

# *Infrared surprises in quantum gravity*

2403.13053, 2412.16149, 2412.16142

& work in progress

to appear & work in progress

with Sangmin Choi & Alok Laddha

with Sangmin Choi & Ameya Kadhe

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**SIM NS**  
FOUNDATION

SIMONS COLLABORATION ON CELESTIAL HOLOGRAPHY

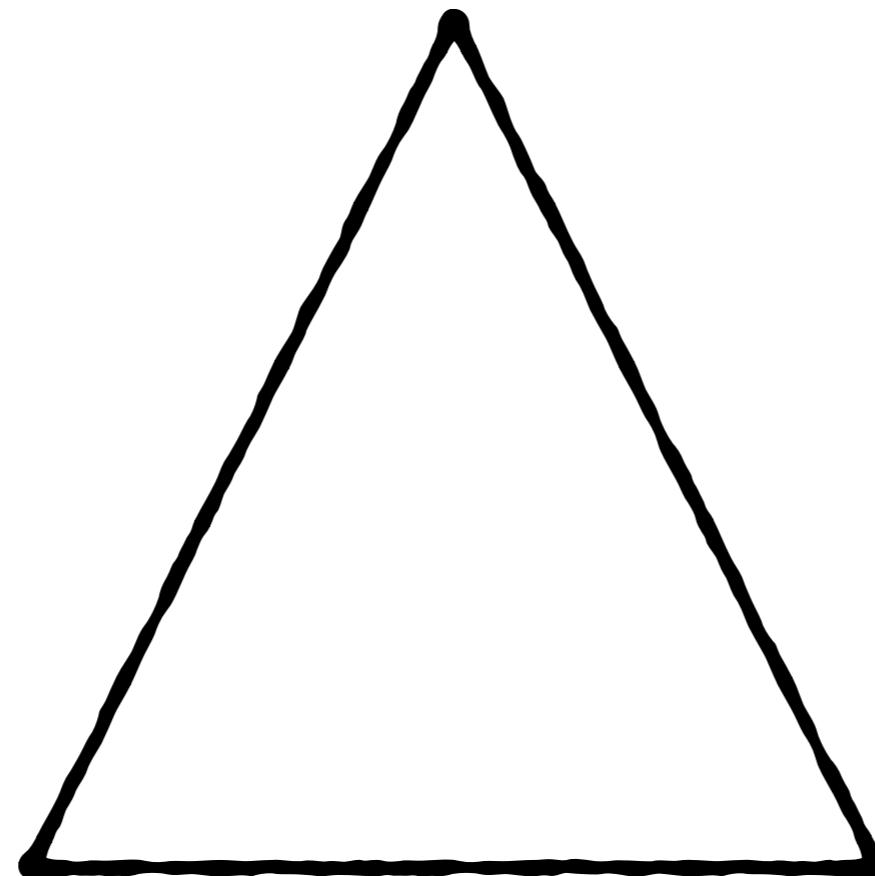
# Universality in the IR

Classical and quantum gravitational scattering in 4D exhibits universal behavior:

**asymptotic symmetry**



approximating  
our world by  
 $\Lambda \approx 0$



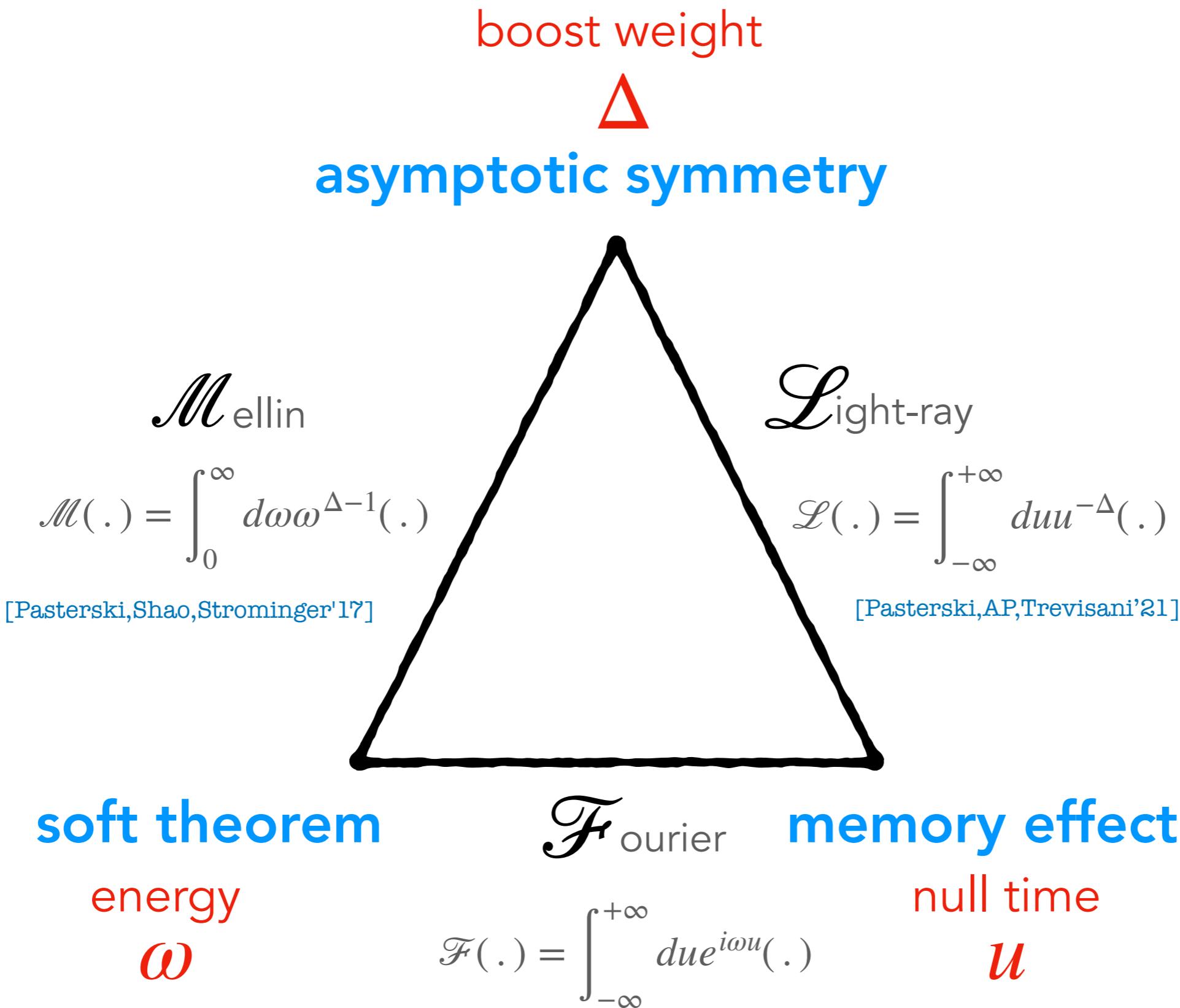
**soft theorem**

low-energy

**memory effect**

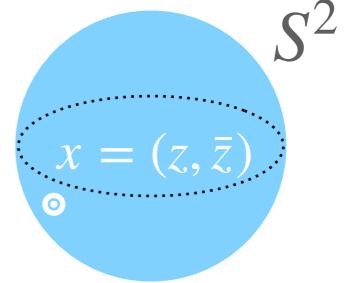
large-distance

# 3 bases for the IR

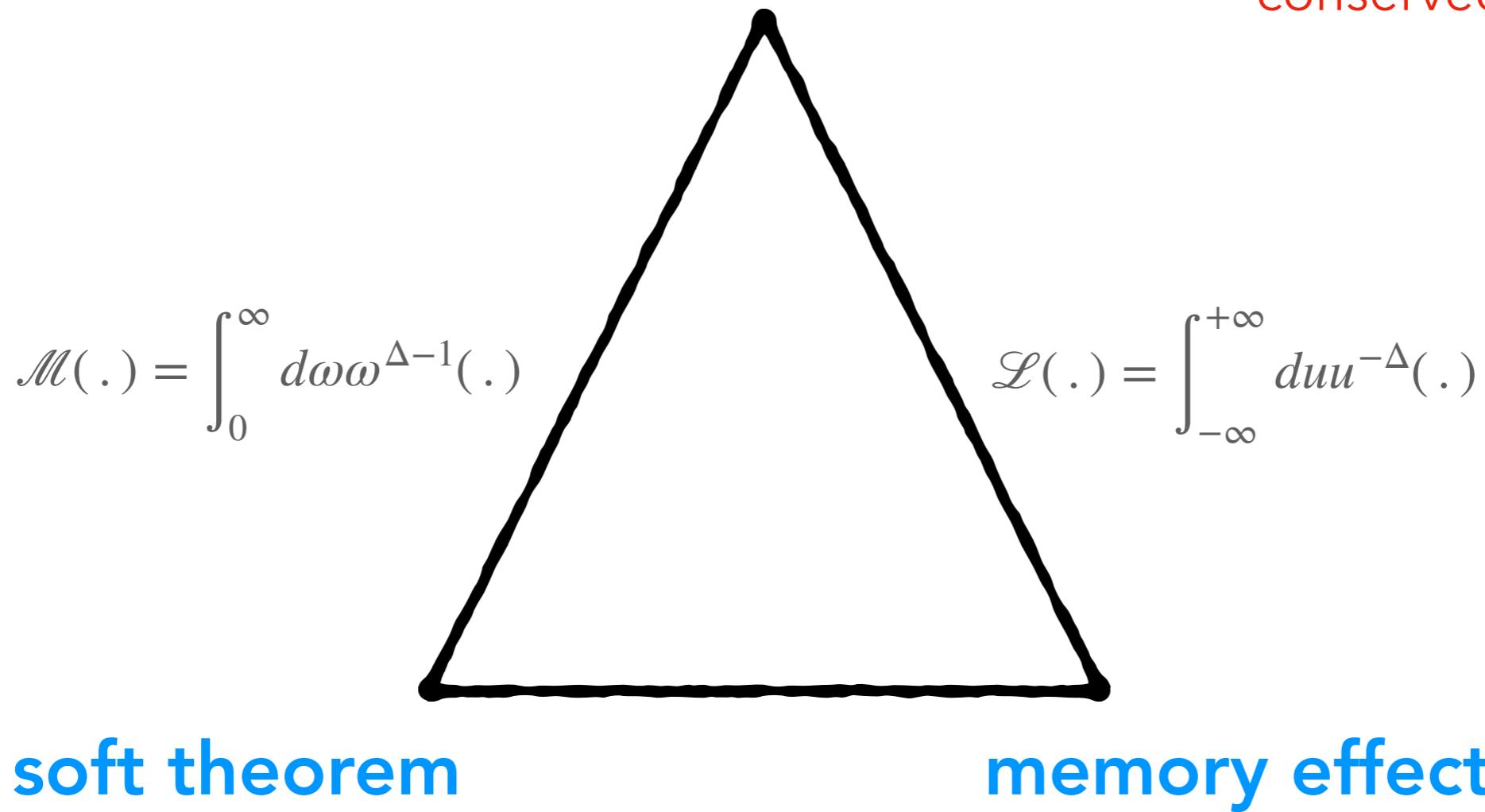


# IR universality is 2D

graviton: directed towards point  $x = (z, \bar{z})$  on the celestial sphere

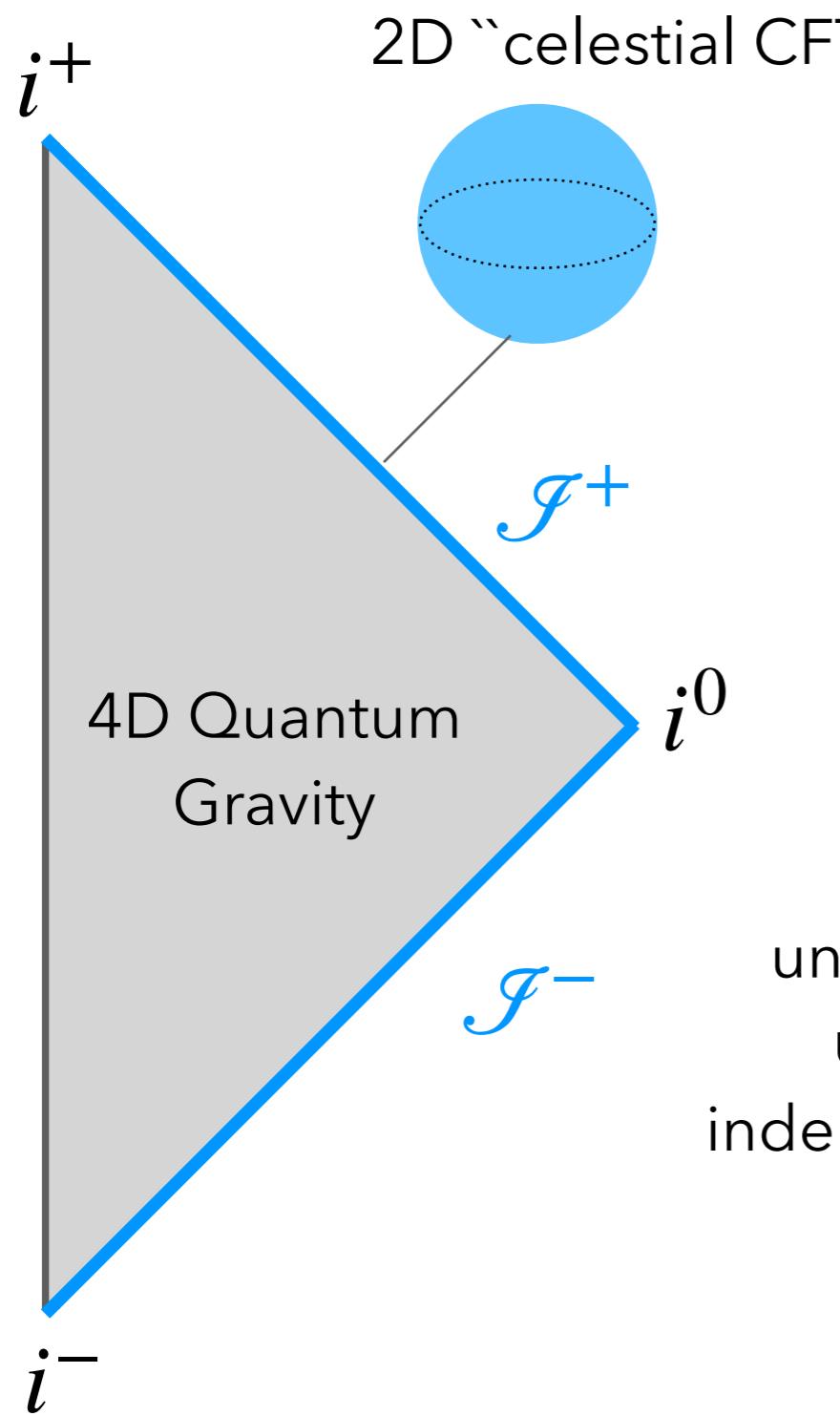


**asymptotic symmetry** → generated by 2D  
conserved operators  $\mathcal{O}_\Delta$



# 2D Hologram for 4D Quantum Gravity ?

(at least for the IR sector)



## Celestial Holography

bottom-up

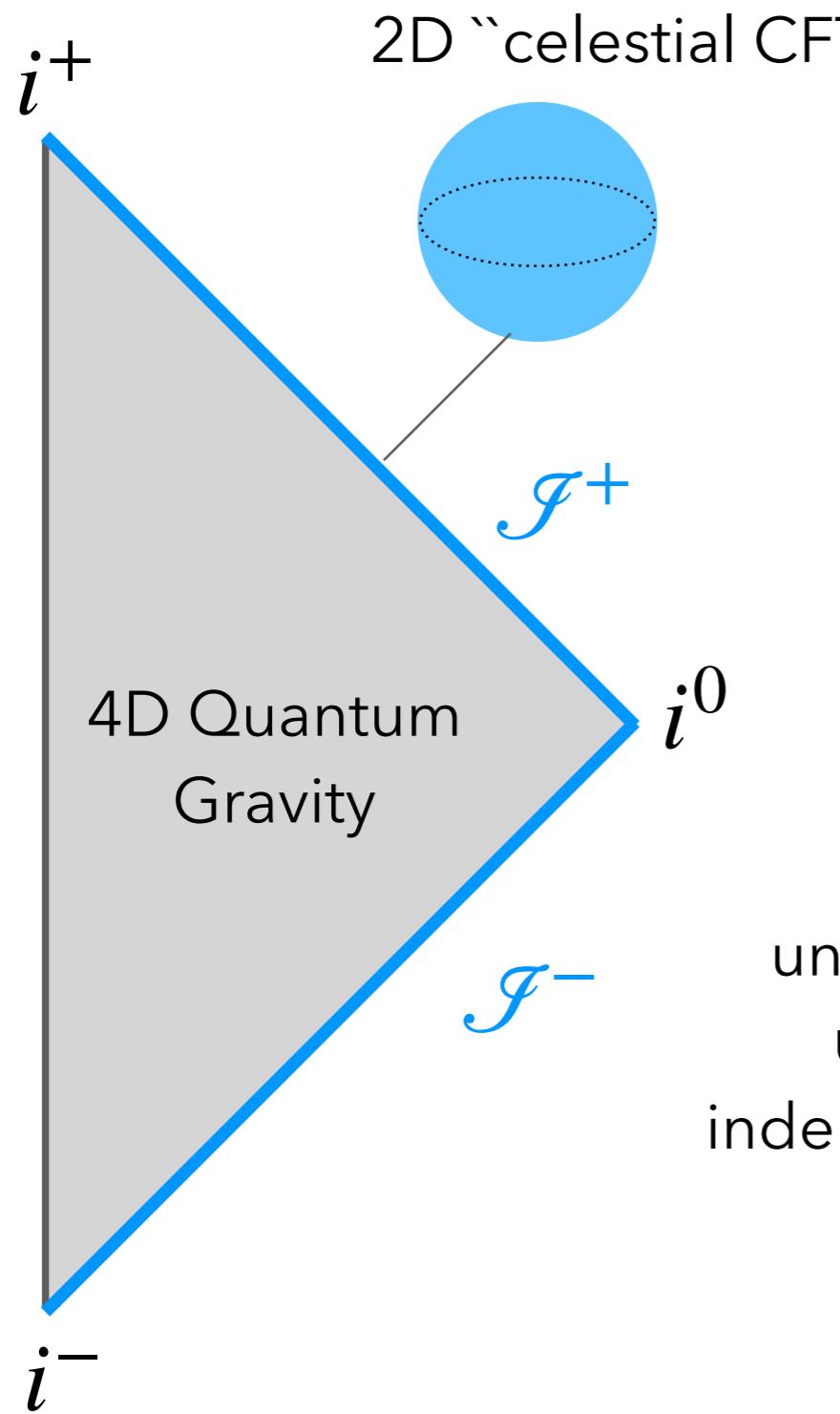
understand long-distance  
**universal properties**  
independent of short distance  
microphysics

top-down

explore toy models to  
**construct dual pairs**

# 2D Hologram for 4D Quantum Gravity ?

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## Celestial Holography

**bottom-up**

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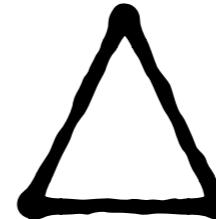
top-down

explore toy models to  
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# Plan: infrared surprises & their connection to “celestial CFT”

