## New Developments in Neutrino (Astro) Physics

### Joachim Kopp (CERN & JGU Mainz) PASCOS 2025 • Durham • 21 July 2025



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solar neutrinos ★ stellar evolution



solar neutrinos ★ stellar evolution

supernova neutrinos ★ nucleosynthesis ★ matter under extreme conditions ★ stellar evolution





solar neutrinos ★ stellar evolution

### high-E neutrinos ★ origin of cosmic rays ★ AGNs, blazars, MW

supernova neutrinos ★ nucleosynthesis ★ matter under extreme conditions ★ stellar evolution





solar neutrinos ★ stellar evolution

### high-E neutrinos ★ origin of cosmic rays ★ AGNs, blazars, MW

cosmology ★ early Universe supernova neutrinos ★ nucleosynthesis ★ matter under extreme conditions ★ stellar evolution





solar neutrinos  $\star$  stellar evolution

### high-E neutrinos ★ origin of cosmic rays ★ AGNs, blazars, MW

cosmology ★ early Universe supernova neutrinos ★ nucleosynthesis ★ matter under extreme conditions ★ stellar evolution

neutron stars ★ cooling common-envelope systems







 supernova neutrinos
 nucleosynthesis
 matter under extreme conditions
 stellar evolution

neutron stars
 \* cooling
 \* common-envelope
 systems





neutron stars
 \* cooling
 \* common-envelope
 systems



### **Neutron Stars**



Image: Gendreau et al.





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□ kinematically forbidden except in the heaviest stars
 □ condition  $p_{Fn} < p_{Fp} + p_{F\ell}$ 





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□ kinematically forbidden except in the heaviest stars
 □ condition  $p_{Fn} < p_{Fp} + p_{F\ell}$ 

Modified Urca Processes



□ allowed in all neutron stars

J Q JG U



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## Neutrinos from Neutron Stars in Chemical Equilibrium



Regular modified Urca (in equilibrium) at 10 kpc: 38 cm<sup>-2</sup> sec<sup>-1</sup> 

■ large flux, but low energy ■ so far undetectable





# in young neutron stars (T ~ yrs): $E_v \sim 100$ keV, $\phi \sim 10^{41}$ erg/sec

# for comparison: diffuse SN neutrinos: ~ 1 cm<sup>-2</sup> sec<sup>-1</sup> at $E_v$ ~ MeV

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### **Neutrinos from Neutron Stars**

neutron stars evolve:
spin-down / spin-up
accretion
expulsion of *B*-fields
tidal deformation



Result:
out-of-equilibrium Urca processes
extra neutrinos

JK Opferkuch arXiv:2312.08457



## **Neutron Stars Away from Thermal Equilibrium**

- neutrino flux can be enhanced by several orders of magnitude
- but low energy still precludes detection so far
- opportunities for large low-threshold DM detectors?















## **Common-Envelope Evolution**

- compact star (neutron star, black hole, white dwarf, ...) enters companion star
- significant friction
- gigantic accretion rates (up to 0.1  $M_{\odot}$ /yr for several months)
- outcome: Thorne–Żytkov object or explosion
- crucial for the formation of gravitational wave sources
- rare (0.01 / century 1 / century in our galaxy)
- never observed





Image: Wikimedia Commons

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- gigantic accretion rates
- only cooling channel is via neutrinos



![](_page_30_Picture_7.jpeg)

temperature / density profile

 solve hydrodynamic equations with appropriate boundary conditions (accretion shock discontinuity)

$$\begin{split} \frac{1}{r^2} \frac{\mathrm{d}(r^2 \rho v)}{\mathrm{d}r} &= 0 \ ,\\ \frac{\mathrm{d}(\rho c^2 + e)}{\mathrm{d}r} - \frac{w}{\rho} \frac{\mathrm{d}\rho}{\mathrm{d}r} &= \frac{\varepsilon_{\mathrm{nuc}} - \mathcal{L}_{\nu}}{v} \ ,\\ v \frac{\mathrm{d}v}{\mathrm{d}r} + \frac{GM_{\mathrm{NS}}}{r^2} + \frac{1}{w} \frac{\mathrm{d}P}{\mathrm{d}r} \left( v^2 + c^2 - \frac{2GM_{\mathrm{NS}}}{r} \right) &= 0 \ , \end{split}$$

### Esteban Beacom JK 2023

![](_page_31_Picture_5.jpeg)

