

BSM MULTI-HIGGS: IMPLICATIONS IN EWPT AND COLLIDERS

Wrishik Naskar

(based on L. Biermann, C. Borschensky, C. Englert, M. Mühlleitner, WN 2408.08043)

18th Dec. 2024 – Young Theorists' Forum 2024

IPPP, Durham

INTRODUCTION

MULTI-HIGGS PROCESSES

 hh has been a cornerstone for physics Beyond the Standard Model (BSM) at the LHC.

(ATLAS 2023; CMS 2023)



MULTI-HIGGS PROCESSES

 hh has been a cornerstone for physics Beyond the Standard Model (BSM) at the LHC.

(ATLAS 2023; CMS 2023)



MULTI-HIGGS PROCESSES

 hh has been a cornerstone for physics Beyond the Standard Model (BSM) at the LHC.

(ATLAS 2023; CMS 2023)



. hhh is gradually gaining more focus.

(Papaefstathiou et al. 2016; Fuks et al. 2016; Chen et al. 2016; Robens et al. 2020; Papaefstathiou et al. 2019; Papaefstathiou et al. 2021; Florian et al. 2020; Papaefstathiou and Tetlalmatzi-Xolocotzi 2023; Stylianou et al. 2023; Delgado et al. 2023; Anisha et al. 2024; Brigljevic et al. 2024; Karkout et al. 2024; ATLAS 2024)

 \Rightarrow Crucial for future collider search strategies.

. SM cross-section for hhh (LO)

 $\left| \sigma_{hhh}^{ggF} = \mathcal{O}(50 \text{ ab}) \right| \sim 4 \text{ (10) events} \text{ (HL-LHC)}$













. SMEFT vs HEFT.

(Delgado et al. 2023; Anisha et al. 2024)





. SMEFT is a linear extension around the SM point.

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}}^{d \leq 4} + \sum_{i,d=5,6,\dots} \frac{C_i^{(d)} \mathcal{O}_i^{(d)}}{\Lambda^{d-4}}.$$

(Henning et al. 2016; Ilaria Brivio et al. 2019) ...



. SMEFT is a linear extension around the SM point.

$$\mathcal{L}^{Higgs,d\geq 6}_{SMEFT}\supset \frac{1}{\Lambda^2}C_6|\Phi|^6+\frac{1}{\Lambda^4}C_8|\Phi|^8, \quad \Phi=\begin{pmatrix}0\\\frac{\nu+h}{\sqrt{2}}\end{pmatrix}$$



. SMEFT is a linear extension around the SM point.

$$\mathcal{L}_{\mathsf{SMEFT}}^{\mathsf{Higgs},d\geq 6} \supset \left(\frac{\mathsf{C}_6}{\mathsf{\Lambda}^2}\frac{\mathsf{5}\mathsf{v}^3}{2} + \frac{\mathsf{C}_8}{\mathsf{\Lambda}^4}\frac{\mathsf{7}\mathsf{v}^5}{2}\right)h^3 + \left(\frac{\mathsf{C}_6}{\mathsf{\Lambda}^2}\frac{\mathsf{15}\mathsf{v}^2}{8} + \frac{\mathsf{C}_8}{\mathsf{\Lambda}^4}\frac{\mathsf{35}\mathsf{v}^4}{8}\right)h^4$$



• SMEFT is a linear extension around the SM point.

$$\mathcal{L}_{\mathsf{SMEFT}}^{\mathsf{Higgs},d\geq 6} \supset \left(\frac{\mathsf{C}_6}{\mathsf{\Lambda}^2}\frac{\mathsf{5}\mathsf{v}^3}{2} + \frac{\mathsf{C}_8}{\mathsf{\Lambda}^4}\frac{\mathsf{7}\mathsf{v}^5}{2}\right)h^3 + \left(\frac{\mathsf{C}_6}{\mathsf{\Lambda}^2}\frac{\mathsf{15}\mathsf{v}^2}{8} + \frac{\mathsf{C}_8}{\mathsf{\Lambda}^4}\frac{\mathsf{35}\mathsf{v}^4}{8}\right)h^4$$

