

Confinement Slingshot: Confined Monopoles in First-Order Phase Transitions

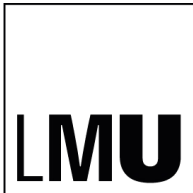
Maximilian Bachmaier

In collaboration with: Gia Dvali

Juan Sebastián Valbuena-Bermúdez

Michael Zantedeschi

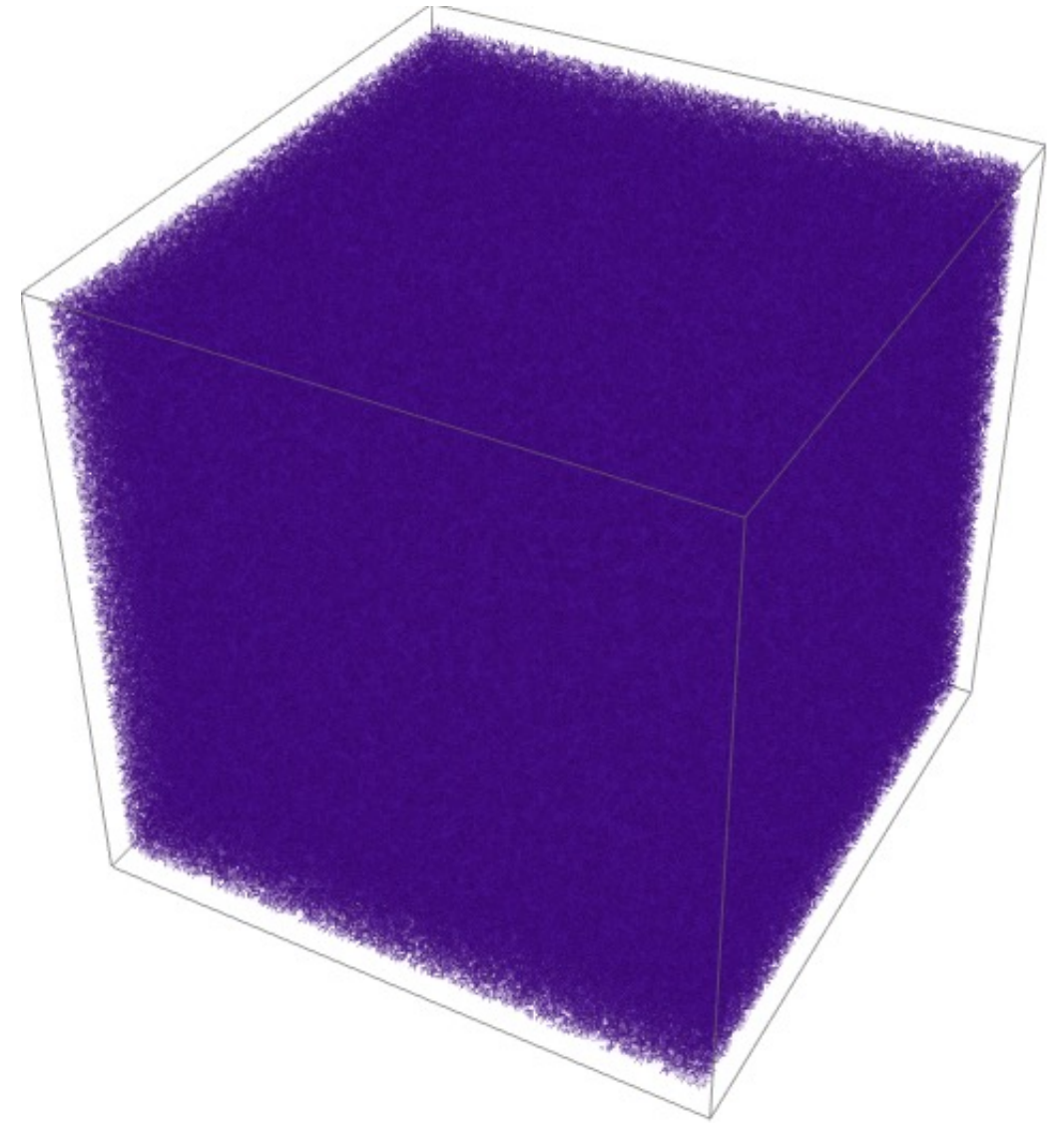
PASCOS 2025 (Durham, July 2025)



The Magnetic Monopole Problem

Grand Unified Theories predict Magnetic Monopoles
(Creation through Kibble Mechanism)

Zeldovich and Khlopov 1978, Preskill 1979:
Magnetic Monopole Problem



The Magnetic Monopole Problem - Solutions

Inflation

BUT: GUT phase transition can happen after inflation¹

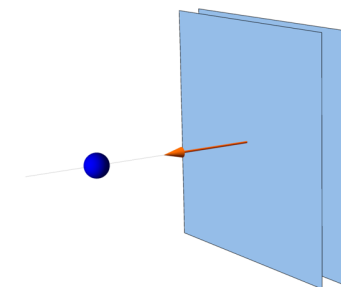
Symmetry Non-Restoration² (Dvali, Melfo and Senjanovic 1995)

Domain Walls erase Magnetic Monopoles² (Dvali, Liu and Vachaspati 1997)

Numerical Studies: Brush, Pogosian, Vachaspati 2015

Dvali, Valbuena-Bermudez 2022

Bachmaier, Dvali, Valbuena-Bermudez 2023

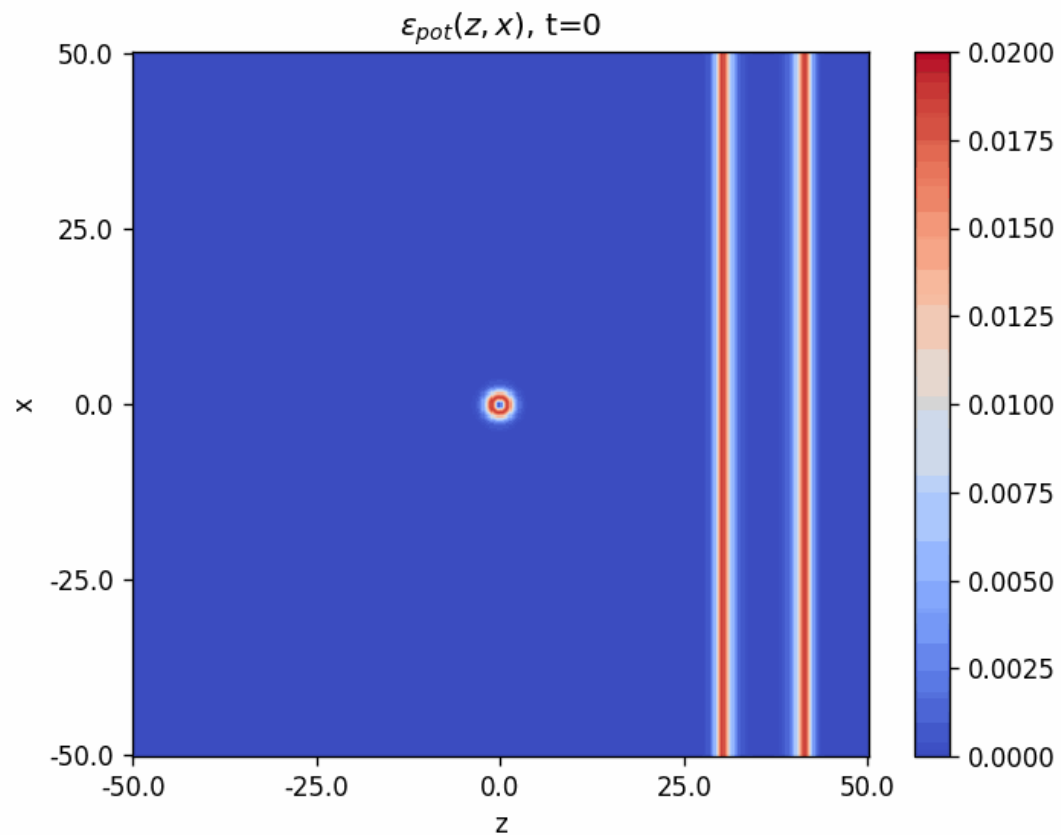


¹see for example G. Dvali, H. Tye – Brane Inflation (1998)

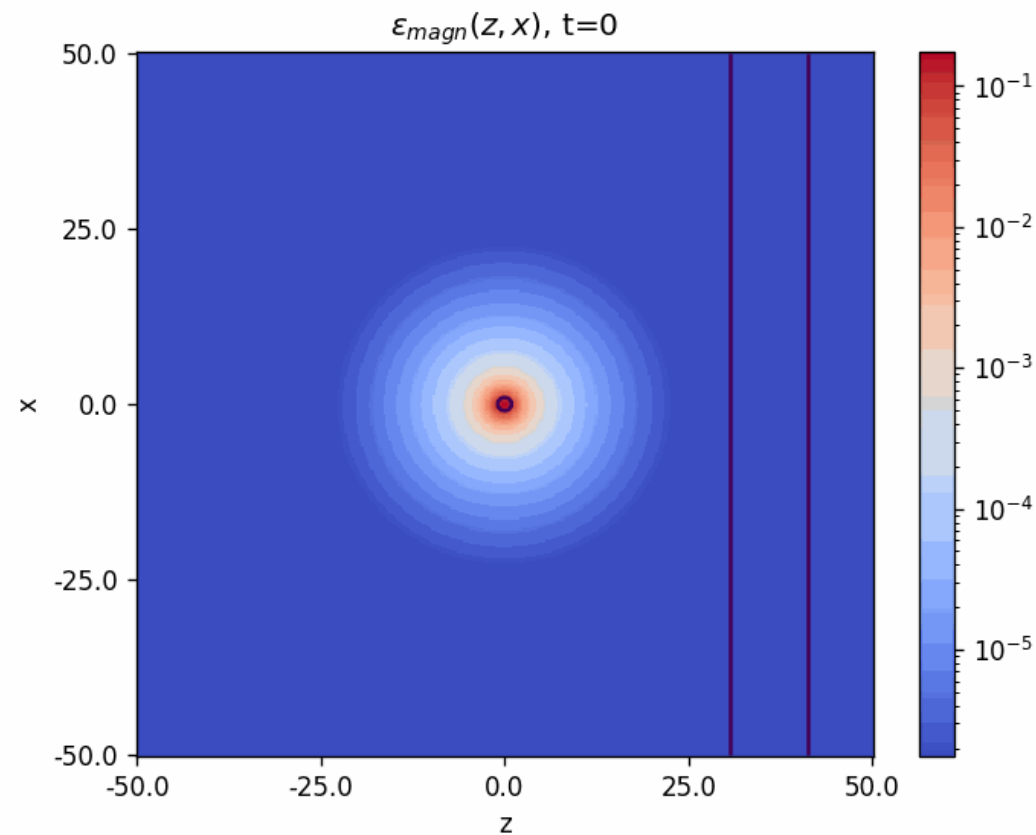
²minimal solutions, since they don't require physics beyond minimal SU(5)

Magnetic Monopole - Domain Wall Collision

Potential Energy Density



Magnetic Energy Density



The Magnetic Monopole Problem - Solutions

Inflation

BUT: GUT phase transition can happen after inflation¹

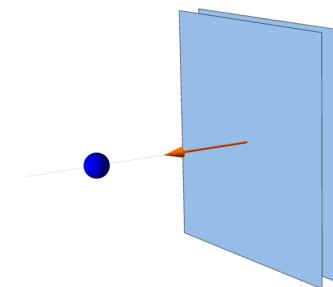
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Magnetic Monopoles connected by Strings (Langacker and Pi 1980)

Numerical Studies: Dvali, Valbuena-Bermudez, Zantedeschi 2022

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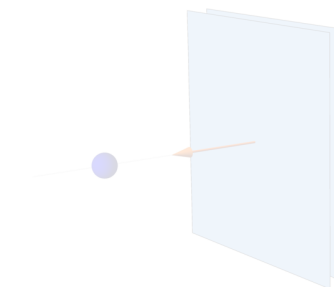
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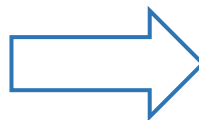
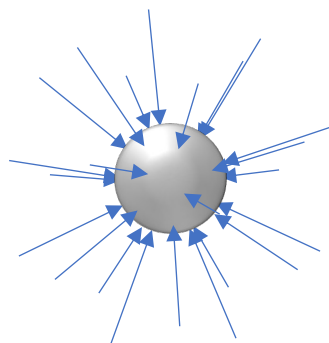
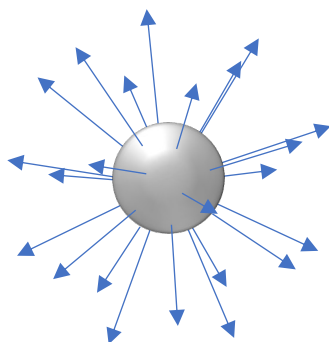


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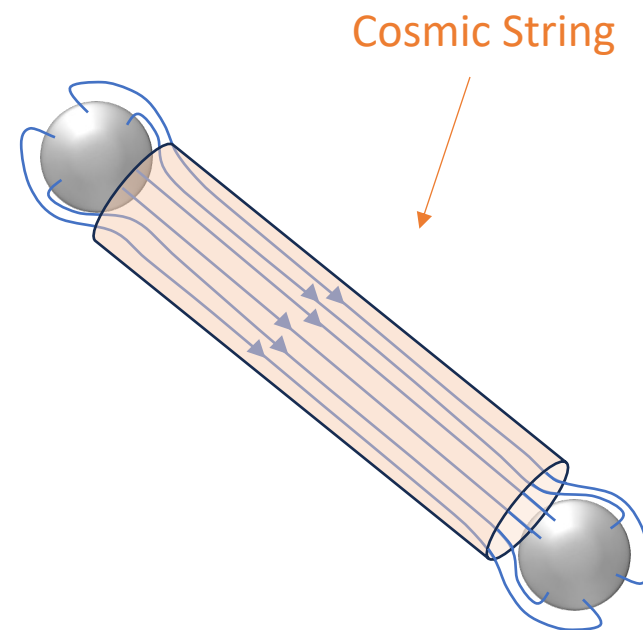
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Magnetic Monopoles Connected by a String

First Phase Transition: $G \rightarrow \dots \times U(1)$



Second Phase Transition: $\dots \times U(1) \rightarrow \dots \times \cancel{U(1)}$