Infrared triangles

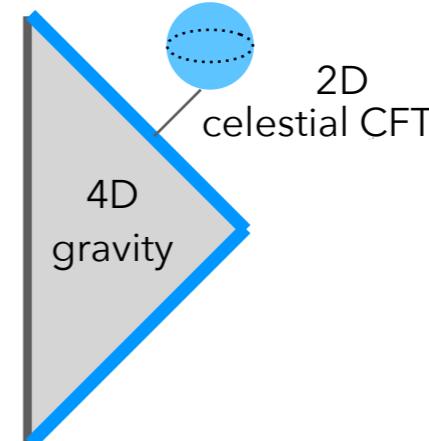
asymptotic symmetry



soft theorem      memory effect

Towers of  $\infty$  symmetries

$4D = 2D$



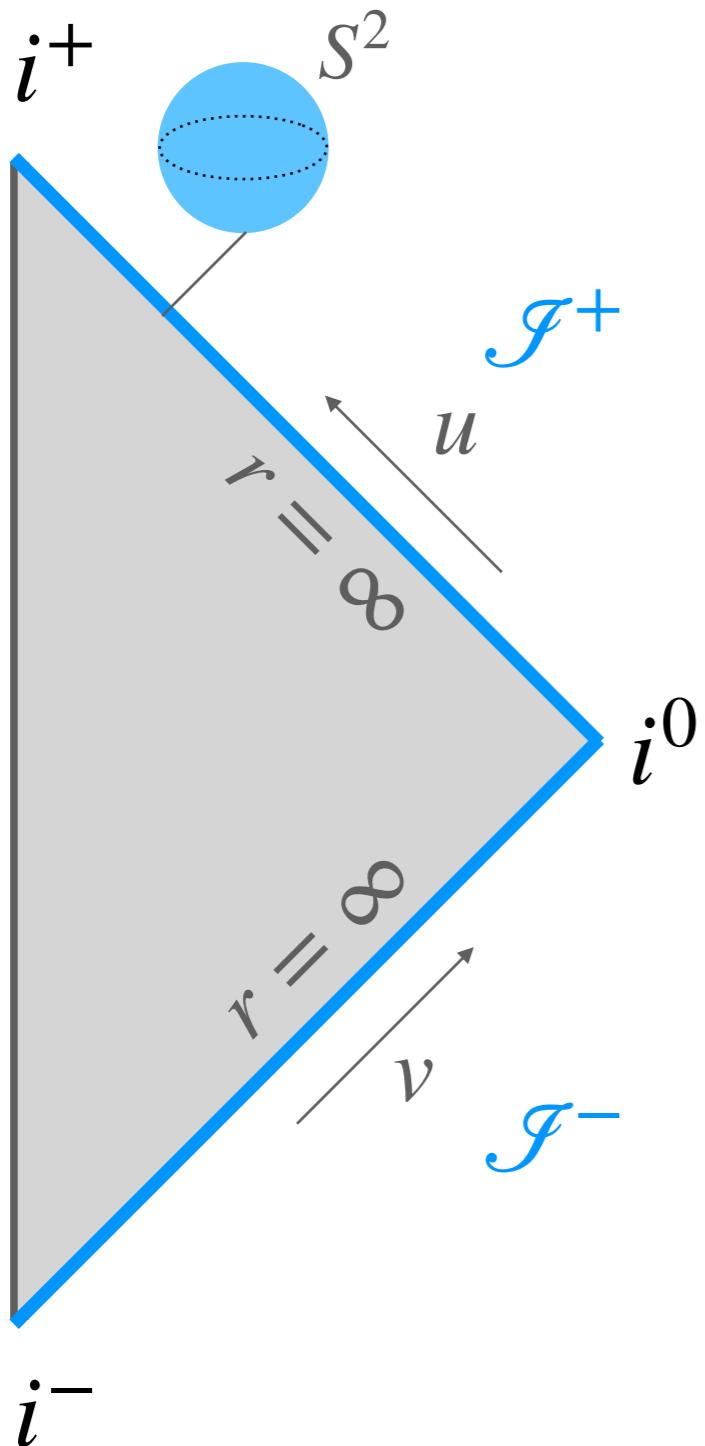
Long-range effects →

**New conservation laws  
in Einstein gravity!**

[Choi,Laddha,AP'24]

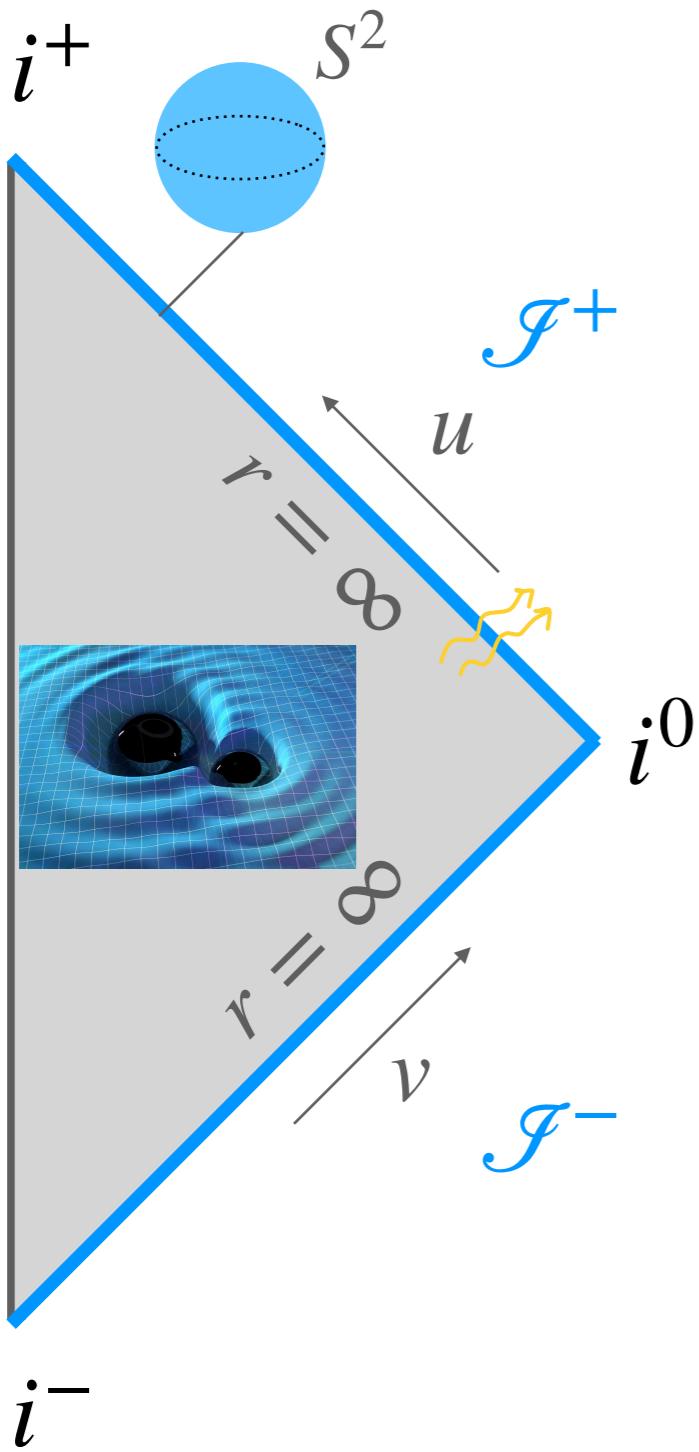
# Infrared triangles

# Asymptotic symmetries



"Large" gauge/diffeo transformations that  
preserve the boundary conditions of the fields,  
i.e. their large-distance fall-offs.

# Asymptotic symmetries



"Large" gauge/diffeo transformations that  
preserve the boundary conditions of the fields,  
i.e. their large-distance fall-offs.

**gravity:**

mass

$$ds^2 = -(1 + \dots)du^2 - (2 + \dots)dudr$$

angular momentum

$$+ (\dots)dudx^A$$

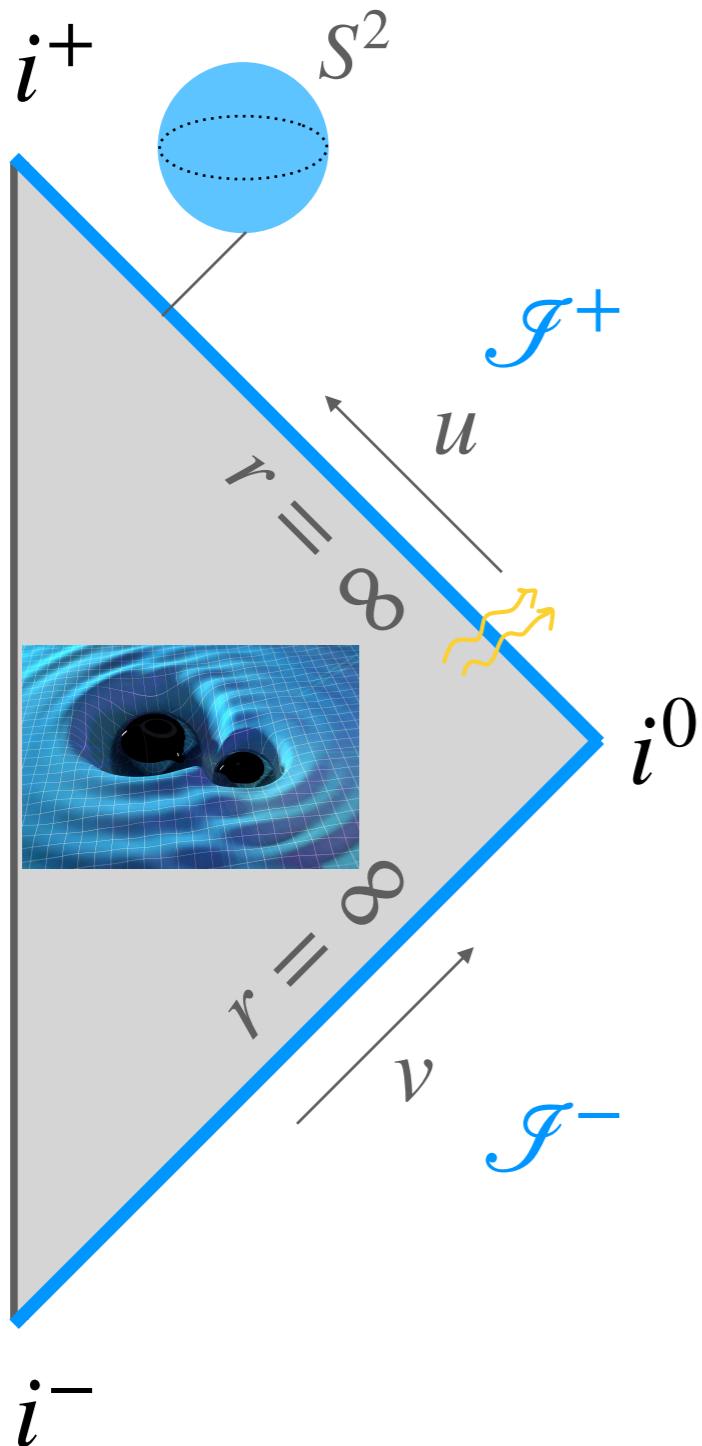
$$+ (r^2\gamma_{AB} + rC_{AB} + \dots)dx^A dx^B$$



shear: gravitational waves

$$\Rightarrow \text{Bondi news } N_{AB} = \partial_u C_{AB}$$

# Asymptotic symmetries



"Large" gauge/diffeo transformations that  
preserve the boundary conditions of the fields,  
i.e. their large-distance fall-offs.

## gravity:

Find  $\xi$  such that  $\mathcal{L}_\xi g_{\mu\nu} \approx "0"$  as  $r \rightarrow \infty$ .

$$\begin{array}{c} \uparrow \\ O(1/r^\#) \end{array}$$

Unlike gauge redundancies, asymptotic  
symmetries act non-trivially on the  
physical data  $\rightarrow$  non-zero charges.

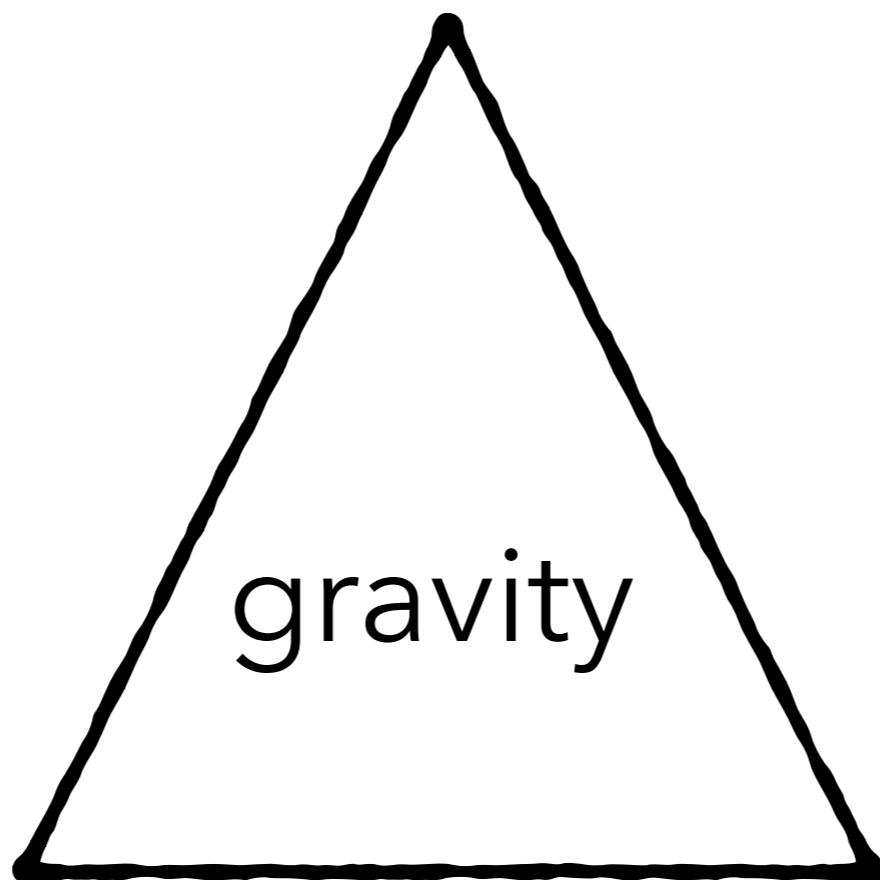
# IR triangle

The symmetries of asymptotically flat space are not just Poincaré but an infinite extension!

