## <span id="page-0-0"></span>Bell inequality violations: the QBist view

### Rüdiger Schack Royal Holloway, University of London

October 2024

つくい

Rüdiger Schack Royal Holloway, University of London [Bell inequality violations: the QBist view](#page-41-0)

## The B in QBism

Rüdiger Schack Royal Holloway, University of London [Bell inequality violations: the QBist view](#page-0-0)

メロメメ 御 メメ きょく きょう

E

### • Bayesian?

Rüdiger Schack Royal Holloway, University of London [Bell inequality violations: the QBist view](#page-0-0)

メタメ メミメ メミメ

そロト

E

### Bayesian? NO

Rüdiger Schack Royal Holloway, University of London [Bell inequality violations: the QBist view](#page-0-0)

 $\sqrt{m}$  )  $\sqrt{m}$  )  $\sqrt{m}$  )

そロト

E

Bayesian? NO

### • Bohr?

 $\sqrt{m}$  )  $\sqrt{m}$  )  $\sqrt{m}$  )

そロト

E

- Bayesian? NO
- Bohr? NO

 $\sqrt{m}$  )  $\sqrt{m}$  )  $\sqrt{m}$  )

そロト

E

- Bayesian? NO
- Bohr? NO
- Bruno de Finetti?

 $\langle \overline{m} \rangle$  and  $\overline{m}$  and  $\overline{m}$  and  $\overline{m}$ 

**ALCOHOL** 

 $299$ 

э

- Bayesian? NO
- Bohr? NO
- **•** Bruno de Finetti? Better

**ALCOHOL** 

**何 ▶ ( ヨ ト ( ヨ ト** 

 $299$ 

∍

- Bayesian? NO
- Bohr? NO
- **•** Bruno de Finetti? Better
- **Bettabilitarian?**

4 17 18

**何 ▶ ( 三 ) ( 三 )** 

 $2990$ 

∍

- Bayesian? NO
- Bohr? NO
- **•** Bruno de Finetti? Better
- Bettabilitarian? Excellent, but it won't catch on...

 $\mathbf{A}$  . The first set of  $\mathbf{A}$ 

- Bayesian? NO
- Bohr? NO
- **•** Bruno de Finetti? Better
- Bettabilitarian? Excellent, but it won't catch on...
- $\bullet$  B?

**ARACTE** 

- Bayesian? NO
- Bohr? NO
- **•** Bruno de Finetti? Better
- Bettabilitarian? Excellent, but it won't catch on...
- B? YES! (QBism is a noun, not an acronym)

## QBism in 2 words

Rüdiger Schack Royal Holloway, University of London [Bell inequality violations: the QBist view](#page-0-0)

メロメメ 御 メメ きょく きょう

E

## The world is bettable.

 $\sim$   $\sim$ 

医尿管的尿管

 $2990$ 

∍

## Personalist decision theory



### Bayes 1755 de Finetti 1931 Savage 1954

# $\left\{ \bigoplus_k k \bigoplus_k k \bigoplus_k k \right\}$

4 17 18

つくへ

Rüdiger Schack Royal Holloway, University of London [Bell inequality violations: the QBist view](#page-0-0)

### QBism, the Perimeter of Quantum Bayesianism

Christopher A. Fuchs Perimeter Institute for Theoretical Physics Waterloo, Ontario N2L 2Y5, Canada

cfuchs@perimeterinstitute.ca



 $\sim$   $\sim$ 

 $\triangleright$   $\rightarrow$   $\exists$   $\triangleright$   $\rightarrow$ 

 $\equiv$ 

 $QQ$ 

### Bell's theorem is the most famous example of what is now often called a no-go theorem.

医间周的

### The assumption of an ontological model:

For any measurement on a physical system, either the outcomes or their probabilities are determined by the system's real properties,  $\lambda$ . (Harrigan and Spekkens, 2007).

(Potentially misleading alternative labels for the same idea: "hidden variables", "realism".)

### Einstein 1927

Assuming  $\lambda$  (elements of physical reality) and locality (no spooky action at a distance) implies that  $\psi$  is not in one-to-one correspondence with  $\lambda$ .

### Einstein 1927

Assuming  $\lambda$  (elements of physical reality) and locality (no spooky action at a distance) implies that  $\psi$  is not in one-to-one correspondence with  $\lambda$ .

Einstein 1935 (letter to Schrödinger, not EPR)

Assuming  $\lambda$  and locality implies  $\psi$  is not determined by  $\lambda$ .

つくい

### Einstein 1927

Assuming  $\lambda$  (elements of physical reality) and locality (no spooky action at a distance) implies that  $\psi$  is not in one-to-one correspondence with  $\lambda$ .

Einstein 1935 (letter to Schrödinger, not EPR)

Assuming  $\lambda$  and locality implies  $\psi$  is not determined by  $\lambda$ .

Recent no-go theorems (e.g., Pusey, Barrett & Rudolph)

Assuming  $\lambda$  plus further assumptions implies  $\psi$  is determined by  $\lambda$ .