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![](_page_32_Figure_1.jpeg)

 solve hydrodynamic equations with appropriate boundary conditions (accretion shock discontinuity)

$$\begin{split} \frac{1}{r^2} \frac{\mathrm{d}(r^2 \rho v)}{\mathrm{d}r} &= 0 \,, \\ \frac{\mathrm{d}(\rho c^2 + e)}{\mathrm{d}r} - \frac{w}{\rho} \frac{\mathrm{d}\rho}{\mathrm{d}r} &= \frac{\varepsilon_{\mathrm{nuc}} - \mathcal{L}_{\nu}}{v} \,, \\ v \frac{\mathrm{d}v}{\mathrm{d}r} + \frac{GM_{\mathrm{NS}}}{r^2} + \frac{1}{w} \frac{\mathrm{d}P}{\mathrm{d}r} \left( v^2 + c^2 - \frac{2GM_{\mathrm{NS}}}{r} \right) &= 0 \,, \end{split}$$

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![](_page_32_Picture_5.jpeg)

continuity equation

energy conservation

Euler equation (in Schwarzschild background)

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temperature / density profile

 solve hydrodynamic equations with appropriate boundary conditions (accretion shock discontinuity)

$$\begin{split} \frac{1}{r^2} \frac{\mathrm{d}(r^2 \rho v)}{\mathrm{d}r} &= 0 \ ,\\ \frac{\mathrm{d}(\rho c^2 + e)}{\mathrm{d}r} - \frac{w}{\rho} \frac{\mathrm{d}\rho}{\mathrm{d}r} &= \frac{\varepsilon_{\mathrm{nuc}} - \mathcal{L}_{\nu}}{v} \ ,\\ v \frac{\mathrm{d}v}{\mathrm{d}r} + \frac{GM_{\mathrm{NS}}}{r^2} + \frac{1}{w} \frac{\mathrm{d}P}{\mathrm{d}r} \left( v^2 + c^2 - \frac{2GM_{\mathrm{NS}}}{r} \right) &= 0 \ , \end{split}$$

### Esteban Beacom JK 2023

![](_page_33_Picture_5.jpeg)

### neutrino emission

- e+e- annihilation (dominant)
  - plasmon decay (subdominant)

![](_page_33_Figure_9.jpeg)

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### **Common-Envelope Evolution**

![](_page_34_Figure_1.jpeg)

Main detection channel is IBD. No directionality Backgrounds:

- Accidental coincidences
- □ Li-9 from spallation
- NC interactions of atmospheric v
- □ reactor v, CC atmospheric

Esteban Beacom JK 2023

![](_page_34_Picture_8.jpeg)

![](_page_35_Picture_1.jpeg)

Esteban Beacom JK 2023

www.esa.int

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Sun

![](_page_35_Picture_5.jpeg)

![](_page_35_Picture_6.jpeg)

![](_page_36_Picture_1.jpeg)

Esteban Beacom JK 2023

existing data (~20% of stars)

Sun

### www.esa.int

![](_page_36_Picture_6.jpeg)

![](_page_36_Picture_7.jpeg)

![](_page_37_Picture_1.jpeg)

Esteban Beacom JK 2023

near future (> 80% of stars)

existing data (~20% of stars)

Sun

www.esa.int

![](_page_37_Picture_8.jpeg)

![](_page_37_Picture_9.jpeg)

- CEE detectable almost anywhere in our galaxy
- novel astrophysical neutrino source
- opportunity for discovery

![](_page_38_Picture_4.jpeg)

Esteban Beacom JK 2023

near future (> 80% of stars)

existing data (~20% of stars)

Sun

www.esa.int

![](_page_38_Picture_11.jpeg)

![](_page_38_Picture_12.jpeg)

![](_page_39_Picture_6.jpeg)

supernova neutrinos ★ nucleosynthesis ★ matter under extreme conditions ★ stellar evolution

![](_page_39_Figure_9.jpeg)

![](_page_39_Picture_10.jpeg)

![](_page_40_Figure_0.jpeg)

### Infalling material produces accretion shock

Neutrino gain region

 $\tau_{\nu} = 1$ 

The convective region must overcome this pressure to launch an explosion

 $P_{\rm shock} = \frac{1}{2} \rho_{\rm s} v_{\rm ff}^2$ 

Image: Young 2021

~100-300 km

![](_page_41_Figure_0.jpeg)

Melson Janka Marek 2015

neutrino density  $> 10^{30}$  cm<sup>-3</sup> each neutrino "feels" the presence of the other neutrinos (via coherent forward scattering)

$$i\mathcal{A} =$$
  $+$ 

![](_page_42_Picture_3.jpeg)

![](_page_42_Figure_5.jpeg)

![](_page_42_Picture_7.jpeg)

![](_page_42_Picture_8.jpeg)

