Electroweak and QCD measurements at e+e- colliders



Chris Hays, Oxford University









ECFA UK meeting 23 September 2024



Proposals for a future e+e- collider





ILC Candidate site in Kitakami, Tohoku

4 years @ \sqrt{s} = 90 GeV: 6 trillion Z bosons 2 years @ \sqrt{s} = 160 GeV: 240 million W bosons 3 years @ \sqrt{s} = 240 GeV: 1.5 million Higgs bosons 5 years @ \sqrt{s} = 365 GeV: 2 million top quarks

2 years @ \sqrt{s} = 90 GeV: 1.4 trillion Z bosons 1 year @ \sqrt{s} = 160 GeV: 80 million W bosons 7 years @ \sqrt{s} = 240 GeV: 2 million Higgs bosons

11 years @ \sqrt{s} = 250 GeV: 0.4 million Higgs bosons 2 years @ \sqrt{s} = 350 GeV: 1.3 million top quarks 9 years @ \sqrt{s} = 500 GeV: 0.2 million Higgs bosons with longitudinal polarization

ECFA physics studies for a future e⁺e⁻ collider

16/09/2024, 16:10

WG1 physics performance · Wiki · ECFA-Study / ECFA HiggsTopEW Factories · GitLab

WG1 physics performance

Identified 'focus topics' for study (arXiv:2401.07564)

Organization

• Coordinators: Jorge de Blas (Univ. Granada), Patrick Koppenburg (Nikhef), Jenny List (DESY), Fabio Maltoni (UC Louvain / Bologna)

Focus Groups

WG1 is organised in five focus groups.

Global interpretations (WG1-GLOB):

- Conveners: Jorge de Blas (Granada), Sven Heinemeyer (IFCA/IFT), Alexander Grohsjean (DESY), Junping Tian (Tokyo), Marcel Vos (Valencia)
- Dedicated wiki page

Precision (WG1-PREC):

- Conveners: Ayres Freitas (Pittsburgh), Paolo Azzurri (Pisa), Adrian Irles (Valencia), Andreas Meyer (DESY)
- Dedicated wiki page

http/09/g02ah.http/2.fr/ecfa-study/ECFA-HiggsTopEW-Factories/-/wikis/WG1-physics-performance · Wiki · ECFA-Study / ECFA HiggsTopEW Factories · GitLab

Higgs/Top/EW (WG1-HTE):

- Conveners: Chris Hays (Oxford), Karsten Köneke (Freiburg), Fabio Maltoni (Louvain)
- Dedicated wiki page

Flavour (WG1-FLAV):

- Conveners: David Marzocca (Trieste), Stephane Monteil (Clermont Ferrand), Pablo Goldenzweig (KIT)
- Dedicated wiki page

Searches (WG1-SRCH):

- Conveners: Roberto Franceschini (Rome III), Rebeca Gonzalez Suarez (Uppsala), Filip Zarnecki (Warsaw)
- Dedicated wiki page

HTE subgroup held series of meetings focussing on each \sqrt{s}

Electroweak and QCD measurements

Observable	Present value \pm error	FCC-ee stat.	FCC-ee syst.	Comment and dominant exp. error
$m_Z (keV/c^2)$	$91,186,700 \pm 2200$	5	100	From Z line shape scan Beam energy calibration
Γ_Z (keV)	$2,495,200 \pm 2300$	8	100	From Z line shape scan beam energy calibration
R^{Z}_{ℓ} (×10 ³)	$20,767\pm25$	0.06	0.2–1	Ratio of hadrons to leptons acceptance for leptons
$\alpha_{s} (m_{Z}) (\times 10^{4})$	1196 ± 30	0.1	0.4–1.6	From R_{ℓ}^{Z} above
$R_b (\times 10^6)$	$216,290 \pm 660$	0.3	< 60	Ratio of bb to hadrons stat. extrapol. from SLD
$\sigma_{\rm had}^0 \ (\times 10^3) \ ({\rm nb})$	$41,541 \pm 37$	0.1	4	Peak hadronic cross-section luminosity measurement
$N_{\nu} (\times 10^3)$	2991 ± 7	0.005	1	Z peak cross sections Luminosity measurement
$\sin^2 \theta_{\rm W}^{\rm eff}$ (×10 ⁶)	$231,480 \pm 160$	3	2–5	From $A_{FB}^{\mu\mu}$ at Z peak Beam energy calibration
$1/\alpha_{\rm QED}~(m_Z)~(\times 10^3)$	$128,952\pm14$	4	Small	From $A_{FB}^{\mu\mu}$ off peak
$A_{FB}^{b,0}$ (×10 ⁴)	992 ± 16	0.02	1–3	b-quark asymmetry at Z pole from jet charge
$A_{FB}^{pol,\tau}$ (×10 ⁴)	1498 ± 49	0.15	< 2	τ Polarisation and charge asymmetry τ decay physics
$m_W (MeV/c^2)$	$80,350 \pm 15$	0.5	0.3	From WW threshold scan Beam energy calibration
Γ_{W} (MeV)	2085 ± 42	1.2	0.3	From WW threshold scan beam energy calibration
$\alpha_{\rm s}~(m_{\rm W})~(\times 10^4)$	1170 ± 420	3	Small	From R^W_ℓ
$N_{\nu} (\times 10^3)$	2920 ± 50	0.8	Small	Ratio of invis. to leptonic in radiative Z returns
m _{top} (MeV/c ²)	$172,740 \pm 500$	17	Small	From tt threshold scan QCD errors dominate
Γ_{top} (MeV)	1410 ± 190	45	Small	From tt threshold scan QCD errors dominate
$\lambda_{top}/\lambda_{top}^{SM}$	1.2 ± 0.3	0.1	Small	From tt threshold scan QCD errors dominate
ttZ couplings	$\pm 30\%$	0.5-1.5%	Small	From $E_{CM} = 365 \text{ GeV run}$