∢何 ▶ ∢ ヨ ▶ ∢ ヨ ▶

つくい

### Einstein 1927

Assuming  $\lambda$  (elements of physical reality) and locality (no spooky action at a distance) implies that  $\psi$  is not in one-to-one correspondence with  $\lambda$ .

Einstein 1935 (letter to Schrödinger, not EPR)

Assuming  $\lambda$  and locality implies  $\psi$  is not determined by  $\lambda$ .

Recent no-go theorems (e.g., Pusey, Barrett & Rudolph)

Assuming  $\lambda$  plus further assumptions implies  $\psi$  is determined by  $\lambda$ .

**ADALER (B)** 

 $QQ$ 

### Bell

Assuming  $\lambda$  and locality contradicts quantum mechanics.

#### Consider the state  $|\psi^{AB}\rangle=\frac{1}{\sqrt{2}}$  $\frac{1}{2}(|0\rangle|0\rangle+|1\rangle|1\rangle),$

where  $|0\rangle$  and  $|1\rangle$  are the eigenstates of the spin Z operator.

伊 ▶ イヨ ▶ イヨ ▶

#### Consider the state  $|\psi^{AB}\rangle=\frac{1}{\sqrt{2}}$  $\frac{1}{2}(|0\rangle|0\rangle+|1\rangle|1\rangle),$

where  $|0\rangle$  and  $|1\rangle$  are the eigenstates of the spin Z operator.

#### Now,  $\frac{1}{\sqrt{2}}$  $\frac{1}{2}(|0\rangle|0\rangle+|1\rangle|1\rangle)=\frac{1}{\sqrt{2}}$  $\frac{1}{2}(|+\rangle|+\rangle + |-\rangle|-\rangle),$

where  $\ket{\pm} = \frac{1}{\sqrt{2}}$  $\frac{1}{2}(|0\rangle\pm|1\rangle)$  are the eigenstates of the spin  $X$ operator.

モー イモン イミン イ野

Consider the state 
$$
|\psi^{AB}\rangle = \frac{1}{\sqrt{2}}(|0\rangle|0\rangle + |1\rangle|1\rangle),
$$

where  $|0\rangle$  and  $|1\rangle$  are the eigenstates of the spin Z operator.

Now, 
$$
\frac{1}{\sqrt{2}}(|0\rangle|0\rangle+|1\rangle|1\rangle) = \frac{1}{\sqrt{2}}(|+\rangle|+\rangle+|-\rangle|-\rangle),
$$

where  $\ket{\pm} = \frac{1}{\sqrt{2}}$  $\frac{1}{2}(|0\rangle\pm|1\rangle)$  are the eigenstates of the spin  $X$ operator.

Let  $\ket{\psi^B}$  be the conditional state after a measurement on A:

す 何 ト す ヨ ト す ヨ ト

つくい

Consider the state 
$$
|\psi^{AB}\rangle = \frac{1}{\sqrt{2}}(|0\rangle|0\rangle + |1\rangle|1\rangle),
$$

where  $|0\rangle$  and  $|1\rangle$  are the eigenstates of the spin Z operator.

Now, 
$$
\frac{1}{\sqrt{2}}(|0\rangle|0\rangle+|1\rangle|1\rangle) = \frac{1}{\sqrt{2}}(|+\rangle|+\rangle+|-\rangle|-\rangle),
$$

where  $\ket{\pm} = \frac{1}{\sqrt{2}}$  $\frac{1}{2}(|0\rangle\pm|1\rangle)$  are the eigenstates of the spin  $X$ operator.

Let  $\ket{\psi^B}$  be the conditional state after a measurement on A:

 $\mathcal{A} \oplus \mathcal{B}$   $\mathcal{B}$   $\mathcal{B}$   $\mathcal{B}$   $\mathcal{B}$   $\mathcal{B}$   $\mathcal{B}$   $\mathcal{B}$ 

 $200$ 

**A** measures Z  $|\psi^B\rangle \in \{|0\rangle,|1\rangle\}$ 

Consider the state 
$$
|\psi^{AB}\rangle = \frac{1}{\sqrt{2}}(|0\rangle|0\rangle + |1\rangle|1\rangle),
$$

where  $|0\rangle$  and  $|1\rangle$  are the eigenstates of the spin Z operator.

Now, 
$$
\frac{1}{\sqrt{2}}(|0\rangle|0\rangle+|1\rangle|1\rangle) = \frac{1}{\sqrt{2}}(|+\rangle|+\rangle+|-\rangle|-\rangle),
$$

where  $\ket{\pm} = \frac{1}{\sqrt{2}}$  $\frac{1}{2}(|0\rangle\pm|1\rangle)$  are the eigenstates of the spin  $X$ operator.

Let  $\ket{\psi^B}$  be the conditional state after a measurement on A:

- **A** measures Z
- $\bullet$  A measures X.

 $|\psi^B\rangle \in \{|0\rangle,|1\rangle\}$  $|\psi^B\rangle \in \{ |+\rangle, |-\rangle \}$ 

す 何 ト す ヨ ト す ヨ ト

つくへ

### Let  $\ket{\psi^B}$  be the conditional state after a measurement on  $A$ :

- A measures Z.
- $\bullet$  A measures X.

