

# Listening for ultra-heavy DM with underwater acoustic detectors

**Damon Cleaver**, Christopher McCabe and Ciaran A.J. O'Hare YTF 2024, 19/12/24

### What is Dark Matter?

### What is Dark Matter?

### $\Omega_{\rm CDM} h^2 \sim 0.120 \pm 0.001$

[Planck, 2018]

### What is Dark Matter?



### $\Omega_{\rm CDM} h^2 \sim 0.120 \pm 0.001$

[Planck, 2018]



• Ultra-heavy dark matter is necessarily composite (if thermally produced) due to s-wave unitarity



- Ultra-heavy dark matter is necessarily composite (if thermally produced) due to s-wave unitarity
- Many different models for UHDM
  - PBH, Nuggets, Blobs, WIMPonium, Q-Balls etc...



- Ultra-heavy dark matter is necessarily composite (if thermally produced) due to s-wave unitarity
- Many different models for UHDM
  - PBH, Nuggets, Blobs, WIMPonium, Q-Balls etc...
- Is there a nice model-independent way to treat them?



- Ultra-heavy dark matter is necessarily composite (if thermally produced) due to s-wave unitarity
- Many different models for UHDM
  - PBH, Nuggets, Blobs, WIMPonium, Q-Balls etc...
- Is there a nice model-independent way to treat them?
- Answer: Yes (for some parts of parameter space)



- Ultra-heavy dark matter is necessarily composite (if thermally produced) due to s-wave unitarity
- Many different models for UHDM
  - PBH, Nuggets, Blobs, WIMPonium, Q-Balls etc...
- Is there a nice model-independent way to treat them?
- Answer: Yes (for some parts of parameter space)



### Macroscopic DM Candidate

\*\*\*\*\*\*\*\*\*\*

# Interaction range



• Consider parameters of models where:

Macroscopic DM Candidate

\*\*\*\*

### **Interaction range**



- Consider parameters of models where:
  - The DM is Planck-mass or larger

Macroscopic DM Candidate

### **Interaction range**

\*\*\*\*\*



- Consider parameters of models where:
  - The DM is Planck-mass or larger
  - DM Radius  $R_{\gamma}$  much larger than interaction length scale



\*

### **Interaction range**

\*\*\*\*\*



- Consider parameters of models where:
  - The DM is Planck-mass or larger
  - DM Radius  $R_{\gamma}$  much larger than interaction length scale
  - Geometric cross section dominates i.e.  $\sigma_{\chi} \approx \pi R_{\chi}^2$



### **Interaction range**

\*\*\*\*\*\*



- Consider parameters of models where:
  - The DM is Planck-mass or larger
  - DM Radius  $R_{\gamma}$  much larger than interaction length scale
  - Geometric cross section dominates
  - Parameterise the interaction in terms of  $R_{\gamma}$ , -> set by the theory -> make experimental statements about multiple models!

s i.e. 
$$\sigma_{\chi} \approx \pi R_{\chi}^2$$

### Macroscopic DM Candidate

### Interaction range



- Consider parameters of models where:
  - The DM is Planck-mass or larger
  - DM Radius  $R_{\gamma}$  much larger than interaction length scale
  - Geometric cross section dominates
  - Parameterise the interaction in terms of  $R_{\gamma}$ , -> set by the theory -> make experimental statements about multiple models!

s i.e. 
$$\sigma_{\chi} \approx \pi R_{\chi}^2$$

### Macroscopic DM Candidate

### Interaction range



### • We parametrise UHDM in grams (g)

• We parametrise UHDM in grams (g)

• DM Flux:  $\phi_{\chi} \approx 6 \left(\frac{1 \text{ g}}{m_{\chi}}\right) \text{ km}^{-2} \text{ yr}^{-1}$ 

• We parametrise UHDM in grams (g)

DM Flux: 
$$\phi_{\chi} \approx 6 \left(\frac{1 \text{ g}}{m_{\chi}}\right) \text{ km}^{-2} \text{ yr}^{-1}$$

number of events.

• Need a very large detector (or very long integration time) to have significant

• We parametrise UHDM in grams (g)

DM Flux: 
$$\phi_{\chi} \approx 6 \left(\frac{1 \text{ g}}{m_{\chi}}\right) \text{ km}^{-2} \text{ yr}^{-1}$$

- number of events.
- No hope for conventional detectors (LUX-ZEPLIN, XENONnT etc.)