[Bondi,van der Burg,Metzner'62]  
[Sachs'62]

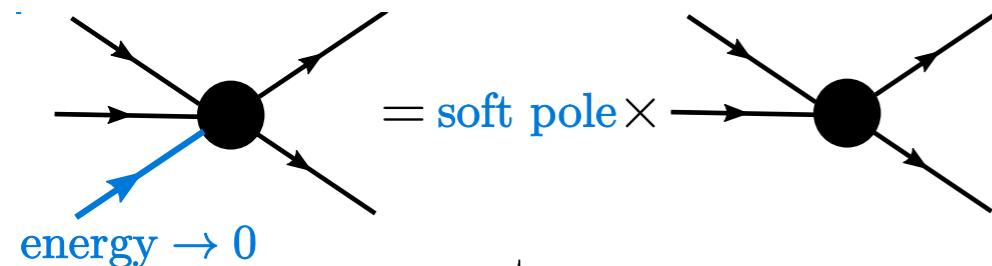
BMS group

**supertranslations**

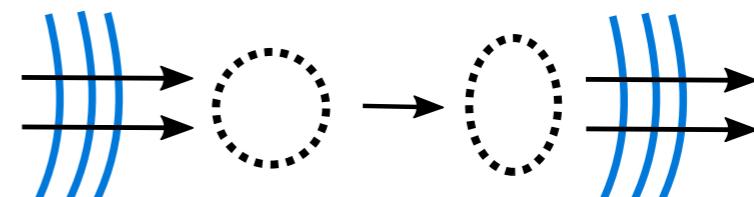


[Weinberg'65]

**soft graviton theorem**

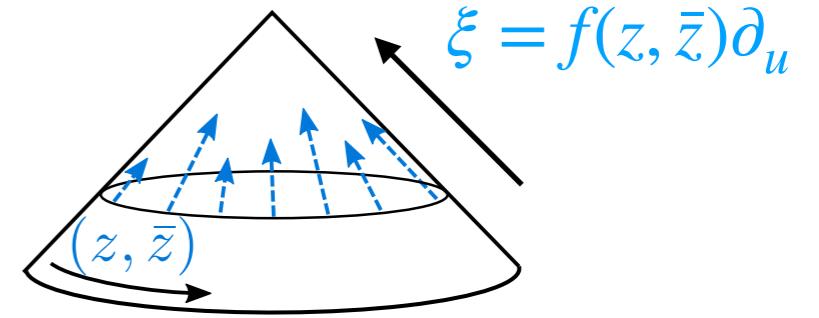


**displacement memory effect**



[He,Lysov,Mitra,Strominger'14]

[Strominger,Zhiboedov'14]



# IR triangle

on  $\mathcal{J}^+$ :

supertranslations  
superrotations

[Barnich,Troessaert'11][Campiglia,Laddha'14]

extended / generalized  
BMS group

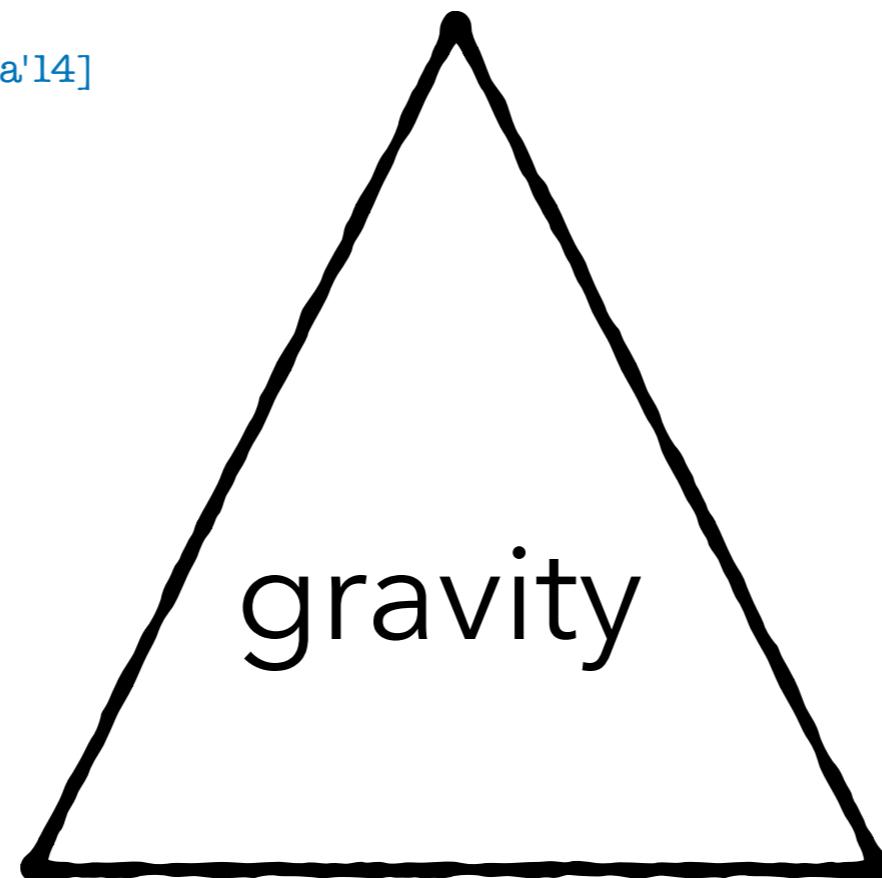
**asymptotic symmetry**

$$\xi^u = f$$

$$\xi^u = -\frac{1}{2}D_A Y^A$$

$$\xi^A = \frac{1}{r}D_A f$$

$$\xi^A = Y^A$$



**soft theorem**

$\omega^{-1}$  leading soft graviton  
 $\omega^0$  subleading soft graviton

[Cachazo,Strominger'14]

**memory effect**

displacement  
spin

[Pasterski,Strominger,Zhiboedov'15]

# IR triangle

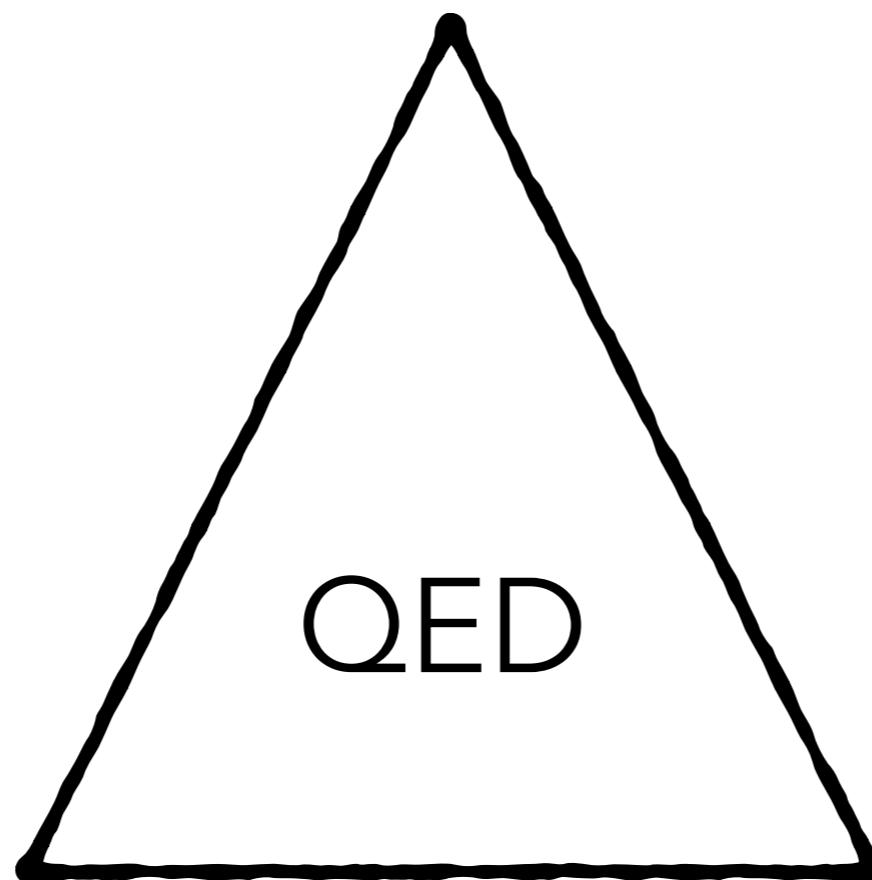
superphaserotation

**asymptotic symmetry**

[He,Mitra,Porfyriadis,Strominger'14]

[Kapac,Pate,Strominger'15]

[Campiglia,Laddha'15]



QED

**soft theorem**

$\omega^{-1}$  leading soft photon

[Weinberg'65]

**memory effect**

electromagnetic kick

[Bieri, Garfinkle'13]

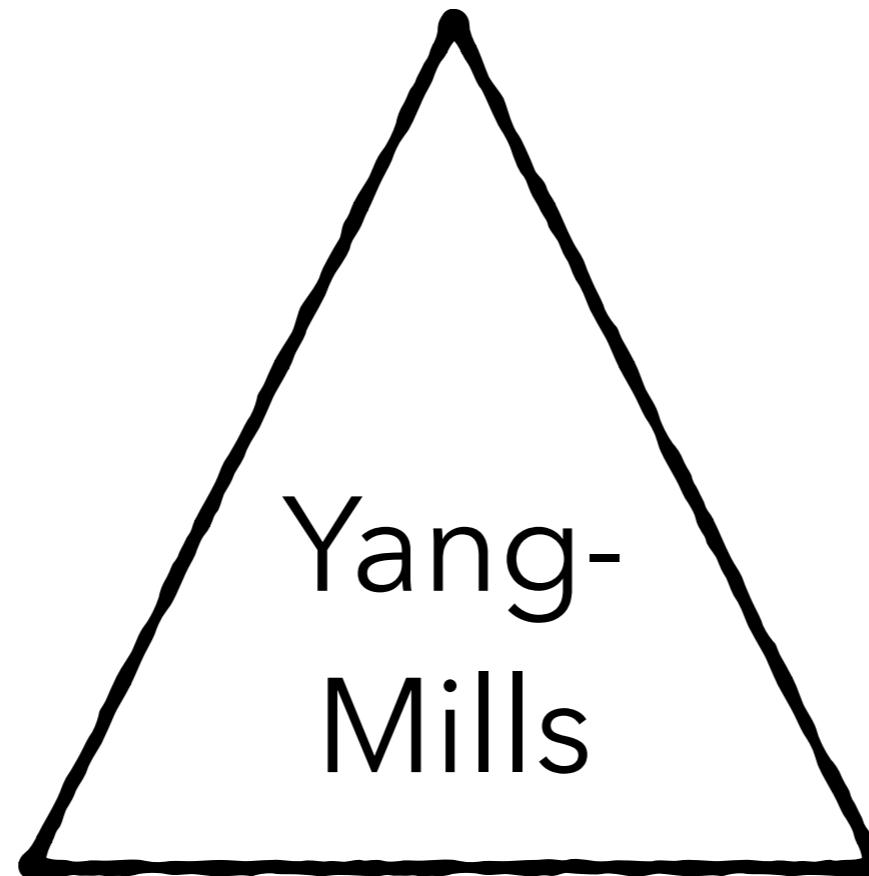
[Pasterski'15]

# IR triangle

superphaserotation

**asymptotic symmetry**

[He,Mitra,Strominger'14]



**soft theorem**

**memory effect**

$\omega^{-1}$  leading soft gluon

[Berends,Giele'89]

color

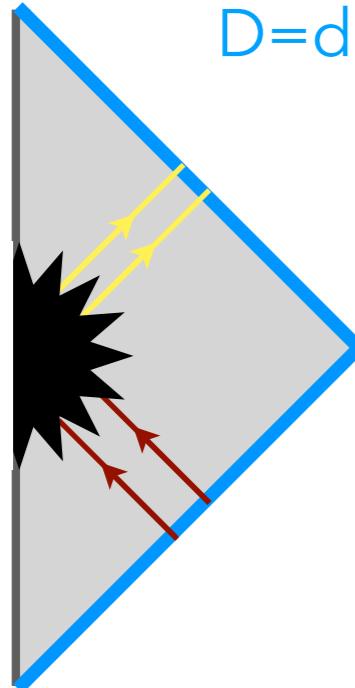
[Pate,Raclariu,Strominger'17]

**4D = 2D**

# Symmetry & observables

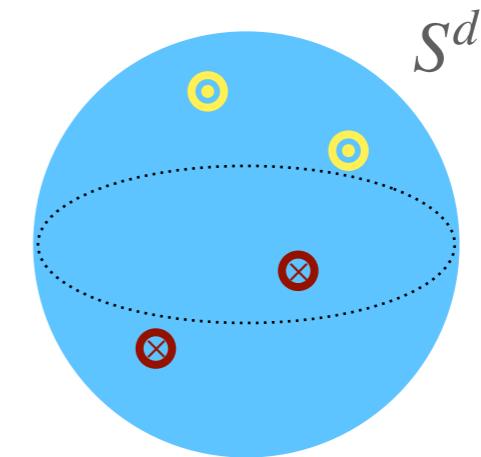
symmetry:

Lorentz group in  
 $D=d+2$  dimensions

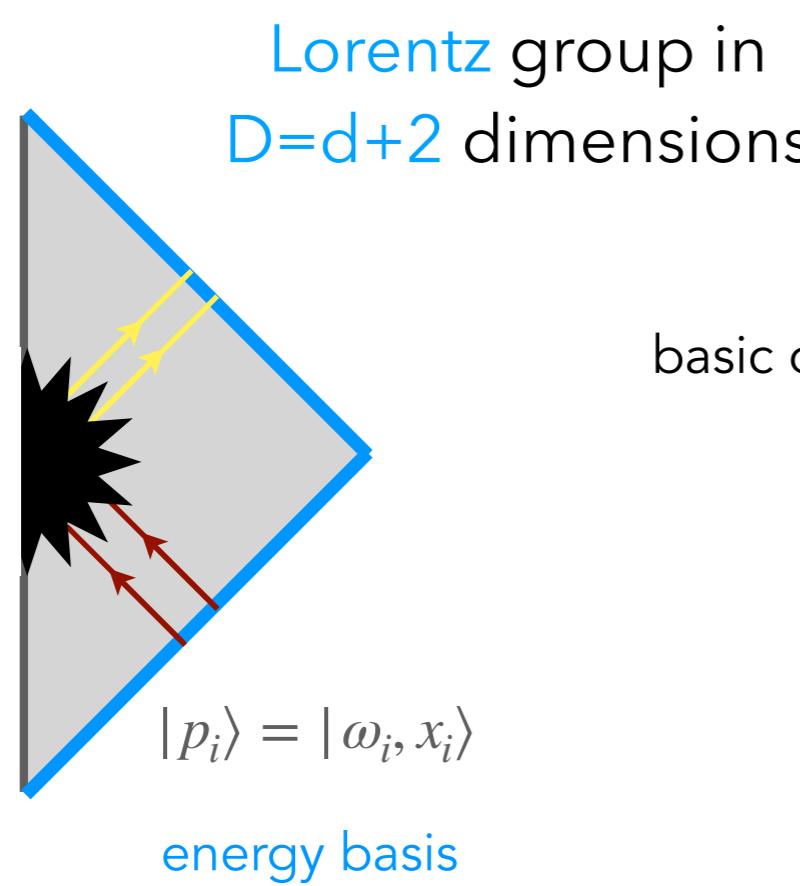


$\approx$

Euclidean **conformal**  
group in  $d$  dimensions



# Symmetry & observables



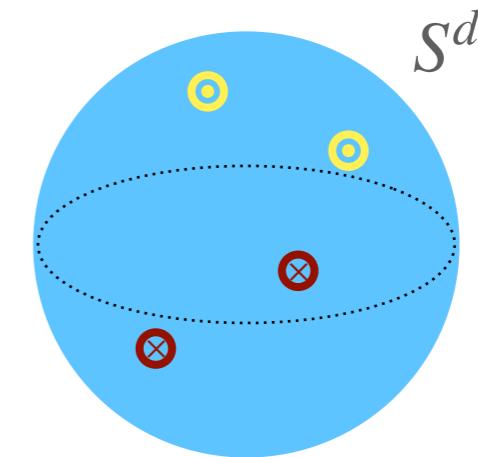
symmetry:

$\approx$

Euclidean **conformal**  
group in  $d$  dimensions

basic observables in flat space:

**S-matrix**

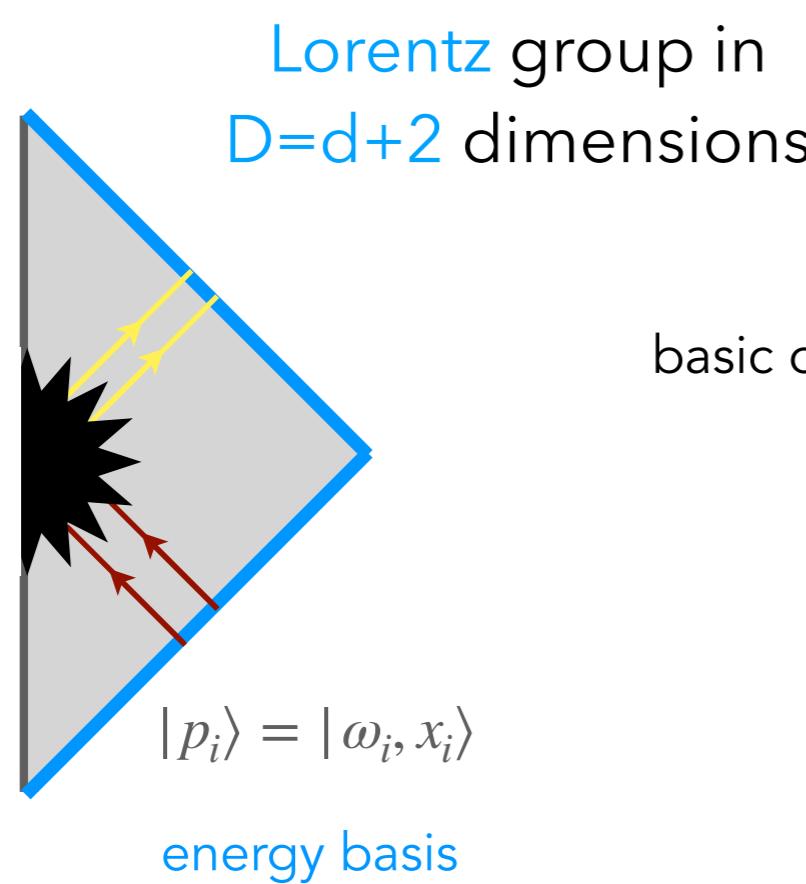


**Standard amplitudes**

$$\langle p_n^{\text{out}} \dots | \mathcal{S} | \dots p_1^{\text{in}} \rangle$$

translation symmetry

# Symmetry & observables



## Standard amplitudes

$$\langle p_n^{\text{out}} \dots | \mathcal{S} | \dots p_1^{\text{in}} \rangle$$

translation symmetry

symmetry:

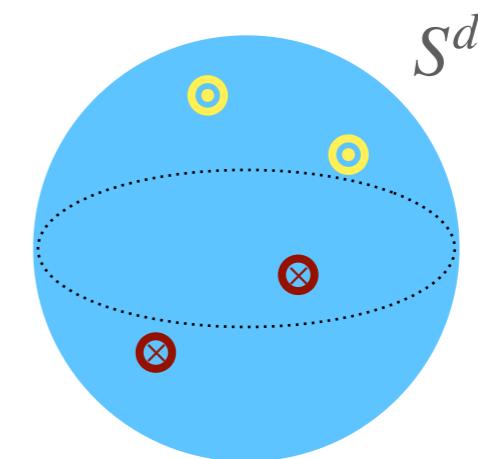
$\approx$

Euclidean conformal  
group in  $d$  dimensions

basic observables in flat space:

## S-matrix

$$\frac{\mathcal{M}_{\text{ellin}}}{\int_0^\infty d\omega \omega^{\Delta-1}}$$



$$|\Delta_i, x_i\rangle$$

boost-weight basis

## Celestial amplitudes

$$\langle \mathcal{O}_{\Delta_1}^-(x_1) \dots \mathcal{O}_{\Delta_n}^+(x_n) \rangle$$

Lorentz symmetry

global conformal symmetry

# From global to local conformal

on  $\mathcal{J}^+$ :

supertranslations,  
**superrotations**, ...