• HEFT is a non-linear realisation of EWSB.

$$\Phi \rightarrow \frac{v+h}{\sqrt{2}} U \begin{pmatrix} 0 \\ 1 \end{pmatrix}, \qquad U = \exp\left(\frac{i\sigma^j \pi_j}{v}\right)$$

(Alonso et al. 2013; I. Brivio et al. 2014; Gerhard Buchalla et al. 2014; G. Buchalla et al. 2016)



• SMEFT is a linear extension around the SM point.

$$\mathcal{L}_{\mathsf{SMEFT}}^{\mathsf{Higgs},d \ge 6} \supset \left(\frac{\mathsf{C}_6}{\mathsf{\Lambda}^2} \frac{\mathsf{5} \mathsf{v}^3}{2} + \frac{\mathsf{C}_8}{\mathsf{\Lambda}^4} \frac{\mathsf{7} \mathsf{v}^5}{2}\right) h^3 + \left(\frac{\mathsf{C}_6}{\mathsf{\Lambda}^2} \frac{\mathsf{15} \mathsf{v}^2}{8} + \frac{\mathsf{C}_8}{\mathsf{\Lambda}^4} \frac{\mathsf{35} \mathsf{v}^4}{8}\right) h^4$$

• HEFT is a non-linear realisation of EWSB.

$$\mathcal{L}_{\text{HEFT}}^{\text{Higgs}} \supset -\left(\frac{1}{2}m_h^2h^2 + \frac{1}{2}\kappa_3\frac{m_h^2}{v}h^3 + \frac{1}{2}\kappa_4\frac{m_h^2}{v^2}h^4\right)$$



• SMEFT is a linear extension around the SM point.

$$\mathcal{L}_{\mathsf{SMEFT}}^{\mathsf{Higgs},d\geq 6} \supset \left(\frac{\mathsf{C}_6}{\mathsf{\Lambda}^2}\frac{\mathsf{5}\mathsf{v}^3}{2} + \frac{\mathsf{C}_8}{\mathsf{\Lambda}^4}\frac{\mathsf{7}\mathsf{v}^5}{2}\right)h^3 + \left(\frac{\mathsf{C}_6}{\mathsf{\Lambda}^2}\frac{\mathsf{15}\mathsf{v}^2}{8} + \frac{\mathsf{C}_8}{\mathsf{\Lambda}^4}\frac{\mathsf{35}\mathsf{v}^4}{8}\right)h^4$$

• HEFT is a non-linear realisation of EWSB.

$$\mathcal{L}_{\text{HEFT}}^{\text{Higgs}} \supset -\left(\frac{1}{2}m_h^2h^2 + \frac{1}{2}\kappa_3\frac{m_h^2}{v}h^3 + \frac{1}{2}\kappa_4\frac{m_h^2}{v^2}h^4\right)$$

 \Rightarrow HEFT can effectively capture any decorrelated new physics effects in hhh that is absent in di-/single Higgs.



• SMEFT vs HEFT.

(Delgado et al. 2023; Anisha et al. 2024)

• **hhh** \Rightarrow Higgs potential (κ_3, κ_4)

COUPLING CONSTRAINTS

• hhh significantly improves constraints on κ_4 .



COUPLING CONSTRAINTS

• hhh significantly improves constraints on κ_4 .





(Fuks et al. 2016; Stylianou et al. 2023)

COUPLING CONSTRAINTS

• hhh significantly improves constraints on κ_4 .



(Fuks et al. 2016; Stylianou et al. 2023)

 Order of magnitude improvement in precision at FCC-hh and high energy lepton colliders.



• SMEFT vs HEFT.

(Delgado et al. 2023; Anisha et al. 2024)

• **hhh** \Rightarrow Higgs potential \Rightarrow EW vacuum structure \Rightarrow EWPTs, baryogenesis, stability of the universe, etc.

