure Ks-beam

- Most rewarding:  $\mathcal{L}^p$  and  $\langle K_L | K_S \rangle \neq 0 \rightarrow \text{No decay channel able to tag either K}_1$  or K<sub>s</sub>
- After 58 years of CPV: this POST-TAG condition in times is the only way to study rare  $K_s$ -decays. Compare with 60 year history of  $K_L$  decays!
- $\triangleright$  Example: Difference of charge Asymmetries  $A_1 A_5 \rightarrow$  Direct test of CPT!

#### CONCLUSION

- $\succ$ Entanglement in particle anti-particle system  $M^0-\overline{M}^0$
- NOVEL EFFECTS Tools for Particle Physics

  Quantum Phenomena
- ➤ Solution for NO GO's 

  TR for Unstable Particles

  K<sub>s</sub> tag
- ➤ POST-TAG of the past-decayed state depending on what and when measurement on the partner in the future.
- ➤In Classical and Quantum Physics, Time is a parameter to describe the evolving definite reality, not an observable.
- With the surviving correlation-in-time, Einstein would claim:
  "A Spooky Action to the Past"

#### NO (UNKNOWN) CAUSAL EFFECT

- CAUSAL INFLUENCE says that the cause must precede the effect according to ALL inertial observers, so that for the Post-Tag effect in the entangled K-mesons system –in which there are both time-like and space- like intervals,
- If the Interval is time-like, future is future for all observers → the future to past "influence" is NOT CAUSAL.
- If the Interval is space-like, there could be observers exhanging future and past, BUT the
  two events could only connect with a signal velocity higher than the speed of light → this
  "influence" is NOT CAUSAL.
- ➤ Then, independent of the space-time interval between the future observation in CM of the second decay and the past state of the partner, "the Post-Tag correlation in time" effect CANNOT BE A CAUSAL INFLUENCE.
- ➤ Whereas the EPR correlation between observables NEEDS a space-like interval to ensure no causal influence, the Post-Tag effect cannot be a causal influence for ALL cases → no loopholes. This is an additional argument, besides the fact that TIME IS NOT AN OBSERVABLE, to skate that the Post-Tag effect goes beyond the EPR correlation.

#### FOR PHILOSOPHERS ..... EPISTEMOLOGY

Physics -> QM correctly describes the behaviour of nature when it is observed

Scientific Methodology

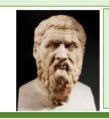
**Philosophy** > What QM says about nature's reality?



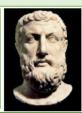
- Spooky Action at a Distance
- EPR Correlation-Bell Theorem
- Lack of Local Realism

- Spooky Action to the Past
- Surviving Correlation-in-time
- Lack of Instant Realism

(x,t) is not a definite, separate event  $\longrightarrow$  Role of time in QM?



TIME versus REALITY
Heraclitus vs. Parmenides



# THANK YOU VERY MUCH FOR YOUR ATTENTION