• Need a very large detector (or very long integration time) to have significant

• We parametrise UHDM in grams (g)

DM Flux: 
$$\phi_{\chi} \approx 6 \left(\frac{1 \text{ g}}{m_{\chi}}\right) \text{ km}^{-2} \text{ yr}^{-1}$$

- number of events.
- No hope for conventional detectors (LUX-ZEPLIN, XENONnT etc.)

• Need a very large detector (or very long integration time) to have significant









6



### Annoying "gap" in constraints



- Annoying "gap" in constraints
- Mica underground too much overburden

6



- Annoying "gap" in constraints
- Mica underground too much overburden
- Radar not sensitive enough - not enough ionisation



- Annoying "gap" in constraints
- Mica underground too much overburden
- Radar not sensitive enough - not enough ionisation
- What phenomena could  $\bullet$ we use to constrain this region?





### • Idea: DM is weakly interacting enough to make it through the atmosphere



- Idea: DM is weakly interacting enough to make it through the atmosphere
- Reaches much more dense medium: the ocean



- Idea: DM is weakly interacting enough to make it through the atmosphere
- Reaches much more dense medium: the ocean
- DM deposits energy into the ocean creating pressure waves



- Idea: DM is weakly interacting enough to make it through the atmosphere
- Reaches much more dense medium: the ocean
- DM deposits energy into the ocean creating pressure waves
- Detect pressure waves using a large hydrophone array



- Idea: DM is weakly interacting enough to make it through the atmosphere
- Reaches much more dense medium: the ocean
- DM deposits energy into the ocean creating pressure waves
- Detect pressure waves using a large hydrophone array



- Idea: DM is weakly interacting enough to make it through the atmosphere
- Reaches much more dense medium: the ocean
- DM deposits energy into the ocean creating pressure waves
- Detect pressure waves using a large hydrophone array


#### Acoustic Detection

- Idea: DM is weakly interacting enough to make it through the atmosphere
- Reaches much more dense medium: the ocean
- DM deposits energy into the ocean creating pressure waves
- Detect pressure waves using a large hydrophone array



# • Propositions for acoustic neutrino experiments with $\mathcal{O}(100 \text{ km}^3)$ hydrophone arrays [Lahmann, 2016]

- arrays [Lahmann, 2016]
- section

#### • Propositions for acoustic neutrino experiments with $\mathcal{O}(100 \text{ km}^3)$ hydrophone

- arrays [Lahmann, 2016]
- section
- Acoustic propagation distance in water much greater than light -> less dense instrumentation

#### • Propositions for acoustic neutrino experiments with $\mathcal{O}(100 \text{ km}^3)$ hydrophone

- arrays [Lahmann, 2016]
- section
- Acoustic propagation distance in water much greater than light -> less dense instrumentation
- Energy deposition comes from hadronic showers

#### • Propositions for acoustic neutrino experiments with $\mathcal{O}(100 \text{ km}^3)$ hydrophone

- arrays [Lahmann, 2016]
- section
- Acoustic propagation distance in water much greater than light -> less dense instrumentation
- Energy deposition comes from hadronic showers

#### • Propositions for acoustic neutrino experiments with $\mathcal{O}(100 \text{ km}^3)$ hydrophone

- arrays [Lahmann, 2016]
- section
- Acoustic propagation distance in water much greater than light -> less dense instrumentation
- Energy deposition comes from hadronic showers

#### • Propositions for acoustic neutrino experiments with $\mathcal{O}(100 \text{ km}^3)$ hydrophone

Pressure waves come from thermo-acoustic heating, which obeys the following DE:

 $\nabla^2 p - \frac{1}{c_s^2} \frac{\partial^2 p}{\partial t^2} = -\frac{\alpha}{c_n} \frac{\partial^2 q(r,t)}{\partial t^2}$ 

Pressure waves come from thermo-acoustic heating, which obeys the following DE:

 $\nabla^2 p - \frac{1}{c_s^2} \frac{\partial^2 p}{\partial t^2} = -\frac{\alpha}{c_p} \frac{\partial^2 q(r,t)}{\partial t^2}$ Acoustic pressure