Electroweak measurements

Precision EW measurements probe multi-TeV physics

Effective field theory used to check consistency and sensitivity across measurements





W boson mass measurement

achieving sub-MeV systematic precision is challenging

$$\Delta m_W(B) = \left(\frac{d\sigma}{dm_W}\right)^{-1} \left(\frac{\Delta\sigma_B}{\varepsilon} \oplus \Delta\sigma_{TH}\right)$$

Background and Theory

$$\Delta \sigma_{TH} < 1 \text{fb} \ (\Delta \sigma_{TH} / \sigma_{TH} < 2 \cdot 10^{-4})$$

$$\Delta \sigma_B / \varepsilon < 1 \text{fb} \ (\Delta \sigma_B / \sigma_B < 4 \cdot 10^{-3})$$

$$\Delta m_W(\varepsilon) = \sigma \left(\frac{d\sigma}{dm_W}\right)^{-1} \left(\frac{\Delta\varepsilon}{\varepsilon} + \frac{\Delta L}{L}\right)$$

Acceptance and Luminosity

$$\left(\frac{\Delta\varepsilon}{\varepsilon} \oplus \frac{\Delta L}{L}\right) < 2 \cdot 10^{-4}$$

$$\Delta m_W(E) = \left(\frac{d\sigma}{dm_W}\right)^{-1} \left(\frac{d\sigma}{dE}\right) \Delta E \le \frac{1}{2} \Delta E$$

ILC: 2.4 MeV statistical uncertainty

ILC polarised collisions : enhance (x4) t-channel WW production or suppress it to control background

Channel	Efficiency (%)	$\sigma^U_{\rm bkgd}$ (fb)	$A^B_{ m LR}$	Eff. syst. (%)	Bkgd syst.	$A^B_{\rm LR}$ syst.
lvlv	87.5	10	0.15	0.1	free	0.025
	av &hift i	n Marc	A391	$d h a^{0.1}$	free	0.012
	ဗေန္မျူး၊	1 12000 C	0.48		free	0.005
caused a polarized	by new s scan	hysics	at w	30 event Velectio	on near thres	hold using 7