(Papaefstathiou and Tetlalmatzi-Xolocotzi 2023; Stylianou et al. 2023; Karkout et al. 2024)

• EW symmetry breaking in the Early Universe at $T \sim \mathcal{O}(100)$ GeV \Rightarrow activation of the Higgs Mechanism.



(Credits: Lisa Biermann)

(Sakharov 1967)

• EW symmetry breaking in the Early Universe at $T \sim \mathcal{O}(100)$ GeV \Rightarrow activation of the Higgs Mechanism.



(Credits: Lisa Biermann)

Strong First Order PTs required to explain EW Baryogenesis
 ⇒ Matter-antimatter asymmetry.

SM predicts smooth crossover for the Higgs mass, NO FOPTs!
 ⇒ insufficient for the required baryogenesis.



SM predicts smooth crossover for the Higgs mass, NO FOPTs!
 ⇒ insufficient for the required baryogenesis.



• Modifications to the scalar potential / couplings \Rightarrow modified EW vacuum structure \Rightarrow Strong FOPTs.



• SMEFT vs HEFT.

(Delgado et al. 2023; Anisha et al. 2024)

• **hhh** \Rightarrow Higgs potential \Rightarrow EW vacuum structure \Rightarrow EWPTs, baryogenesis, stability of the universe, etc.

(Papaefstathiou and Tetlalmatzi-Xolocotzi 2023; Stylianou et al. 2023; Karkout et al. 2024)

• Enhancements in σ^{ggF}_{hhh}

hhh enhancements



⁽ATLAS 2024; Karkout et al. 2024)

• Upper limit on SM hhh cross-section = 59 fb.



• SMEFT vs HEFT.

(Delgado et al. 2023; Anisha et al. 2024)

• **hhh** \Rightarrow Higgs potential \Rightarrow EW vacuum structure \Rightarrow EWPTs, baryogenesis, stability of the universe, etc.

(Papaefstathiou and Tetlalmatzi-Xolocotzi 2023; Stylianou et al. 2023; Karkout et al. 2024)

• Enhancements in $\sigma_{hhh}^{ggF} \sim \mathcal{O}(100)$ events at HL-LHC, can be relevant at FCC-*hh*.

(Papaefstathiou et al. 2019; Papaefstathiou et al. 2021; Papaefstathiou et al. 2016; Fuks et al. 2016; ATLAS 2024)

METHODOLOGY

$$\begin{split} V_{\text{2HDM}} &= m_{11}^2 \Phi_1^{\dagger} \Phi_1 + m_{22}^2 \Phi_2^{\dagger} \Phi_2 - m_{12}^2 \left(\Phi_1^{\dagger} \Phi_2 + \text{h.c.} \right) + \frac{\lambda_1}{2} \left(\Phi_1^{\dagger} \Phi_1 \right)^2 \\ &+ \frac{\lambda_2}{2} \left(\Phi_2^{\dagger} \Phi_2 \right)^2 + \lambda_3 \left(\Phi_1^{\dagger} \Phi_1 \right) \left(\Phi_2^{\dagger} \Phi_2 \right) + \lambda_4 \left(\Phi_1^{\dagger} \Phi_2 \right) \left(\Phi_2^{\dagger} \Phi_1 \right) \\ &+ \frac{\lambda_5}{2} \left[\left(\Phi_1^{\dagger} \Phi_2 \right)^2 + \text{h.c.} \right] \qquad \text{(softly broken } \mathbb{Z}_2 \text{ symmetry}) \end{split}$$