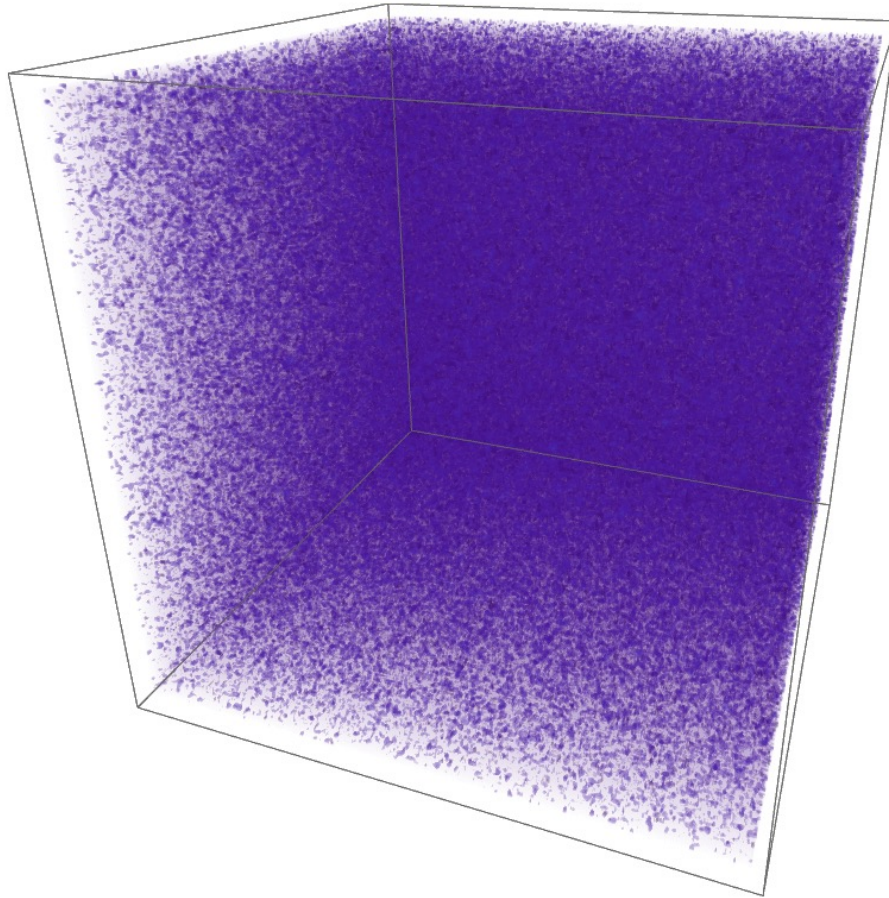


P. Langacker, S. Pi (1980)

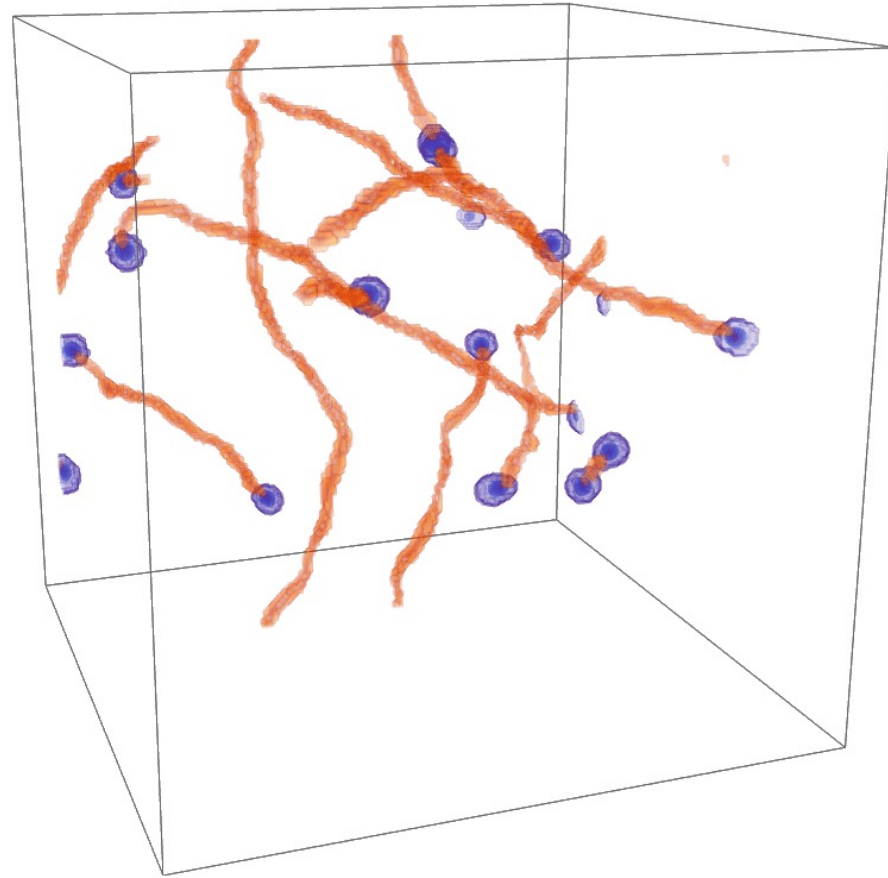
X. Martin, A. Vilenkin (1997)

G. Dvali, J. S. Valbuena-Bermúdez, M. Zantedeschi (2022)

Second-Order Phase Transition

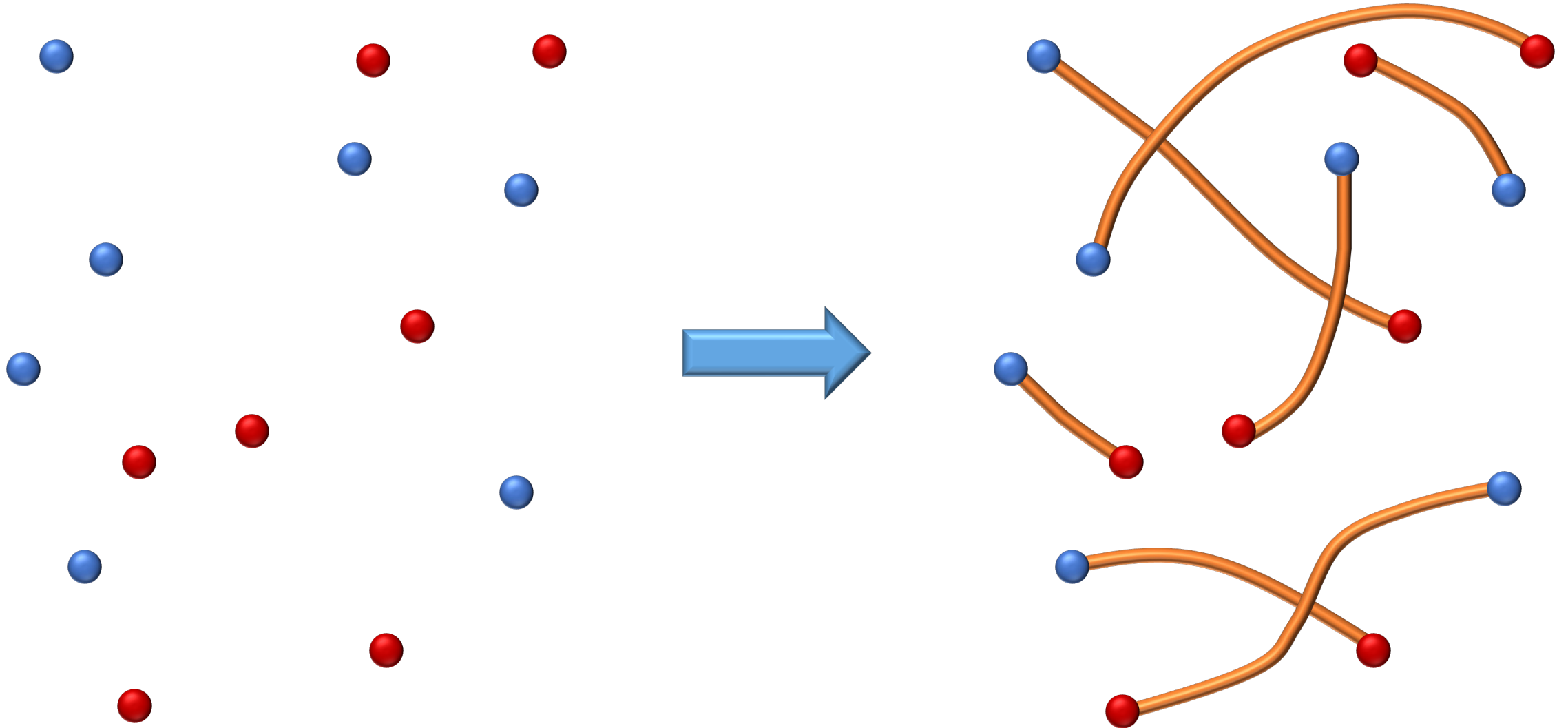


Second-Order Phase Transition

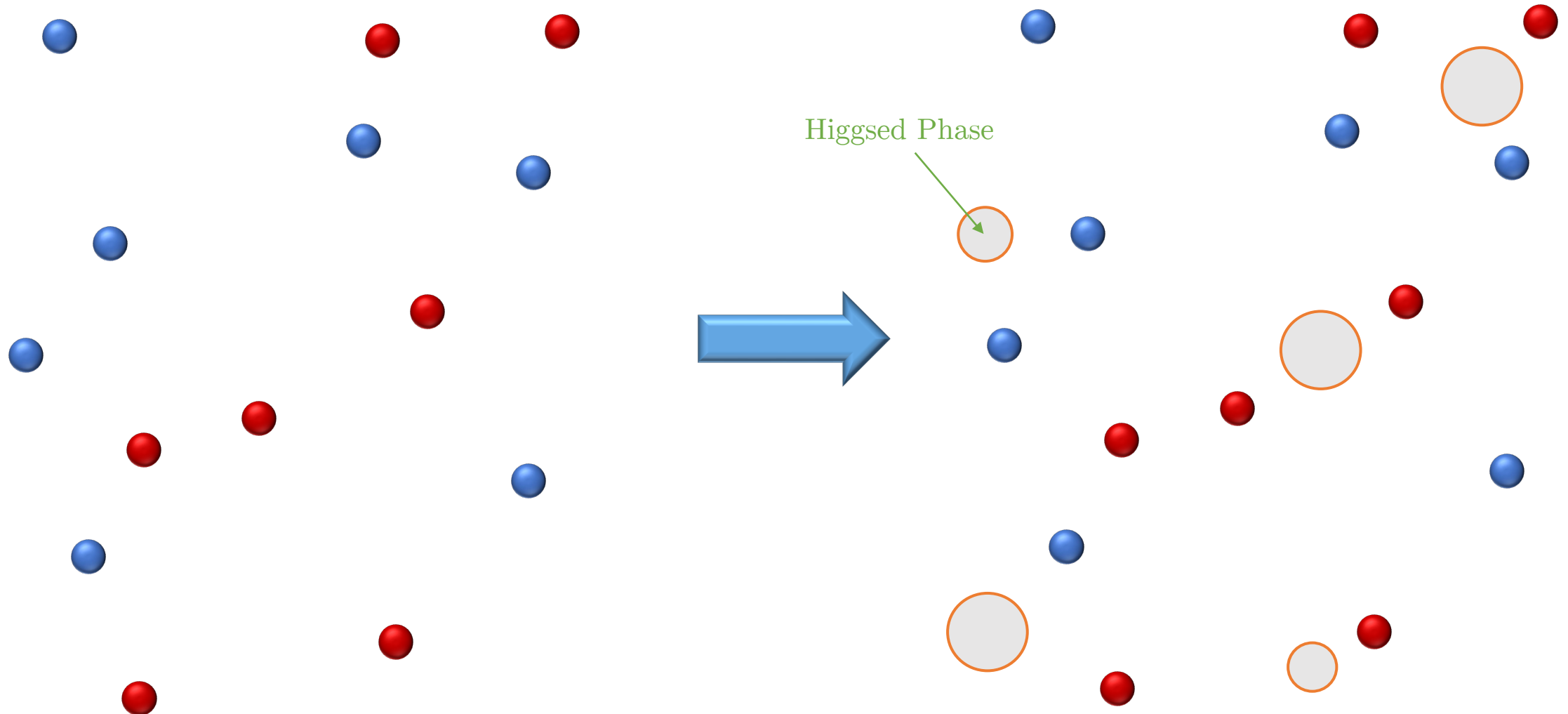


What happens in a first-order phase transition?

Second-Order Phase Transition



First-Order Phase Transition



The Model

We consider an $SU(2)$ gauge theory with the following potential

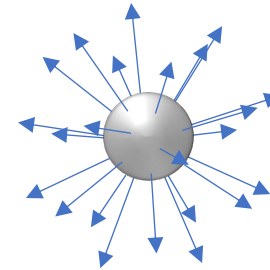
$$V(\phi) = \lambda_\phi (\phi^a \phi^a - v_\phi^2)^2 + \lambda_\psi (\psi^\dagger \psi - v_\psi^2)^2 (\psi^\dagger \psi) + \beta \psi^\dagger \phi \psi$$

ϕ : $SU(2)$ adjoint

ψ : $SU(2)$ fundamental

$\langle \phi^a \phi^a \rangle = v_\phi^2 \rightarrow SU(2)$ breaks down to $U(1)$
 $\rightarrow$ Magnetic Monopoles

