

“Hairy Black Holes by Spontaneous Symmetry Breaking in ESGB theory”

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PASCOS 2025

July 25, 2025 @ Durham, UK

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In this talk, I will consider **Einstein-Scalar-Gauss-Bonnet Theory (ESGB)**.

$$S = \int d^4x \sqrt{-g} \left[\frac{R}{2\kappa^2} - \frac{1}{2} \nabla_\alpha \varphi \nabla^\alpha \varphi + f(\varphi) \mathcal{G} \right],$$

where $\mathcal{G}$ is the GB term

$$\mathcal{G} = R_{\mu\nu\rho\sigma} R^{\mu\nu\rho\sigma} - 4R_{\mu\nu} R^{\mu\nu} + R^2,$$

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- ESGB theory evades **the no-hair theorems** which indicates the existence of hairy black holes.

Evasion of (old/novel) No-Hair theorem

Rectifying the no-hair theorem in EsGB ^{1 2 3}

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No restriction to evade the no-hair theorem for any coupling function of $f(\varphi)$

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We suggest the formation of hairy black holes by spontaneous symmetry breaking ⁴

⁴Boris Latosh, Miok Park, Phys.Rev.D 110 (2024) 2, 024012, "Hairy black holes by spontaneous symmetry breaking"

Our Lagrangian

$$S = \int d^4x \sqrt{-g} \left[\frac{1}{2\kappa^2} R - \nabla_\alpha \varphi^* \nabla^\alpha \varphi + f(\varphi) \mathcal{G} \right], \quad (1)$$

$$\mathcal{L}_\varphi = -\nabla_\alpha \varphi^* \nabla^\alpha \varphi + f(\varphi) \mathcal{G} = T - V, \quad V = -f(\varphi) \mathcal{G} \quad (2)$$

where $\mathcal{G}$ is the Gauss-Bonnet term and the scalar field coupling function is

$$f(\varphi) = \alpha \varphi^*(r) \varphi(r) - \lambda (\varphi^*(r) \varphi(r))^2, \quad (\lambda > 0)$$

Our metric ansatz is

$$ds^2 = -A(r) dt^2 + \frac{1}{B(r)} dr^2 + r^2 (d\theta^2 + \sin^2 \theta d\phi^2). \quad (3)$$

- This Lagrangian respects the global $U(1)$ symmetry

$$\varphi(r) \rightarrow e^{i\chi} \varphi(r)$$

- This action allows to have Schwarzschild black holes ($\varphi = \text{const}$).

Schwarzschild BH are always stable in ESGB?

Scalar field perturbation

$$\delta\varphi_1(t, r, \theta, \phi) = \sum_{l, m} \frac{\Phi(r) Y_{lm}(\theta, \phi)}{r} e^{-i\omega t} \quad (4)$$

$$\Phi''(r_*) - (V_{\text{eff}} - \omega^2)\Phi(r_*) = 0, \quad dr_* = \frac{1}{\sqrt{AB}} dr, \quad (5)$$

$$V_{\text{eff}}(r) = \frac{l(l+1)A}{r^2} + \frac{1}{2r} \left(A'B + AB' \right) - \frac{1}{2} f_{\varphi_1 \varphi_1} A \mathcal{G}, \quad (6)$$

where l is the angular momentum.

$$\int_{r_h}^{\infty} dr \frac{1}{\sqrt{AB}} V_{\text{eff}}(r) < 0 \quad \rightarrow \quad \alpha > \alpha_{\text{Sch.}} = \frac{5}{24} \approx 0.2083, \quad (7)$$

$$A(r) = B(r) = 1 - \frac{2M}{r}, \quad l = 0 \quad (8)$$

Schwarzschild black holes in ESGB

become unstable beyond the critical value of α ($= \alpha_{\text{Sch.}}$)

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Hairy Black Holes by SSB

$$V = -f(\varphi)\mathcal{G},$$

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- Black holes in the symmetric phase :

$$\langle \varphi \rangle = 0, \quad \varphi(r) = \frac{1}{\sqrt{2}} (\varphi_1(r) + i \varphi_2(r)) \quad (9)$$