[Barnich,Troessaert'11]

**local** conformal  
symmetry on  $S^2$  !

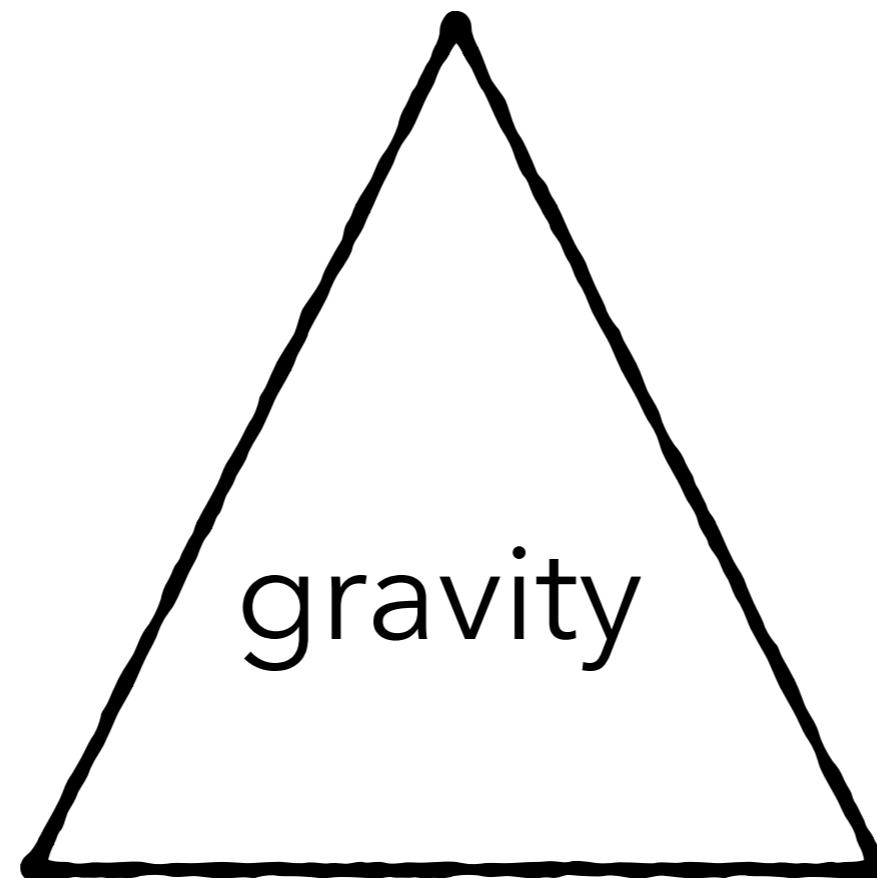
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**soft theorem**

$\omega^{-1}$  leading soft graviton,

$\omega^0$  subleading soft graviton, ...

[Cachazo,Strominger'14]

**memory effect**

displacement,

spin, ...

[Pasterski,Strominger,Zhiboedov'15]

Tower of  $\infty$  symmetries

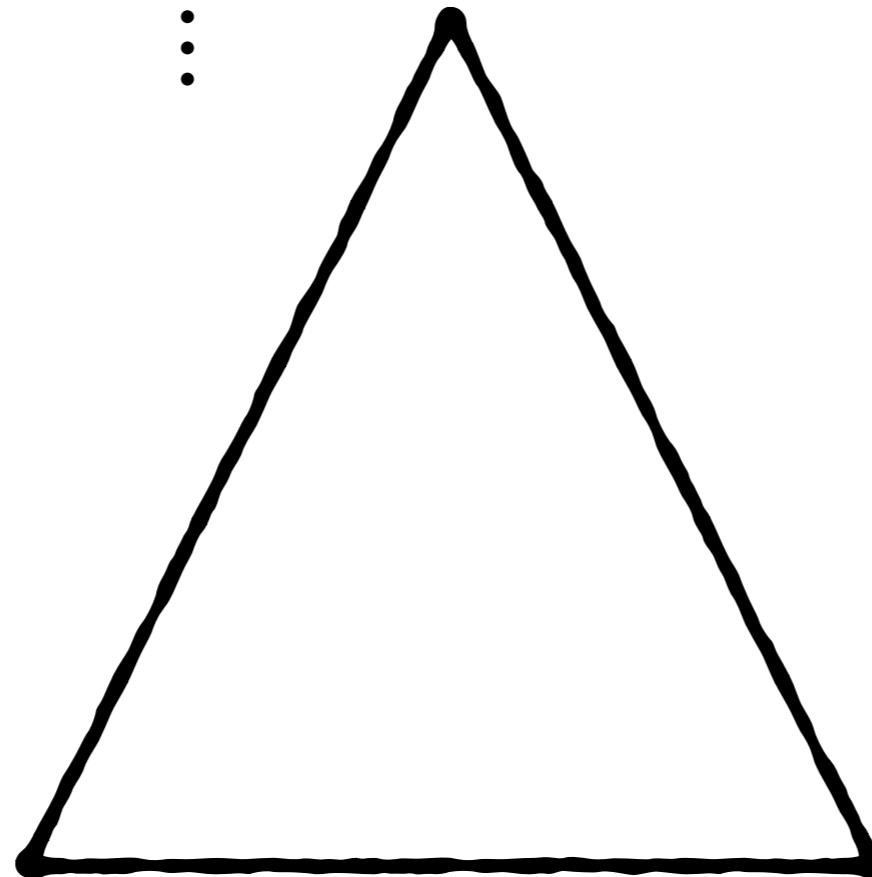


# IR triangle

(for projected S-matrix)

**asymptotic symmetry**

[Hamada,Shiu'18]  
[Li,Lin,Zhang'18]



**soft theorem**

$\omega^n$

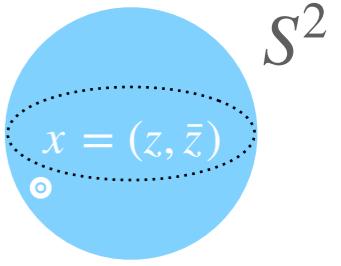
:

$n = -1, 0, 1, \dots$

**memory effect**

:

# Celestial symmetries



*celestial diamonds:*

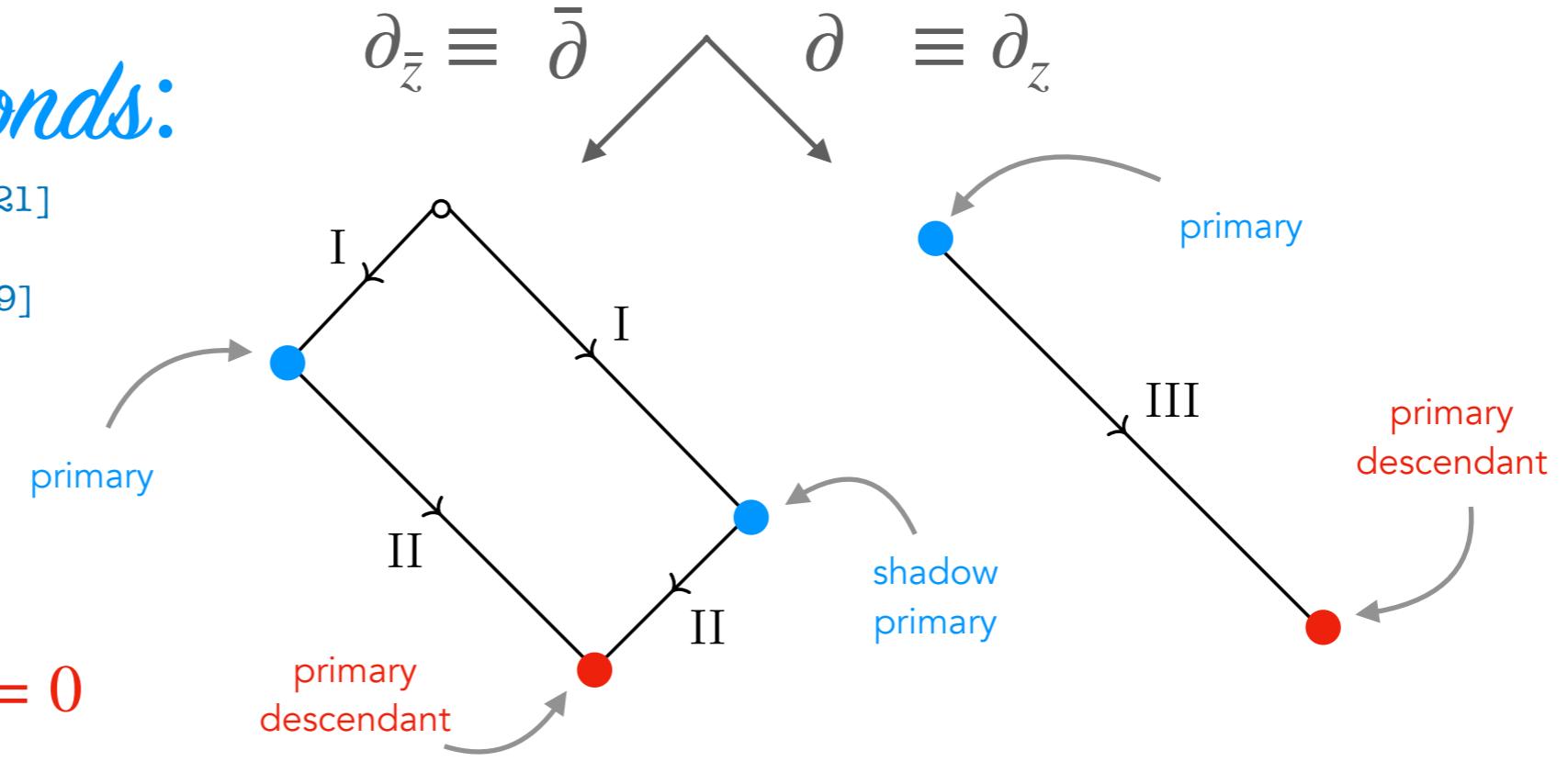
[Pasterski,AP,Trevisani'21]

also: [Banerjee,Pandey,Paul'19]

conservation  
equation

$$\bar{\partial}^\# \mathcal{O}(z, \bar{z}) = 0$$

# depends on types I, II, III: spin of descendant >,<,= spin of parent primary



Conformally soft operator = primary

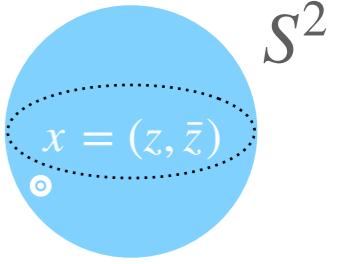
$$\downarrow \bar{\partial}^\#$$

primary descendant =

conservation equation for  
**topological operator**

Noether current:  $\mathcal{J} = \sum_{m=0}^{\#-1} (-1)^m \bar{\partial}^m \epsilon(z, \bar{z}) \bar{\partial}^{\#-m-1} \mathcal{O}(z, \bar{z})$

# Celestial symmetries



*celestial diamonds:*

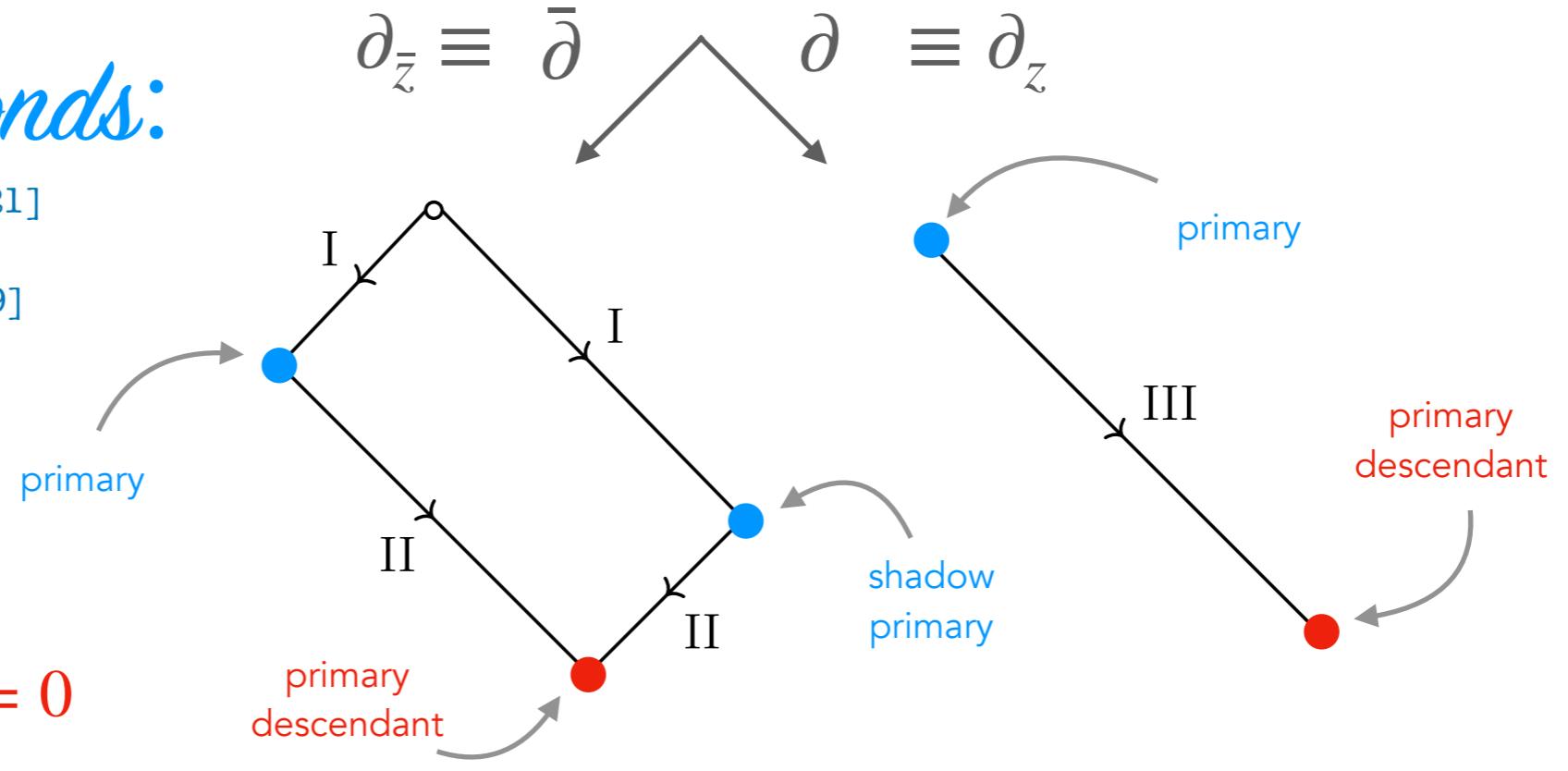
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**Classification of all symmetries in celestial CFT**  $d \geq 2$

[Pasterski,AP,Trevisani'21]  $d = 2$   
[Pano,AP,Trevisani'23]  $d > 2$

$$\mathcal{O}_\Delta : \quad \Delta = 1, 0, -1, \dots$$

supertranslations      |       $\infty$  tower

superrotation

# Towers of $\infty$ symmetries

Define a **discrete family** of conformally soft positive-helicity gravitons

$$H^k = \lim_{\epsilon \rightarrow 0} \epsilon \mathcal{O}_{k+\epsilon, +2}$$

$k = 2, 1, 0, -1, -2, \dots$

with weights  $(h, \bar{h}) = \left( \frac{k+2}{2}, \frac{k-2}{2} \right)$

and a **consistently-truncated** antiholomorphic mode expansion

$$H^k(z, \bar{z}) = \sum_{n=\frac{k-2}{2}}^{\frac{2-k}{2}} \frac{H_n^k(z)}{\bar{z}^{n+\frac{k-2}{2}}}$$

`wedge condition' = conservation equation for topological operator

$w_n^p = \frac{1}{\kappa} (p-n-1)! (p+n-1)! H_n^{-2p+4}$        $p, q = 1, \frac{3}{2}, 2, \frac{5}{2}, \dots$

$$[w_m^p, w_n^q] = [m(q-1) - n(p-1)] w_{m+n}^{p+q-2}$$

Arises already in Penrose's twistor construction!

This is the  $w_{1+\infty}$  algebra.