't Hooft-Polyakov
Magnetic Monopole



Coulomb (unconfined)
Vacuum

The Model

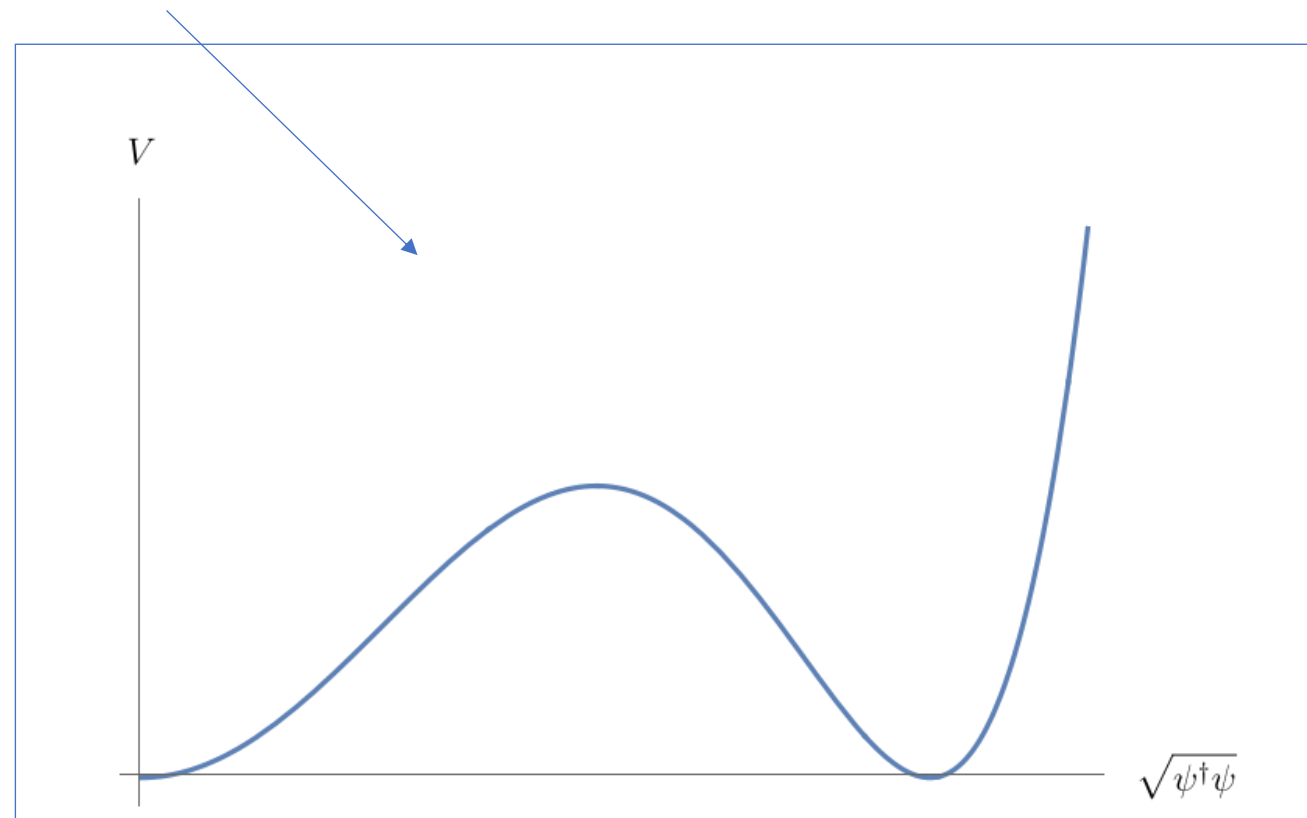
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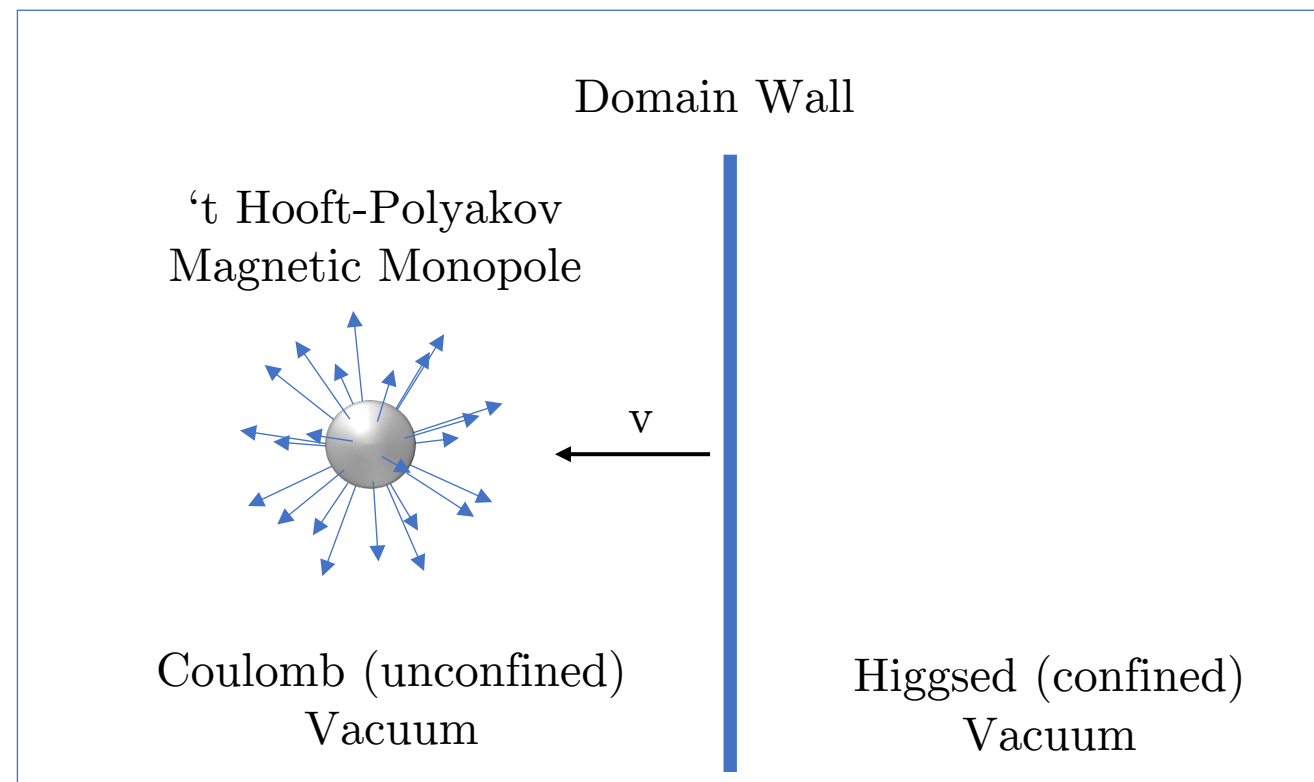
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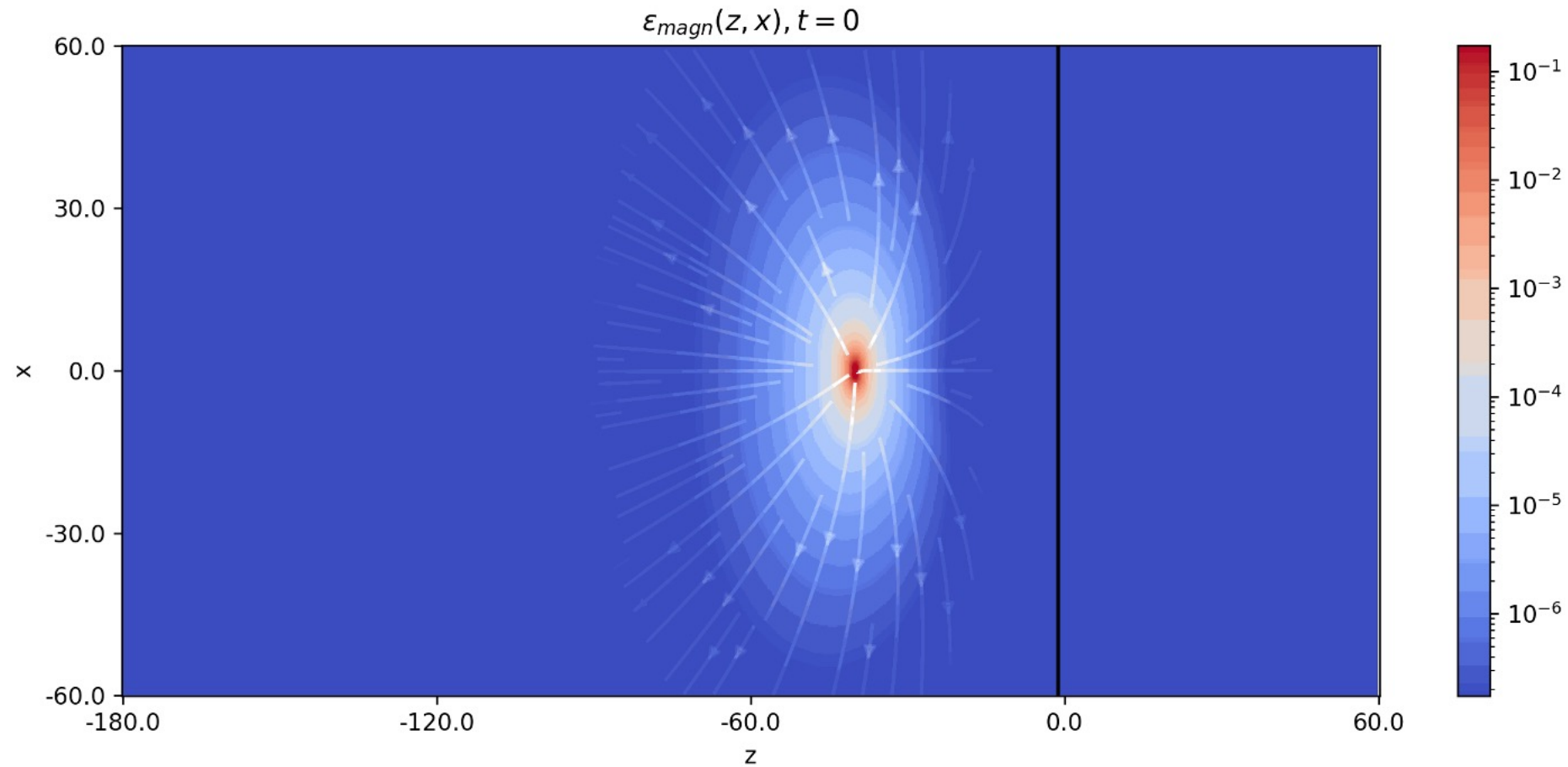
Disconnected Vacuum Manifold for ψ
 $\rightarrow$ Domain Walls

$\langle \psi^\dagger \psi \rangle = v_\psi^2 \rightarrow U(1)$ breaks down to 1
 $\rightarrow$ Cosmics Strings