Pressure waves come from thermo-acoustic heating, which obeys the following DE:

**Acoustic pressure** 

**Energy Deposition Density** 



Pressure waves come from thermo-acoustic heating, which obeys the following DE:



#### General Solution: p(r, t)

**Energy Deposition Density** 

$$\frac{p}{2} = \frac{\alpha}{c_p} \frac{\partial^2 q(r,t)}{\partial t^2}$$

$$= \frac{\alpha}{4\pi c_p} \int \frac{\mathrm{d}^3 r'}{|r - r'|} \frac{\partial^2 q(r', t')}{\partial t^2}$$
$$t' = t - \frac{|r - r'|}{r} \frac{\partial^2 q(r', t')}{\partial t^2}$$

Taking a geometric cross section only and taking the number of scatters to infinity:

$$\frac{\mathrm{d}E_{\chi}}{\mathrm{d}z} = -\rho_{\mathrm{sea}}\sigma_{\chi}v_{\chi}^{2}\exp\left(-\frac{z}{\ell_{\mathrm{sea}}}\right)$$

Taking a geometric cross section only and taking the number of scatters to infinity:

$$\frac{\mathrm{d}E_{\chi}}{\mathrm{d}z} = -\rho_{\mathrm{sea}}\sigma_{\chi}v_{\chi}^{2}\exp\left(-\frac{z}{\ell_{\mathrm{sea}}}\right)$$

Where  $\ell_{sea}$  is the characteristic length of the energy deposition:

$$\mathscr{\ell}_{\text{sea}} = \frac{m_{\chi}}{2\rho_{\text{sea}}\sigma_{\chi}} \simeq 480 \,\text{km} \times \left(\frac{m_{\chi}}{10^{-2} \,\text{g}}\right) \left(\frac{10^{-10} \,\text{cm}^2}{\sigma_{\chi}}\right)$$

Taking a geometric cross section only and taking the number of scatters to infinity:

$$\frac{\mathrm{d}E_{\chi}}{\mathrm{d}z} = -\rho_{\mathrm{sea}}\sigma_{\chi}v_{\chi}^{2}\exp\left(-\frac{z}{\ell_{\mathrm{sea}}}\right)$$

Where  $\ell_{sea}$  is the characteristic length of the energy deposition:

$$\ell_{\text{sea}} = \frac{m_{\chi}}{2\rho_{\text{sea}}\sigma_{\chi}} \simeq 480 \,\text{km} \times \left(\frac{m_{\chi}}{10^{-2} \,\text{g}}\right) \left(\frac{10^{-10} \,\text{cm}^2}{\sigma_{\chi}}\right)$$

 $\ell_{\rm sea}$  can be very long, in this case:

$$\frac{\mathrm{d}E_{\chi}}{\mathrm{d}z} \simeq -\rho_{\mathrm{sea}}\sigma_{\chi}v_{\chi}^2 = \mathrm{const}$$



Model energy deposition rate as Gaussian cylinder:

$$q(r) = \sum_{A = \{H,O\}} \frac{1}{2\pi} \frac{dE_A}{dz} \frac{1}{\sigma_A^2} \exp\left(-\frac{\rho^2}{2\sigma_A^2}\right)$$



Model energy deposition rate as Gaussian cylinder:

$$q(r) = \sum_{A = \{H,O\}} \frac{1}{2\pi} \frac{dE_A}{dz} \frac{1}{\sigma_A^2} \exp\left(-\frac{\rho^2}{2\sigma_A^2}\right)$$

Where  $\sigma_A$  is the characteristic scattering length of species A.



Model energy deposition rate as Gaussian cylinder:

$$q(r) = \sum_{A = \{H,O\}} \frac{1}{2\pi} \frac{dE_A}{dz} \frac{1}{\sigma_A^2} \exp\left(-\frac{\rho^2}{2\sigma_A^2}\right)$$

Where  $\sigma_A$  is the characteristic scattering length of species A.