$\begin{array}{c c c c c c c c c c c c c c c c c c c $			Channel	Effi	cie	ncy (%	b) $\mid \sigma_1$	_{okgd} (f	b)	$A_{\rm LR}^{\rm B}$	$_{l} \mid Eff.$	syst.	(%)	Bkgd syst.	A
$\begin{array}{c c c c c c c c c c c c c c c c c c c $	Collicion		lvlv		8	7.5		10		0.15	5	0.1		free	
$\frac{qqqq}{1} = \frac{83.5}{2.43} = \frac{200}{0.48} = 0.1 \text{free}$ Table 3: Experimental assumptions for the WW event selection near thresh a polarized scan $(Oxford)$ $= \frac{0.1704 + -}{0.0254 + +} = \frac{20}{19} = \frac{67}{27} = \frac{158}{6661} = \frac{139932}{6661}$ $= \frac{0.1704 + -}{161.2} = 21.739 = 0.7789 - + 16096 = 67610 = 73538 = 4635245$ $= \frac{20}{0002} = \frac{697141}{100} = \frac{222}{222} = \frac{242832}{222} = \frac{20}{222} = \frac{242832}{222} = \frac{20}{2230} = \frac{22}{22} = \frac{242832}{222} = \frac{20}{230} = \frac{22}{9} = \frac{242832}{222} = \frac{22}{22} = \frac{242832}{222} = \frac{20}{230} = \frac{22}{9} = \frac{242832}{222} = \frac{22}{22} = \frac{242832}{222} = \frac{24}{23} = \frac{2}{2} = \frac{242832}{22} = \frac{2}{2} = $	COMISION	IE	qqlv		8	7.5		40		0.30)	0.1		free	
Table 3: Experimental assumptions for the WW event selection near thresholds a polarized scan $ 0.1704 + - 20 & 67 & 158 & 139932 \\ 0.0254 + + 2 & 19 & 27 & 6661 \\ 0.0254 + + 2 & 19 & 27 & 6661 \\ 0.0254 & 21 & 100 & 102 & 8455 \\ 161.2 & 21.739 & 0.7789 - + & 16096 & 67610 & 73538 & 4635245 \\ 0.0254 & 21 & 100 & 102 & 8455 \\ 0.0254 & 21 & 100 & 102 & 8455 \\ 0.0254 & 16096 & 67610 & 73538 & 4635245 \\ 0.025 & 20 & 697141 \\ 0.0254 & & 16096 & 67610 & 73538 & 4635245 \\ 0.025 & 22 & 42832 \\ 0.025 & 22 & 42832 \\ 0.025 & 0.025 & 0.788 & 61 & 42979 \\ 0.025 & 0.012 & 0.788 & 0.80 & 03 & 697851 \\ 0.025 & 0.025 & 0.78 & 61 & 42979 \\ 0.005 & 0.025 & 0.78 & 61 & 42979 \\ 0.005 & 0.025 & 0.78 & 61 & 42979 \\ 0.005 & 0.025 & 0.78 & 61 & 42979 \\ 0.005 & 0.025 & 0.78 & 61 & 42979 \\ 0.005 & 0.025 & 0.78 & 61 & 42979 \\ 0.005 & 0.025 & 0.78 & 61 & 42689 \\ 0.1704 + + & 10 & 25 & 43 & 6633 \\ 0.1704 + + & 24 & 957 & 147 & 838233 \\ 0.1704 + & 106 & 451 & 466 & 41196 \\ 0.1704 + & 106 & 451 & 466 & 41196 \\ 0.1704 + & 106 & 451 & 466 & 41196 \\ 0.1704 + & 106 & 4$			qqqq		8	3.5		200		0.48	3	0.1		free	
$\frac{\begin{vmatrix} 0.1704 & +- & 20 & 67 & 158 & 139932 \\ 0.0254 & ++ & 2 & 19 & 27 & 6661 \\ 0.0254 & & 21 & 100 & 102 & 8455 \\ \hline 161.2 & 21.739 & 0.7789 & -+ & 16096 & 67610 & 73538 & 4635245 \\ \hline 161.2 & 21.739 & 0.7789 & -+ & 16096 & 67610 & 73538 & 4635245 \\ \hline 10012 & Background & 3.20 & 2.30 & 22 & 42832 \\ \hline Dolarization & 3.73 & 1.27 & 70 & 697459 \\ \hline Polarization & 3.73 & 1.27 & 70 & 697459 \\ \hline Polarization & 3.76 & 0.78 & 61 & 42979 \\ \hline Polarization & 3.76 & 0.78 & 61 & 42979 \\ \hline Luminosity & 3.76 & 0.78 & 61 & 42979 \\ \hline A_{LR}^B & 3.86 & 0.80 & 03 & 697851 \\ \hline Statistical & 2.43 & 74 & 33271 \\ \hline Systematic & 3.94 & 74 & 33271 \\ \hline 1005 \\ \hline 1 using \\ 7 & \hline 170.0 + 10007 & 10000 & 1000 & 1000 \\ \hline 0.0254 & & 46 & 135 & 141 & 8463 \\ \hline 0.0254 & & 46 & 135 & 141 & 8463 \\ \hline 0.0254 & & 46 & 135 & 141 & 8463 \\ \hline 0.0254 & & 46 & 135 & 141 & 8463 \\ \hline 0.0254 & & 46 & 135 & 141 & 8463 \\ \hline 0.0254 & ++ & 100 & 457 & 1447 & 838233 \\ \hline 0.012 & 0.012 & 0.012 & 0.012 & 0.012 & 0.012 & 0.012 & 0.012 & 0.012 & 0.012 & 0.012 & 0.0254 & & 466 & 135 & 141 & 8463 \\ \hline 10005 & 0.012 & 0.025 & 0.1700 & +++ & 100 & 25 & 43 & 6633 \\ \hline 10005 & 0.0254 & & 466 & 135 & 141 & 8463 \\ \hline 10005 & 0.0254 & & 466 & 135 & 141 & 8463 \\ \hline 0.0254 & ++ & 1000 & 451 & 264869 & 270577 & 5500286 \\ \hline 0.1704 & +- & 224 & 957 & 1447 & 838233 \\ \hline 0.0254 & ++ & 1000 & 451 & 466 & 40196 \\ \hline \end{array}$			Table 3: E a polarized	xper l sca	ime n	ental a	FC	C be	a fo y	m e r the Wilk	nerg kinso	v gro vent se n (O)	up lectio xfor	ed by n near thre d)	shol
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							0.0254			106	451	466	401	96	