Breaking Pattern: $SU(2) \xrightarrow{\phi} U(1) \xrightarrow{\psi} 1$

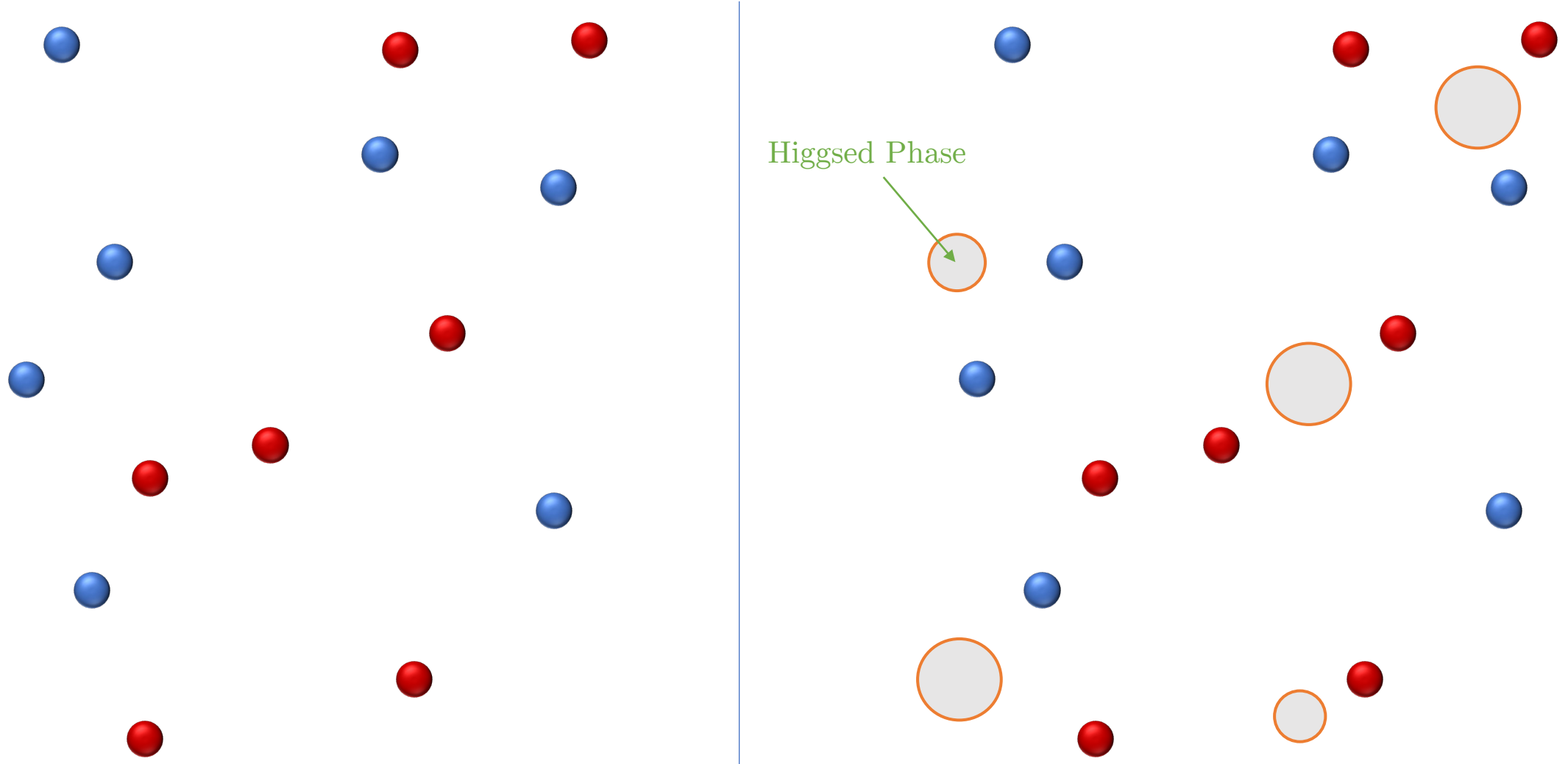


Confinement Slingshot

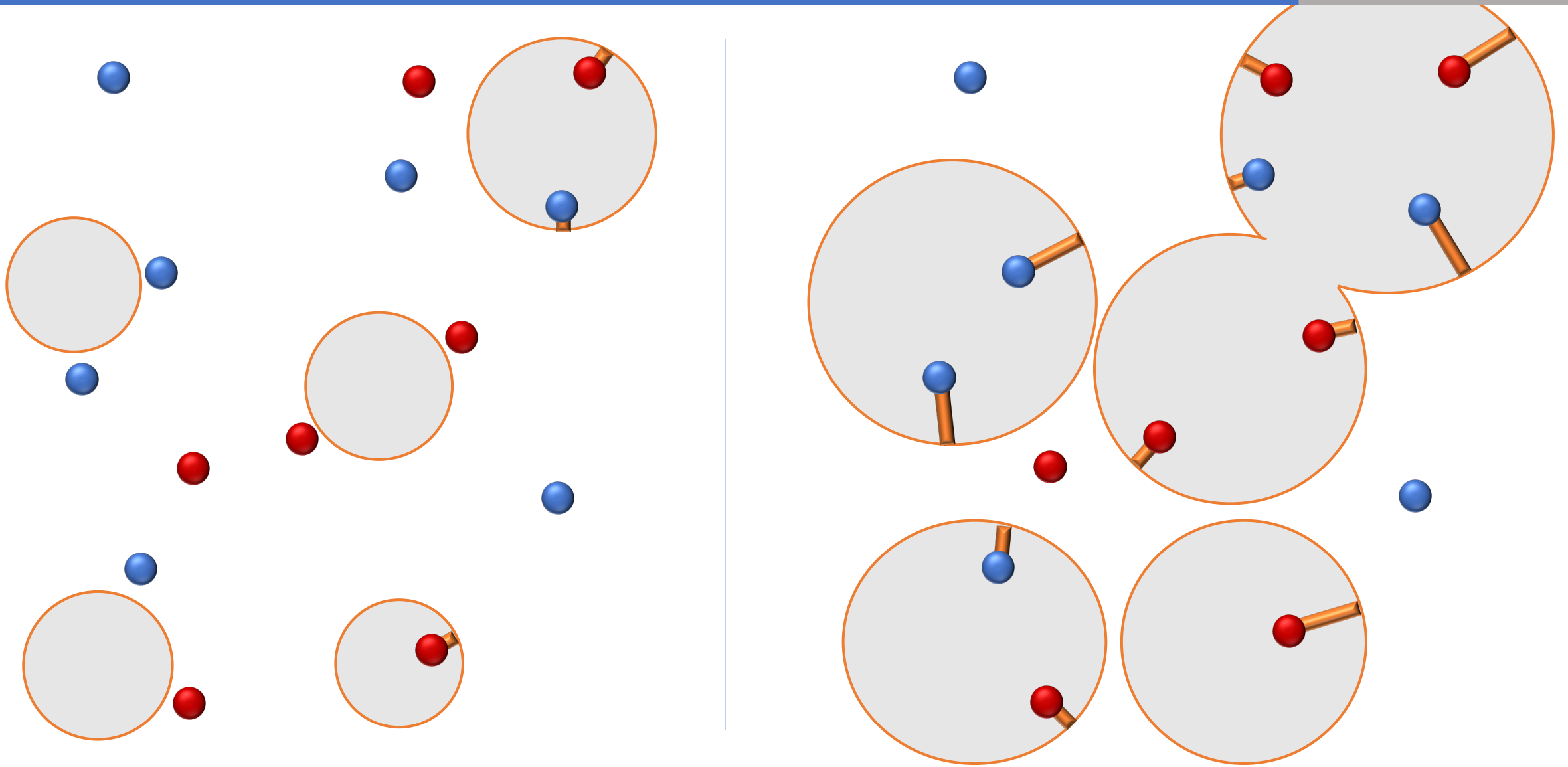


M. B., G. Dvali, J. S. Valbuena-Bermúdez, M. Zantedeschi (2023)

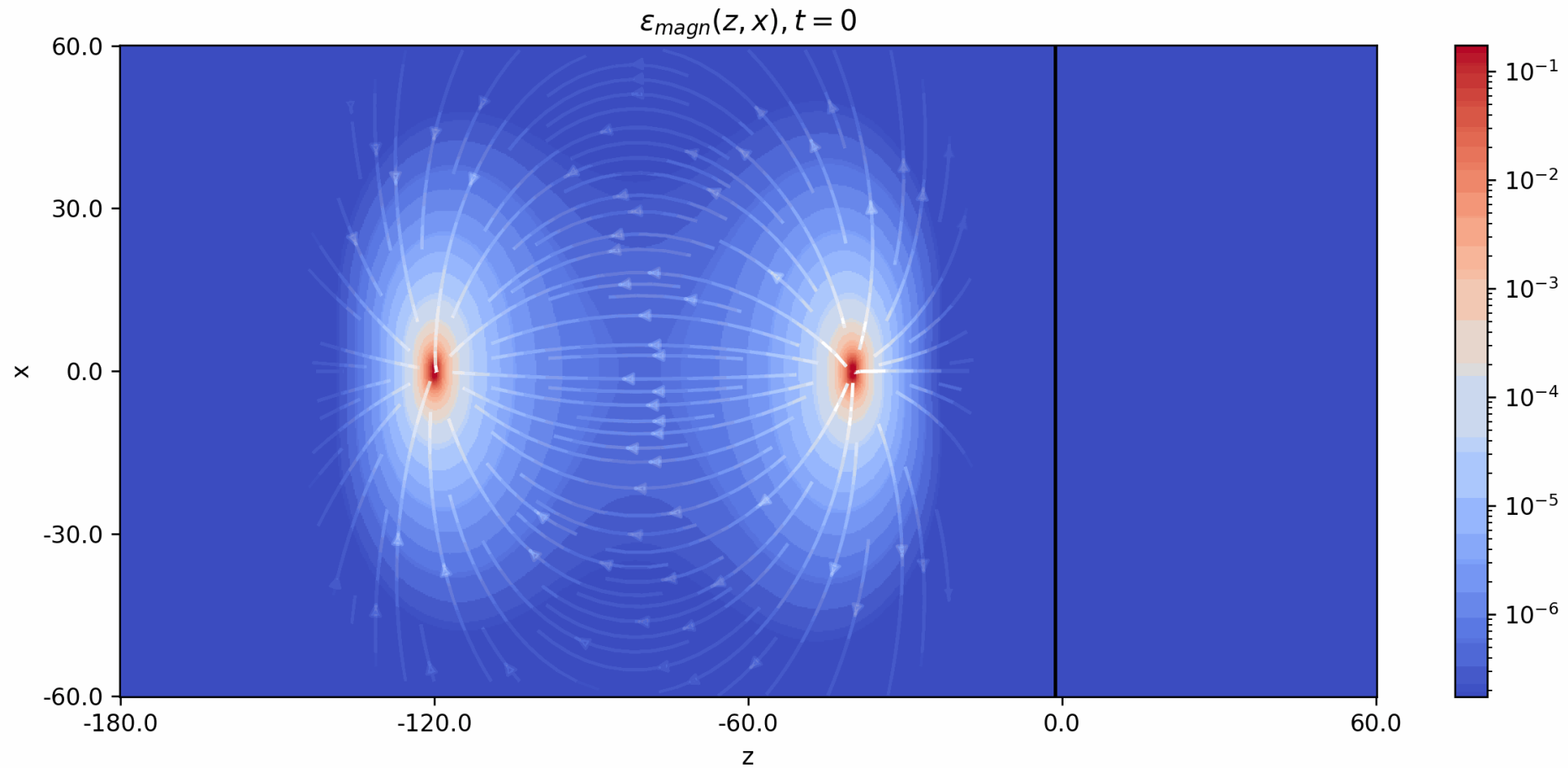
First-Order Phase Transition



First-Order Phase Transition

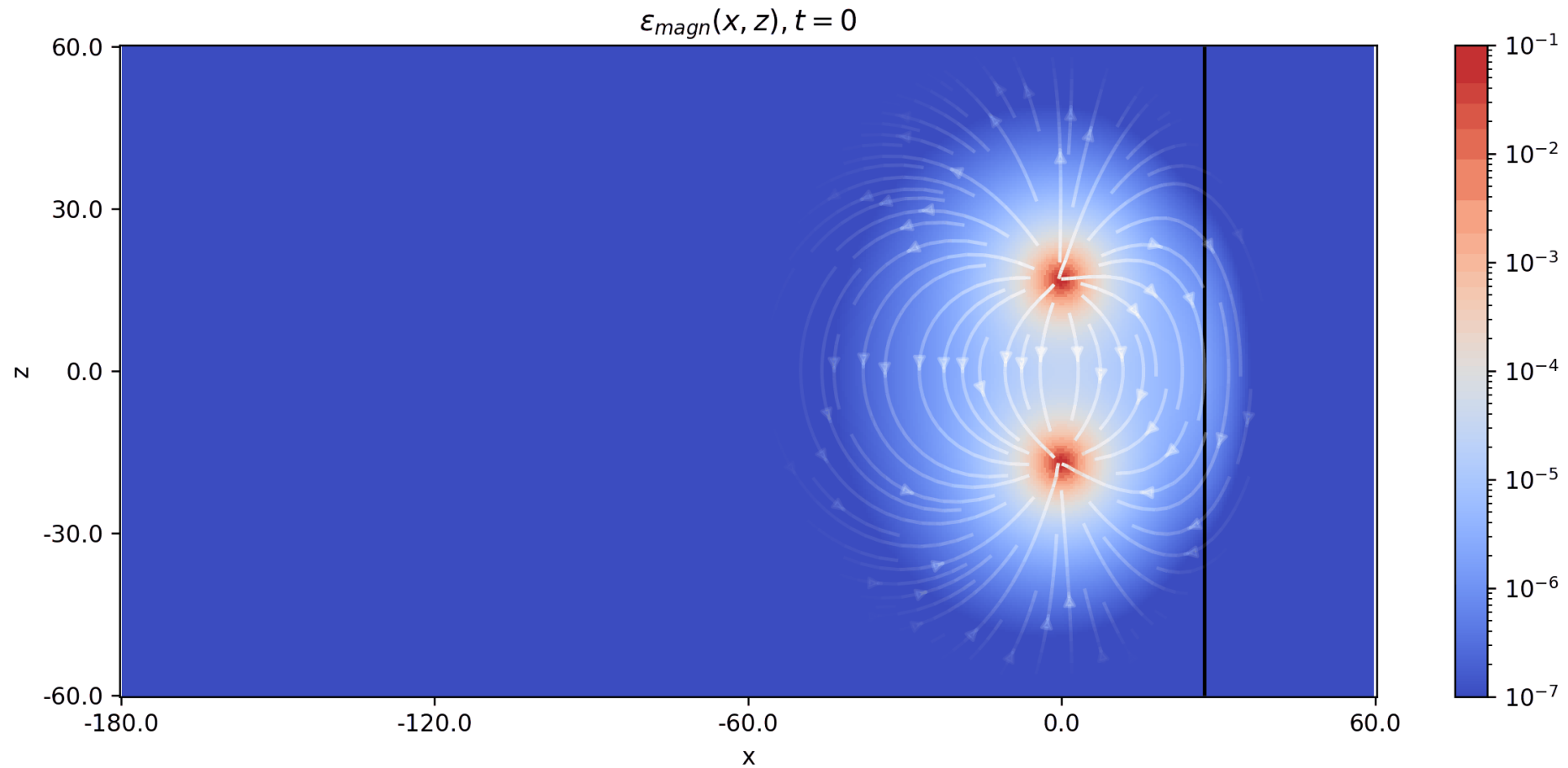


Confinement Slingshot – Two Monopoles



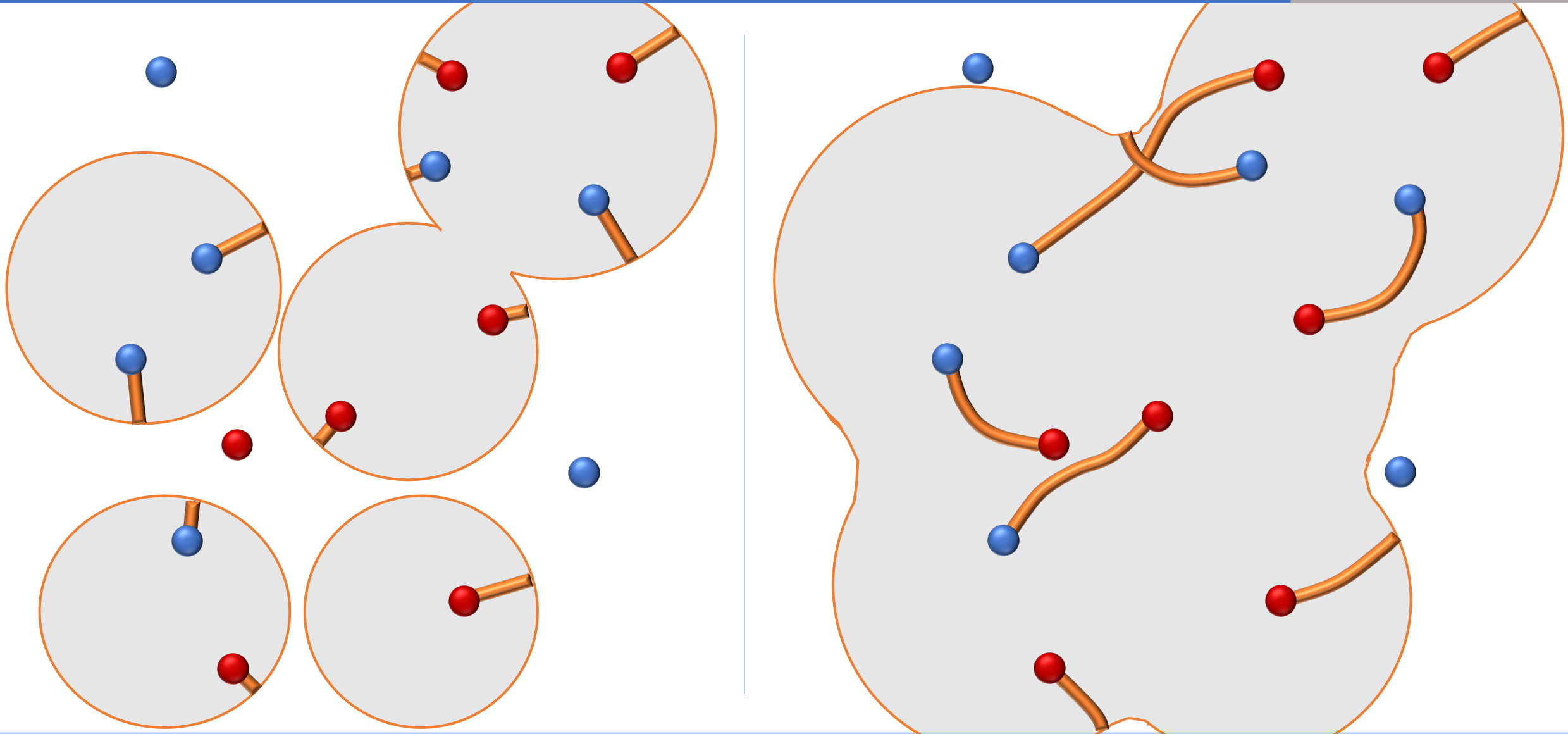
M. B., G. Dvali, J. S. Valbuena-Bermúdez, M. Zantedeschi (2023)