Gaussian allows us to find analytic solutions for the pressure - turns out to be enough to capture the physics



#### What is the pressure solution?

#### What is the pressure solution?

and take width much smaller than detection distance  $\sigma \ll \rho$ :

$$p(r,t;\sigma \ll \rho) = \frac{\alpha}{2\pi c_p} \frac{dE}{dz} \frac{c_s^2}{\sqrt{2\pi\sigma^3}} \frac{1}{\sqrt{\rho}} I_p\left(\frac{t-\rho/c_s}{\sigma/c_s}\right)$$

Take infinitely long line track (good for large  $\ell_{\rm DM}$ ), instantaneous energy deposition



#### What is the pressure solution?

and take width much smaller than detection distance  $\sigma \ll \rho$ :

$$p(r,t;\sigma \ll \rho) = \frac{\alpha}{2\pi c_p} \frac{dE}{dz} \frac{c_s^2}{\sqrt{2\pi\sigma^3}} \frac{1}{\sqrt{\rho}} I_p\left(\frac{t-\rho/c_s}{\sigma/c_s}\right)$$

$$I_p(A) = \int_0^\infty dY \sqrt{Y} \exp\left(-\frac{1}{2}\right) dY \sqrt{$$

$$= -\frac{\pi A}{4\sqrt{2}(A^2)^{1/4}} \exp\left(-\frac{A^2}{4}\right) \left[ \left(A + \sqrt{A^2}\right) \left(I_{1/4}\left(\frac{A^2}{4}\right) - I_{3/4}\left(\frac{A^2}{4}\right)\right) + \frac{\sqrt{2}}{\pi} \left(\sqrt{A^2}K_{1/4}\left(\frac{A^2}{4}\right) - AK_{3/4}\left(\frac{A^2}{4}\right)\right) \right]$$

Take infinitely long line track (good for large  $\ell_{\rm DM}$ ), instantaneous energy deposition

$$\frac{Y^2}{2}\right)\cos\left(A\,Y + \frac{\pi}{4}\right)$$



 $p(r,t;\sigma \ll \rho) = \frac{\alpha}{2\pi c_p} \frac{dE}{dz} \frac{c_s^2}{\sqrt{2\pi\sigma^3}} \frac{1}{\sqrt{\rho}} I_p\left(\frac{t-\rho/c_s}{\sigma/c_s}\right)$ 



$$p(r,t;\sigma \ll \rho) = \frac{\alpha}{2\pi c_p} \frac{dE}{dz} \frac{c_s^2}{\sqrt{2\pi\sigma^3}} \frac{1}{\sqrt{\rho}} I_p \left(\frac{t-\rho/\sigma}{\sigma/c_s}\right)$$

Shape determined by  $I_p \sim \mathcal{O}(1)$ . Solution is **bi-polar** 



$$p(r,t;\sigma \ll \rho) = \frac{\alpha}{2\pi c_p} \frac{dE}{dz} \frac{c_s^2}{\sqrt{2\pi\sigma^3}} \frac{1}{\sqrt{\rho}} I_p \left(\frac{t-\rho/\sigma}{\sigma/c_s}\right)$$

Shape determined by  $I_p \sim \mathcal{O}(1)$ . Solution is **bi-polar** 

Large MPa signal for UHDM in target parameter regions - **determined by prefactor** 



$$p(r,t;\sigma \ll \rho) = \frac{\alpha}{2\pi c_p} \frac{dE}{dz} \frac{c_s^2}{\sqrt{2\pi\sigma^3}} \frac{1}{\sqrt{\rho}} I_p \left(\frac{t-\rho/\sigma}{\sigma/c_s}\right)$$

Shape determined by  $I_p \sim \mathcal{O}(1)$ . Solution is **bi-polar** 

Large MPa signal for UHDM in target parameter regions - **determined by prefactor** 

Full pressure solution is sum of O and H contributions.