R2HDM : $m_{11}^2, m_{22}^2, m_{12}^2, \lambda_{1,...,4}, \lambda_5 \in \mathbb{R}$ (CP-Conserving)**C2HDM** : $m_{11}^2, m_{22}^2, \lambda_{1,...,4} \in \mathbb{R}, \ m_{12}^2, \lambda_5 \in \mathbb{C}$ (CP-Violating)

$$\begin{split} V_{\text{2HDM}} &= m_{11}^2 \Phi_1^{\dagger} \Phi_1 + m_{22}^2 \Phi_2^{\dagger} \Phi_2 - m_{12}^2 \left(\Phi_1^{\dagger} \Phi_2 + \text{h.c.} \right) + \frac{\lambda_1}{2} \left(\Phi_1^{\dagger} \Phi_1 \right)^2 \\ &+ \frac{\lambda_2}{2} \left(\Phi_2^{\dagger} \Phi_2 \right)^2 + \lambda_3 \left(\Phi_1^{\dagger} \Phi_1 \right) \left(\Phi_2^{\dagger} \Phi_2 \right) + \lambda_4 \left(\Phi_1^{\dagger} \Phi_2 \right) \left(\Phi_2^{\dagger} \Phi_1 \right) \\ &+ \frac{\lambda_5}{2} \left[\left(\Phi_1^{\dagger} \Phi_2 \right)^2 + \text{h.c.} \right] \qquad \text{(softly broken } \mathbb{Z}_2 \text{ symmetry}) \end{split}$$

R2HDM : $m_{11}^2, m_{22}^2, m_{12}^2, \lambda_{1,...,4}, \lambda_5 \in \mathbb{R}$ (CP-Conserving)**C2HDM** : $m_{11}^2, m_{22}^2, \lambda_{1,...,4} \in \mathbb{R}, \ m_{12}^2, \lambda_5 \in \mathbb{C}$ (CP-Violating)

Next-to-Minimal 2HDM (**N2HDM**): R2HDM + $\Phi_S \leftarrow$ Singlet

$$V_{\text{N2HDM}} = V_{\text{R2HDM}} + \frac{1}{2}m_{\text{S}}^2\Phi_{\text{S}}^2 + \frac{\lambda_6}{8}\Phi_{\text{S}}^4 + \frac{\lambda_7}{2}\left(\Phi_1^{\dagger}\Phi_1\right)\Phi_{\text{S}}^2 + \frac{\lambda_8}{2}\left(\Phi_2^{\dagger}\Phi_2\right)\Phi_{\text{S}}^2$$

$$\begin{split} V_{2\text{HDM}} &= m_{11}^2 \Phi_1^{\dagger} \Phi_1 + m_{22}^2 \Phi_2^{\dagger} \Phi_2 - m_{12}^2 \left(\Phi_1^{\dagger} \Phi_2 + \text{h.c.} \right) + \frac{\lambda_1}{2} \left(\Phi_1^{\dagger} \Phi_1 \right)^2 \\ &+ \frac{\lambda_2}{2} \left(\Phi_2^{\dagger} \Phi_2 \right)^2 + \lambda_3 \left(\Phi_1^{\dagger} \Phi_1 \right) \left(\Phi_2^{\dagger} \Phi_2 \right) + \lambda_4 \left(\Phi_1^{\dagger} \Phi_2 \right) \left(\Phi_2^{\dagger} \Phi_1 \right) \\ &+ \frac{\lambda_5}{2} \left[\left(\Phi_1^{\dagger} \Phi_2 \right)^2 + \text{h.c.} \right] \\ V_{\text{N2HDM}} &= V_{\text{R2HDM}} + \frac{1}{2} m_5^2 \Phi_5^2 + \frac{\lambda_6}{8} \Phi_5^4 + \frac{\lambda_7}{2} \left(\Phi_1^{\dagger} \Phi_1 \right) \Phi_5^2 + \frac{\lambda_8}{2} \left(\Phi_2^{\dagger} \Phi_2 \right) \Phi_5^2 \end{split}$$

• R2HDM: 2 CP-Even neutral mass eigenstates (h, H)