Multiple Confinement Slingshots

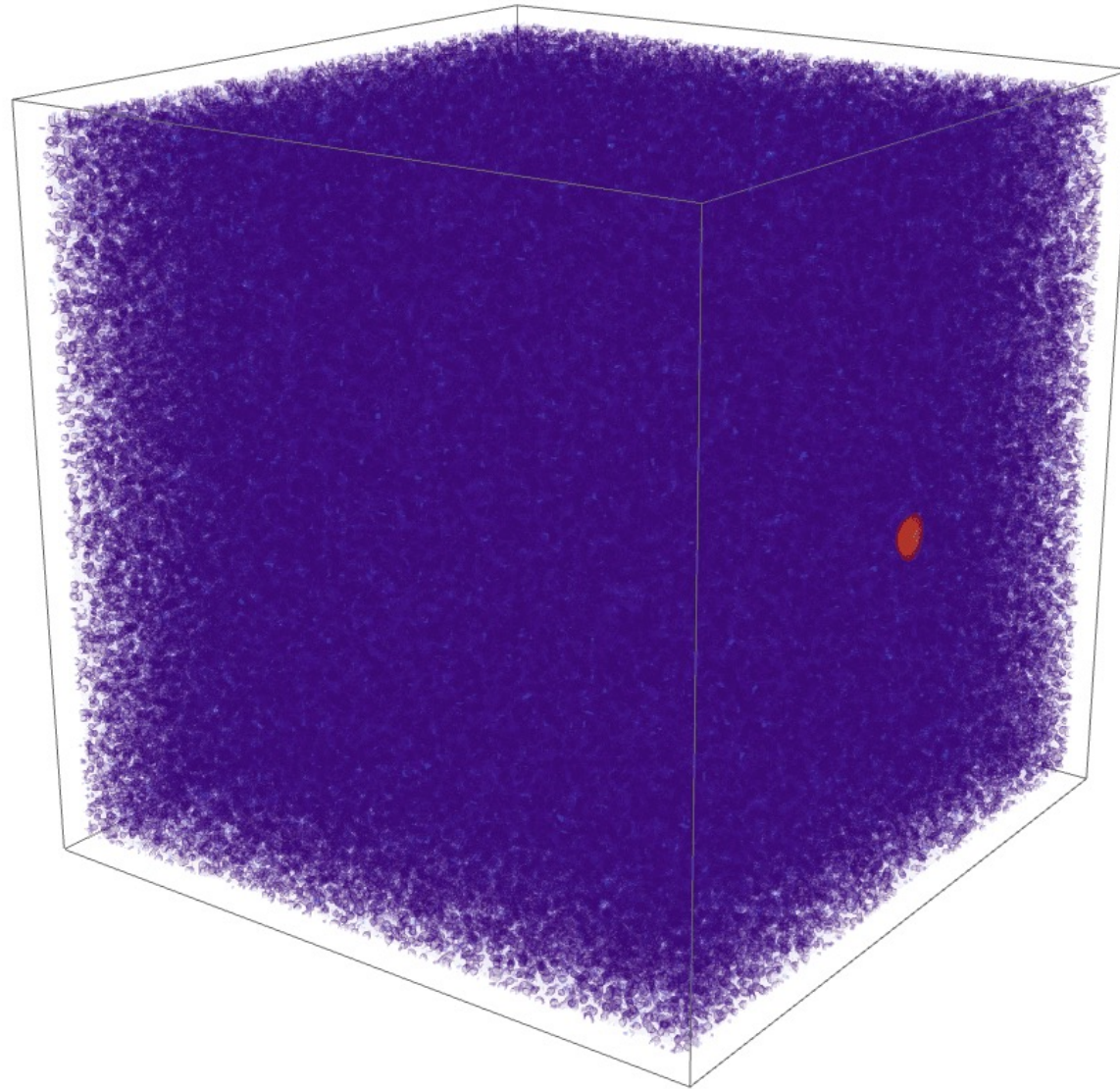


M. B., G. Dvali, J. S. Valbuena-Bermúdez, M. Zantedeschi (to be published)

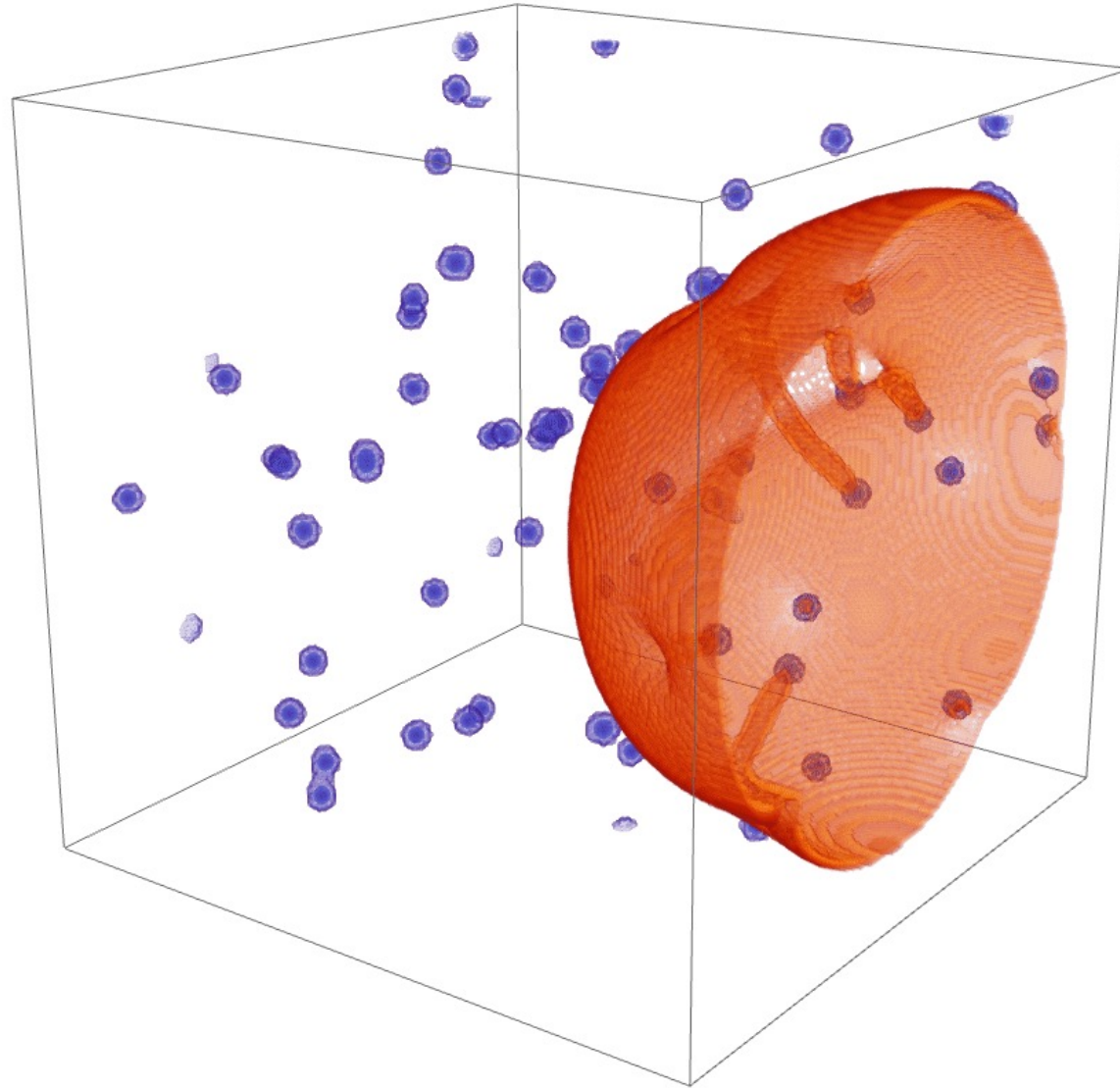
First-Order Phase Transition



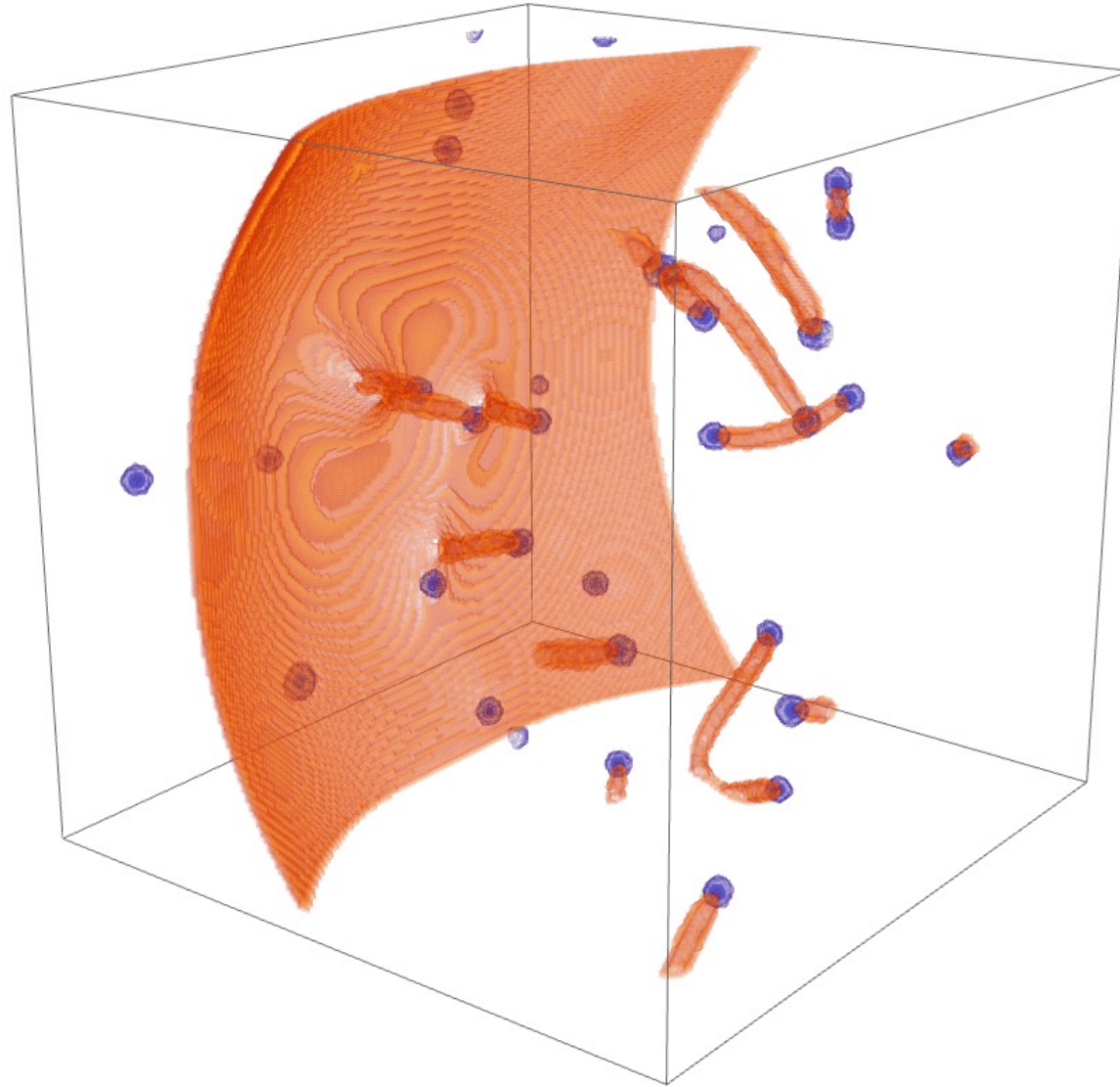
First-Order Phase Transition



First-Order Phase Transition



First-Order Phase Transition

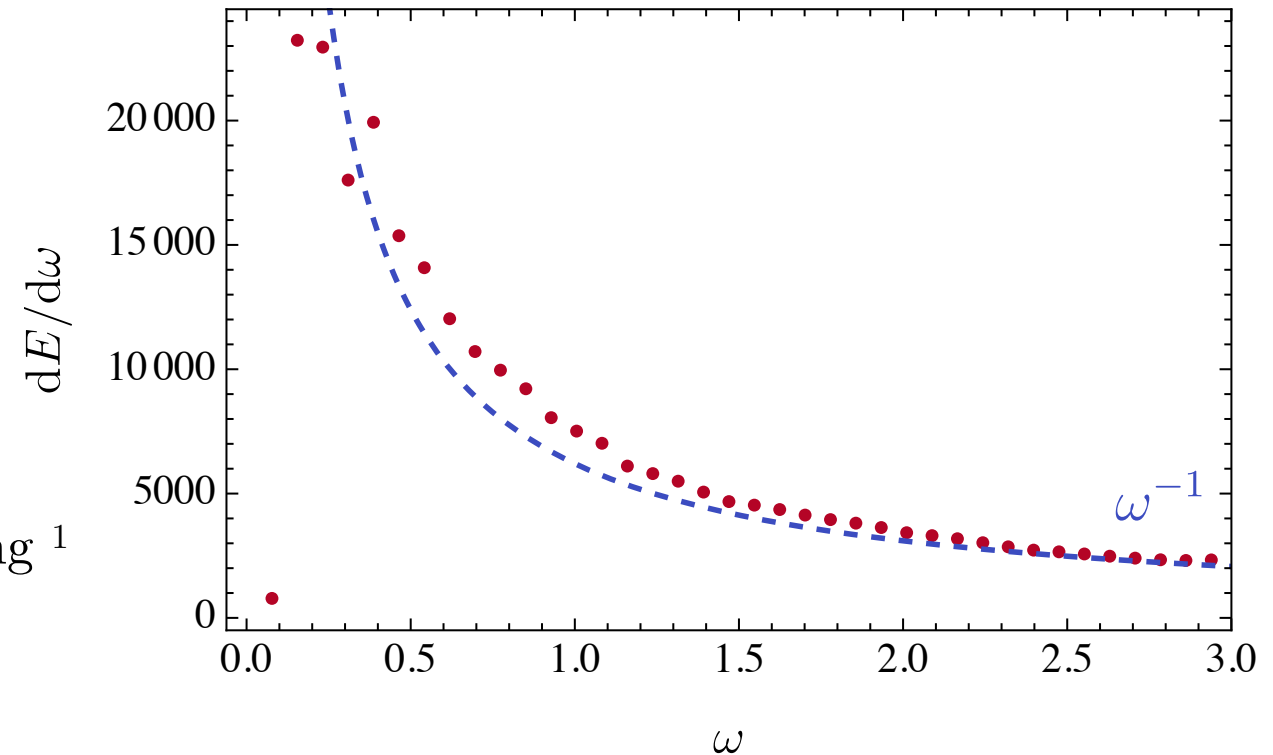


Confinement Slingshot – Gravitational Waves

The slingshot effect leads to the emission of gravitational waves

The energy spectrum decays with ω^{-1}

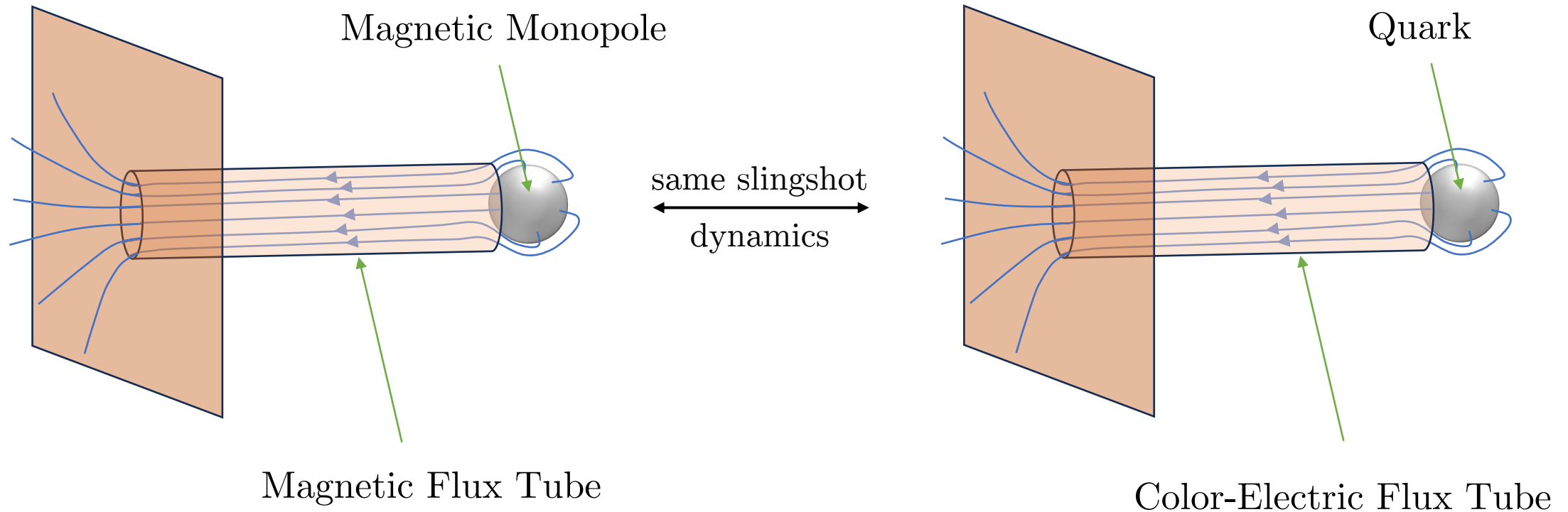
→ Similar to two monopoles connected by a string ¹



¹ X. Martin, A. Vilenkin (1997) / G. Dvali, J. S. Valbuena-Bermúdez, M. Zantedeschi (2022)

Confinement Slingshot – Quark Confinement

The similar slingshot effect is expected in a “dual” picture when a heavy quark crosses into a confined vacuum of QCD.