W boson mass $M_{\rm Z}^2 = s \frac{\beta_1 \sin \theta_1 + \beta_2 \sin \theta_2 - \beta_1 \beta_2 |\sin(\theta_1 + \theta_2)|}{\beta_1 \sin \theta_1 + \beta_2 \sin \theta_2 + \beta_1 \beta_2 |\sin(\theta_1 + \theta_2)|}$

e⁺e⁻-collider measurements can also used reconstructed W-boson pairs

 $M_{\rm Z}^2 = s \frac{\beta_1 \sin \theta_1 + \beta_2 \sin \theta_2 - \beta_1 \beta_2 |\sin(\theta_1 + \theta_2)|}{\beta_1 \sin \theta_1 + \beta_2 \sin \theta_2 + \beta_1 \beta_2 |\sin(\theta_1 + \theta_2)|}$ E_{CM} is again a main ingredient: sets jet energy scale

 $(\sqrt{s} = 160 \text{ and } 240 \text{ GeV})$

other main ingredients are the jets (and lepton) angles secondary ingredients are the jet velocities ($\beta = p/E$)







W boson mass measurement



Focus topic team also considering single-W production at $\sqrt{s} = 91$ GeV

1 million W bosons expected

Graham Wilson, Josh Bendavid, Juergen Reuter, Keisho Hidaka, Martin Beneke, Raimund Strohmer, Simon Platzer, Stefan Dittmaier

Forward backward asymmetries

FCC	Observable	Present value \pm error	FCC-ee stat.	FCC-ee syst.	Comment and dominant exp. error
Z pole	$\sin^2 \theta_{\rm W}^{\rm eff}$ (×10 ⁶)	$231,480 \pm 160$	3	2–5	From $A_{FB}^{\mu\mu}$ at Z peak Beam energy calibration
	$A_{FB}^{b,0}$ (×10 ⁴)	992 ± 16	0.02	1–3	b-quark asymmetry at Z pole from jet charge
	$A_{FB}^{pol,\tau}~(\times 10^4)$	1498 ± 49	0.15	< 2	τ Polarisation and charge asymmetry τ decay physics

aggressive targets: systematic uncertainties are key

A_{FB}^b systematics studies ongoing

We know that statistical uncertainty will not be an issue

- LEP combination has ~equal stat and syst contributions
- We expect ~10⁵ times more statistics at FCC-ee \Rightarrow ~300 times smaller stat. uncertainty

Systematic uncertainties expected to be dominant

- Modelling b-fragmentation
 - Affecting B-hadron kinematics
- Final-state QCD radiation effects
 - Affecting jet shapes, distribution of charge, B-hadron kinematics...
- b-tagging efficiency:
 - Uncertainty on mis-tag rate affecting background prediction
 - $p_{\rm T}$ and η dependency of b-tagging eff. for signal

Giovanni Guerrieri, FCC workshop Krakow

Forward backward asymmetries

Per mil level statistical uncertainties reachable for

• Fragmentation, angular correlations → minimized

the nominal ILC250-500 program

Smaller exp syst. Uncertainties

ILC studying sensitivity to new physics off the Z-pole

- At least 4 observables for AFB at ILC250 per energy point
 - 2 quarks and 2 polarisations (eLpR, eRpL)



Forward backward asymmetries

Expansion to light flavours possible with particle id in the ParticleNet jet tagger



FIG. 1. Event display of an $e^+e^- \rightarrow \nu\bar{\nu}H \rightarrow \nu\bar{\nu}gg$ ($\sqrt{s} = 240 \text{ GeV}$) event simulated and reconstructed with the CEPC baseline detector [17]. Different particles are depicted with colored curves and straight lines: red for e^{\pm} , cyan for μ^{\pm} , blue for π^{\pm} , orange for photons, and magenta for neutral hadrons.