#### Can also find full frequency solution (must be solved numerically)





Can also find full frequency solution (must be solved numerically)

Frequency cut-off set by  $c_s/\sigma_A$ 





Can also find full frequency solution (must be solved numerically)

Frequency cut-off set by  $c_s/\sigma_A$ 

Can "integrate out" width to get an analytic approximation at lower freq:





Can also find full frequency solution (must be solved numerically)

Frequency cut-off set by  $c_s/\sigma_A$ 

Can "integrate out" width to get an analytic approximation at lower freq:

$$\tilde{p}_A(\rho,\omega) \approx \frac{\omega\alpha}{2\pi c_p} \frac{dE_A}{dz} \frac{\pi}{2} H_0^{(2)} \left(\frac{\rho\omega}{c_s}\right)$$





# Is this the full story?

# Is this the full story?

Need to account for other attenuation effects! Packaged into an absorption coefficient  $\tilde{a}(\omega)$ :

$$\tilde{a}(\omega) = \frac{\omega^2}{\omega_0 c_s} + \frac{2}{\lambda_1} \frac{i\omega}{\omega_1 + i\omega} + \frac{2}{\lambda_2} \frac{i\omega}{\omega_2 + i\omega}$$

# Is this the full story?

Need to account for other attenuation effects! Packaged into an absorption coefficient  $\tilde{a}(\omega)$ :


Need to account for other attenuation effects! Packaged into an absorption coefficient  $\tilde{a}(\omega)$ :



The pressure in frequency space becomes:

Need to account for other attenuation effects! Packaged into an absorption coefficient  $\tilde{a}(\omega)$ :



The pressure in frequency space becomes:

$$\tilde{p}_a(\rho,\omega) = \exp\left(\right.$$



 $-\frac{\tilde{a}(\omega)\rho}{2}\right) \tilde{p}(\rho,\omega)$ 





#### Takes frequency cut-off from $\mathcal{O}(10^{11} \,\mathrm{Hz})$ to $\mathcal{O}(10^5 \,\mathrm{Hz})$



#### Takes frequency cut-off from $\mathcal{O}(10^{11} \,\mathrm{Hz})$ to $\mathcal{O}(10^5 \,\mathrm{Hz})$

Cut-off profile is not Gaussian at certain characteristic distances





Nano-second pulse of MPa amplitude has become micro-second pulse at Pa amplitude



Nano-second pulse of MPa amplitude has become micro-second pulse at Pa amplitude

non-Gaussian frequency cut-off = different shaped pulses. **Distance dependent** 



Nano-second pulse of MPa amplitude has become micro-second pulse at Pa amplitude

non-Gaussian frequency cut-off = different shaped pulses. **Distance dependent** 



#### we can see.

• We need to understand the experimental noise sources to understand what signals



- we can see.
- state noise.

• We need to understand the experimental noise sources to understand what signals

• In the 10-50 kHz band, the noise is dominated by sea surface agitation called sea-



- we can see.
- state noise.
- **lowest sea state noise.** We are always sea-state noise limited

• We need to understand the experimental noise sources to understand what signals

• In the 10-50 kHz band, the noise is dominated by sea surface agitation called sea-

• Self noise of hydrophones designed for UHE neutrino detection is comparable to



- we can see.
- state noise.
- **lowest sea state noise.** We are always sea-state noise limited
- average to be approx. 5 mPa we take this as a baseline value.

• We need to understand the experimental noise sources to understand what signals

• In the 10-50 kHz band, the noise is dominated by sea surface agitation called sea-

• Self noise of hydrophones designed for UHE neutrino detection is comparable to

• Ocean Noise Experiment ( $O\nu NE$ ) has measured this noise is Mediterranean on



- we can see.
- state noise.
- **lowest sea state noise.** We are always sea-state noise limited
- average to be approx. 5 mPa we take this as a baseline value.
- using machine learning algorithms

• We need to understand the experimental noise sources to understand what signals

• In the 10-50 kHz band, the noise is dominated by sea surface agitation called sea-

• Self noise of hydrophones designed for UHE neutrino detection is comparable to

• Ocean Noise Experiment ( $O\nu NE$ ) has measured this noise is Mediterranean on

• Transient noise sources exist (e.g. dolphins) we assume these can be taken care of





- we can see.
- state noise.
- **lowest sea state noise.** We are always sea-state noise limited
- average to be approx. 5 mPa we take this as a baseline value.
- using machine learning algorithms

• We need to understand the experimental noise sources to understand what signals

• In the 10-50 kHz band, the noise is dominated by sea surface agitation called sea-

• Self noise of hydrophones designed for UHE neutrino detection is comparable to

• Ocean Noise Experiment ( $O\nu NE$ ) has measured this noise is Mediterranean on