(In the case of light quarks, quark-antiquark pairs can emerge and break the string.)

Summary

We analyzed the slingshot effect that can appear in first-order phase transitions involving magnetic monopoles/quarks

The slingshot can also happen for vortices/strings in $2+1/3+1$ dimensions

The slingshot effect leads to the emission of gravitational waves
→ observable? bounds on monopoles, GUT's, etc.?

Summary

We analyzed the slingshot effect that can appear in first-order phase transitions involving magnetic monopoles/quarks

The slingshot can also happen for vortices/strings in $2+1/3+1$ dimensions

The slingshot effect leads to the emission of gravitational waves
→ observable? bounds on monopoles, GUT's, etc.?

Thank you!

More Videos on Youtube:



Maximilian Bachmaier



Backup

Numerical Simulation



Python Package Numba:



- Translates Python and NumPy code into fast machine code
- Easy implementation of parallelization

Numerical Simulation – Axial Symmetry

Axial Symmetry:

$$\phi^1 = x f_1 + y f_2$$

$$\phi^2 = y f_1 - x f_2$$

$$\phi^3 = f_3$$

$$W_x^1 = x y f_4 + y^2 f_5 + f_6$$

$$W_x^2 = -x^2 f_4 - x y f_5 + f_7$$

$$W_x^3 = x f_8 + y f_9$$

$$W_y^1 = y^2 f_4 - x y f_5 - f_7$$

$$W_y^2 = -x y f_4 + x^2 f_5 + f_6$$

$$W_y^3 = y f_8 - x f_9$$

$$W_z^1 = x f_{10} + y f_{11}$$

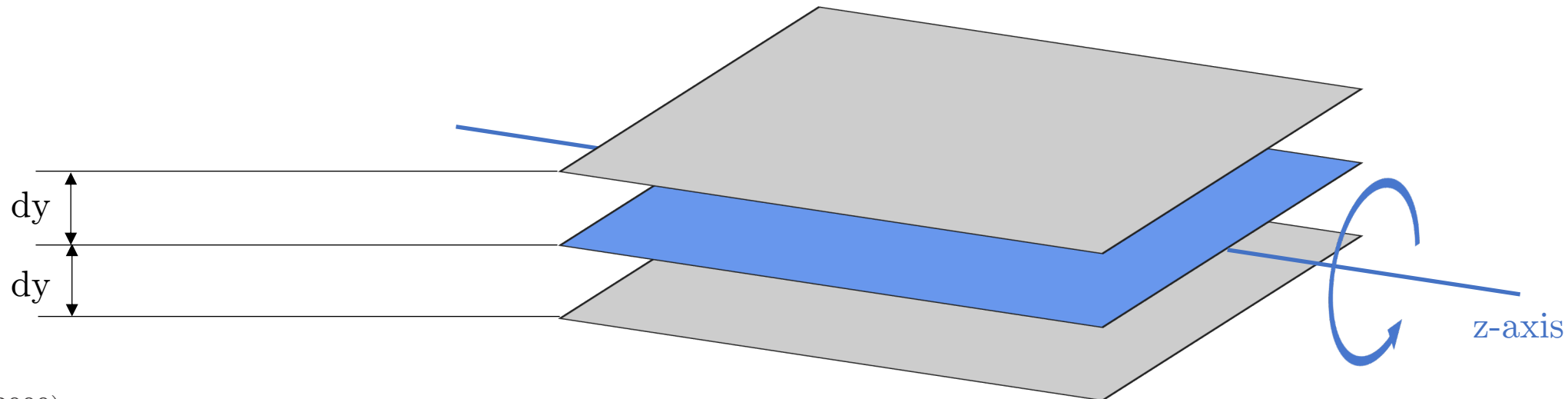
$$W_z^2 = -x f_{11} + y f_{10}$$

$$W_z^3 = 0$$

$$W_t^1 = x f_{12} + y f_{13}$$

$$W_t^2 = -x f_{13} + y f_{12}$$

$$W_t^3 = 0$$

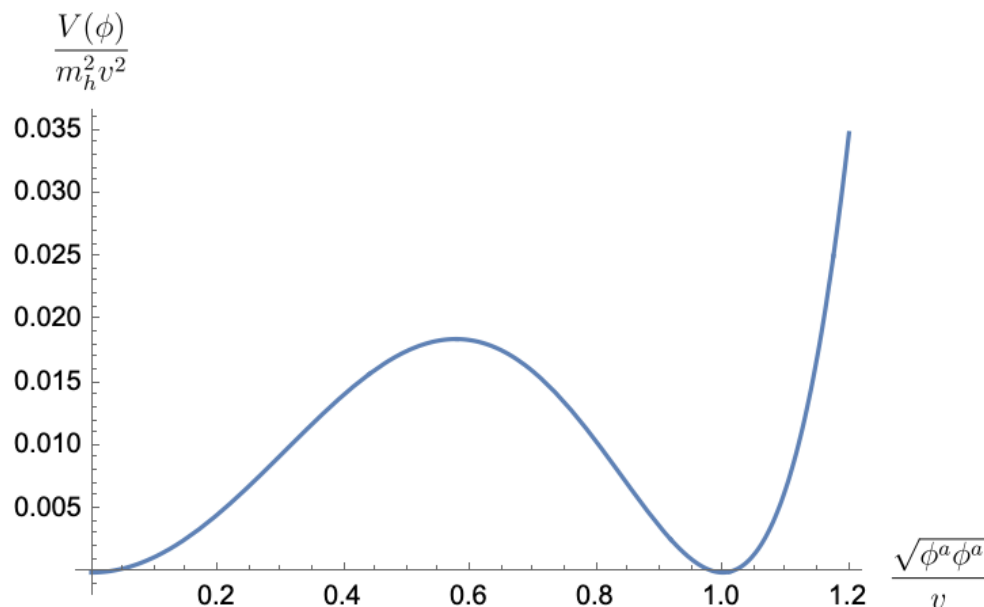


The Model

We consider an $SU(2)$ gauge theory with sextic potential

$$V(\phi) = \frac{\lambda}{8} (\phi^a \phi^a - v^2)^2 (\phi^a \phi^a)$$

↖
adjoint



$$\langle \phi^a \phi^a \rangle = \begin{cases} 0 \rightarrow SU(2) \text{ invariant vacuum} \\ v^2 \rightarrow U(1) \text{ invariant vacuum} \end{cases}$$

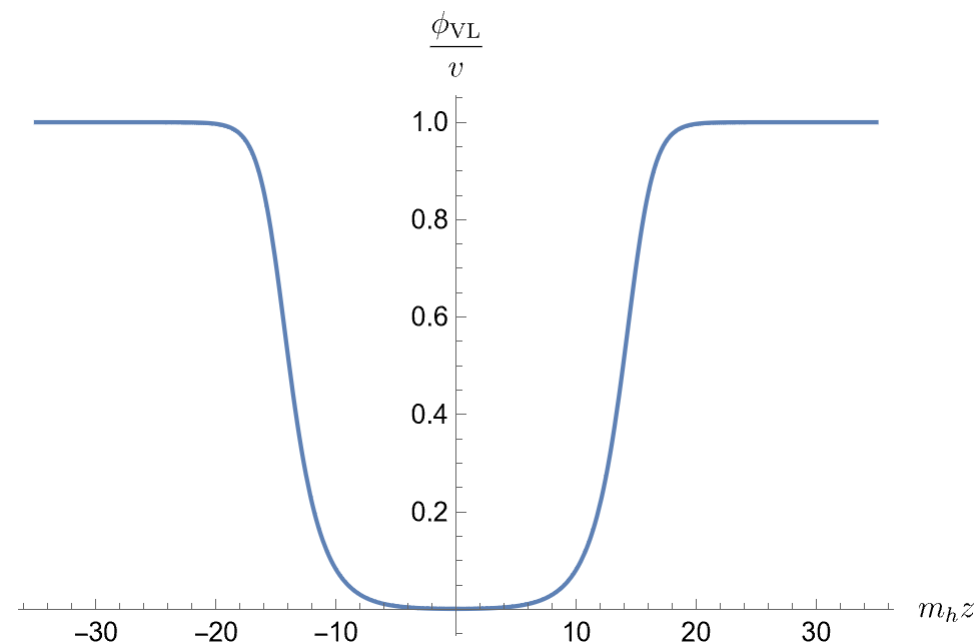
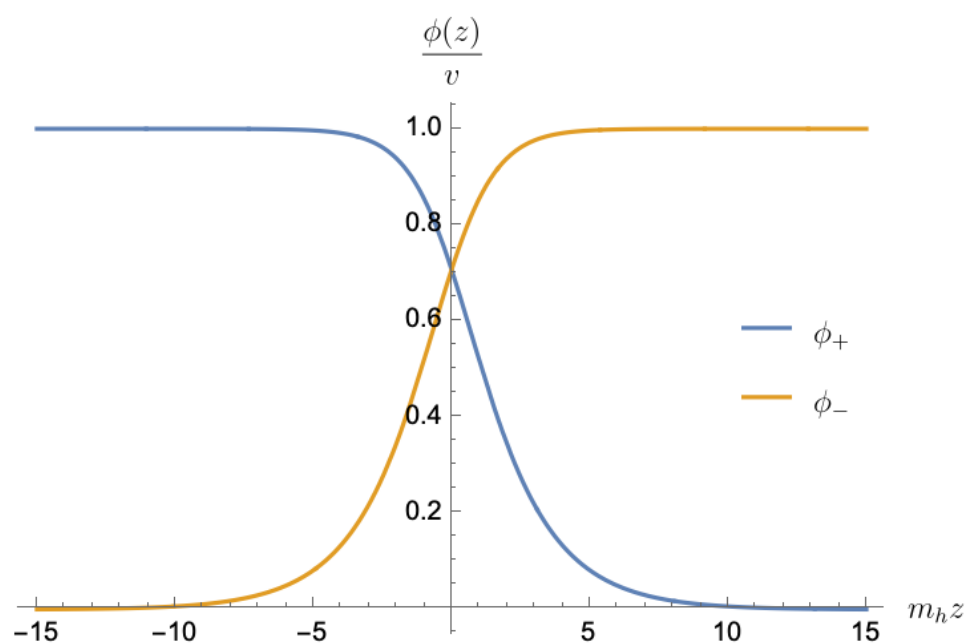
Disconnected Vacuum Manifold
→ Domain Walls

$\langle \phi^a \phi^a \rangle = v^2 \rightarrow SU(2)$ breaks down to $U(1)$
→ Magnetic Monopoles

Domain Wall Solution

Two domain wall solutions are

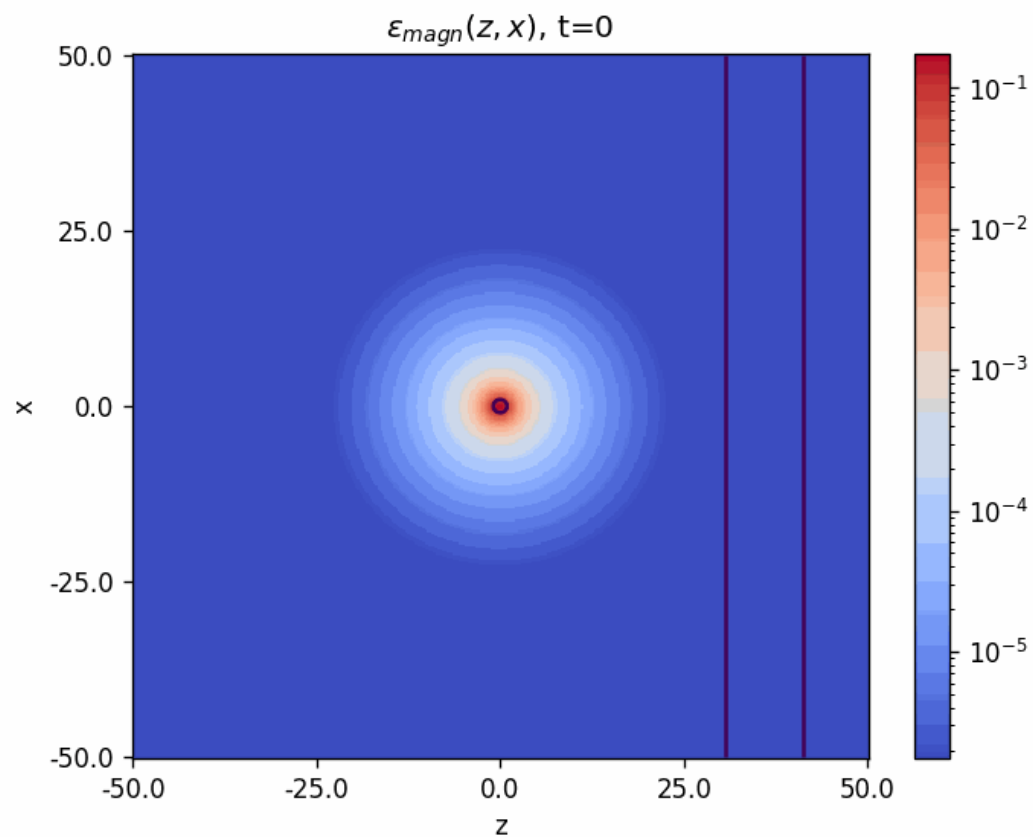
$$\phi_{\pm} = \frac{v}{\sqrt{1 + e^{\pm m_h z}}}$$