[Guevara, Himwich, Pate, Strominger '21]  
[Strominger '21]

# 2D soft actions

A toy model that captures the features of the diamond structure is the higher derivative Gaussian theory with action

[Pasterski,AP,Trevisani'21]

$$S = \int d^2z \left[ \partial^k \mathcal{O}_{\Delta,J}^s \bar{\partial}^{\bar{k}} \mathcal{O}_{\Delta,J}^s + \partial^k \mathcal{O}_{\Delta,-J}^s \bar{\partial}^{\bar{k}} \mathcal{O}_{\Delta,-J}^s \right]$$

and with conservation equation

$$\partial_z^k \partial_{\bar{z}}^{\bar{k}} \mathcal{O}_{\Delta,J}^s = 0$$

$$\Delta = 1 - \frac{k + \bar{k}}{2} \quad J = \frac{k - \bar{k}}{2}$$
$$k, \bar{k} \in \mathbb{Z}_>$$

QED

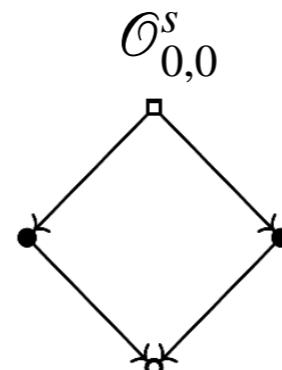
Simplest example:  
**free boson** ( $k = \bar{k} = 1$ )

Yang-Mills

[Cheung,de La Fuente, Sundrum'15]  
[Magnea'21] [González,Rojas'21]

gravity

[Nguyen,Salzer'20] [Nguyen'21]  
[Kalyanapuram'20+'21]

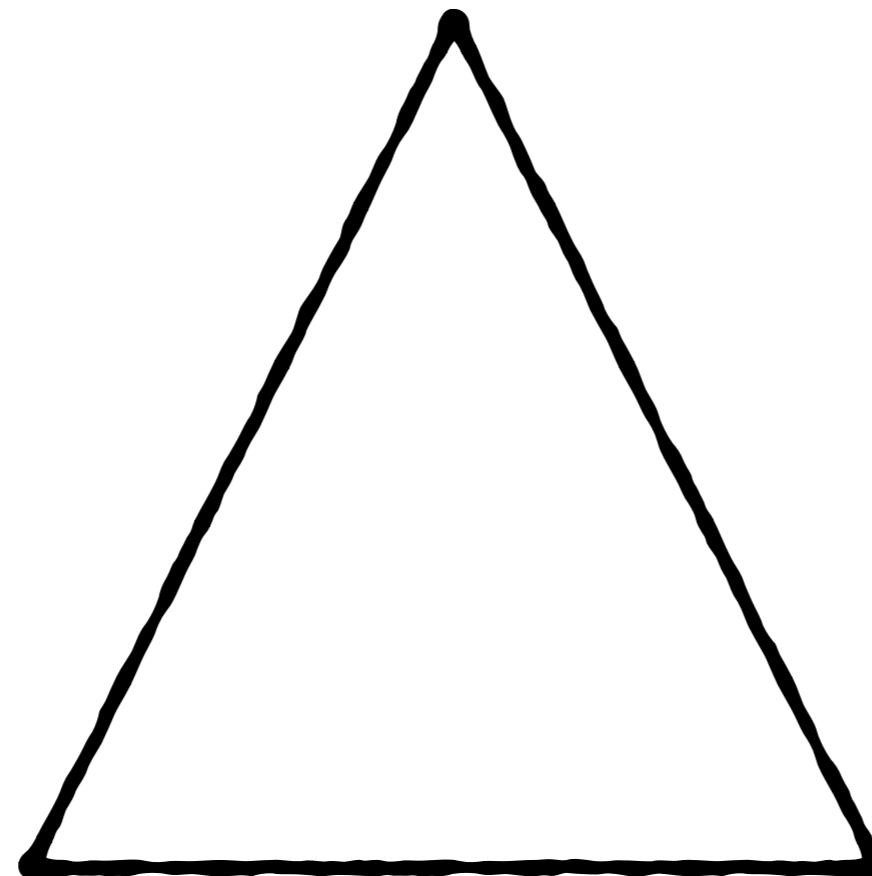


# Long-range effects

# IR triangle @ tree !

$\infty$ -dimensional  
symmetry algebra  
 $\supset$  local conformal  
symmetry on  $S^2$  !

**asymptotic symmetry**



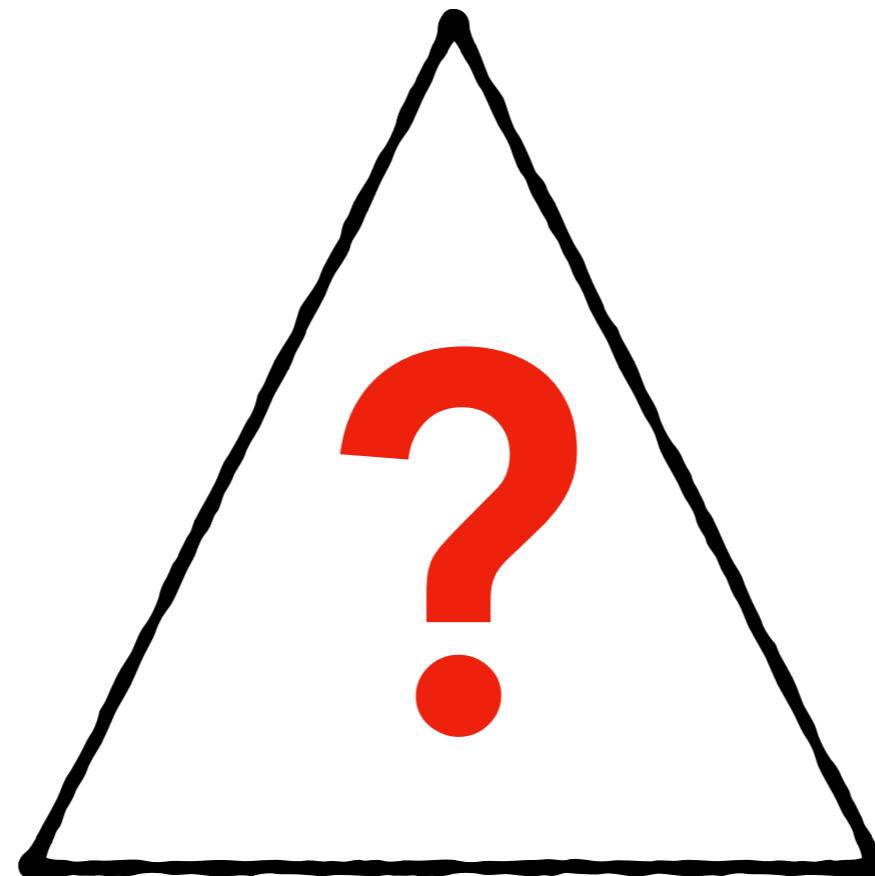
**soft theorem**

**memory effect**

# IR triangle @ loop ?

$\infty$ -dimensional  
symmetry algebra  
 $\supset$  local conformal  
symmetry on  $S^2$  ?

**asymptotic symmetry**



**soft theorem**

**memory effect**

# IR triangle

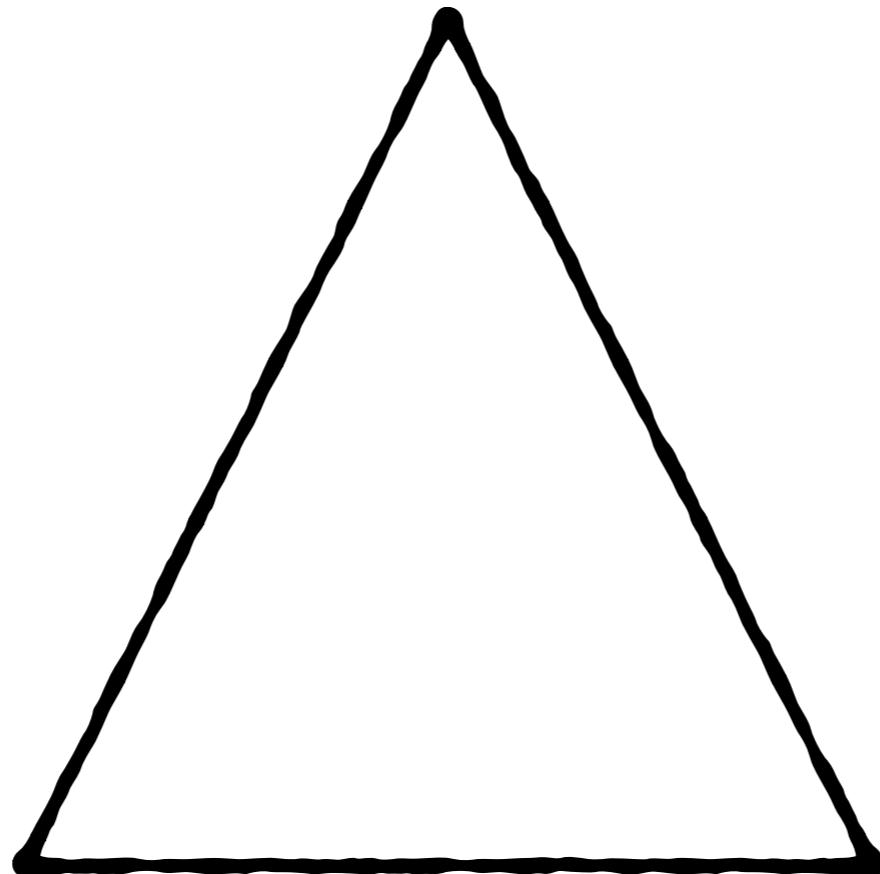
@

long-range  
IR effects

classical and quantum

$\infty$ -dimensional  
symmetry algebra  
 $\hookrightarrow$  local conformal  
symmetry on  $S^2$  ?

asymptotic ? symmetry



logarithmic soft theorem

tail memory effect

[Laddha,Sen'18]  
[Sahoo,Sen'18]  
[Saha,Sahoo,Sen'19]

[Sahoo'20]  
[Sahoo,Sen'21]  
[Sahoo,Krishna'23]

# IR triangle

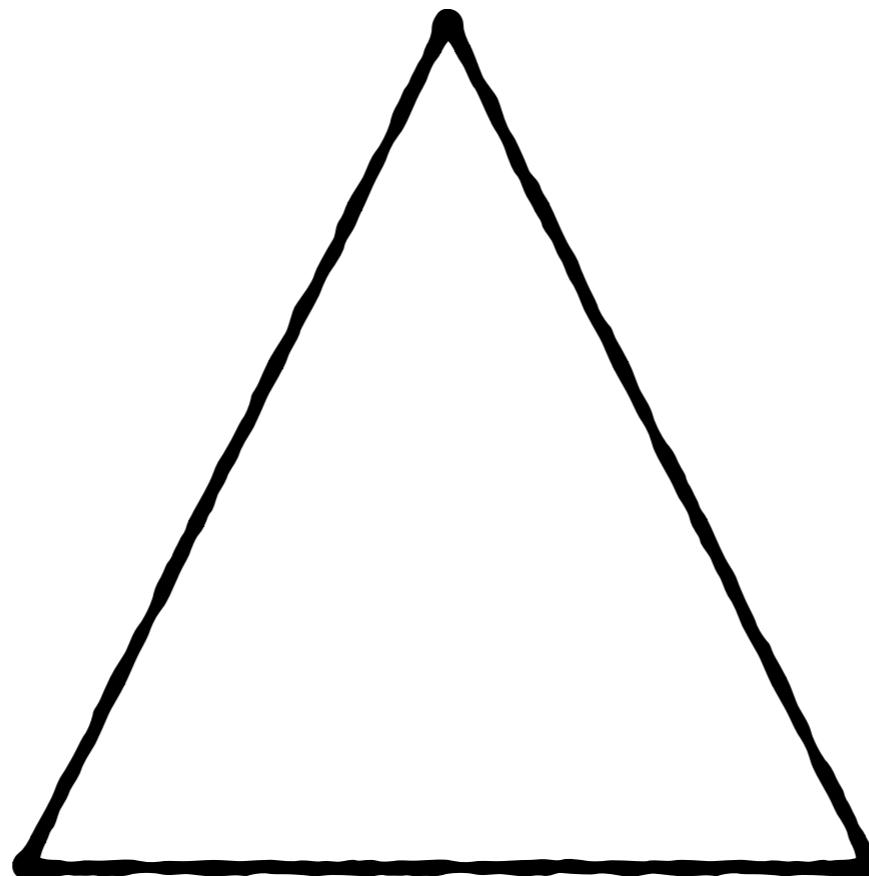
@

long-range  
IR effects

classical and quantum

$\infty$ -dimensional  
symmetry algebra  
 $\hookrightarrow$  local conformal  
symmetry on  $S^2$  ?

asymptotic ? symmetry



logs render ambiguous  
all subleading tree-level  
soft theorems



**logarithmic soft theorem**

$\omega^{-1}$  leading soft  
 $\log \omega$  subleading soft  
⋮

[Laddha,Sen'18]  
[Sahoo,Sen'18]  
[Saha,Sahoo,Sen'19]

**tail memory effect**

[Sahoo'20]  
[Sahoo,Sen'21]  
[Sahoo,Krishna'23]

# Power-law soft theorems

**Tree-level** amplitudes admit a soft expansion:

[Low'58] [Weinberg'65]

[Cachazo,Strominger'14]

[Hamada,Shiu'18] [Li,Lin,Zhang'18]

$$\mathcal{M}_{N+1}(p_1, \dots, p_N; (\omega, q, \ell)) = \sum_{n=-1}^{\infty} \text{[blue box]} \omega^n S_n(\{p_i\}, (q, \ell)) \mathcal{M}_N(p_1, \dots, p_N) + \dots$$

↑

hard momenta      helicity  
 soft momentum

$p^\mu = \omega q^\mu(z, \bar{z})$

$n = -1$	$S_{-1}$	Weinberg (leading) soft factor	tree exact & universal
$n = 0$	$S_0$	subleading tree soft factor	
$n > 0$	$S_{n>0}$	sub $n+1$ leading tree soft factors	

↑  
⇒ non-universal \*

\* see [Elvang,Jones,Naculich'16] for classification in EFT of massless particles

[Weinberg'65]

[Cachazo, Strominger'14]

# Soft graviton factor

Leading soft factor  $\sim \frac{1}{\omega}$ :

$$S_{-1} = \frac{\kappa}{2} \sum_{i=1}^N \frac{\epsilon_{\mu\nu} p_i^\mu p_i^\nu}{q \cdot p_i}$$

$p_i^\mu$  ... hard momenta

$p^\mu = \omega q^\mu$  ... soft momentum

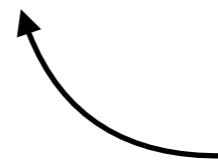
$\epsilon^{\mu\nu}$  ... soft graviton polarization

$\kappa = \sqrt{32\pi G_N}$  ... coupling

Subleading soft factor  $\sim \omega^0$ :

$$S_0 = -\frac{i\kappa}{2} \sum_{i=1}^N \frac{\epsilon_{\mu\nu} p_i^\mu q_\lambda J_i^{\lambda\nu}}{q \cdot p_i}$$

ambiguous if long-range IR effects



# Logarithmic Soft Theorems

**Long-range effects** yield **novel soft theorems**:

[Sahoo,Sen'18]  
 [Saha,Sahoo,Sen'19]

$$\mathcal{M}_{N+1}(p_1, \dots, p_N; (\omega, q, \ell)) = \sum_{n=-1}^{\infty} \text{hard momenta} \quad \text{helicity} \quad \text{soft momentum}$$

$$\omega^n (\ln \omega)^{n+1} S_n^{(\ln \omega)}(\{p_i\}, (q, \ell)) \mathcal{M}_N(p_1, \dots, p_N) + \dots$$

$\uparrow$

$\supset$  non-universal  
 $\sim \omega^n (\ln \omega)^{m \neq n+1}$

$n = -1$	$S_{-1}^{(\ln \omega)} \equiv S_{-1}$	Weinberg (leading) soft factor	<b>tree exact &amp; universal</b>
$n = 0$	$S_0^{(\ln \omega)} \neq S_0$	leading log soft factor	<b>1-loop exact &amp; universal</b>
$n > 0$	$S_{n>0}^{(\ln \omega)} \neq S_{n>0}$	sub $^n$ leading log soft factor	<b>(n+1)-loop exact &amp; universal ?</b>

Is the **universality** of the loop exact logarithmic soft theorems a consequence of **asymptotic symmetries of the S-matrix**?

[Weinberg'65]  
 [Sahoo,Sen'18]  
 [Saha,Sahoo,Sen'19]

# Soft graviton factor

Leading soft factor  $\sim \frac{1}{\omega}$ :

$$S_{-1} = \frac{\kappa}{2} \sum_{i=1}^N \frac{\epsilon_{\mu\nu} p_i^\mu p_i^\nu}{q \cdot p_i}$$

$p_i^\mu$  ... hard momenta  
 $p^\mu = \omega q^\mu$  ... soft momentum  
 $\epsilon^{\mu\nu}$  ... soft graviton polarization

Subleading soft factor  $\sim \ln \omega$ :

$$S_0^{(\ln \omega)} = S_{0,\text{classical}}^{(\ln \omega)} + S_{0,\text{quantum}}^{(\ln \omega)}$$

classical: late time gravitational radiation from  
 particle acceleration via long-range  
 gravitational interaction

$$S_{0,\text{classical}}^{(\ln \omega)} = \frac{i(\frac{\kappa}{2})^3}{8\pi} \sum_{i=1}^N \frac{\epsilon_{\mu\nu} p_i^\nu q_\rho}{q \cdot p_i} \sum_{j \neq i, \eta_i \eta_j = 1} \frac{(p_i \cdot p_j) \left[ p_i^\mu p_j^\rho - p_j^\mu p_i^\rho \right] \left[ 2(p_j \cdot p_j)^2 - 3p_i^2 p_j^2 \right]}{\left[ (p_i \cdot p_j)^2 - p_i^2 p_j^2 \right]^{3/2}}$$

quantum:  
 (1-loop)

$$S_{0,\text{quantum}}^{(\ln \omega)} = -\frac{(\frac{\kappa}{2})^3}{16\pi^2} \sum_{i=1}^N \frac{\epsilon_{\mu\rho} p_i^\rho q_\nu}{q \cdot p_i} \left( p_i^\mu \partial_{p_i}^\nu - p_i^\nu \partial_{p_i}^\mu \right) \sum_{j \neq i} \frac{2(p_i \cdot p_j)^2 - p_i^2 p_j^2}{\sqrt{(p_i \cdot p_j)^2 - p_i^2 p_j^2}} \ln \left( \frac{p_i \cdot p_j + \sqrt{(p_i \cdot p_j)^2 - p_i^2 p_j^2}}{p_i \cdot p_j - \sqrt{(p_i \cdot p_j)^2 - p_i^2 p_j^2}} \right)$$

# Soft graviton factor

[Weinberg'65]  
 [Sahoo,Sen'18]  
 [Saha,Sahoo,Sen'19]

Leading soft factor  $\sim \frac{1}{\omega}$ :

$$S_{-1} = \frac{\kappa}{2} \sum_{i=1}^N \frac{\epsilon_{\mu\nu} p_i^\mu p_i^\nu}{q \cdot p_i}$$

$p_i^\mu$  ... hard momenta  
 $p^\mu = \omega q^\mu$  ... soft momentum  
 $\epsilon^{\mu\nu}$  ... soft graviton polarization

