Scotogenic mechanism from a 3221 symmetry [2507.XXXX] (soon[™])

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From ancient greek

Σκότος : Darkness Γένος : Kin, generation

• Take the SM and add a Z₂ symmetry

 $SU(3)_C \times SU(2)_L \times U(1)_Y \times Z_2$





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- New scalar doublet and two to three $\eta \sim (2, 1/2, -)$ $N^{\sigma} \sim (1, 0, -)$ generations of a neutral fermionic singlet



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- All **SM fields are even** under Z₂, new **BSM** $L, H, e_R \dots \sim +$ fields will be odd
- $\eta \sim (2, 1/2, -) \quad N^{\sigma} \sim (1, 0, -)$ New scalar doublet and two to three generations of a **neutral fermionic singlet**
- The η_0 VEV is zero, so that \mathbf{Z}_2 is exact

 $\langle \eta_0 \rangle = 0 \rightarrow Z_2$ does not break

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Fermionic: Lightest mass eigenstate of N^{σ}

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Bosonic: Lightest $\operatorname{Re}(\eta_0)$ or $\operatorname{Im}(\eta_0)$ **Fermionic:** Lightest mass eigenstate of N^{σ} Radiative Majorana neutrino masses at 1-loop LO



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 $m_{\nu}^{\alpha\beta} \propto \frac{\lambda_5 \ \nu^2 \ Y^2}{16\pi^2 M}$



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$$\lambda_5 Y^2 \sim 10^{-8}$$



Neutrino masses!

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$$m \relton$$
 But why's λ_5 small?



There is always a bigger fish

 We want to build a SM extension that includes the Scotogenic at low energies...



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- ... with the added benefit of generating λ_5 's smallness naturally



There is always a bigger fish

- We want to build a SM extension that includes the **Scotogenic at low energies**...
- ... with the added benefit of generating λ_5 's smallness naturally
- We will also generate Z₂ as an accidental symmetry





$SU(3)_C \times SU(2)_1 \times SU(2)_2 \times U(1)_Y$





$SU(3)_{C} \times SU(2)_{1} \times SU(2)_{2} \times U(1)_{Y}$ $L \qquad 1 \qquad 2 \qquad 1 \qquad -1/2$ $H \qquad 1 \qquad 2 \qquad 1 \qquad +1/2$





$SU(3)_C \times SU(2)_1 \times SU(2)_2 \times U(1)_Y$

L	1	2	1	-1/2
Н	1	2	1	+1/2
η	1	1	2	+1/2

+1/22











Going down the ladder

 How do we return to the scotogenic's diagram?







Why is v_{Δ} the smallest?

 The smallness of λ₅ comes from it being an effective operator after Δ is integrated out.



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Why is v_{Δ} the smallest?

- The smallness of λ_5 comes from it being an effective operator after Δ is integrated out.
- Therefore we need $m_{\Delta} \gg$ any other mass scale
- From the tadpole equations, this forces small v_{Δ}

$$v_{\Delta} \propto \frac{v_{H}^{2} v_{\Omega}}{m_{\Delta}^{2}}$$

Effective theory!





- The first symmetry breaking step comes from v_Ω and v_ξ

$$v_{\Omega}, v_{\xi} \gg v_H \gg v_L$$





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$$v_\Omega$$
 , $v_\xi \gg v_H \gg v_L$

$$SU(2)_1 \times SU(2)_2 \xrightarrow{v_\Omega, v_\xi} SU(2)_L$$





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 $\eta!$

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- The scotogenic fields always appear in pairs within the lagrangian

$$\eta \sim (\mathbf{1}, \mathbf{2}) \xrightarrow{SU(2)_L} \mathbf{2}$$

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*Artist's rendition



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• It is indeed!

Session

Particle Phenomenology

() 21 Jul 2025, 16:00**♥** Durham

Scotogenic mechanism from an extended \$SU(2)_1 \times SU(2)_2 \times U(1)_Y\$ electroweak symmetry

③ 17:20 - 17:40

Presenter Javier Perez-Soler

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 $\rho = 1.00031 \pm 0.00019$

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These limits also suppress gauge boson mixing

 $\theta_c, \theta_n \leq 10^{-3}$

erez-Soler

Secoi

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- We have Majorana neutrino masses... LFV decays?
- $\mu \rightarrow e + \gamma$ is a classic choice
- BR($\mu \to e + \gamma$) < 1.5 · 10⁻¹³

(All particles mass eigenstates)

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Final point: Dark matter

• We have a bunch of stable DM candidates thanks to Z_2

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Final point: Dark matter

- We have a bunch of stable DM candidates thanks to Z_2
- For scalars, it's the scotogenic guys η_R and η_I
- For fermions, we have two, χ_{t1} and χ_{s1} ; coming from the neutral components of

 $\chi \rightarrow 3 \oplus 1$

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What to do if i have theories?

Question

Vote

I contacted a college and they ignored me

67

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