- Black holes in the symmetry-broken phase : the stable minima are determined by

$$\langle \varphi \rangle = v e^{i\beta}, \quad v = \sqrt{\frac{\alpha}{2\lambda}} \quad (10)$$

we expand a field around the ground state v by reparameterizing it as follows

$$\varphi(r) = \left(v + \frac{\sigma(r)}{\sqrt{2}} \right) e^{i\theta(r)} \quad (11)$$

Flux near the black hole horizon ⁵

The conserved current is given by

$$\partial_\alpha J^\alpha = 0, \quad J_\alpha = i g (\varphi^* \partial_\alpha \varphi - \varphi \partial_\alpha \varphi^*).$$

- In the symmetric phase :

$$\int_{\Sigma} J_\alpha n^\alpha \sqrt{-h} d^3 y = \int_{\Sigma} \left[g(\varphi_2 \partial_r \varphi_1 - \varphi_1 \partial_r \varphi_2) \right] \left[\sqrt{A(r)B(r)} r^2 \sin \theta d\theta d\phi dt \right] = 0 \quad (12)$$

- In the symmetry-broken phase :

$$\theta'(r) = \frac{c_2}{4r^2 \sqrt{A(r)B(r)}} \left(v + \frac{\sigma(r)}{\sqrt{2}} \right)^{-2}, \quad (13)$$

$$\begin{aligned} \int_{\Sigma} J_\alpha n^\alpha \sqrt{-h} d^3 y &= \int_{\Sigma} \left[-2g \left(v + \frac{\sigma(r)}{\sqrt{2}} \right)^2 \theta'(r) \right] \left[\sqrt{A(r)B(r)} r^2 \sin \theta d\theta d\phi dt \right] \\ &= -8\pi g c_2 \end{aligned} \quad (14)$$

$c_2 = 0$ is required. The Goldstone boson is trivial.

⁵special thanks to Prof. Seong Chan Park

Hairy Black Holes by SSB in EsGB with global $U(1)$

We are interested in the situation that

“Scalar fields are about to grow from black holes without scalar hair. Finally it evolves to hairy black holes.”

ex. Schwarzschild black hole $\longrightarrow$ Hairy Black holes

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- “ $V = -f(\varphi)\mathcal{G}$ ” as an “interacting potential” :
 - effective near the black hole horizon
 - not effective at infinity ($V \rightarrow 0$ as $r \rightarrow \infty$, since $\mathcal{G} \rightarrow 0$ as $r \rightarrow \infty$)

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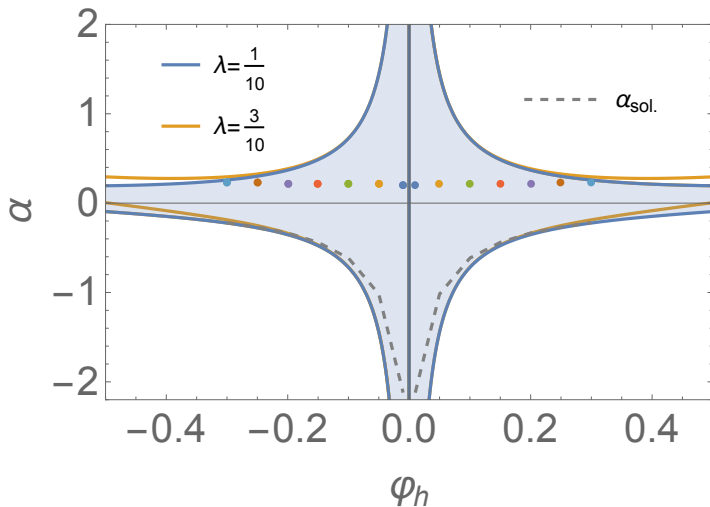
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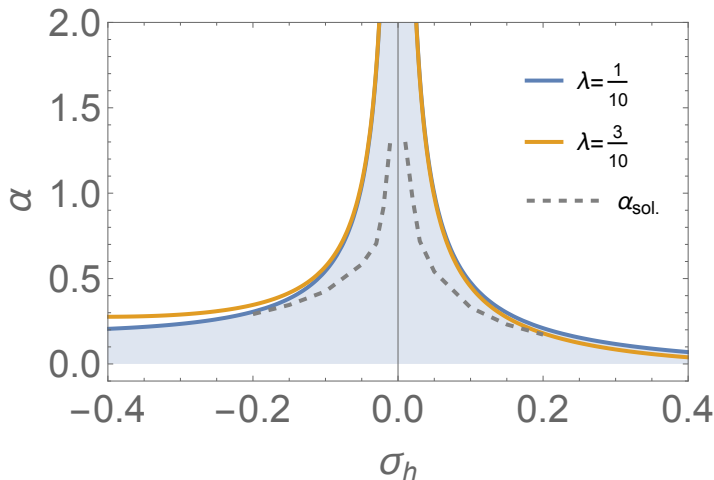
Based on these assumption,