H Liang, M Ruan et al, 2310.03440

	Predicted											
		b	$\frac{1}{b}$	C	$\frac{1}{C}$	S	$\frac{1}{S}$	u	$\frac{1}{u}$	d	$\frac{1}{d}$	Ġ
	G -	0.015	0.014	0.024	0.024	0.052	0.052	0.043	0.041	0.034	0.034	0.667
	d -	0.003	0.003	0.020	0.012	0.092	0.112	0.219	0.076	0.079	0.272	0.113
	d -	0.003	0.003	0.012	0.019	0.112	0.092	0.082	0.207	0.277	0.079	0.112
	u -	0.003	0.003	0.011	0.019	0.132	0.043	0.062	0.356	0.178	0.081	0.111
	u -	0.002	0.003	0.020	0.011	0.044	0.131	0.367	0.055	0.080	0.174	0.111
בכו	<u>-</u>	0.003	0.003	0.018	0.020	0.102	0.542	0.084	0.028	0.045	0.062	0.094
	s -	0.003	0.002	0.020	0.018	0.543	0.102	0.030	0.080	0.063	0.045	0.092
	- 7	0.016	0.015	0.056	0.739	0.032	0.037	0.009	0.026	0.017	0.010	0.043
	с -	0.015	0.014	0.743	0.055	0.036	0.031	0.025	0.009	0.009	0.018	0.043
	b -	0.170	0.737	0.026	0.033	0.003	0.004	0.003	0.002	0.002	0.003	0.018
	b -	0.745	0.163	0.033	0.025	0.004	0.003	0.002	0.003	0.002	0.002	0.017

TwoF team to report at ECFA Paris workshop E Bagnaschi, A Irles, D Jeans, A Vicini

Differential WW measurements

Snowmass study with HEPfit used optimal observables to estimate aTGC sensitivity





Electron couplings using transverse spin asymmetry

(a)

Zh

📫 W+W

(c)

 $e[\Gamma_{\gamma}^{e}] \times 10^{2}$

 $\operatorname{Im}[\Gamma_{\gamma}^{e}] \times 10^{2}$

 $\frac{1}{-2}$

Opposite Spin

-1

Opposite Spin

0

 $\text{Re}[\Gamma_Z^e] \times 10^3$

 $\mu^+\mu^-$

Zh Zγ

W⁺W

μ⁺μ⁻ Zh

Recent study proposes transversely polarised beams to constrain dipole interactions

Unpolarized beams: dominant sensitivity to new physics decreases as $1/\Lambda^4$ Due to helicity flip in interference term

Transversely polarized beams: sensitive to interference with SM ($1/\Lambda^2$)

2 Aligned Spin

 $\mathrm{Re}[\Gamma_{\gamma}^{e}]{\times}10^{2}$

 $\mathrm{Im}[\Gamma_{\gamma}^{e}]{\times}10^{2}$

-1

Aligned Spin

0

 $\text{Re}[\Gamma_Z^e] \times 10^3$

Spin dependent amplitude square:

 $(0 \downarrow)$

1 1

U

$$|\mathcal{M}|^2 = \rho_{\alpha_1 \alpha_1'}(\boldsymbol{s}) \rho_{\alpha_2 \alpha_2'}(\bar{\boldsymbol{s}}) \mathcal{M}_{\alpha_1 \alpha_2}(\phi) \mathcal{M}^*_{\alpha_1' \alpha_2'}(\phi)$$

$$oldsymbol{s} = (b_1, b_2, \lambda) = (b_{\mathrm{T}} \cos \phi_0, b_{\mathrm{T}} \sin \phi_0, \lambda)$$

$$\rho = \frac{1}{2} \left(1 + \boldsymbol{\sigma} \cdot \boldsymbol{s} \right) = \frac{1}{2} \begin{pmatrix} 1 + \lambda & b_{\mathrm{T}} e^{-i\phi_0} \\ b_{\mathrm{T}} e^{i\phi_0} & 1 - \lambda \end{pmatrix}$$

$$\mathcal{M}_{\lambda_1,\lambda_2}(\theta,\phi) = e^{-\frac{1}{2}\phi} \mathcal{T}_{\lambda_1,\lambda_2}(\theta)$$

