Birational Transformations and Mass Deformations of 2d SUSY gauge theories

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Brane Brick Models : brane construction

•	World-volume theory	of D1-brane whose	e transverse geomet	ry is a tor	ric Calabi-Yau 4-folo
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We call its T-dual Type IIA brane construction as a 'brane brick'



Birational Transformations



Mass-Deformation of Gauge Theory



Gauge Theory

Brane Brick

Fig. 2. Toric diagram (left), Periodic quiver (middle) and brane brick (right)

KEY : Mass Deformation of Brane Brick models corresponds to the birational transformation on Newton polynomial of probed CY 4-folds.

Mass deformations of (0,2) quiver gauge theories

Holomorphic interaction of 2d (0,2) gauge theories is given by J- and Eulletterms, each of which is associated with each Fermi multiplet.

$$\Lambda = \lambda_{-} - \theta^{+} G - i \theta^{+} \overline{\theta}^{+} D_{+} \lambda_{-} - \overline{\theta}^{+} E, \quad \overline{\mathcal{D}} \Lambda = E(\Phi_{i})$$
$$L_{J} = -\int d^{2}x d\theta^{+} \sum_{a} (\Lambda_{a} |J_{a}(\Phi_{i})|_{\overline{\theta}^{+}=0}) - h.c.$$

- J- and E-terms that are **linear** in chiral multiplet turns on the mass → A chiral-Fermi pair in the same (or opposite) gauge representation can trigger mass deformations.
 - → Massive fields are replaced by binomials of surviving chirals

$$(\Lambda_{ij}, X_{ij}) \in Q : \quad J'_{ji} = J_{ji} , E'_{ij} = E_{ij} \pm m X_{ij}$$

$$(\overline{\Lambda_{ij}}, X_{ji}) \in Q : \quad J'_{ii} = J_{ii} \pm m X_{ii} , E'_{ij} = E_{ij}$$

$$X_{ij} = \mp \frac{1}{m} E_{ij}$$

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Birational transformation as Algebraic mutation

Toric Diagram

- Newton polynomial from toric diagram (0, -1, 2)(0, -1, 1) $_{q_1, q_2}$ $P(x, y, z) = \sum_{\mathbf{v} \in \Delta} c_{\mathbf{v}} \mathbf{x}^{\mathbf{v}} = \frac{1}{y} (1 + 2z + z)$ (1, 0, 0)(0, -1, 0)Fig 3. Mapping between terms in Newton polynomial and vertices of toric diagram
- Birational transformation and transform of Newton polynomial
 - 1. SL(3, Z) rotation

$$x \mapsto x^{M_{11}} y^{M_{12}} z^{M_{13}}, y \mapsto x^{M_{21}} y^{M_{22}} z^{M_{23}}, z \mapsto x^{M_{31}} y^{M_{32}} z^{M_{33}}$$
$$M = (M_{ij}) \in \mathrm{SL}(3, \mathbf{Z})$$

- 2. rational transformation
 - $x \mapsto x, \quad y \mapsto y, \quad z \mapsto A(x,y)z$

$\Lambda_{ji} = + \frac{1}{m} J_{ji}$

Example : Mass deformation of Brane Brick Models

For an example of mass deformation of brane brick models, we consider C_{++} theory of which quiver is constructed by 3*d* printing of 4*d* N=1 conifold theory. RG flow is triggered by the deformation colored in blue.



 $X_{21} \cdot Y_{12} \cdot Y_{21} - Y_{21} \cdot Y_{12} \cdot X_{21} - Y_{12} \cdot X_{21} - Y_{12} + Z_{13} \cdot X_{32} - X_{14} \cdot Z_{42}$ Λ_{12} : $Y_{34} + Z_{31} \cdot X_{14} - X_{32} \cdot Z_{24}$ $X_{43} \cdot Y_{34} \cdot Y_{43} - Y_{43} \cdot Y_{34} \cdot X_{43}$ Λ_{34} : $Z_{13} \cdot Y_{34} - Y_{12} \cdot Z_{24}$ $Y_{43} \cdot X_{32} \cdot X_{21} - X_{43} \cdot X_{32} \cdot Y_{21}$ Λ_{14} : $Y_{21} \cdot X_{14} \cdot X_{43} - X_{21} \cdot X_{14} \cdot Y_{43}$ $Z_{31} \cdot Y_{12} - Y_{34} \cdot Z_{42}$ Λ_{32} : $\Lambda^{1}_{23}:$ $Z_{24} \cdot X_{43} - X_{21} \cdot Z_{13}$ $Y_{34} \cdot Y_{43} \cdot X_{32} - X_{32} \cdot Y_{21} \cdot Y_{12}$ $\Lambda^{2}_{23}:$ $X_{32} \cdot X_{21} \cdot Y_{12} - Y_{34} \cdot X_{43} \cdot X_{32}$ $Z_{24} \cdot Y_{43} - Y_{21} \cdot Z_{13}$ Λ^1_{41} : $Y_{12} \cdot Y_{21} \cdot X_{14} - X_{14} \cdot Y_{43} \cdot Y_{34}$ $Z_{42} \cdot X_{21} - X_{43} \cdot Z_{31}$ $\Lambda^{2}_{41}:$ $X_{14} \cdot X_{43} \cdot Y_{34} - Y_{12} \cdot X_{21} \cdot X_{14}$ $Z_{42} \cdot Y_{21} - Y_{43} \cdot Z_{31}$