- neutrino density  $> 10^{30}$  cm<sup>-3</sup> each neutrino "feels" the presence of the other neutrinos

$$i(\partial_t + \vec{v} \cdot \vec{\nabla}_{\vec{r}})\rho_{\vec{r},\vec{p}} =$$

![](_page_43_Picture_4.jpeg)

# flavour evolution described by von Neumann equation (mean field approach)

 $= \left[ H_{\text{vac}} + H_{\text{MSW}} + H_{\nu\nu}, \rho_{\vec{r},\vec{p}} \right]$ 

![](_page_43_Figure_8.jpeg)

- neutrino density > 10<sup>30</sup> cm<sup>-3</sup>
   each neutrino "feels" the presence of the other neutrinos
- □ flavour evolution described by von N

$$i(\partial_t + \vec{v} \cdot \vec{\nabla}_{\vec{r}})\rho_{\vec{r},\vec{p}} =$$

### vacuum oscillations

$$H_{\rm vac} = \frac{1}{2E} U_{\rm PMNS} M^2 U_{\rm PMNS}^{\dagger}$$

$$H_{\rm MSW} = \sqrt{2}G_{F}$$

![](_page_44_Picture_8.jpeg)

![](_page_44_Figure_9.jpeg)

![](_page_44_Picture_11.jpeg)

- neutrino density  $> 10^{30}$  cm<sup>-3</sup> each neutrino "feels" the presence of the other neutrinos

$$i(\partial_t + \vec{v} \cdot \vec{\nabla}_{\vec{r}})\rho_{\vec{r},\vec{p}} =$$

non-linear equation dynamics highly non-trivial computationally intractable so far a pure Standard Model problem possible quantum entanglement? 

![](_page_45_Picture_5.jpeg)

# flavour evolution described by von Neumann equation (mean field approach)

 $= \left[ H_{\text{vac}} + H_{\text{MSW}} + H_{\nu\nu}, \rho_{\vec{r},\vec{p}} \right]$ 

### solution will be crucial for understanding the next Galactic supernova

![](_page_45_Figure_11.jpeg)

![](_page_46_Picture_0.jpeg)

### Cartoon: Reddit u/TheVeryNearFuture

## Supernova Neutrinos on a Quantum Computer

(in 2-flavour approximation)

states: 
$$|\psi\rangle = |q_1\rangle \otimes |q_2\rangle \otimes \ldots \otimes$$

- However: in the large-N limit, fully entangled N-qubit system should reduce to the standard mean-field picture Friedland Lunardini 2003

![](_page_47_Picture_6.jpeg)

highly entangled quantum system calls for simulation on a quantum system basic idea: flavour state of each neutrino mode represented by qubit  $q_i$ 

Hall et al. 2021, Amitrano et al. 2022, Siwach et al. 2023

$$|q_N\rangle$$

time-evolution via Trotterization (discretisation in t + low-order expansion of  $S = e^{i\hat{H}\delta t}$ )

our goal here is to demonstrate this explicitly on a quantum computer

![](_page_47_Picture_13.jpeg)

### Neutrino Qubit Hamiltonian

![](_page_48_Figure_1.jpeg)

![](_page_48_Picture_2.jpeg)

![](_page_48_Picture_3.jpeg)

### Neutrino Qubit Hamiltonian

![](_page_49_Figure_1.jpeg)

![](_page_49_Figure_2.jpeg)

![](_page_49_Figure_3.jpeg)

 $\Leftrightarrow$ 

$$H_{\text{int}} \propto \begin{pmatrix} 2 & & \\ & 1 & 1 & \\ & & 1 & 1 & \\ & & & 2 \end{pmatrix} \begin{vmatrix} \nu_e \nu_e \rangle \\ |\nu_\mu \nu_e \rangle \\ |\nu_\mu \nu_\mu \rangle$$

![](_page_49_Picture_6.jpeg)

![](_page_49_Picture_7.jpeg)

$$H = \sum_{k=1}^{N} \vec{b} \cdot \vec{\sigma}_k + \sum_{p < q}^{N} J_{pq} \vec{\sigma}_q$$

![](_page_49_Picture_9.jpeg)

![](_page_49_Picture_10.jpeg)

### Neutrino Qubit Hamiltonian

![](_page_50_Figure_1.jpeg)

![](_page_50_Figure_2.jpeg)

![](_page_50_Figure_3.jpeg)

 $\Leftrightarrow$ 

$$H_{\text{int}} \propto \begin{pmatrix} 2 & & \\ & 1 & 1 & \\ & & 1 & 1 & \\ & & & 2 \end{pmatrix} \begin{vmatrix} \nu_e \nu_e \rangle \\ |\nu_\mu \nu_e \rangle \\ |\nu_\mu \nu_\mu \rangle$$