Term linear in b_T interferes with NF

 $i(\lambda_1 - \lambda_2)\phi \boldsymbol{\tau}$

(n)

			F 🥻 🛛 🖊 🕹
U	L	$T = \frac{2\gamma}{W^+W^-}$	
$ \mathcal{M} ^2_{UU} \to 1$	$ \mathcal{M} ^2_{UL} \to 1$	$ \mathcal{M} ^{2}_{UT^{\mathrm{Im}}\overline{[\Gamma_{Z}]} imes} \cos \phi, \sin^{2} \phi = 1$	$\begin{bmatrix} 2 & -1 & 0 \\ & & \text{Im}[\Gamma_Z^e] \times 10^3 \end{bmatrix}$
$ \mathcal{M} _{LU}^2 o 1$	$ \mathcal{M} ^2_{LL} ightarrow 1$	$ \mathcal{M} _{LT}^2 o \cos \phi, \sin \phi$]
			1





 $egin{aligned} \lambda & b_1 - ib_2 \ ib_2 & 1 - \lambda \end{pmatrix} &= rac{1}{2} egin{pmatrix} 1+\lambda & b_{\mathrm{T}}e^{-i\phi_0} \ b_{\mathrm{T}}e^{i\phi_0} & 1-\lambda \end{pmatrix} \end{pmatrix} &= rac{1}{2} egin{pmatrix} 1 \ b_{\mathrm{T}} \end{pmatrix} \end{aligned}$

X-K Wen, B Yan, Z Yu, C-P Yuan, 2307.05236

Neutrino anomalous magnetic moment

Z-pole run could improve ν_{τ} magnetic moment constraints by two orders of magnitude

$$\mathscr{L}_{V} = \mu_{B} \sum_{i=e,v,\tau} k_{i} [v_{i} \sigma^{\mu v} v_{i}] F_{\mu v}$$

Bohr magneton

$$\Gamma(Z \to v \bar{v} \gamma) = \frac{\mu_B^2 \alpha m_Z^3 \left(\sum_i |k_i|^2\right)}{512\pi^2 c_W^2 s_W^2}$$

From PDG

$$k_e < 2.9 \times 10^{-11}$$

 $k_\mu < 6.8 \times 10^{-10}$
 $k_\tau < 3.9 \times 10^{-7}$

$$BR(Z \to v \bar{v} \gamma) = 2.3 \times 10^5 k^2$$

Improving bounds on MDM of neutrino-tau



Emidio Gabrielli, HTE meeting

Neutrino four-fermion interactions



QCD measurements

QCD measurements crucial to achieving precision physics goals

e.g. fragmentation & jet shapes

Heavy flavour production and fragmentation identified as a focus topic (BCfrag and Gsplit)

• anticipated (perturbative) theory precision: $\leq 1\%$

(personal take!)

- at least $\mathcal{O}(\alpha_{s}^{3})$ corrections for QCD event shapes
- at least $\mathcal{O}(\alpha_5^2)$ corrections for QCD final states with up to four jets

(I think we'll see the first of such calculations in the next 2-3 years)

- systematic inclusion of $\mathcal{O}(\alpha_{EW})$ in a multiplicative scheme and a lot of mixed $\mathcal{O}(\alpha_{S}\alpha_{EW})$ corrections
- NNLL parton shower matched to NNLO QCD

soft physics effects may dominate theory uncertainties: no first-principles theory \longrightarrow must measure!

> Frank Krauss (IPPP), ECFA workshop DESY

Heavy flavour production and fragmentation

b quark fragmentation function $f(x_{B}^{\text{weak}})$ *b*-quark fragmentation function $f(x_B^{weak})$ *b* quark fragmentation function $f(x_B^{\text{weak}})$ $\frac{1}{N} \frac{dN}{N} \frac{dx_B}{dN} \frac{dx_B}{dx_B}$ ation function $f(x_n^v)$ Analysis Analysis Analysis Analysis nalvsis Analysis nalysis nalvsis Analysis Analysis Analysis Σ 0.2 0.3 0.4 0.5 0.7 0.2 0.3 0.4 0.5 0.7 0.9 0.4 0.5 0.7 0.6 0.8 0.6 0.8 0.3 0.6 0.8 0.9 χ_{R} x_B

 ${\scriptstyle \bullet}$ disagreement in $b\mbox{-}quark$ fragmentation measurements

(look at results from ALEPH, OPAL, SLD \longrightarrow need to chose one!)