we numerically generated the hairy black holes in symmetric and symmetry-broken phase

Hairy black holes in *symmetric phase*



Hairy black holes in *symmetry broken phase*



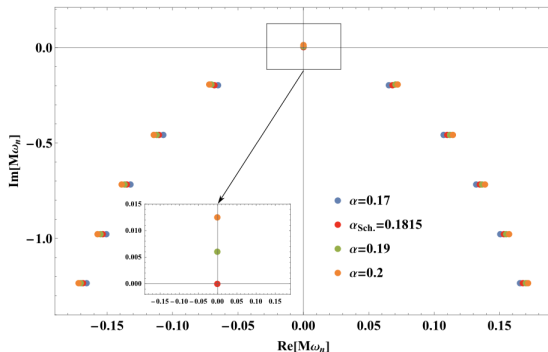
Hairy black holes + Scalar field perturbation

→ **Quasinormal modes (QNM)** ⁶

⁶Young-Hwan Hyun, Boris Latosh, Miok Park, JHEP 08 (2024) 163, "Scalar field perturbation of hairy black holes in EsGB theory"

G_{QNM} : Schwarzschild black holes

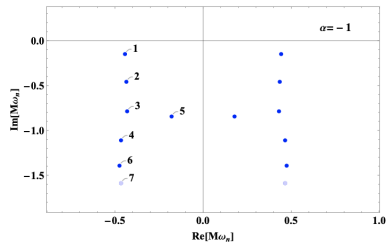
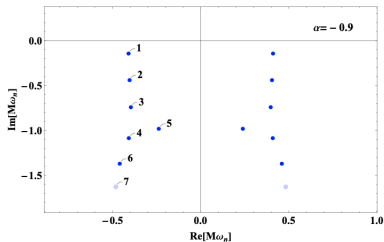
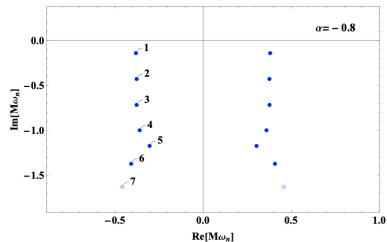
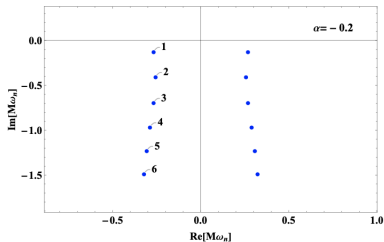
$$\delta\varphi(v, z, \theta, \phi) = \sum_{l,m} \Phi(z) Y_{lm}(\theta, \phi) e^{-i\omega v}, \quad \omega = \omega_R + i\omega_I$$