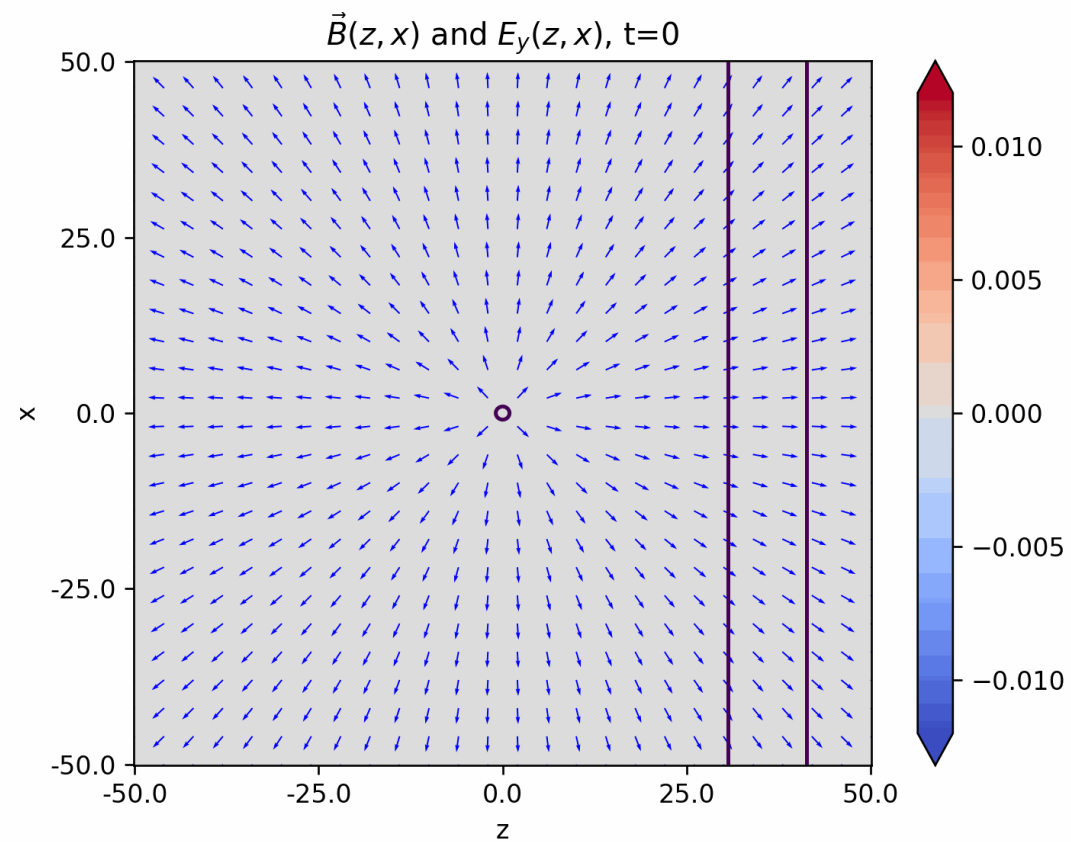
SU(2) Invariant Vacuum Layer

Erasure of the Monopole

Magnetic Energy Density

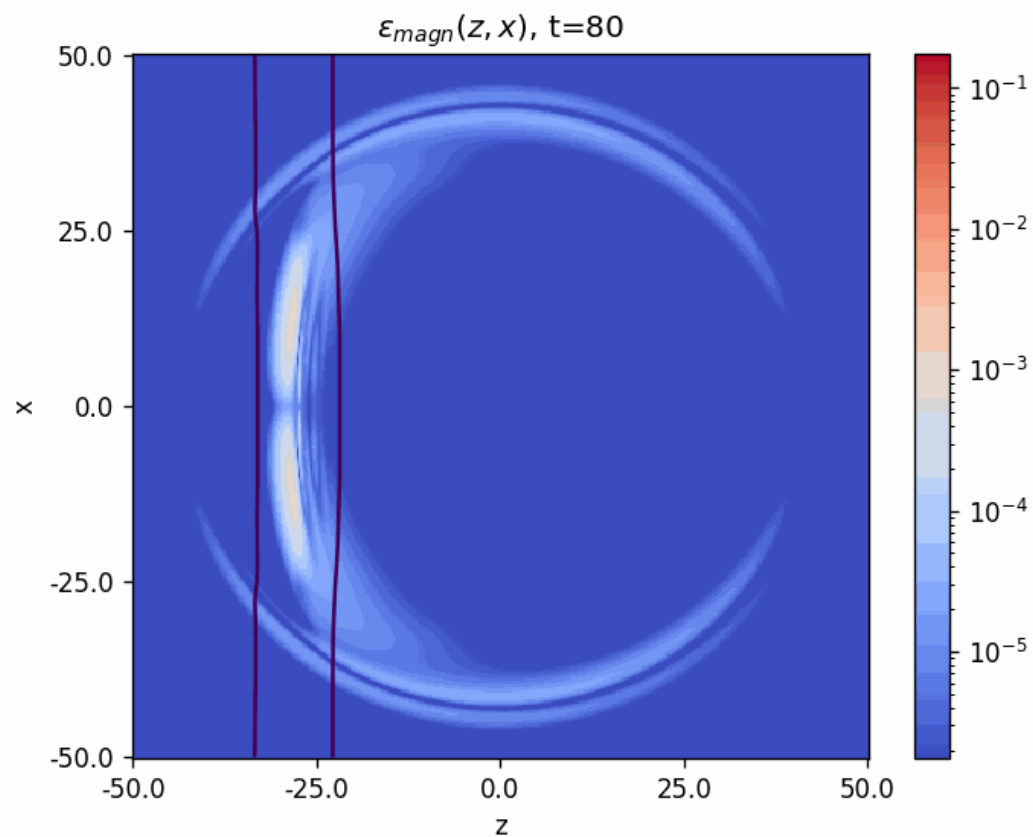


Magnetic and Electric Field

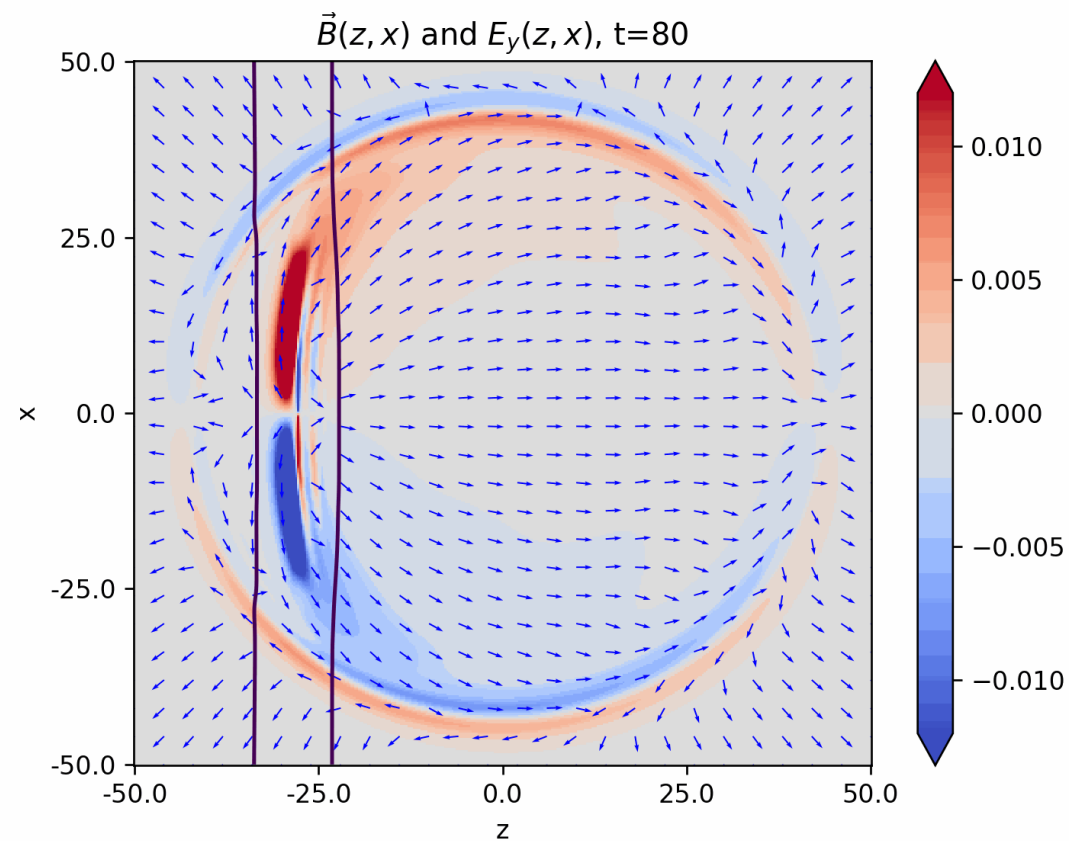


Erasure of the Monopole

Magnetic Energy Density



Magnetic and Electric Field



Initial Configuration – Domain Wall

We consider an $SU(2)$ gauge theory with the following potential

$$V(\phi) = \lambda_\phi (\phi^a \phi^a - v_\phi^2)^2 + \lambda_\psi (\psi^\dagger \psi - v_\psi^2)^2 (\psi^\dagger \psi) + \beta \psi^\dagger \phi \psi$$

ϕ : $SU(2)$ adjoint

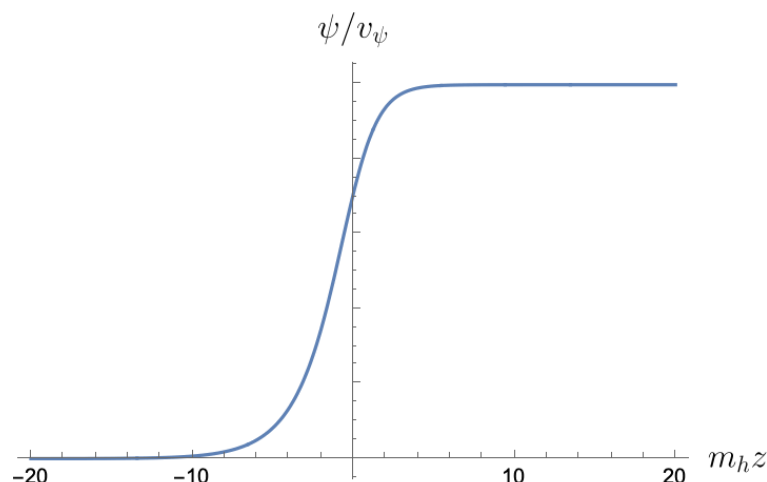
ψ : $SU(2)$ fundamental

The profile of the domain wall is

$$|\psi| = \frac{v_\psi}{\sqrt{1 + e^{-m_h z}}}$$

This term needs to be minimized

$$\rightarrow \psi^\dagger \sigma^a \psi \sim \phi^a$$



Initial Configuration – Magnetic Monopole

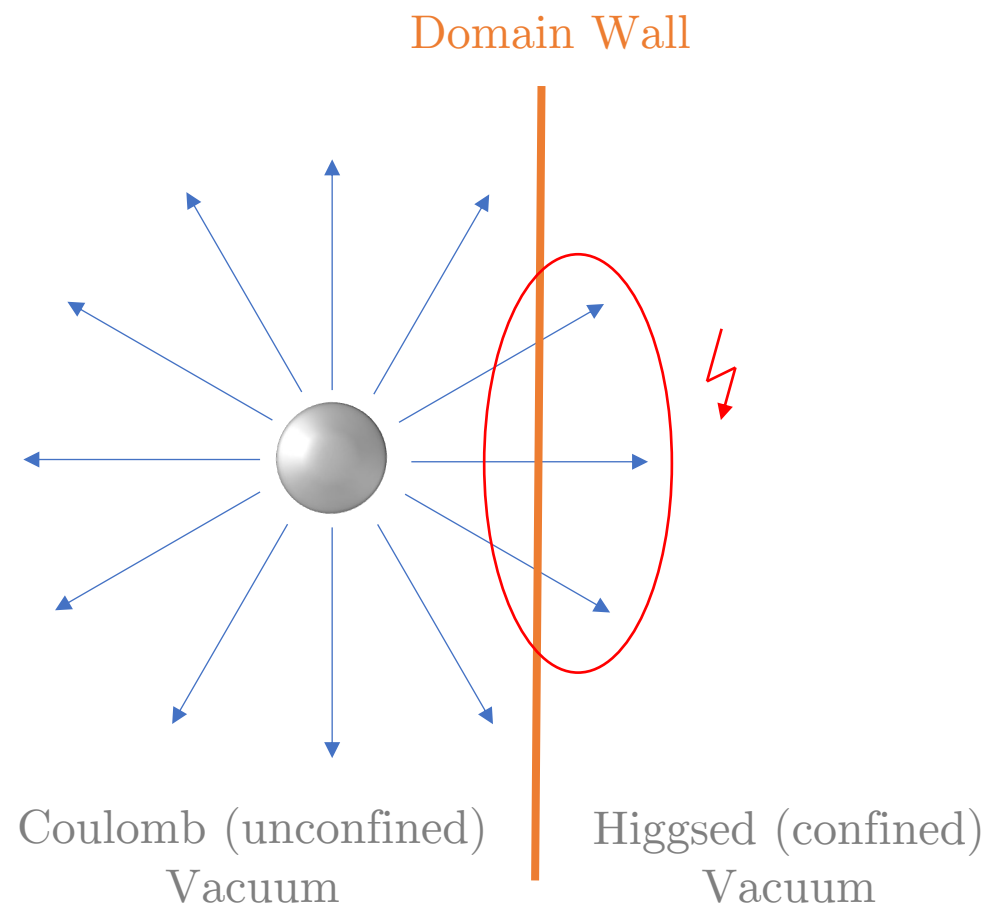
Ansatz for 't Hooft-Polyakov magnetic monopole:

$$W_i^a = \varepsilon_{aij} \frac{r_j}{r^2} \frac{1}{g} (1 - K(r))$$

$$W_t^a = 0$$

$$\phi^a = \frac{r^a}{r^2} \frac{1}{g} H(r)$$

Magnetic field needs to be repelled from the domain wall!



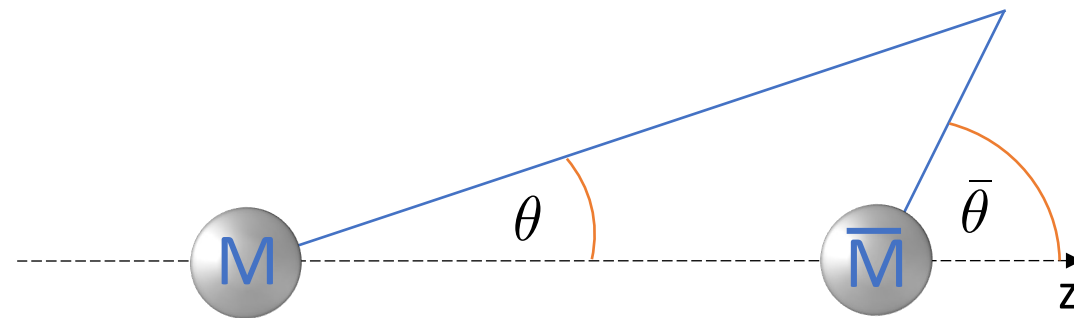
Initial Configuration – Magnetic Monopole

Ansatz for a monopole-antimonopole configuration¹:

$$\hat{\phi}_1 = (\sin \bar{\theta} \cos \theta - \sin \theta \cos \bar{\theta} \cos \alpha) \cos (\varphi - \alpha/2) \\ + \sin \theta \sin \alpha \sin (\varphi - \alpha/2)$$

$$\hat{\phi}_2 = (\sin \bar{\theta} \cos \theta - \sin \theta \cos \bar{\theta} \cos \alpha) \sin (\varphi - \alpha/2) \\ - \sin \theta \sin \alpha \cos (\varphi - \alpha/2)$$

$$\hat{\phi}_3 = -\cos \theta \cos \bar{\theta} - \sin \theta \sin \bar{\theta} \cos \alpha$$



α : twist parameter

The gauge field is given by

$$W_\mu^a \sim \varepsilon_{abc} \hat{\phi}^b \partial_\mu \hat{\phi}^c$$

¹T. Vachaspati, G. B. Field (1994) / A Saurabh, T. Vachaspati (2018)