• g
ightarrow Q ar Q splitting tricky in parton showers

Frank Krauss (IPPP), ECFA workshop DESY

- measurement strategy:
 - "Mercedes star" with two id'd heavy quark jets
 - \longrightarrow third jet is gluon jet
 - jet-shape measurements: sub-jettiness & friends
 - hadron yields inside jet
 - leading hadron identity $/x_p$
 - di-baryon/di-strange correlations inside jet



Heavy flavour production and fragmentation

List of important observables produced by Torbjorn Sjostrand and the focus topic team

Observable	e^+e^-	pp
Event shapes and angular distributions		
Inclusive B/D production cross section	primary production is well known from theory, so any "excess" is from gluon split- ting	combines primary production, gluon splitting, and MPI (multiparton interactions) contributions, each with signific- ant theoretical uncertainties
Flavour composition as far back in decay chains as can be traced (even equal D^{*0} and D^{*+} rates gives unequal D^0 and D^+ ones)	we do not expect sizeable momentum dependence, but interesting to contrast mesons and baryons for smaller ones	significant $p_{\rm T}$ dependence observed and to be studied fur- ther, also high- vs. low-multiplicity events, rapidity,, which is important for development/tuning of colour recon- nection models
Particle-antiparticle production asymmetries	none expected, except tiny from CP- violation in oscillations	asymmetries expected and observed from p flavour content, increasing at larger rapidities; relates to how string (and cluster?) fragmentation connects central rapidities to beam remnants
Momentum spectra	dn/dx_E with $x_E = 2E_{had}/E_{cm}$; basic distribution for tuning of "fragmentation function"	$dn/dp_{\rm T}$ and dn/dy give basic production kinematics, but the many production channels give less easy interpretation
Energy flow around B/D hadrons, excluding the hadron itself, as a test that dead cone effects are correctly described	$dE/d\theta$ where θ is the distance from B/D on the sphere	$dp_{\rm T}/dR$ where R is the distance in (η,ϕ) or (y,ϕ) space, only applied for B/D above some $p_{\rm T}$ threshold
B/D hadron momentum fraction of total E or $p_{\rm T}$ in a jet, with $x = p_{\rm T}^{\rm had}/p_{\rm T}^{\rm jet}$, as a test of the fragmentation function combined with almost collinear radiation, suitably for some slices of $p_{\rm T}$ (and in addition with a veto that no other B/D should be inside the jet cone, so as to suppress the gluon splitting contribution)	draw a jet cone in θ around B/D and measure x	draw a jet cone in R around B/D and measure x
B/D hadron multiplicity, as a measure of how often several pairs are produced		
Separation inside B/D pairs, where large sep- aration suggests back-to-back primary production, while small separation suggests gluon splitting	separation in θ	separation both in ϕ and in R , since for primary produc- tion $\phi = \pi$ is hallmark with η/y separation less interesting, while gluon splitting means R is small while ϕ and y/η in- dividually are less interesting
Hardness difference within (reasonably hard) pairs, $\Delta = (p_{\rm T}^{\rm max} - p_{\rm T}^{\rm min})/(p_{\rm T}^{\rm max} + p_{\rm T}^{\rm min})$, where for gluon splitting $r^2 + (1 - r)^2$ translates to $1 + \Delta^2$	separately for small or large θ	separately for large or small ϕ

Eli Ben Haim, Loukas Gouskos, Simon Platzer, Andrzej Siodmok, Maria Ubiali (Cambridge)

Top-quark measurements

Top quark measurements can enter new realm of precision (TTthresh focus topic)



Top-quark measurements

Ongoing FCC effort to update experimental fit

Study detector-level distributions, use MVA to suppress backgrounds



- A simultaneous fit of m_t, total width, and y_t seems possible based on a threshold scan of [-4,+1] GeV around the threshold
- 30 (50) MeV shift in mass (width) induce a 4% shift in the xsec
- 10% shift in Yukawa produce a ~1% effect just above threshold
- Limited impact from aS assuming expected precision at Z pole
- Some residual sensitivity to y_t above threshold -> we will investigate the possibility of one additional scan point well above threshold (continuum)
- Presented studies do not include impact from <u>ISR</u> and beam energy resolution, which can be significant -> will be included at next update



Wh

WAN

Global EFT constraints



Eugenia Clelada, Alejo Rossia, Marion Thomas, Eleni Vryonidou (Manchester), Luca Mantani (Cambridge), Tommaso Giani, Jaco ter Hoeve, Juan Rojo

Summary

The proposed e⁺e⁻ colliders can provide a step-change in precision More than an order of magnitude in many cases