The integration-out of massive fields induce the following substitution of chiral fields in the J- and E-terms.

> $Y_{12} = Z_{13} \cdot X_{32} - X_{14} \cdot Z_{42}$ $Y_{34} = -Z_{31} \cdot X_{14} + X_{32} \cdot Z_{24}$

The resulting quiver gauge theory is given by H_4 theory. The toric diagrams of two gauge theories are described below. You could check \bullet

3. SL (3,Z) rotation $x \mapsto x^{N_{11}} y^{N_{12}} z^{N_{13}}, y \mapsto x^{N_{21}} y^{N_{22}} z^{N_{23}}, z \mapsto x^{N_{31}} y^{N_{32}} z^{N_{33}}$ $N = (N_{ij}) \in \mathrm{SL}(3, \mathbf{Z})$

• Birational transform of Newton polynomial for
$$C_{++}$$
 theory
- step 1.
 $M = \begin{pmatrix} 0 & 1 & 1 \\ 0 & 0 & 1 \\ 1 & 0 & 0 \end{pmatrix}$ $P(x, y, z) = (1 + 2x + x^2) \frac{1}{z} + (1 + y) + yz$
- step 2. $A(x, y) = (1 + x)$ $P(x, y, z) = (1 + x) \frac{1}{z} + (1 + y) + (y + xy)z$
- step 3. $N = M^{-1}$ $P^{\vee}(x, y, z) = \frac{1}{y}(1 + z) + 1 + \frac{x}{y} + (1 + z)x$

Birational transformation as Combinatorial mutation • Height assignment and decomposition of toric diagram $H_{+1} = G_{+1} + F$ w = (1, 1, 0) $= \{(1,0,0)\} + \{(0,0,0), (0,0,1)\}$ $F = \{(0, 0, 0), (0, 0, 1)\}$ $H_0 = G_0 = \{(0, 0, 0), (-1, -1, 0)\}$ $H_{-1} = G - 1 = \{(0, -1, 0), (0, -1, 1)\}$ $G_{+1} = \{(1,0,0)\}$ q_1, q_2 $G_0 = \{(0, 0, 0), (-1, -1, 0)\}$ $G_{-1} + F = \{(0, -1, 0), (0, -1, 1)\}$ h=+1 $+ \{(0,0,0), (0,0,1)\}$ h=+1

- the vertex in move and where it moves •
- the massive vertex which contains chiral fields to become massive \bullet



h = -1 $p_1 h=0$ h = -1

Fig 4. Mutation : moving a factor *F* in underground layers h < 0 to overground layers h > 0

Reference

S. Franco, D. Ghim, G. P. Goulas, R. -K. Seong, Mass Deformations of Brane Brick Models, JHEP 09 (2023) 176 [arXiv:2307.03220].

[2] D. Ghim, M. Kho and R.-K. Seong, *Combinatorial and algebraic mutations of toric Fano 3-folds and* mass deformations of 2d (0,2) quiver gauge theories, Phys. Rev. D 110 (2024) 8, 08600 [arXiv:2407.19924].

[3] D. Ghim, M. Kho and R.-K. Seong, Birational Transformations and 2d (0,2) Quiver Gauge Theories beyond Toric Fano 3-folds, JHEP 06 (2025) 032 [arXiv:2502.08741].

[4] S. Franco, D. Ghim, S. Lee, R.-K. Seong and D. Yokoyama, 2d (0,2) Quiver Gauge Theories and D-Branes, JHEP 09 (2015) 072, [arXiv:1506.03818].



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