## Using $W \rightarrow \ell \nu_{\ell}$ Processes to Probe Sterile Neutrinos at Colliders New- $\nu$ Physics: From Colliders to Cosmology 10th April 2025

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#### Overview What I'll cover:

- The model
- Effect of  $W \rightarrow eN$  on W-boson mass
- Angular observables' sensitivity to sterile neutrinos

## Formalism

- We work in vSMEFT.
  - Operators contributing to  $W \rightarrow Ne$ :  $\Delta \mathcal{L} = \frac{c_{LNH}^{i}}{\Lambda^2} \bar{L}_i \nu_R \tilde{H} H^{\dagger} H + \frac{c_{HNe}^{i}}{\Lambda^2} \bar{\nu}_I$
  - kinematics of the W decays

 $\Gamma(W -$ 

#### arXiv:1909.04665

$$\bar{\nu}_R \gamma^\mu e_{iR} H^\dagger i D_\mu H + \frac{c_{NW}^i}{\Lambda^2} \bar{L}_i \sigma^{\mu\nu} \nu_R \sigma_I \tilde{W}^I_{\mu\nu},$$

•  $c_{LNH}^{i}$  and  $c_{NW}^{i}$  give rise to parity conserving interactions, while  $c_{HNe}^{i}$  gives rise to a parity-violating V+A interaction - only this operator changes the

• Thus we assume only  $c_{HNe}^{\iota} \neq 0$  which gives rise to a W decay of width:

$$\stackrel{}{} \rightarrow e\nu_R) = \frac{m_w^3 v^2}{48\pi \Lambda^4} (c_{HNe}^e)^2$$



### Motivation **W** Mass Discrepancy

- Most precise W mass measurement to date much higher than other measurements
- LHC measurements closer to EW SM prediction



#### CMS

Electroweak fit PRD 110 (2024) 030001 LEP combination Phys. Rep. 532 (2013) 119 PRL 108 (2012) 151804 Science 376 (2022) 6589 LHCb JHEP 01 (2022) 036 ATLAS arXiv:2403.15085 CMS This work



2412.13872



## W Mass Measurement at Hadron Colliders

- W-mass measured in  $W \to \ell \nu_\ell$  processes at colliders
- Observables used are  $p_T^{\ell}$ ,  $p_T^{miss}$  and  $m_T$ , where  $p_T^{miss}$  is an approximation of the neutrino momentum
  - Defined by  $\vec{p}_T^{\text{miss}} = -(\vec{u}_T + \vec{p}_T^{\ell})$  where the recoil  $\vec{u}_T$  is the vector sum of the transverse energy of all other deposits in calorimeters
- Distributions of  $p_T^{\ell}$ ,  $p_T^{\text{miss}}$  have Jacobian endpoints at m/2 while  $m_T$  drops off at m
  - *m* is invariant mass of charged lepton-neutrino system.

## W Mass Measurement at Hadron Colliders

• Invariant mass m related to  $m_W$  via the **Breit-Wigner resonance** 

$$\frac{d\sigma}{dm} \propto \frac{m^2}{(m^2 - m_W^2)^2 + m^4 \Gamma_W^2}$$

 Simulate expected final state distributions for many  $m_{W}$  and then fit to experimental data using  $\chi^2$  test





- $x_u > x_d \Rightarrow W$  boosted in direction of u or  $\overline{u}$
- $\Rightarrow$  charged lepton more boosted in decays to RHN than SM decays

 $\Rightarrow \nu_R$  decays must be more transverse to enter detector  $m_W$ 

- $\Rightarrow$  W transverse decays lead to a harder  $p_T$  spectrum, leading to a higher



- Valence quarks higher  $p_7 \Rightarrow$  W boosted in direction of valence quark
- $\Rightarrow$  in RHN decays, charged lepton

 $\Rightarrow$ in  $\nu_R$  decays,  $\nu_R$  decays more transverse to enter detector for  $W^+$ less transverse to enter detector for  $W^ \Rightarrow$  W mass artificially raised in  $W^+$  decays, artificially lowered in  $W^-$  decays

more boosted in  $W^+$  case less boosted in  $W^-$  case

## Motivation





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## **Event Generation**

• Given our assumptions, neutrino decays with width

$$\Gamma_{\nu_R} = \frac{m_N^5}{(2\pi)^3} \frac{G_F^2}{192} \frac{v^4 \sum_i (c_{HN}^i)}{\Lambda^4}$$

- Standard W selection requires  $\nu_R$  to decay outside detector
  - of typical detector ( $\gamma c \tau \approx 200 \text{m}$ )
- We find existing W-boson measurements are sensitive to  $c_{HNe}/\Lambda^2 \approx 0.1/v^2$  and masses  $m_N \lesssim 1 \text{ GeV}$ .

 $_{Ve})^2$ 

• With  $\Lambda$  at the lower bound of v = 246 GeV and  $c_{HNe}^{\iota}$  of  $\mathcal{O}(1)$  we find  $m_N \lesssim 100 \,\mathrm{MeV}$  for almost all events to decay outside fiducial volume

#### **Results** CDF Mass Fit

- Simulated at CDF, 1.96TeV
- Likelihood fit for  $m_W$  gives  $|c_{HNe}| = 1.2 \pm 0.2$  to account for CDF theory-experiment discrepancy of  $m_{W,\text{CDF}} - m_{W,\text{SM}} = 76 \text{ MeV},$  $(\Lambda = v)$