Broad range of physics topics within electroweak and QCD physics LEP programme produced more than 1700 papers with no Higgs or top measurements Truncated ECFA studies are only the tip of the iceberg

Active FCC and ILC physics working groups



Michael Peskin, Aidan Robson (Glasgow), Junping Tan, ILC physics group conveners

https://agenda.linearcollider.org/event/9154/

FCCeePhysicsPerformance

Welcome to the FCC-ee Physics Performance Documentation

Table of Contents

1. Organisation

- 2. Towards the definition of detector requirements
- 3. List of Active Case studies (evolving)
- 4. General information for FCC-ee analyses
- 5. LOIs submitted to Snowmass
- 6. Software

Organisation

Coordinators

- Patrizia Azzi (INFN Padova) Patrizia.Azzi@cern.ch
- Emmanuel Perez (CERN) Emmanuel.Perez@cern.ch
- Michele Selvaggi (CERN) michele.Selvaggi@cern.ch

Physics Performance meetings

O(monthly) meetings: Mondays, 3pm-5pm, CERN time. Usually the third Monday of each month.

Exotic Z decays



Landau-Yang theorem forbids Z → 2 photons -> amplitude vanishes
 avoided due to distinguishability of photon and dark-photon interaction (blob)

Top quark FCNCs at the LHC and e^+e^- colliders



top physics opportunities at a new e⁺e⁻ collider – CP violation in the top sector

Marcel Vos,

IFIC, CSIC/UV, Valencia, Spain

CP-odd (imaginary parts of) operators at the LHC



	C _t	W	C_{itW}		
	68% CL	95% CL	68% CL	95% CL	
All terms	[-0.3, 0.8]	[-0.9, 1.4]	[-0.5, -0.1]	[-0.8, 0.2]	
Order $1/\Lambda^4$	[-0.3, 0.8]	[-0.9, 1.4]	[-0.5, -0.1]	[-0.8, 0.2]	
Order $1/\Lambda^2$	[-0.3, 0.8]	[-0.8, 1.5]	[-0.6, -0.1]	[-0.8, 0.2]	

CP violating analysis: construct dedicated CP-odd triple-product observables

$$\mathcal{O}_{+}^{Re} = (\hat{\mathbf{q}}_{\bar{X}} \times \hat{\mathbf{q}}_{+}^{*}) \cdot \hat{\mathbf{p}}_{+},$$

$$\mathcal{O}_{+}^{Im} = -[1 + (\frac{\sqrt{s}}{2m_{t}} - 1)(\hat{\mathbf{q}}_{\bar{X}} \cdot \hat{\mathbf{p}}_{+})^{2}]\hat{\mathbf{q}}_{+}^{*} \cdot \hat{\mathbf{q}}_{\bar{X}} + \frac{\sqrt{s}}{2m_{t}}} \hat{\mathbf{q}}_{\bar{X}} \cdot \hat{\mathbf{p}}_{+} \hat{\mathbf{q}}_{+}^{*} + \hat{\mathbf{p}}_{+}.$$
ECCA users Lines/top/EW factories, Oct '22 9
leptons from t - Wb - Wb decay, and qX that of the hadronic top quark

$$\Gamma_{\mu}^{t\bar{t}X}(k^{2}, q, \bar{q}) = ie \left\{ \gamma_{\mu} \left(F_{11Y}^{X}(k^{2}) + \gamma_{5} r^{*X} \cdot 2\lambda - \sigma_{\mu\nu} \cdot \dots \cdot \lambda\nu \cdot r^{*X} \cdot r^{2}\lambda - r^{*X} \cdot r^{2}\lambda - \sigma_{\mu\nu} \cdot \dots \cdot \nu\nu \cdot r^{*X} \cdot r^{2}\lambda - \sigma_{\mu\nu} \cdot \dots \cdot \nu\nu \cdot r^{*X} \cdot r^{2}\lambda - \sigma_{\mu\nu} \cdot (r^{*X}(k^{2})) - \frac{\sigma_{\mu\nu}}{2m_{t}}(q + \bar{q})^{\nu} (iF_{2Y}^{X}(k^{2}) + \gamma_{5}F_{2X}^{*A}(k^{2}))$$
ECFA wisp Higgstop/EW factories, Oct '
ECFA wisp Higgstop/EW factories, Oct '

$$\int_{0}^{\pi} \frac{1}{\sigma_{\mu\nu}} \frac{1}{\sigma$$

Bernreuther & Chen, in arXiv:1710.06737

ECFA wksp Higgs/top/EW factories, Oct '22 13