• Transient noise sources exist (e.g. dolphins) we assume these can be taken care of





- we can see.
- state noise.
- **lowest sea state noise.** We are always sea-state noise limited
- average to be approx. 5 mPa we take this as a baseline value.
- using machine learning algorithms

• We need to understand the experimental noise sources to understand what signals

• In the 10-50 kHz band, the noise is dominated by sea surface agitation called sea-

• Self noise of hydrophones designed for UHE neutrino detection is comparable to

• Ocean Noise Experiment ( $O\nu NE$ ) has measured this noise is Mediterranean on

• Transient noise sources exist (e.g. dolphins) we assume these can be taken care of







## **Preliminary Sensitivities**

- Assuming proposed acoustic neutrino experiment parameters, could constrain the gap!
- Array Geometry: 10km x 10km x 1km with 13 x 13 x 10 hydrophone distribution. 2km depth.
- Requirements p = 5 mPa,  $\rho = 300$ m,  $N_{\rm events} \ge 100/{
  m yr}$
- Complementary to Humans, Mica, Ohya and Cosmological Bounds



# Punchline:

Future acoustic neutrino experiments could have the power to constrain **UHDM** candidates

# Thank you for listening! Any Questions?

# Backup slides





• DM is galactic in origin. Neutrinos are extragalactic.





- DM is galactic in origin. Neutrinos are extragalactic.
- Direction of dark matter in the direction of the Cygnus constellation.





- DM is galactic in origin. Neutrinos are extragalactic.
- Direction of dark matter in the direction of the Cygnus constellation.
- So UHDM flux modulates





- DM is galactic in origin. Neutrinos are extragalactic.
- Direction of dark matter in the direction of the Cygnus constellation.
- So UHDM flux modulates
- Varies throughout the year due to following the sidereal day over solar day





- DM is galactic in origin. Neutrinos are extragalactic.
- Direction of dark matter in the direction of the Cygnus constellation.
- So UHDM flux modulates
- Varies throughout the year due to following the sidereal day over solar day



# Angular Entry to Array







• Pulse gets more asymmetry from non-Gaussian cut-off shape in freq space.





- Pulse gets more asymmetry from non-Gaussian cut-off shape in freq space.
- In pure water, freq. cut-off is Gaussian.
   Shape is the same as unattenuated case!



- Pulse gets more asymmetry from non-Gaussian cut-off shape in freq space.
- In pure water, freq. cut-off is Gaussian.
   Shape is the same as unattenuated case!
- Maximal asymmetry near the characteristic absorption scale of Magnesium Sulphate ( $\lambda_2 = 152.7$ m)



- Pulse gets more asymmetry from non-Gaussian cut-off shape in freq space.
- In pure water, freq. cut-off is Gaussian.
   Shape is the same as unattenuated case!
- Maximal asymmetry near the characteristic absorption scale of Magnesium Sulphate ( $\lambda_2 = 152.7$ m)
- Never quite becomes same as pure water due to imaginary components to the absorption coefficient



- Pulse gets more asymmetry from non-Gaussian cut-off shape in freq space.
- In pure water, freq. cut-off is Gaussian.
   Shape is the same as unattenuated case!
- Maximal asymmetry near the characteristic absorption scale of Magnesium Sulphate ( $\lambda_2 = 152.7$ m)
- Never quite becomes same as pure water due to imaginary components to the absorption coefficient



- Pulse gets more asymmetry from non-Gaussian cut-off shape in freq space.
- In pure water, freq. cut-off is Gaussian.
   Shape is the same as unattenuated case!
- Maximal asymmetry near the characteristic absorption scale of Magnesium Sulphate ( $\lambda_2 = 152.7$ m)
- Never quite becomes same as pure water due to imaginary components to the absorption coefficient


## Finding a good hydrophone distance





• Calculate range using slowing down approximation range of H and O using SRIM software package.