Subleading soft factor  $\sim \ln \omega$ :

$$S_0^{(\ln \omega)} = S_{0,\text{classical}}^{(\ln \omega)} + S_{0,\text{quantum}}^{(\ln \omega)}$$

classical: late time gravitational radiation from  
 particle acceleration via long-range  
 gravitational interaction

$$S_{0,\text{classical}}^{(\ln \omega)} = \frac{i(\frac{\kappa}{2})^3}{8\pi} \sum_{i=1}^N \frac{\epsilon_{\mu\nu} p_i^\nu q_\rho}{q \cdot p_i} \sum_{j \neq i, \eta_i \eta_j = 1} \frac{(p_i \cdot p_j) \left[ p_i^\mu p_j^\rho - p_j^\mu p_i^\rho \right] \left[ 2(p_j \cdot p_j)^2 - 3p_i^2 p_j^2 \right]}{\left[ (p_i \cdot p_j)^2 - p_i^2 p_j^2 \right]^{3/2}} \quad + \text{drag}$$

quantum:  
 (1-loop)

$$S_{0,\text{quantum}}^{(\ln \omega)} = -\frac{(\frac{\kappa}{2})^3}{16\pi^2} \sum_{i=1}^N \frac{\epsilon_{\mu\rho} p_i^\rho q_\nu}{q \cdot p_i} \left( p_i^\mu \partial_{p_i}^\nu - p_i^\nu \partial_{p_i}^\mu \right) \sum_{j \neq i} \frac{2(p_i \cdot p_j)^2 - p_i^2 p_j^2}{\sqrt{(p_i \cdot p_j)^2 - p_i^2 p_j^2}} \ln \left( \frac{p_i \cdot p_j + \sqrt{(p_i \cdot p_j)^2 - p_i^2 p_j^2}}{p_i \cdot p_j - \sqrt{(p_i \cdot p_j)^2 - p_i^2 p_j^2}} \right) \quad + \text{drag}$$

$\omega_{\text{IR}} \ll \omega_{\text{loop}} \ll \omega_{\text{soft}}$

# Soft graviton factor

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quantum:  
 (1-loop)

$$S_{0,\text{quantum}}^{(\ln \omega)} = -\frac{(\frac{\kappa}{2})^3}{16\pi^2} \sum_{i=1}^N \frac{\epsilon_{\mu\rho} p_i^\rho q_\nu}{q \cdot p_i} \left( p_i^\mu \partial_{p_i}^\nu - p_i^\nu \partial_{p_i}^\mu \right) \sum_{j \neq i} \frac{2(p_i \cdot p_j)^2 - p_i^2 p_j^2}{\sqrt{(p_i \cdot p_j)^2 - p_i^2 p_j^2}} \ln \left( \frac{p_i \cdot p_j + \sqrt{(p_i \cdot p_j)^2 - p_i^2 p_j^2}}{p_i \cdot p_j - \sqrt{(p_i \cdot p_j)^2 - p_i^2 p_j^2}} \right) \quad + \text{drag}$$

# Conservation laws

To establish a symmetry interpretation for a soft theorem from *first principles*: for asymptotic symmetry transformations  $\delta$  compute charges  $Q^\pm$  from the symplectic structure

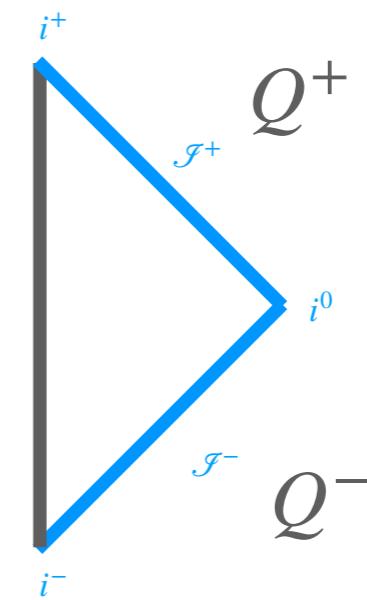
$$\Omega_{i^\pm \cup \mathcal{J}^\pm}(\delta, \delta') = \delta' Q^\pm$$

in the **covariant phase space formalism** and show that the **charge conservation law**

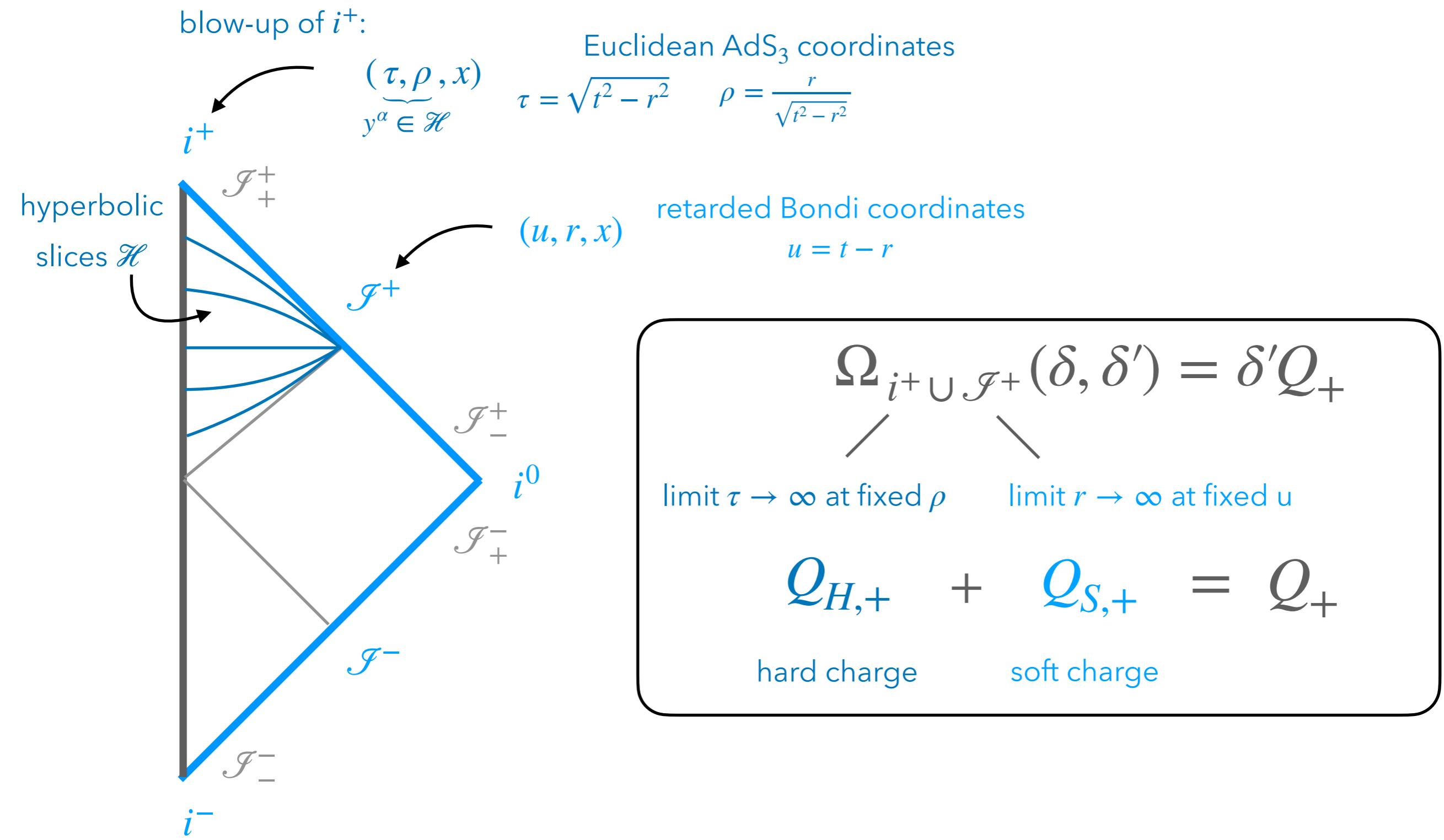
Upon identifying the fields and symmetry parameter antipodally:

$$Q^+ = Q^-$$

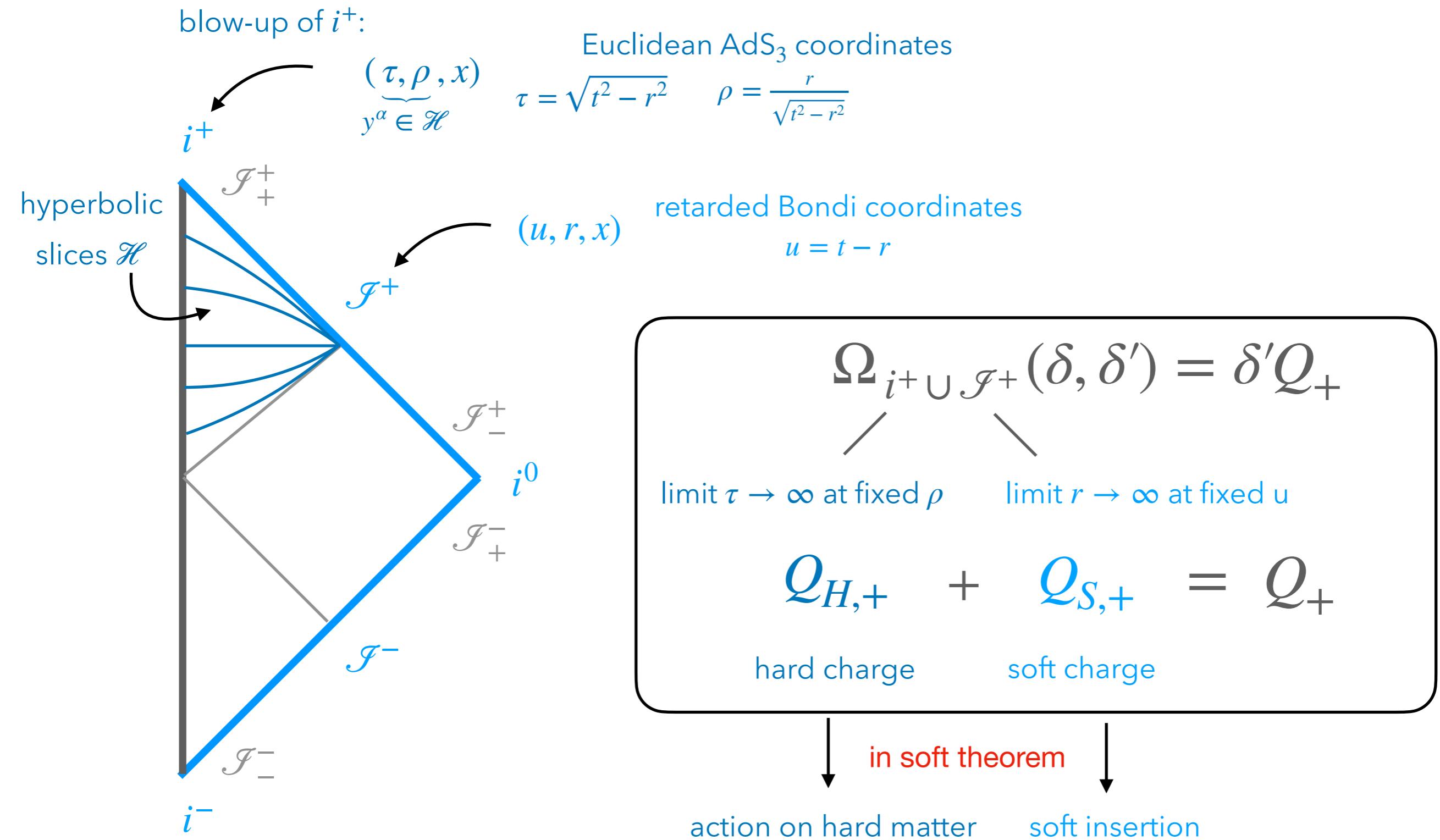
corresponds to the **soft theorem**.



# Structure at $\infty$

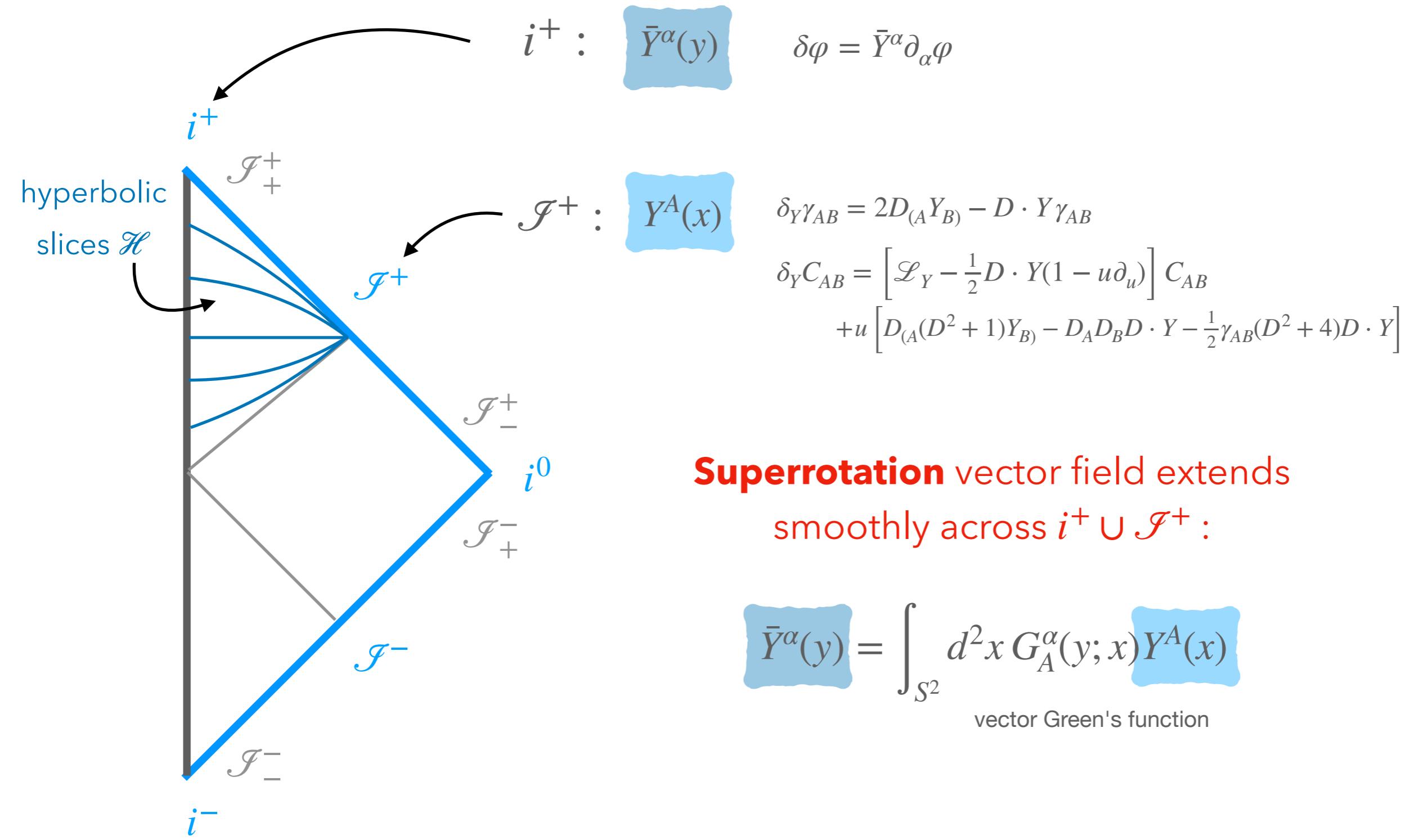


# Structure at $\infty$



# Symmetry at $\infty$

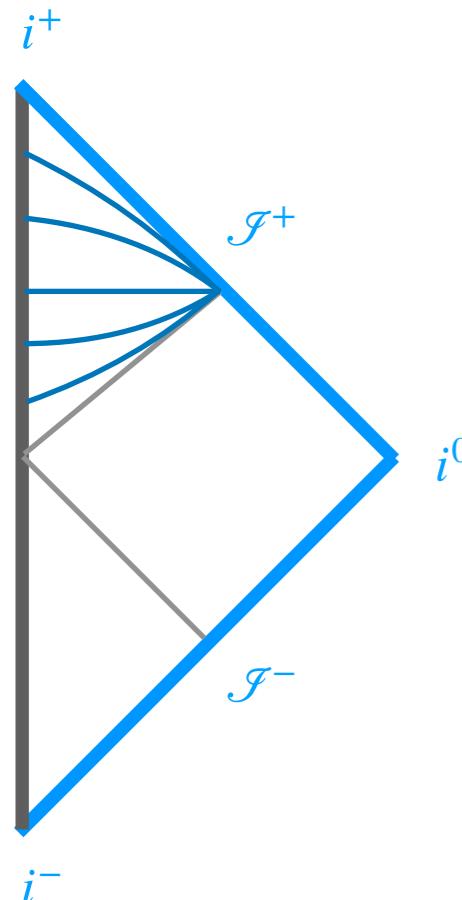
[Choi,Laddha,AP'24]



# Effect of long-range interactions

[Choi,Laddha,AP'24]

Matter is **not free** at late times but gravitationally dressed by **logs**:



$$i^+ : h_{\tau\tau}(\tau, y) \stackrel{\tau \rightarrow \infty}{=} \frac{1}{\tau} h_{\tau\tau}(y) + \dots$$

'Coulombic' mode sourced by matter stress tensor

$$\varphi(\tau, y) = \frac{\sqrt{m}}{2(2\pi)^{3/2}} \sum_{n=0}^{\infty} \frac{e^{-im\tau}}{\tau^{\frac{3}{2}+n}} \left( b_n(y) \ln \tau + b_n(y) \right) + \text{c.c.} + \dots$$

Late-time **tails** with observable consequences:

$$\mathcal{J}^+ : C_{AB}(u, x) \stackrel{u \rightarrow +\infty}{=} C_{AB}^{(0),+}(x) + \frac{1}{u} C_{AB}^{(1),+}(x) + \dots$$

↑                              ↑                              ↑

displacement memory      tail to the memory

$$= \lim_{r \rightarrow \infty} \frac{1}{r} h_{AB}(u, r, x)$$

# Gravity: hard and soft charges

[Choi,Laddha,AP'24]

$$\Omega_{i^\pm \cup \mathcal{J}^\pm} = \Omega_{i^\pm} + \Omega_{\mathcal{J}^\pm}$$



$$Q_\pm = Q_{H,\pm} + Q_{S,\pm}$$

hard charge      soft charge

$$\Omega_{i^\pm \cup \mathcal{J}^\pm}(\delta, \delta_Y) = \delta Q_\pm$$

**superrotation**

# Gravity: hard and soft charges

diverges  $\longrightarrow$  logarithmically as  $\tau \rightarrow \infty, u \rightarrow \infty$

$$\Omega_{i^\pm \cup \mathcal{J}^\pm} = \Omega_{i^\pm} + \Omega_{\mathcal{J}^\pm}$$
$$\downarrow \qquad \downarrow$$
$$Q_\pm = Q_{H,\pm} + Q_{S,\pm}$$