#### **Results** ATLAS Mass Fits

 Simulated at ATLAS, 7TeV: compared MC to ATLAS measurement

• 
$$m_{W^+} - m_{W^-} = -29.2 \pm 28.0 \,\mathrm{MeV}$$

• Fit for  $p_T^{\ell}$ ,  $p_T^{\text{miss}}$  and  $m_T$  in ranges used in ATLAS measurements

#### Eur.Phys.J.C 78 (2018) 2, 110



<b>Results</b> ATLAS Mass Fits			Joir 2.00	
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•	Best constraint given by $p_T^\ell$		1.50 -	
	fit to $m_{W^+} - m_{W^-}$ , which		1.25 -	
	corresponds to $ c_{HNe} v^2/\Lambda^2 < 0.13$ at	- $\Delta \ln \mathcal{L}$	1.00 -	
	68% CL		0.75 -	
•	Additional constraints from cross-sections are small.		0.50 -	
			0.25 -	



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#### **Results** CMS Mass Fits

- Compare to 2024 CMS measurement
- Equivalent value at CMS is  $m_{W^+} m_{W^-} = 57 \pm 30.3 \text{ MeV}$ 
  - Corresponds to  $|c_{HNe}|v^2/\Lambda^2 = 0.28 \pm 0.09$

- Combining ATLAS and CMS gives  $|c_{HNe}| \in [0, 0.27], 95\% \text{ CL}$  assuming  $\Lambda = v$ 

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## Angular Observables (CDF)

- An ideal observable for detecting the difference between the SM and BSM decays would be  $\cos \theta_{CS}$
- $\theta_{CS}$  is the final angle of one of the leptons in the Collins-Soper frame - a proxy for the W rest frame.
- However, our lack of knowledge of  $p_Z$ of the W means we may only reconstruct  $\cos \theta_{CS}$  up to a sign.





## Angular Observables (CDF)

- After reconstruction the stark difference between these distributions goes away, as we can only plot  $|\cos\theta_{CS}|$ .
- $e^+e^-$  measurements would allow full reconstruction.





# **Angular Observables (Atlas)**



- An alternative reference frame is the helicity frame, a different proxy for the W rest frame corresponding to a rotation of axes, in which we define  $\cos \theta_H$  in a similar manner to the CS frame.
- Although this is a different frame, the principles which motivated the use of the CS frame still apply.

A Sparrow (2012), PhD thesis; arXiv:1203.2165



# **Angular Observables (Atlas)**

$$L_p = \frac{\vec{p}_T^{\ell} \cdot \vec{p}_T^W}{|\vec{p}_T^W|^2}$$

- quantities.
  - Can be shown

high  $p_{_{T}}^{W}$ 

#### A Sparrow (2012), PhD thesis

• There is a well documented variable,  $L_p$ , such that  $2L_p - 1$  is highly correlated to  $\cos \theta_H$  at high  $p_T^W$ .

• This allows us to observe the decay angle in the W rest frame at high  $p_T^W$  using only transverse, lab frame

in that 
$$\cos \theta_H = \frac{2 |\vec{p}_\ell| - E_W}{|\vec{p}_W|} \approx 2L_p - 1$$
 a

at

## Angular Observables (Atlas) W+, SM



### **Angular Observables (Atlas)** W+, SM



## **Angular Observables (Atlas)** Sensitivity

- Despite the looser correlation, the highest sensitivity is achieved for  $p_T^W > 50 \,\mathrm{GeV}$ , which gives sensitivities of (for  $\Lambda = v$ )
  - $c_{HNe} = 0.090 (0.072)$  for  $W^-(W^+)$ at 95% CL
- Sensitivities are calculated by statistical uncertainty given by fitting SM pseudodata to SM+BSM simulations for different values of µ at integrated luminosity of 4.7 fb<sup>-1</sup>.



## Summary

- New constraints on sterile neutrinos from  $m_W$  measurements
  - Sensitivity to parity violating V+A interactions caused by C<sub>HNP</sub> Wilson coefficient
  - Cases where standard W-boson selection may be employed
    - Masses  $\sim MeV$  up to  $\approx 1 \, GeV$
  - Could obtain better sensitivity by directly fitting to W-mass data
- Angular variables like  $\cos \theta_H$  can hold significant sensitivity



#### **Formalism** vSMEFT



$$\overline{I}_{i}\tilde{H}N + rac{1}{2}m_{N}\overline{N^{c}}N + ext{h.c.}
ight],$$

## **Mass Measurement Fits**

as in  $m_W$  measurements

$$\frac{\left(m^{2}-m_{W}^{2}\right)^{2}+m^{4}\Gamma_{W}^{2}/m_{W}^{2}}{\left(m^{2}-m_{W}^{ref^{2}}\right)^{2}+m^{4}\Gamma_{W}^{2}/m_{W}^{ref^{2}}}$$



#### • We performed a likelihood fit, reweighting $m_T$ , $p_T^{\ell}$ and $p_T^{\text{miss}}$ by a Breit-Wigner

• Where  $m_W$  is a candidate mass, and  $m_W^{ref} = 79.824 \,\text{MeV}$  is the SM value.

### Results **CDF PDF Reweighting**

- Lepton Asymmetry affected by introduction of right-handed neutrino, which is used to set PDFs at CDF.
- Estimated impact of this by reweighting SM signal in order for A to match that of BSM model with signal strength  $\mu$ .

 $\mathcal{A}(\eta_{\ell}) \equiv \frac{dN^+/d\eta_{\ell} - dN^-/d\eta_{\ell}}{dN^+/d\eta_{\ell} + dN^-/d\eta_{\ell}}$ 





### **Results** CDF PDF Reweighting

• We find after reweighting, effect of sterile neutrinos on W mass reduced by  $\sim 30\%$ .



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