 $m_h pprox$ 125 GeV $< m_H$

$$\begin{split} V_{2\text{HDM}} &= m_{11}^2 \Phi_1^{\dagger} \Phi_1 + m_{22}^2 \Phi_2^{\dagger} \Phi_2 - m_{12}^2 \left(\Phi_1^{\dagger} \Phi_2 + \text{h.c.} \right) + \frac{\lambda_1}{2} \left(\Phi_1^{\dagger} \Phi_1 \right)^2 \\ &+ \frac{\lambda_2}{2} \left(\Phi_2^{\dagger} \Phi_2 \right)^2 + \lambda_3 \left(\Phi_1^{\dagger} \Phi_1 \right) \left(\Phi_2^{\dagger} \Phi_2 \right) + \lambda_4 \left(\Phi_1^{\dagger} \Phi_2 \right) \left(\Phi_2^{\dagger} \Phi_1 \right) \\ &+ \frac{\lambda_5}{2} \left[\left(\Phi_1^{\dagger} \Phi_2 \right)^2 + \text{h.c.} \right] \\ V_{\text{N2HDM}} &= V_{\text{R2HDM}} + \frac{1}{2} m_5^2 \Phi_5^2 + \frac{\lambda_6}{8} \Phi_5^4 + \frac{\lambda_7}{2} \left(\Phi_1^{\dagger} \Phi_1 \right) \Phi_5^2 + \frac{\lambda_8}{2} \left(\Phi_2^{\dagger} \Phi_2 \right) \Phi_5^2 \end{split}$$

• R2HDM: 2 **CP-Even** neutral mass eigenstates (*h*, *H*)

 $m_h pprox$ 125 GeV $< m_H$

• C2HDM and N2HDM : $\frac{3}{2}$ neutral mass eigenstates (H_1 , H_2 , H_3).

$$m_{H_1} \cong m_h pprox$$
 125 GeV $< m_{H_2} < m_{H_3}$

. R2HDM: 2 **CP-Even** neutral mass eigenstates (*h*, *H*).

 $m_h pprox$ 125 GeV $< m_H$

• C2HDM and N2HDM : 3 neutral mass eigenstates (H₁, H₂, H₃).

 $m_{H_1} \cong m_h pprox$ 125 GeV $< m_{H_2} < m_{H_3}$

. R2HDM: 2 **CP-Even** neutral mass eigenstates (*h*, *H*).

 $m_h pprox$ 125 GeV $< m_H$

• C2HDM and N2HDM : 3 neutral mass eigenstates (H_1, H_2, H_3) .

 $m_{H_1} \cong m_h pprox 125 \text{ GeV} < m_{H_2} < m_{H_3}$

Parameter scans:

- ⇒ **ScannerS**: Vary exotic Higgs masses, mixing angles.
- ⇒ HiggsTools: Apply theoretical, experimental constraints

(Mühlleitner et al. 2022; Bechtle, Dercks, et al. 2020; Bechtle, Heinemeyer, et al. 2021; Bahl et al. 2023)

. R2HDM: 2 **CP-Even** neutral mass eigenstates (*h*, *H*).

 $m_h pprox$ 125 GeV $< m_H$

• C2HDM and N2HDM : 3 neutral mass eigenstates (H_1, H_2, H_3) .

 $m_{H_1} \cong m_h pprox$ 125 GeV $< m_{H_2} < m_{H_3}$

Parameter scans:

- ⇒ ScannerS: Vary exotic Higgs masses, mixing angles.
- ⇒ HiggsTools: Apply theoretical, experimental constraints

(Mühlleitner et al. 2022; Bechtle, Dercks, et al. 2020; Bechtle, Heinemeyer, et al. 2021; Bahl et al. 2023)

 FeynRules ⇒ MadGraph_aMC@NLO: Model implementations, generate hh(h) cross-sections.

(Alloul et al. 2014; Degrande et al. 2012; Darmé et al. 2023; Alwall et al. 2014)

 Model-independent implementation of the one-loop daisy-resummed effective potential at finite temperature.



 Model-independent implementation of the one-loop daisy-resummed effective potential at finite temperature.