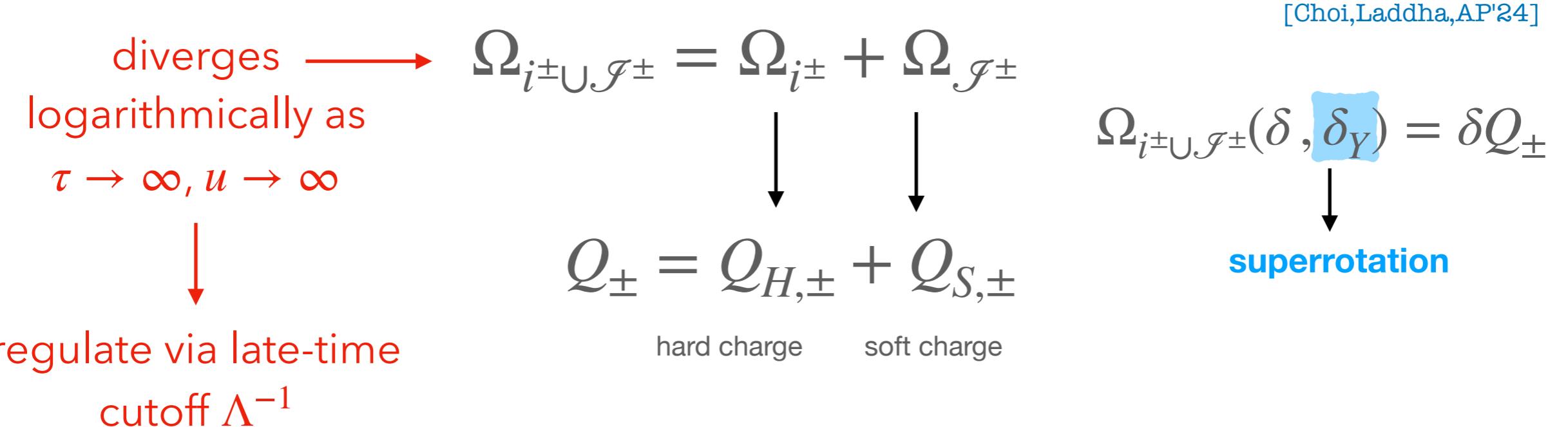
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[Choi,Laddha,AP'24]

$$\Omega_{i^\pm \cup \mathcal{J}^\pm}(\delta, \delta_Y) = \delta Q_\pm$$
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superrotation

Regularized Noether charge:

$$Q_\pm^\Lambda = \ln \Lambda^{-1} \left( Q_{H,\pm}^{(\ln)} + Q_{S,\pm}^{(\ln)} \right) + \left( Q_{H,\pm}^{(0)} + Q_{S,\pm}^{(0)} \right) + \dots$$

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[Choi,Laddha,AP'24]

$$\Omega_{i^\pm \cup \mathcal{J}^\pm} = \Omega_{i^\pm} + \Omega_{\mathcal{J}^\pm}$$



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Receives **loop-corrections**  
and is rendered  
**ambiguous** by logs.

$$Q_{H,+}^{(0)}[\bar{Y}] = \int_{i^+} d^3y \bar{Y}^\alpha \bar{T}_{\tau\alpha}^3$$

free  
stress  
tensor

$$Q_{S,+}^{(0)}[Y] = -\frac{2}{\kappa} \int_{\mathcal{J}^+} du d^2x D_z^3 Y^z u \partial_u C^{zz} + \text{c.c.}$$

$D_A \dots$  covariant  
derivative on  $S^2$

sub tree soft  
projector

# $Q^{(0)}$ charge

$$Q^\Lambda = \ln \Lambda^{-1} \left( Q_H^{(\ln)} + Q_S^{(\ln)} \right) + \left( Q_H^{(0)} + Q_S^{(0)} \right) + \dots$$

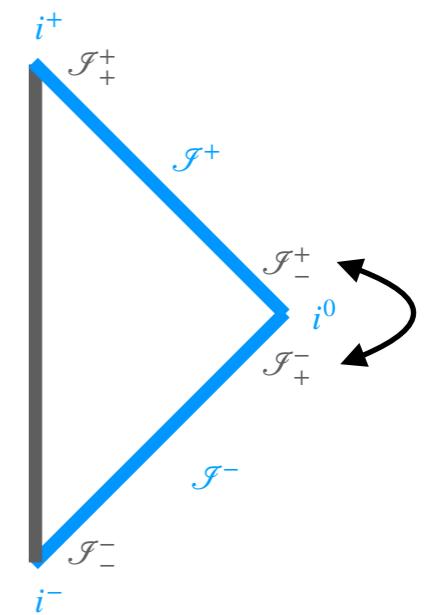
$\equiv Q^{(0)}$

Conservation law:

[Campiglia,Laddha'15]

$$Q_+^{(0)} = Q_-^{(0)}$$

Upon identifying the fields and gauge parameter antipodally:



# $Q^{(0)}$ charge

$$Q^\Lambda = \ln \Lambda^{-1} \left( Q_H^{(\ln)} + Q_S^{(\ln)} \right) + \left( Q_H^{(0)} + Q_S^{(0)} \right) + \dots$$

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[Campiglia,Laddha'15]

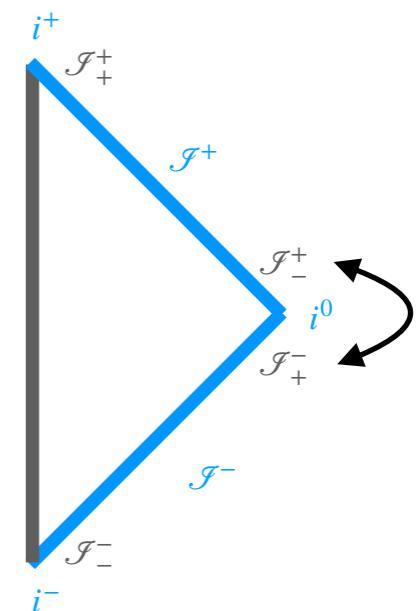
judicious choice of  
**superrotation**  $Y^A(x)$

$$Q_+^{(0)} = Q_-^{(0)}$$



[Kapec,Lysov,Pasterski,Strominger'14]

Upon identifying the fields and  
gauge parameter antipodally:



Tree-level subleading soft **graviton theorem**:

$$\mathcal{M}_{N+1} = (\omega^{-1} S_{-1} + \omega^0 S_0) \mathcal{M}_N + \dots$$

tree-level soft expansion

[Cachazo,Strominger'14]

# $Q^{(0)}$ charge

$$Q^\Lambda = \ln \Lambda^{-1} \left( Q_H^{(\ln)} + Q_S^{(\ln)} \right) + \left( Q_H^{(0)} + Q_S^{(0)} \right) + \dots$$

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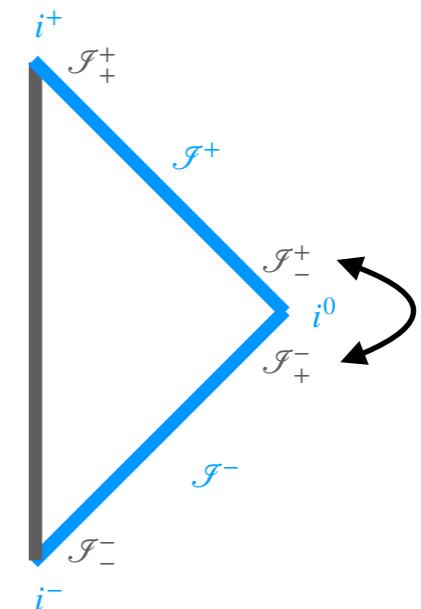
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[

Kapec,Lysov,Pasterski,Strominger'14]

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$$\mathcal{M}_{N+1} = (\omega^{-1} S_{-1} + \omega^0 S_0) \mathcal{M}_N + \dots$$

tree-level soft expansion

[Cachazo,Strominger'14]

Recall: IR effects render **subleading** soft theorem at **tree-level ambiguous**.

# Gravity: hard and soft charges

[Choi,Laddha,AP'24]

$$\Omega_{i^\pm \cup \mathcal{J}^\pm} = \Omega_{i^\pm} + \Omega_{\mathcal{J}^\pm}$$



$$Q_\pm = Q_{H,\pm} + Q_{S,\pm}$$

hard charge      soft charge

$$\Omega_{i^\pm \cup \mathcal{J}^\pm}(\delta, \delta_Y) = \delta Q_\pm$$

superrotation

Regularized Noether charge:

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**superrotation**

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$$Q_{H,+}^{(\ln)}[\bar{Y}] = \int_{i^+} d^3y \bar{Y}^\alpha {}^{3,\ln} T_{\tau\alpha}$$

leading  
interacting  
stress tensor

$$Q_{S,+}^{(\ln)}[Y] = -\frac{2}{\kappa} \int_{\mathcal{J}^+} du d^2x D_z^3 Y^z \partial_u (u^2 \partial_u C^{zz}) + \text{c.c.}$$

log soft  
projector

This is **exact** in the  
**gravitational coupling  $\kappa$  !**

# $Q^{(\ln)}$ charge

$$Q^\Lambda = \ln \Lambda^{-1} \left( Q_H^{(\ln)} + Q_S^{(\ln)} \right) + \left( Q_H^{(0)} + Q_S^{(0)} \right) + \dots$$

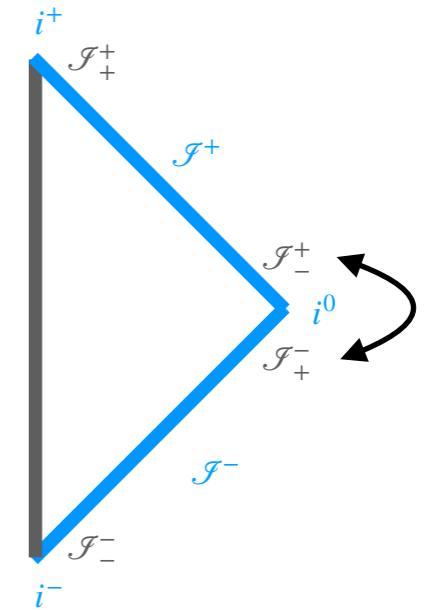
$\equiv Q^{(\ln)}$

Conservation law:

[Choi,Laddha,AP'24]

$$Q_+^{(\ln)} = Q_-^{(\ln)}$$

Upon identifying the fields and gauge parameter antipodally:



# $Q^{(\ln)}$ charge

$$Q^\Lambda = \ln \Lambda^{-1} \left( Q_H^{(\ln)} + Q_S^{(\ln)} \right) + \left( Q_H^{(0)} + Q_S^{(0)} \right) + \dots$$

$\equiv Q^{(\ln)}$

Conservation law:

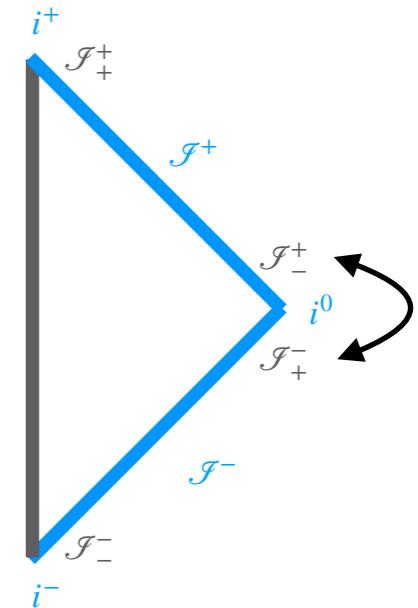
[Choi,Laddha,AP'24]

judicious choice of  
superrotation  $Y^A(x)$

$$Q_+^{(\ln)} = Q_-^{(\ln)}$$



Upon identifying the fields and  
gauge parameter antipodally:



Logarithmic soft graviton theorem:

$$\mathcal{M}_{N+1} = \left( \omega^{-1} S_{-1} + \omega^0 \ln \omega S_0^{(\ln \omega)} \right) \mathcal{M}_N + \dots$$

log soft expansion

[Sahoo,Sen'18] [Saha,Sahoo,Sen'19]

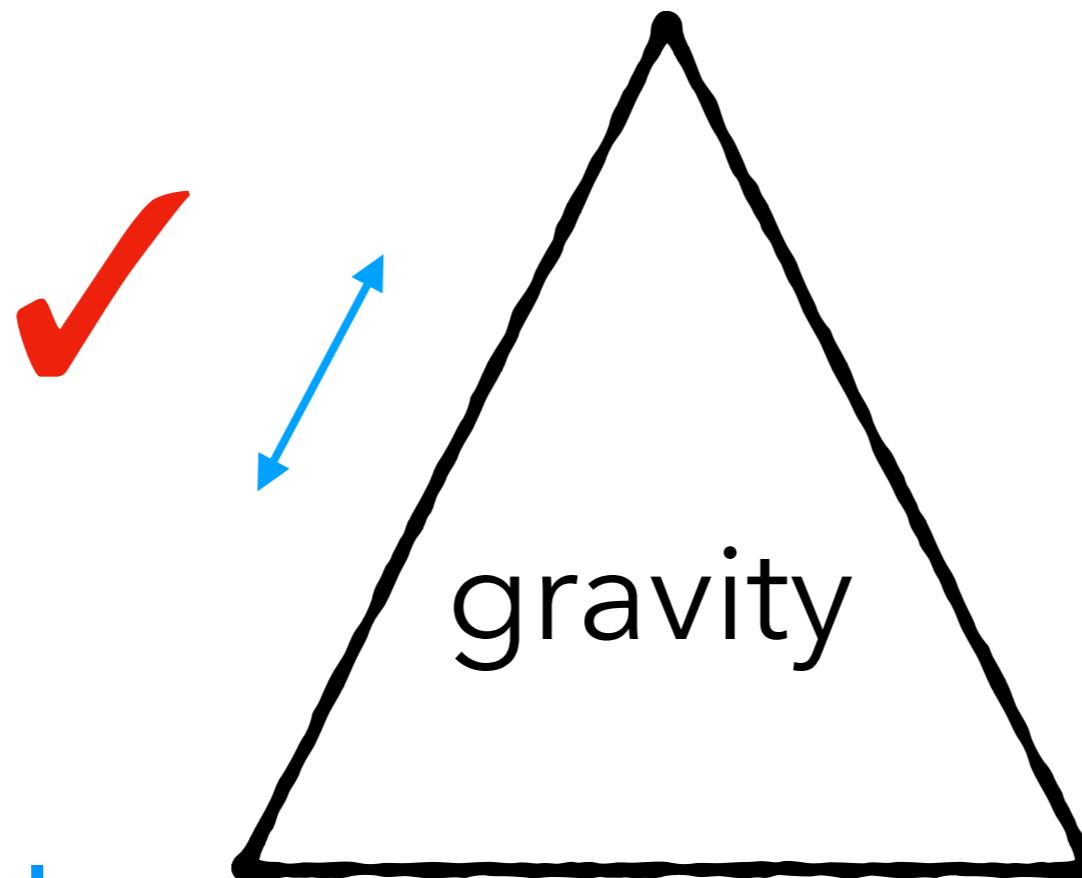
This establishes the symmetry interpretation of the  
classical logarithmic soft graviton theorem.\* [Choi,Laddha,AP'24]

\* drag term  $\simeq$  phase

# Classical superrotation IR triangle

[Choi,Laddha,AP'24]

**superrotation**



**classical log  
soft theorem**

[Laddha,Sen'18]

[Sahoo,Sen'18]

[Saha,Sahoo,Sen'19]

# Memory and its tail

Shear @ late times:  $C_{AB}(u, x) \xrightarrow{u \rightarrow \pm\infty} C_{AB}^{(0),+}(x) + \frac{1}{u} C_{AB}^{(1),+}(x) + \dots$

**linear displacement memory**  
sourced by matter field:

$$\Delta C_{AB}^{(0)} = C_{AB}^{(0),+} - C_{AB}^{(0),-}$$

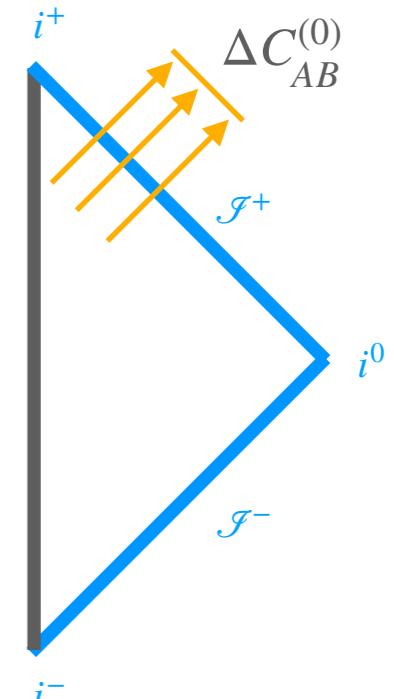
$$C_{AB}^{(0),\pm} = -\frac{\kappa^2}{8\pi} \int_{i^\pm} d^3y \frac{(\partial_A q \cdot \mathcal{Y})(\partial_B q \cdot \mathcal{Y}) + \frac{1}{2}\gamma_{AB}}{q \cdot \mathcal{Y}} \overset{3}{T}_{\tau\tau}^{\text{matt}}$$