Initial Configuration – Magnetic Monopole

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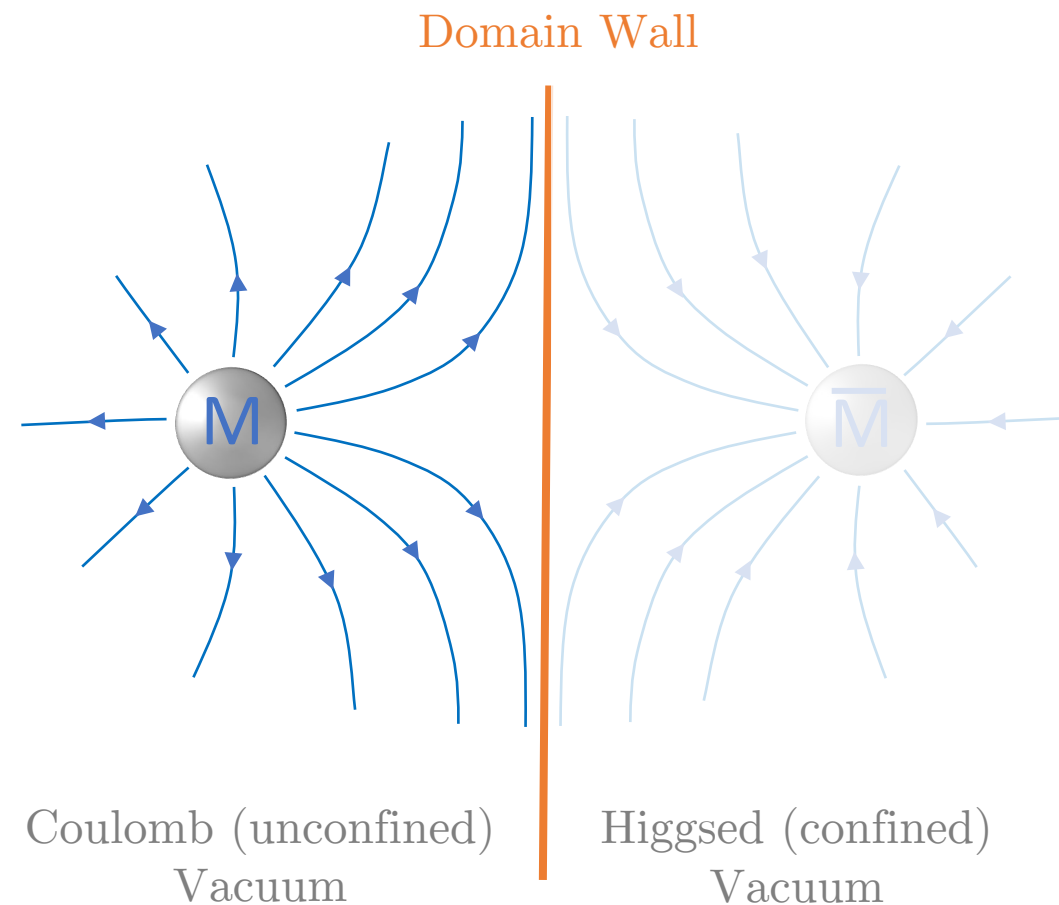
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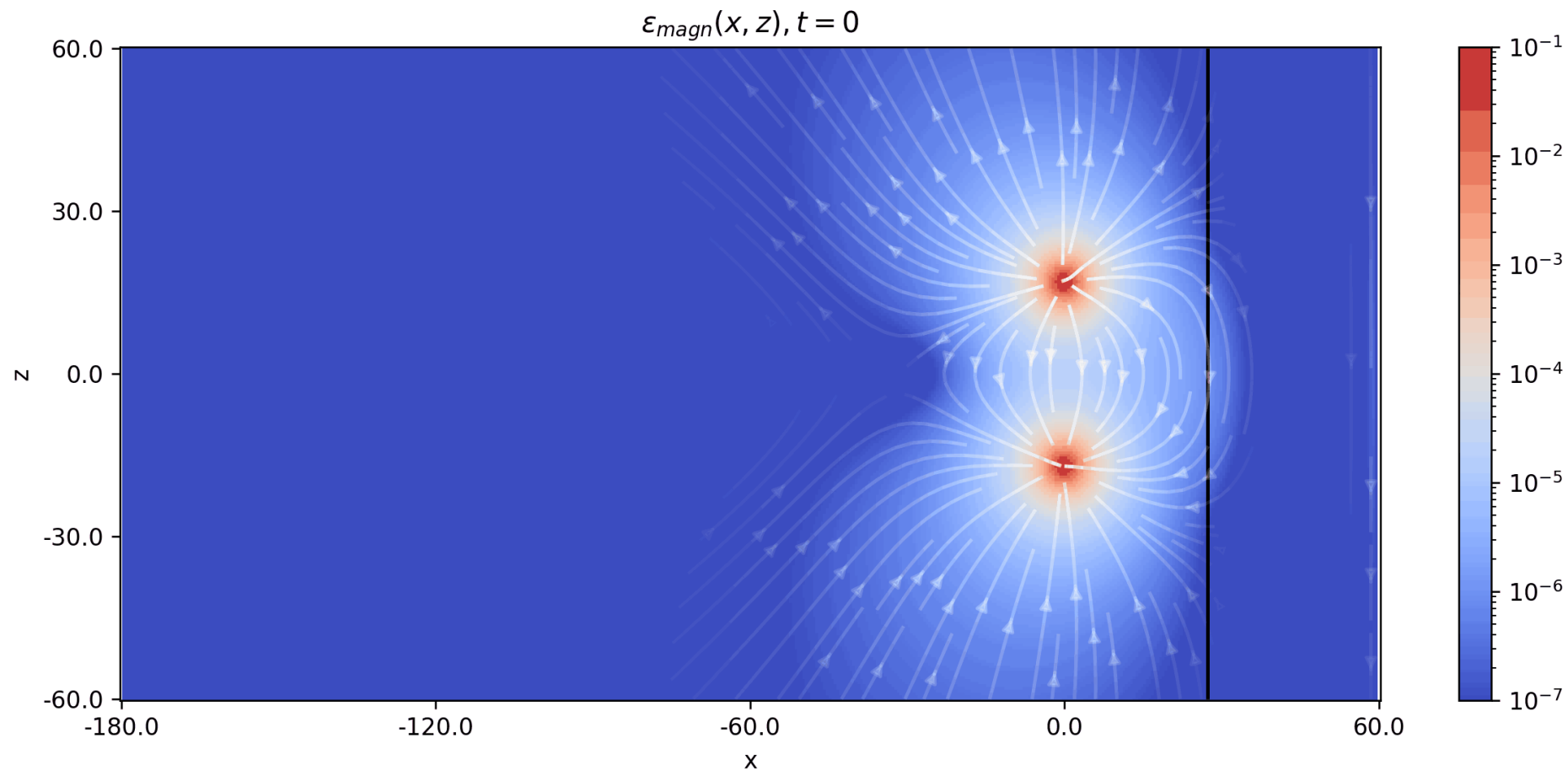
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Multiple Confinement Slingshots (maximal Twist)



M. B., G. Dvali, J. S. Valbuena-Bermúdez, M. Zantedeschi (to be published)

Confinement Slingshot – Gravitational Waves

The gravitational radiation spectrum can be calculated by Weinberg's formula:

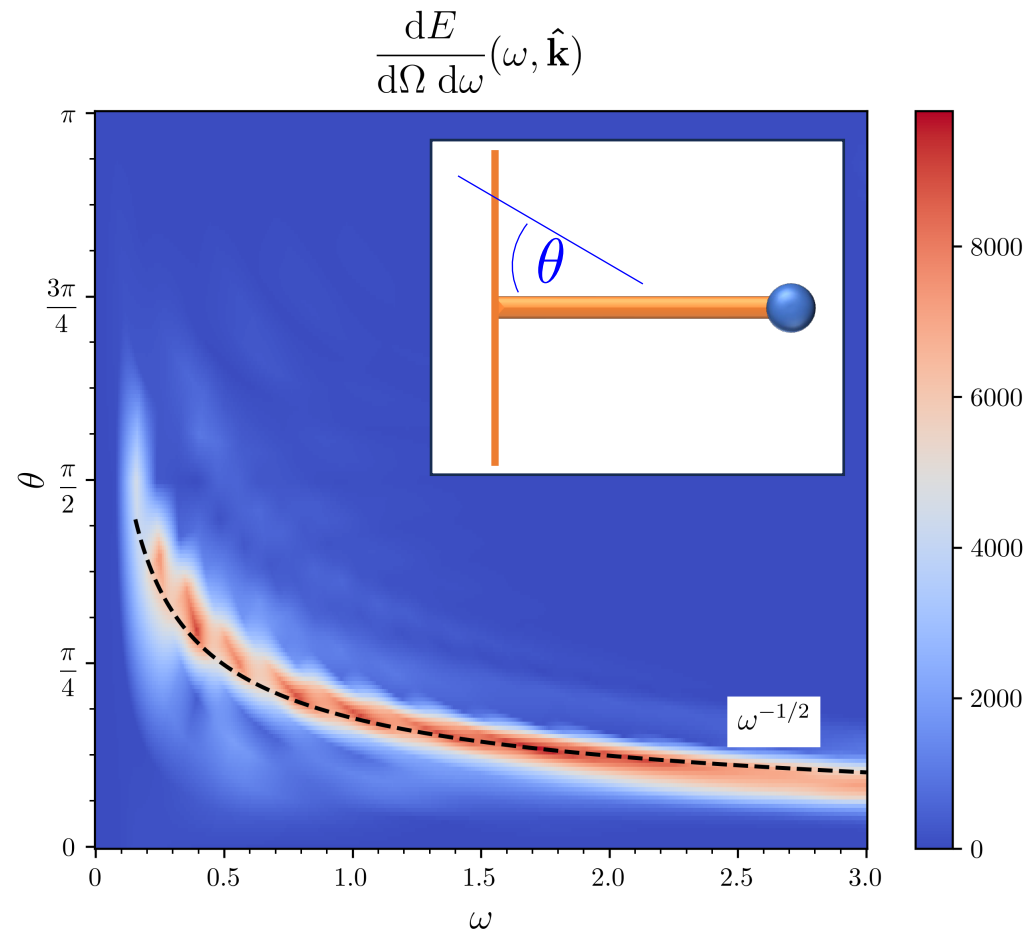
$$\frac{dE}{d\Omega d\omega} = \frac{G \omega^2}{2\pi^2} \Lambda_{ij,lm}(\hat{\mathbf{k}}) T^{ij*}(\mathbf{k}, \omega) T^{lm}(\mathbf{k}, \omega), \quad (1)$$

with

$$\Lambda_{ij,lm}(\hat{\mathbf{k}}) \equiv P_{il}(\hat{\mathbf{k}}) P_{jm}(\hat{\mathbf{k}}) - \frac{1}{2} P_{ij}(\hat{\mathbf{k}}) P_{lm}(\hat{\mathbf{k}}), \quad (2)$$

where $P_{ij}(\hat{\mathbf{k}}) = \delta_{ij} - \hat{k}_i \hat{k}_j$.

Confinement Slingshot – Gravitational Waves



The angle of emission depends on the frequency

Most of the radiation is emitted in the direction of acceleration

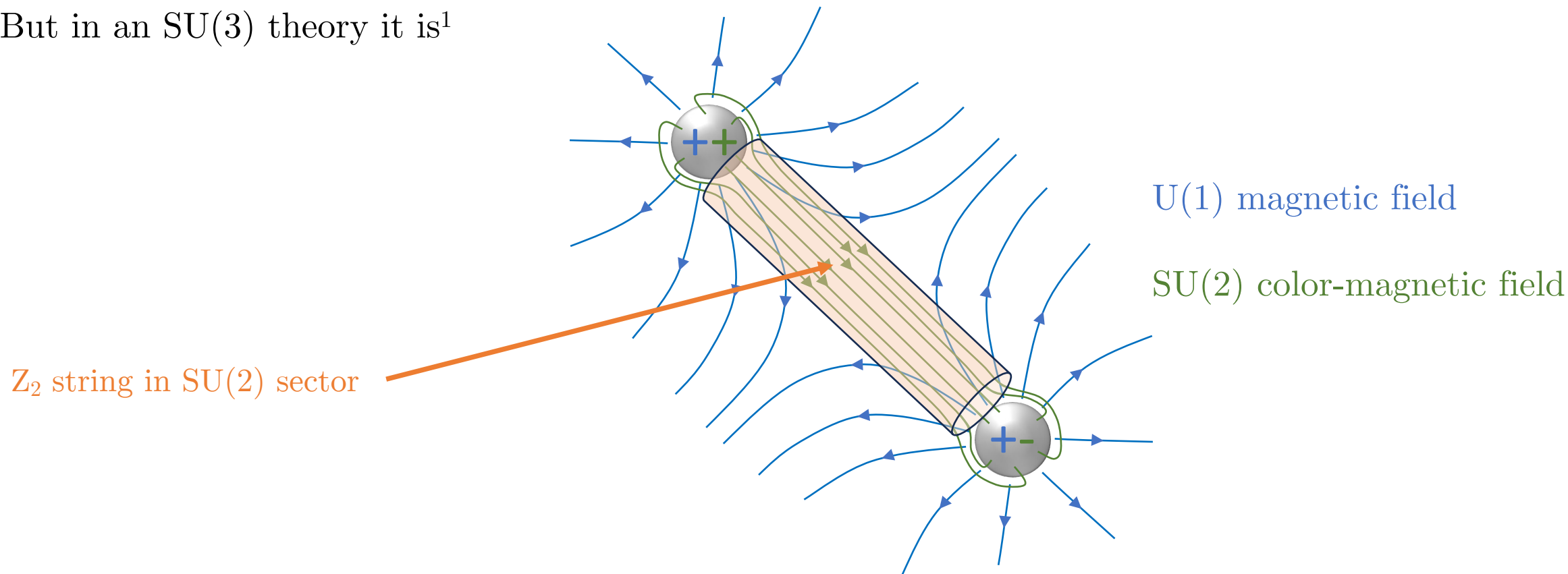
→ Again similar to two monopoles connected by a string ¹

¹ L. Leblond, B. Shlaer, X. Siemens (2009)

Monopole Confinement – Outlook

In $SU(2)$ it's not possible to connect two monopoles of same charge

But in an $SU(3)$ theory it is¹



¹ Y. Ng, T.W.B. Kibble, T. Vachaspati (2008)

Slingshot for Vortices/Strings

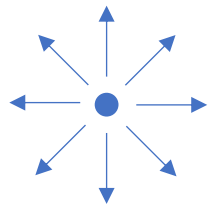
We consider an $U(1)$ gauge theory with two complex fields ϕ, χ and the following potential

$$U(\phi, \chi) = \lambda_\phi (|\phi|^2 - v_\phi^2)^2 + \lambda_\chi (|\chi|^2 - v_\chi^2)^2 |\chi|^2 + \beta \phi^* \chi^2 + c.c.$$

Breaking Pattern: $U(1) \xrightarrow{\phi} Z_2 \xrightarrow{\chi} 1$

2+1 Dimensions

Nielsen-Olesen Vortex



Z_2



$\mathbb{Z}_2$

3+1 Dimensions

String



Z_2



$\mathbb{Z}_2$