 \rightarrow Strength of FOPTs:

$$\xi(T) = \frac{v_{\rm EW}(T)}{T}$$

 $v_{EW}(T=0) = 246 \text{ GeV}$

 Model-independent implementation of the one-loop daisy-resummed effective potential at finite temperature.



$$\rightarrow$$
 Strength of FOPTs:

 $\rightarrow\,$ Strong FOPTs:

$$\xi(T) = \frac{v_{\rm EW}(T)}{T}$$

$$\xi(T_p) = \frac{v_{\rm EW}(T_p)}{T_p} \gtrsim 1$$

 T_p : percolation stage $< T_c$

$$v_{\rm EW}(T=0)=246~{\rm GeV}$$

 Model-independent implementation of the one-loop daisy-resummed effective potential at finite temperature.



$$ightarrow$$
 Strength of FOPTs:

 \rightarrow Strong FOPTs:

$$\xi(T) = \frac{v_{\text{EW}}(T)}{T} \qquad \qquad \xi(T_p) = \frac{v_{\text{EW}}(T_p)}{T_p} \gtrsim 1$$

 $v_{EW}(T = 0) = 246 \text{ GeV}$ T_p : percolation stage $< T_c$

. Minimising $V_{\text{eff.}} \Rightarrow (g_{hhh}, g_{hhhh}, ...) \Rightarrow MadGraph$

RESULTS

R2HDM AT THE LHC AND FCC-hh



- . hhh enhanced compared to hh, although clearly correlated!
- . Enhancements generalise to FCC-hh!

ENHANCING hh/hhh (R2HDM)

Resonant contributions to hh



ENHANCING hh/hhh (R2HDM)

Resonant contributions to hhh



ENHANCEMENTS IN R2HDM



• BP-1: largely enhanced $(m_H \simeq 2m_h)$.

BP-2: SM-like.

BPs	$\sigma_{hh}/\sigma_{hh}^{ m SM}$	$\sigma_{hhh}/\sigma_{hhh}^{ m SM}$	M _H [GeV]	g _{hhH} [GeV]	g _{hhhH}
BP-1	3.24	15.26	274.29	75.28	0.203
BP-2	1.02	1.02	469.30	-7.11	-0.011

MASS SPECTRA OF R2HDM



• EWPTs driven by the physics of light dofs.

⇒ Stronger phase transitions proceed via lighter spectra!



⇒ Neutral Higgs rates alone NOT indicative of the strength of EWPTs.

ENHANCING hh/hhh (3 DOFS)

Resonant contributions to hh



ENHANCING hh/hhh (3 DOFS)

Resonant contributions to hhh



R2HDM + C2HDM



- Additional dof \Rightarrow hhh more enhanced.
- Stringent EDMs \Rightarrow Minimal CP admixture (\lesssim 10%), thus no dramatic changes.

N2HDM



- . The additional dof enhances hhh, like C2HDM.
- . Enhancements \sim 10 25 in hhh, within HL-LHC hh sensitivity
 - \Rightarrow can be accessible in FCC-hh!

 hhh-production enhanced over the SM cross-section for extended scalar sectors.

- hhh-production enhanced over the SM cross-section for extended scalar sectors.
- Enhancements driven by $H \rightarrow hh$, $H \rightarrow hhh$, additional dofs, largest increases resulting from their combination.

- hhh-production enhanced over the SM cross-section for extended scalar sectors.
- Enhancements driven by $H \rightarrow hh$, $H \rightarrow hhh$, additional dofs, largest increases resulting from their combination.
- The lightness of the exotics affects EWPT, multi-Higgs production rates.

- hhh-production enhanced over the SM cross-section for extended scalar sectors.
- Enhancements driven by $H \rightarrow hh$, $H \rightarrow hhh$, additional dofs, largest increases resulting from their combination.
- The lightness of the exotics affects EWPT, multi-Higgs production rates.
- . hhh sensitivity at HL-LHC is challenging, but still motivated!

- hhh-production enhanced over the SM cross-section for extended scalar sectors.
- Enhancements driven by $H \rightarrow hh$, $H \rightarrow hhh$, additional dofs, largest increases resulting from their combination.
- The lightness of the exotics affects EWPT, multi-Higgs production rates.
- . hhh sensitivity at HL-LHC is challenging, but still motivated!

Thank you!