**free**

**matter stress tensor**

$$x^\mu = \tau \mathcal{Y}(y)$$

unit vector in Minkowski



# Memory and its tail

Shear @ late times:  $C_{AB}(u, x) \xrightarrow{u \rightarrow +\infty} C_{AB}^{(0),+}(x) + \frac{1}{u} C_{AB}^{(1),+}(x) + \dots$

encodes gravitational waves

**linear displacement memory**

sourced by matter field:

$$\Delta C_{AB}^{(0)} = C_{AB}^{(0),+} - C_{AB}^{(0),-}$$

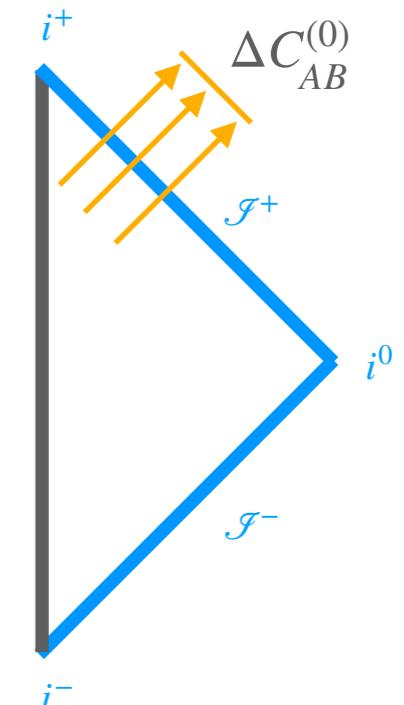
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**free**

**matter stress tensor**

$$x^\mu = \tau \mathcal{Y}(y)$$

unit vector in Minkowski



**tail to the memory**

sourced by matter field:

$$\Delta C_{AB}^{(1)} = C_{AB}^{(1),+} - C_{AB}^{(1),-}$$

**long-range interaction**

$$C_{AB}^{(1),\pm} = -\frac{\kappa^2}{8\pi} (\partial_A q^\mu)(\partial_B q^\nu) \int_{i^\pm} d^3y \left[ \frac{(q \cdot \mathcal{Y}) \mathcal{D}^\alpha (\mathcal{Y}_\mu \mathcal{Y}_\nu) - (\mathcal{Y}_\mu \mathcal{Y}_\nu + \frac{1}{2}\eta_{\mu\nu}) \mathcal{D}^\alpha (q \cdot \mathcal{Y})}{q \cdot \mathcal{Y}} \overset{3,\ln}{T}_{\tau\alpha}^{\text{matt}} \right]$$

**matter stress tensor**

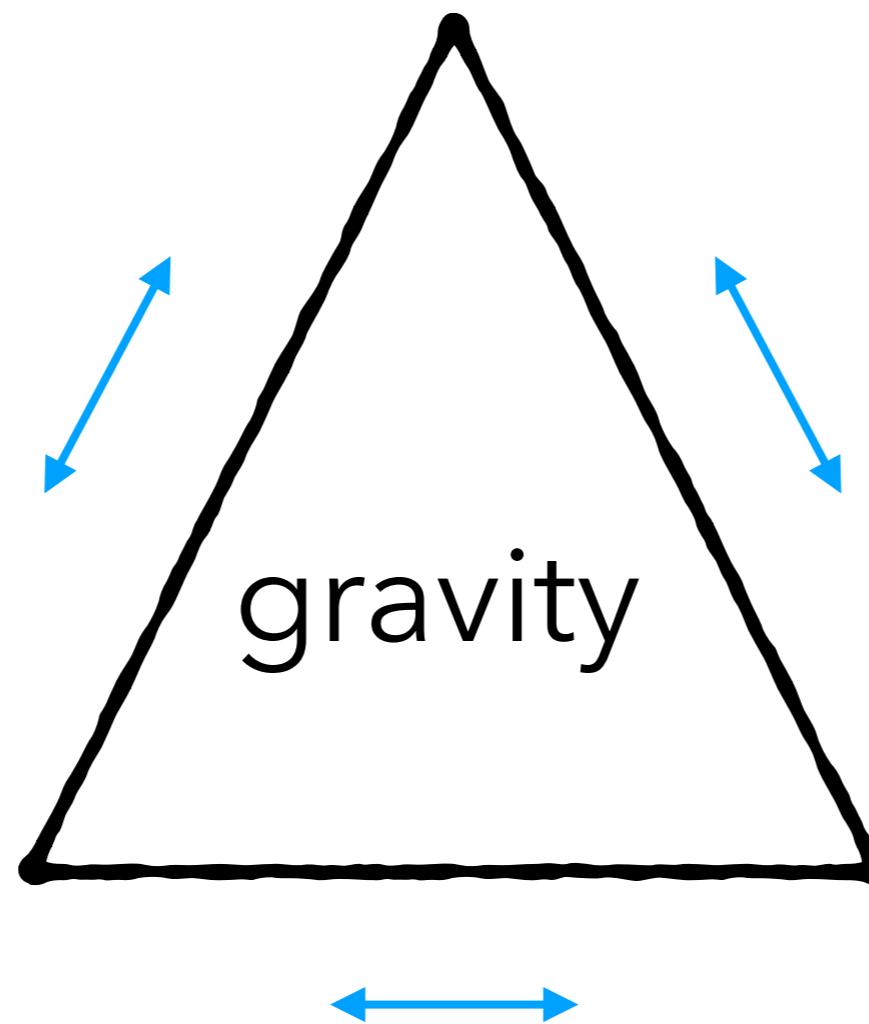
$$- \left( (\mathcal{Y}_\mu \mathcal{Y}_\nu + \frac{1}{2}\eta_{\mu\nu}) k^{\alpha\beta} + (\mathcal{D}^\alpha \mathcal{Y}_\mu)(\mathcal{D}^\beta \mathcal{Y}_\nu) - \frac{1}{2}\eta_{\mu\nu} (\mathcal{D}^\alpha \mathcal{Y}_\sigma)(\mathcal{D}^\beta \mathcal{Y}^\sigma) \right) \overset{2}{T}_{\alpha\beta}^h$$

**gravity stress tensor**

# Classical superrotation IR triangle

[Choi,Laddha,AP'24]

**superrotation**



**classical log  
soft theorem**

[Laddha,Sen'18]

[Sahoo,Sen'18]

[Saha,Sahoo,Sen'19]

**tail to the  
memory effect**

[Saha,Sahoo,Sen'19]

[Choi,Laddha,AP'24]

particles

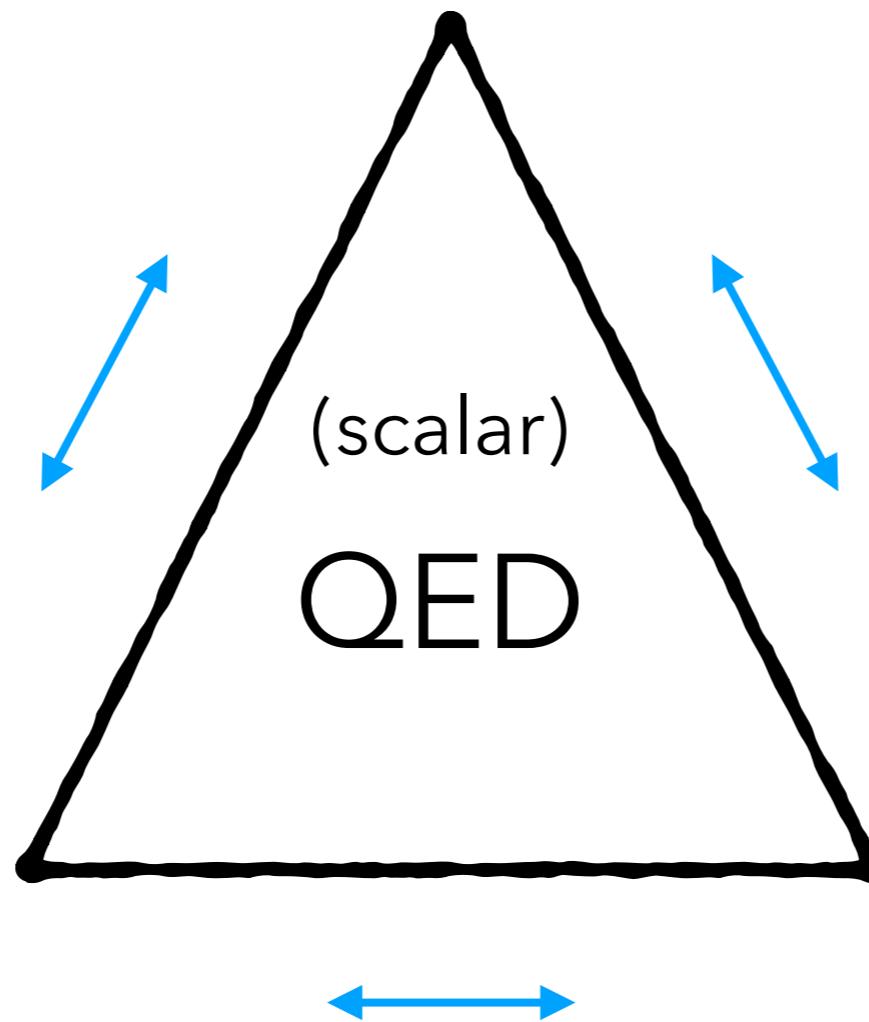
**fields**



# Classical superphaserotation IR triangle

[Choi,Laddha,AP'24]

**superphaserotation**



**classical log  
soft theorem**

[Laddha,Sen'18]  
[Sahoo,Sen'18]  
[Saha,Sahoo,Sen'19]

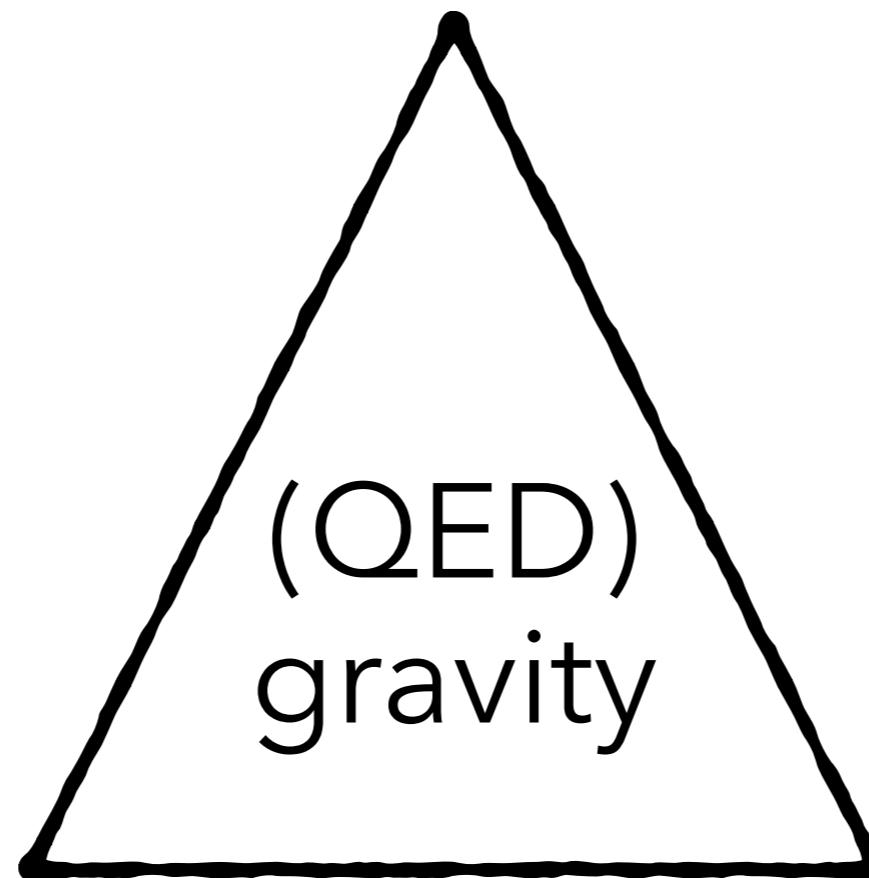
**tail to the  
memory effect**

[Saha,Sahoo,Sen'19] particles  
[Choi,Laddha,AP'24] fields

# Classical IR triangle with **massive** matter

[Choi,Laddha,AP'24]

**super(phase)rotation**



**classical log  
soft theorem**

[Laddha,Sen'18]

[Sahoo,Sen'18]

[Saha,Sahoo,Sen'19]

**tail to the  
memory effect**

[Saha,Sahoo,Sen'19]

[Choi,Laddha,AP'24]

particles

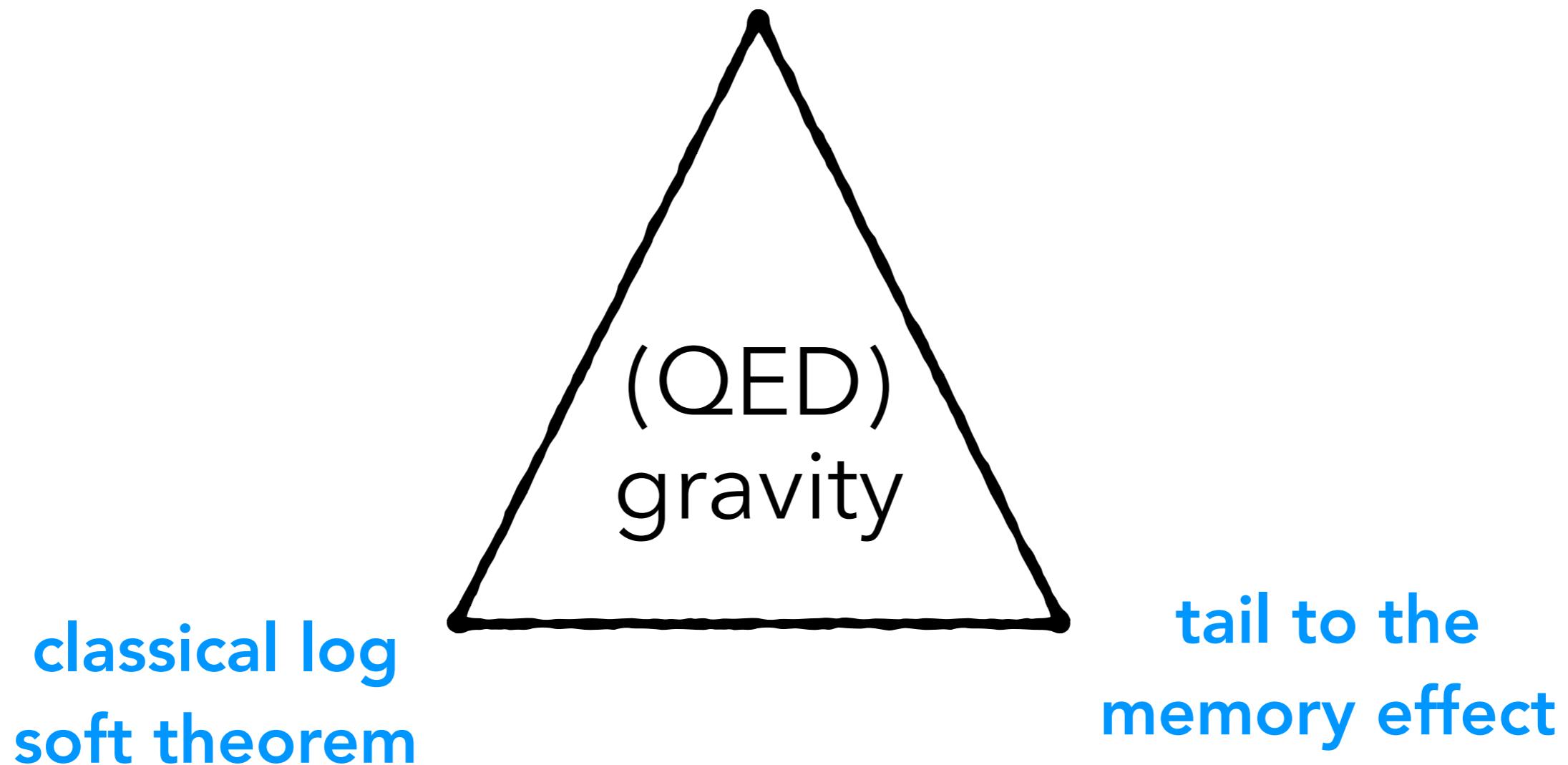
fields

# Classical IR triangle with **massless** matter

(more subtle)

[Choi,Kadhe,AP - to appear] [Choi,Kadhe,AP - work in progress]

## super(phase)rotation



# New conservation laws

Noether charge from long-range IR effects:

$$Q^\Lambda = \ln \Lambda^{-1} \left( Q_H^{(\ln)} + Q_S^{(\ln)} \right) + \left( Q_H^{(0)} + Q_S^{(0)} \right) + \dots$$

physical IR scale

- ▶ super (phase) rotation symmetry ✓
- ▶ associated charge  $Q^\Lambda$  from first principles & exact in the coupling !
- ▶ regulator  $\Lambda^{-1}$  arises from the relevant infrared scale: large  $|\tau|$  &  $|u|$
- ▶  $\Lambda$  can be removed in the end since  $Q_\pm^{(\ln)}$  are finite
- ▶ log charge conservation law = log soft theorem ✓

**Exact in coupling !**

**Exact in coupling !**

**Long-range interactions lead to new conservation laws!**

# Infrared surprises

## Infrared triangles

Complete?

## 4D = 2D

What are the axioms of celestial CFT?

Exact celestial duals? [\[Costello,Paquette'22\]](#)

[\[Costello,Paquette,Sharma'22\]](#)

## Towers of $\infty$ symmetries

Symmetries of what theories?

How powerful constraints?

## Long-range effects

Beyond gravity & QED ?   Beyond leading log ?

see also [\[Donnay,Nguyen,Ruzziconi'22\]](#),  
[\[Agrawal,Donnay,Nguyen,Ruzziconi'23\]](#)

Quantum log soft factor?

$\log(u)$ ? [\[Campiglia,Laddha'19\]](#)

logarithmic CFT? [\[Bissi,Donnay,Valsesia'24\]](#)

# In the spirit of PASCOS

## PArticles

symmetries

scattering amplitudes

infrared divergences

soft physics

...

$\exists \Lambda$  deformation of  $w_{1+\infty}$

## COSmology

## Strings

string world sheet pinned  
to celestial sphere:

$$\sum_i \text{Im}(\Delta_i) \rightarrow \infty$$

quantum gravity  
NOT in a box  
(unlike AdS/CFT)

time is emergent in  
celestial & dS holography

**Math**  
**Gravitational waves**

Thank you!