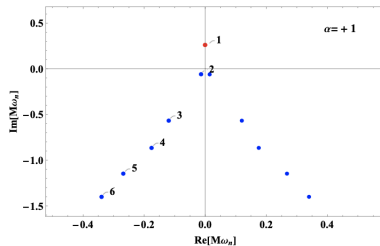
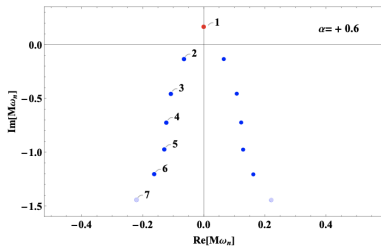
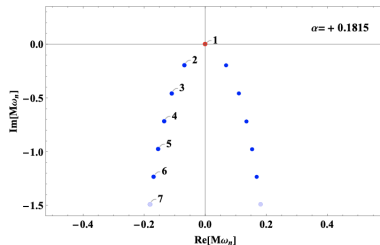
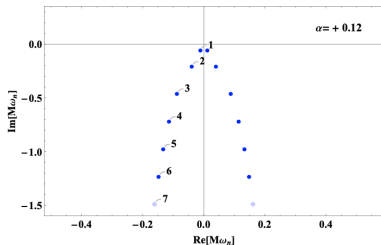
	$\text{Re}[M\omega_n]$	$\text{Im}[M\omega_n]$
$\alpha=0.17$	± 0.06528	-0.19795
	± 0.10718	-0.45845
	± 0.13196	-0.71851
	± 0.15043	0.97682
$\alpha_{Sch.} \approx 0.1815$	0	0.00002
	± 0.06826	-0.19618
	± 0.11012	-0.45776
	± 0.13483	-0.71820
	± 0.15324	-0.97669
$\alpha=0.19$	0	0.00600
	± 0.07013	-0.19489
	± 0.11208	-0.45728
	± 0.13677	-0.71798
	± 0.15513	-0.97662
$\alpha=0.20$	0	0.01254
	± 0.07198	-0.19340
	± 0.11415	-0.45675
	± 0.13881	-0.71775
	± 0.15712	-0.97654

Figure 4 & Table 1: QNMs of Schwarzschild black holes for several values α

G_{QNM} : Symmetric Phase with $\alpha < 0$



G_{QNM} : Symmetric Phase with $\alpha > 0$



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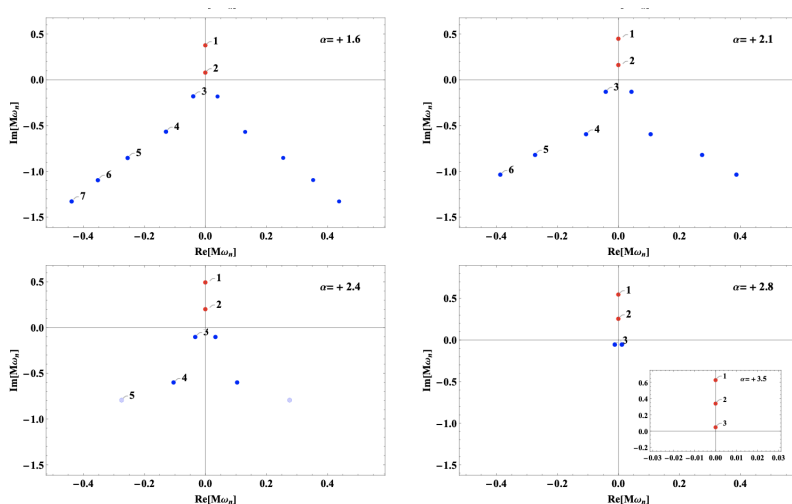
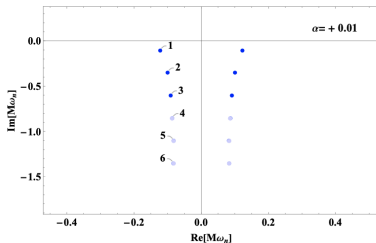
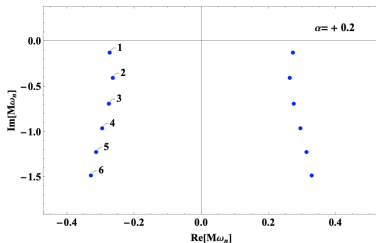


Figure 6: QNMs with positive values of α for $\varphi_h = 0.01$. Blue dots are tolerable less than 10^{-3} , whereas lighter blue dots have a tolerance of 10^{-2} .