- Calculate range using slowing down approximation range of H and O using SRIM software package.
- Find at typical recoil energy of  $\overline{E}_O = 30.2$  keV and  $\overline{E}_H = 1.9$  keV, the fit widths are  $\sigma_O = 0.14 \ \mu m$  and  $\sigma_H = 0.082 \ \mu m$



- Calculate range using slowing down approximation range of H and O using SRIM software package.
- Find at typical recoil energy of  $\overline{E}_O = 30.2$  keV and  $\overline{E}_H = 1.9$  keV, the fit widths are  $\sigma_O = 0.14 \ \mu m$  and  $\sigma_H = 0.082 \ \mu m$
- While Gaussian is not an excellent fit, the true nature of the distribution is irrelevant after attenuation



- Calculate range using slowing down approximation range of H and O using SRIM software package.
- Find at typical recoil energy of  $\overline{E}_O = 30.2$  keV and  $\overline{E}_H = 1.9$  keV, the fit widths are  $\sigma_O = 0.14 \ \mu m$  and  $\sigma_H = 0.082 \ \mu m$
- While Gaussian is not an excellent fit, the true nature of the distribution is irrelevant after attenuation



- Calculate range using slowing down approximation range of H and O using SRIM software package.
- Find at typical recoil energy of  $\overline{E}_O = 30.2$  keV and  $\overline{E}_H = 1.9$  keV, the fit widths are  $\sigma_O = 0.14 \ \mu m$  and  $\sigma_H = 0.082 \ \mu m$
- While Gaussian is not an excellent fit, the true nature of the distribution is irrelevant after attenuation



 $p(\rho, z, t) = \frac{\alpha}{4\pi c_p} \frac{dE}{dz'}$ 

 $q(\rho', z') = \frac{1}{2\pi} \frac{\delta(\rho')}{\rho'} C e^{az'}$ 

$$\left| \frac{\partial}{\partial t} \left( \frac{e^{a(c_s \sqrt{t^2 - t_0^2})}}{\sqrt{t^2 - t_0^2}} \Theta(t - t_0) \right). \right.$$

If we let the DM energy evolve but with no track width:

 $q(\rho', z') =$ 

 $p(\rho, z, t) = \frac{\alpha}{4\pi c_p} \frac{dE}{dz'}$ 

$$= \frac{1}{2\pi} \frac{\delta(\rho')}{\rho'} C e^{az'}$$

$$\left| \frac{\partial}{\partial t} \left( \frac{e^{a(c_s \sqrt{t^2 - t_0^2})}}{\sqrt{t^2 - t_0^2}} \Theta(t - t_0) \right). \right.$$

If we let the DM energy evolve but with no track width:

 $q(\rho', z') =$ 

We find a pressure solution:

$$p(\rho, z, t) = \frac{\alpha}{4\pi c_p} \frac{dE}{dz'} \bigg|_{z'=z} \frac{\partial}{\partial t} \left( \frac{e^{a(c_s \sqrt{t^2 - t_0^2})}}{\sqrt{t^2 - t_0^2}} \Theta(t - t_0) \right).$$

$$= \frac{1}{2\pi} \frac{\delta(\rho')}{\rho'} C e^{az'}$$

If we let the DM energy evolve but with no track width:

$$q(\rho',z')$$
 =

#### We find a pressure solution:

$$p(\rho, z, t) = \frac{\alpha}{4\pi c_p} \frac{dE}{dz'} \bigg|_{z'=z} \frac{\partial}{\partial t} \left( \frac{e^{a(c_s \sqrt{t^2 - t_0^2})}}{\sqrt{t^2 - t_0^2}} \Theta(t - t_0) \right).$$

Signal amplitude determined by the energy deposition rate **perpendicular to the hydrophone and the track - makes sense for cylindrical plane waves!** 

$$= \frac{1}{2\pi} \frac{\delta(\rho')}{\rho'} C e^{az'}$$

If we let the DM energy evolve but with no track width:

 $q(\rho', z') =$ 

We find a pressure solution:

$$p(\rho, z, t) = \frac{\alpha}{4\pi c_p} \frac{dE}{dz'} \bigg|_{z'=z} \frac{\partial}{\partial t} \left( \frac{e^{a(c_s \sqrt{t^2 - t_0^2})}}{\sqrt{t^2 - t_0^2}} \Theta(t - t_0) \right).$$

Signal amplitude determined by the energy deposition rate **perpendicular to the hydrophone and the track - makes sense for cylindrical plane waves!** 

$$= \frac{1}{2\pi} \frac{\delta(\rho')}{\rho'} C e^{az'}$$