![](_page_50_Picture_6.jpeg)

![](_page_50_Picture_7.jpeg)

### vacuum oscillations

self-interactions

![](_page_50_Picture_10.jpeg)

![](_page_50_Picture_11.jpeg)

## **Emergence of the Mean Field Picture**

in the large-N limit, fully entangled N-qubit system
 reduces to the standard mean-field picture Friedland Lunardini 2003

![](_page_51_Figure_2.jpeg)

![](_page_51_Figure_3.jpeg)

![](_page_51_Picture_5.jpeg)

## **Emergence of the Mean Field Picture**

in the large-N limit, fully entangled N-qubit system
 reduces to the standard mean-field picture Friedland Lunardini 2003

![](_page_52_Figure_2.jpeg)

![](_page_52_Figure_3.jpeg)

![](_page_52_Picture_5.jpeg)

## Implementation on Quantum Hardware

- deployment to IBM QPUs (superconducting transmons)
- severely noise-limited
  - more challenging than Heisenberg / Ising models due to all-to-all interactions
  - smart algorithm design allows implementation on linear qubit chain

![](_page_53_Picture_5.jpeg)

![](_page_53_Figure_8.jpeg)

IBM Heron r2 (156 qubits)

![](_page_53_Picture_11.jpeg)

![](_page_54_Figure_1.jpeg)

![](_page_54_Picture_2.jpeg)

background

beam (noiseless simulation)

![](_page_54_Figure_5.jpeg)

### IBM Heron r2 (156 qubits)

![](_page_54_Picture_8.jpeg)

![](_page_55_Figure_1.jpeg)

![](_page_55_Picture_2.jpeg)

background

![](_page_55_Figure_5.jpeg)

### IBM Heron r2 (156 qubits)

![](_page_55_Picture_8.jpeg)

![](_page_55_Picture_9.jpeg)

![](_page_56_Figure_1.jpeg)

![](_page_56_Picture_2.jpeg)

background

beam (IBM Heron r2, Pauli twirling + readout mitigation + zero-noise extrapolation

![](_page_56_Figure_5.jpeg)

### IBM Heron r2 (156 qubits)

![](_page_56_Picture_8.jpeg)

![](_page_56_Picture_9.jpeg)

![](_page_57_Figure_1.jpeg)

![](_page_57_Picture_2.jpeg)

background

beam (IBM Heron r2, probabilistic error cancellation)

![](_page_57_Figure_5.jpeg)

### IBM Heron r2 (156 qubits)

![](_page_57_Figure_8.jpeg)

![](_page_57_Picture_9.jpeg)

![](_page_57_Picture_10.jpeg)

![](_page_58_Picture_6.jpeg)

supernova neutrinos ★ oscillations of SN neutrinos poorly understood ★ playground for quantum computing

neutron stars neutrinos may enable first discovery of common-envelope evolution

![](_page_58_Figure_9.jpeg)

![](_page_58_Picture_10.jpeg)

## Thank You!

![](_page_59_Picture_1.jpeg)

![](_page_59_Picture_2.jpeg)

![](_page_59_Picture_3.jpeg)

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![](_page_59_Picture_5.jpeg)

### **Bonus Slides**

## **Departure from Equilibrium**

Strategy for calculating rates

- apply Feynman rules + phase space integral
   + Pauli-blocking factors
- phenomenological parameterisation of nuclear matrix element
- neglect angular dependence of hadronic + leptonic matrix element
- □ treat nucleons as non-relativistic
- □ all momenta close to Fermi surfaces
- □ carry out angular integrals
  - carry out energy integrals

(multiple applications of residue theorem)

Friman Maxwell 1979 Yakovlev Levenfish 1995 Yakovlev Kaminker Gnedin Haensel 2000 Shapiro Teukolsky 1983

![](_page_62_Picture_11.jpeg)

![](_page_62_Figure_12.jpeg)

![](_page_62_Picture_14.jpeg)

## Neutron Stars Away from Thermal Equilibrium

- very strong dependence on
   departure from equilibrium and on T
  - For muons:

- diffusion (over O(yr) time scales) + decay
- potential source of MeV neutrinos
- would require a mechanisms that drives all NSs in the MW towards
   lower muon abundances
   (~1% over 10<sup>9</sup> yrs)
- all mechanisms known to us do the opposite

no known mechanism that does this

![](_page_63_Picture_8.jpeg)

![](_page_63_Figure_9.jpeg)

![](_page_63_Picture_11.jpeg)