G_{QNM} : Symmetry-Broken Phase

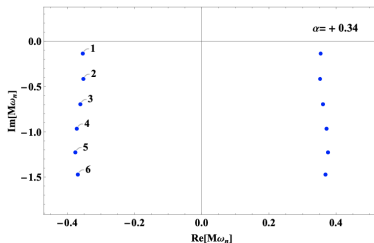


	$\text{Re}[M\omega_n]$	$\text{Im}[M\omega_n]$
1	± 0.12229	-0.10701
2	± 0.10025	-0.35077
3	± 0.09104	-0.60312
4	± 0.08662	-0.85448
5	± 0.08226	-1.10205
6	± 0.08295	-1.35246

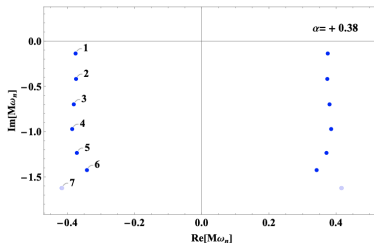


	$\text{Re}[M\omega_n]$	$\text{Im}[M\omega_n]$
1	± 0.27261	-0.13084
2	± 0.26325	-0.40969
3	± 0.27500	-0.69474
4	± 0.29492	-0.96587
5	± 0.31297	-1.22847
6	± 0.32848	-1.4865

G_{QNM} : Symmetry-Broken Phase



	$\text{Re}[M\omega_n]$	$\text{Im}[M\omega_n]$
1	± 0.35393	-0.13610
2	± 0.35217	-0.41623
3	± 0.36116	-0.69570
4	± 0.37179	-0.96582
5	± 0.37610	-1.22716
6	± 0.36866	-1.47193

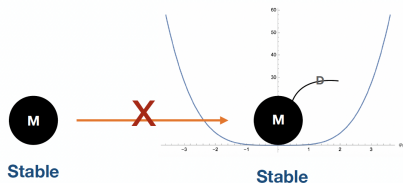


	$\text{Re}[M\omega_n]$	$\text{Im}[M\omega_n]$
1	± 0.37538	-0.13731
2	± 0.37399	-0.41842
3	± 0.38069	-0.69878
4	± 0.38549	-0.97137
5	± 0.37156	-1.23482
6	± 0.34159	-1.42465
7	± 0.41624	-1.62298

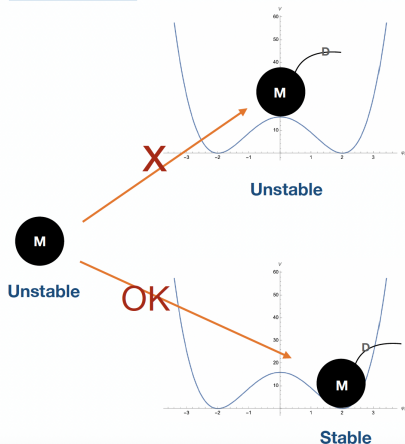
Figure 11 & Table 3: QNMs with various values of α for $\sigma_h = 0.01$. Blue dots are tolerable less than 10^{-3} , whereas lighter blue dots have a tolerance of 10^{-2} .

Summary

When $\alpha < 0$

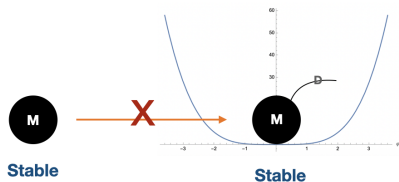


When $\alpha > 0$



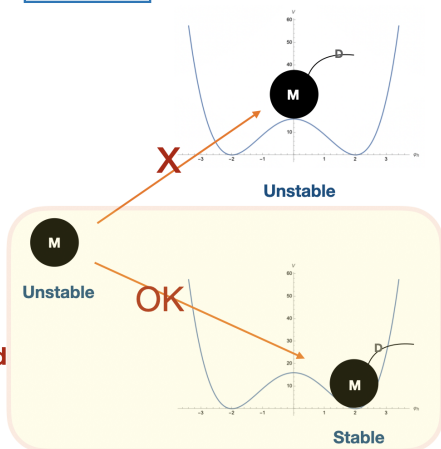
Summary

When $\alpha < 0$



Our Claim :
Phase transition will occur and
form black hole's hairs
in symmetry-broken phase!

When $\alpha > 0$



Future works

- 1 quasinormal modes with metric perturbations
- 2 dynamical evolution from non-hairy black holes to hairy black holes
- 3 hairy black holes by SSB with $U(1)$
- 4 thermodynamic stability
- 5 relation with Love number

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Thank you!