 $|\psi^B\rangle \in \{|0\rangle,|1\rangle\}$  $|\psi^B\rangle \in \{ |+\rangle, |-\rangle \}$ 

つくへ

### Let  $\ket{\psi^B}$  be the conditional state after a measurement on  $A$ :

- A measures Z.
- $\bullet$  A measures X.

 $|\psi^B\rangle \in \{|0\rangle,|1\rangle\}$  $|\psi^B\rangle \in \{ |+\rangle, |-\rangle \}$ 

### Einstein:

"[...] the real state of  $(AB)$  consists precisely of the real state of  $A$  and the real state of  $B$ , which two states have nothing to do with one another. The real state of  $B$  thus cannot depend upon the kind of measurement I carry out on A."

### Let  $\ket{\psi^B}$  be the conditional state after a measurement on  $A$ :

- A measures Z.
- $\bullet$  A measures X.

 $|\psi^B\rangle \in \{|0\rangle,|1\rangle\}$  $|\psi^B\rangle \in \{ |+\rangle, |-\rangle \}$ 

つくへ

### Einstein:

"[...] the real state of  $(AB)$  consists precisely of the real state of  $A$  and the real state of  $B$ , which two states have nothing to do with one another. The real state of  $B$  thus cannot depend upon the kind of measurement I carry out on A."

### Implication, assuming locality (Caves,Fuchs,RS 2002):

 $|\psi^{B}\rangle$  is not a function of "the real state at  $B$ ", i.e.,  $|\psi^{B}\rangle$  is not a real property of the system at  $B$ .

## A choice: do you give up locality or  $\lambda$ ?

If you accept the validity of quantum mechanics, you have to give up either locality or  $\lambda$ , i.e., the assumption of an ontological model.

(There are many good reasons to accept the validity of quantum mechanics. For instance, loophole-free Bell tests.)

## A choice: do you give up locality or  $\lambda$ ?

If you accept the validity of quantum mechanics, you have to give up either locality or  $\lambda$ , i.e., the assumption of an ontological model.

(There are many good reasons to accept the validity of quantum mechanics. For instance, loophole-free Bell tests.)

QBism rejects  $\lambda$ , i.e., in QBism,

- quantum states
- measurement outcomes
- **•** probabilities

are not determined by a system's real properties.

Quantum mechanics is a theory of the world. It is concerned with properties of physical systems.

つくへ

Quantum mechanics is a theory of the world. It is concerned with properties of physical systems.

QBism:

Quantum mechanics is a decision theory. It guides agents in their actions. (But its mathematical form tells us about the character of the world. QBism is a form of "participatory realism".)



### Agents are entities that

• can take actions freely on parts of the world external to themselves

so that

• the consequences of their actions matter to them.

つくへ



### Agents are entities that

• can take actions freely on parts of the world external to themselves

so that

• the consequences of their actions matter to them.

Users of quantum mechanics are agents

capable of applying the quantum formalism normatively.

A measurement is modeled by unitary interaction between a system and a meter,

 $\rho \otimes \vert 0 \rangle \langle 0 \vert \longrightarrow U(\rho \otimes \vert 0 \rangle \langle 0 \vert) U^{\dagger} \; ,$ 

followed by a readout of the meter. The outcome is objective.

A measurement is modeled by unitary interaction between a system and a meter,

 $\rho \otimes \vert 0 \rangle \langle 0 \vert \longrightarrow U(\rho \otimes \vert 0 \rangle \langle 0 \vert) U^{\dagger} \; ,$ 

followed by a readout of the meter. The outcome is objective.

QBism:

A measurement is an action an agent takes on a system. The meter is an extension of the agent. Outcomes as well as outcome probabilities are personal to the agent.

Quantum mechanics describes the world from an agent-independent perspective. Third person.

QBism:

The quantum formalism is a tool that I can use to make decisions regarding the consequences for me of my measurement actions. First person.

## Quantum dynamics

The mainstream approach:

Unitary evolution is fundamental and well understood, but there is a "measurement problem".

つくへ

## Quantum dynamics

The mainstream approach:

Unitary evolution is fundamental and well understood, but there is a "measurement problem".

QBism:

Measurement is fundamental. Unitary (and non-unitary) dynamics can be understood by analysing an agent's current decisions regarding future measurements.

## <span id="page-41-0"></span>Quantum dynamics

The mainstream approach:

Unitary evolution is fundamental and well understood, but there is a "measurement problem".

QBism:

Measurement is fundamental. Unitary (and non-unitary) dynamics can be understood by analysing an agent's current decisions regarding future measurements.

#### **Accepted Paper**

QBism's account of quantum dynamics and decoherence Phys. Rev. A

 $\mathcal{A}$  and  $\mathcal{A}$  . The set of  $\mathcal{B}$  is a set of  $\mathcal{B}$  is a set of  $\mathcal{B}$ 

 $\Omega$ 

John B. DeBrota, Christopher A. Fuchs, and Rúdiger Schack

Accepted 13 September 2024

# Thank you!

 $\leftarrow$ 

 $\mathbf{h}$ a. ∍  $\mathbf{h}$ э  $\mathbf{h}$  E