

Perturbative uncertainties on the p_T Z/W boson spectrum

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Durham, England

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CLUSTER OF EXCELLENCE
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Outline

- » Color singlet production p_T spectrum
- » Meaningful theory uncertainty and perturbative uncertainty
 - » Correlations
 - » Scale variations
 - » Theory Nuisance Parameters (TNPs)
- » Applications with TNPs
 - » $\alpha_s(m_Z)$ from $Z p_T$ spectrum
 - » W mass determination

Color singlet production p_T spectrum

» Wide-ranging applications, many precise measurements:

[ATLAS '20](#), [ATLAS '24](#), [CMS '17](#), [CMS '19](#), [LHCb '16](#), ...

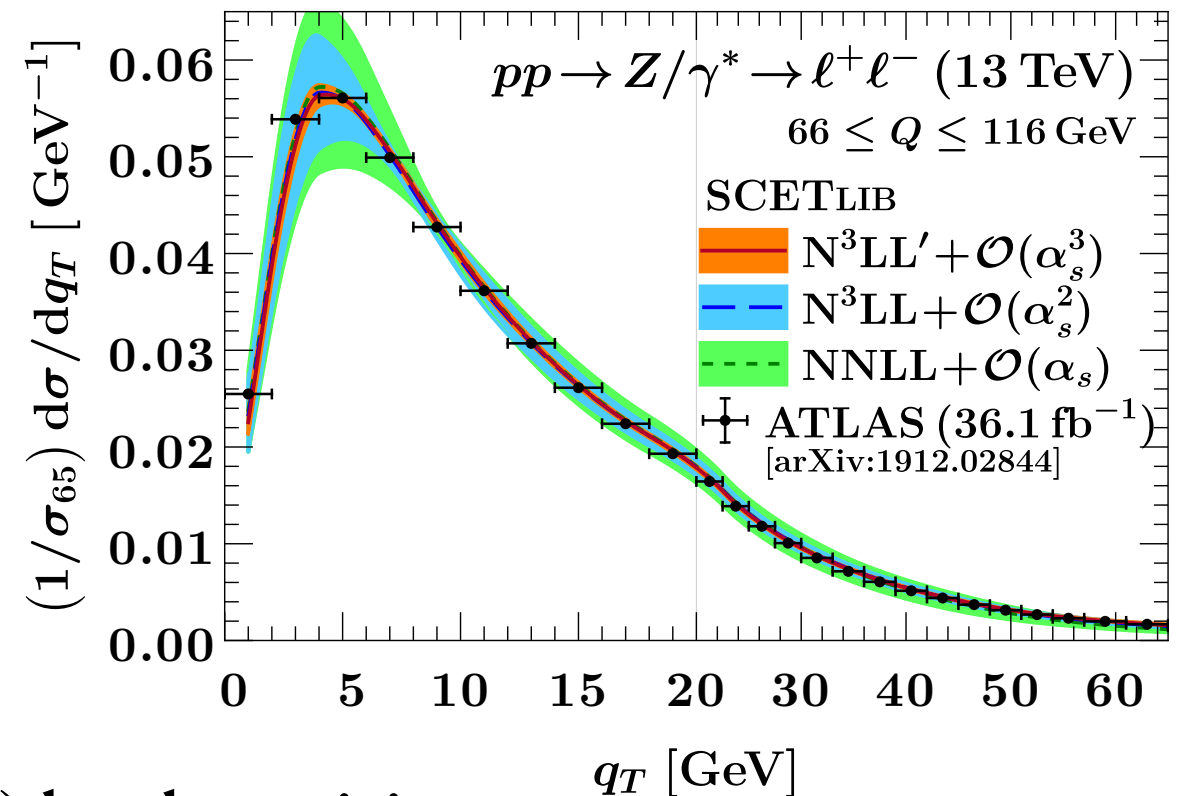
- » determination of the strong coupling α_s
- » W mass measurement
- » weak mixing angle
- » determination of PDFs at full N³LO
- » Higgs Yukawa couplings constraining to light quarks

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» Many theory requirements to reach $\mathcal{O}(1\%)$ level precision:

- | | | |
|---------------------------------|---|---|
| » resummation | $\mathcal{O}(\log^{2n}(p_T/m_Z)) \rightarrow \text{N}^3\text{LL}'/\text{approx N}^4\text{LL}$ | Billis, Michel, Tackmann '25 ,
Moos, Scimeni, Vladimirov, Zurita '24 ,
Camarda, Cieri, Ferrera '23 ,
... |
| » perturbative corrections | $\mathcal{O}(p_T^2/Q^2)$ | |
| » nonperturbative modeling | $\mathcal{O}(\Lambda_{\text{NP}}^2/p_T^2)$ | |
| » finite quark mass corrections | $\mathcal{O}(m_q^2/p_T^2)$ | |
| » electroweak corrections | $(\alpha_{\text{em}} \sim \alpha_s^2)$ | |

Perturbative uncertainty

Consider a series expansion in a small parameter α :

$$f(\alpha) = f_0 + \alpha f_1 + \alpha^2 f_2 + \alpha^3 f_3 + \alpha^4 f_4 + \mathcal{O}(\alpha^5)$$

$$\text{LO} : f(\alpha) = \hat{f}_0 \pm \Delta f$$

$$\text{NLO} : f(\alpha) = \hat{f}_0 + \alpha \hat{f}_1 \pm \Delta f$$

$$\text{NNLO} : f(\alpha) = \hat{f}_0 + \alpha \hat{f}_1 + \alpha^2 \hat{f}_2 \pm \Delta f$$

Δf is due to the series of the unknown true values $\hat{f}_n \rightarrow$ missing higher orders (MHOs)

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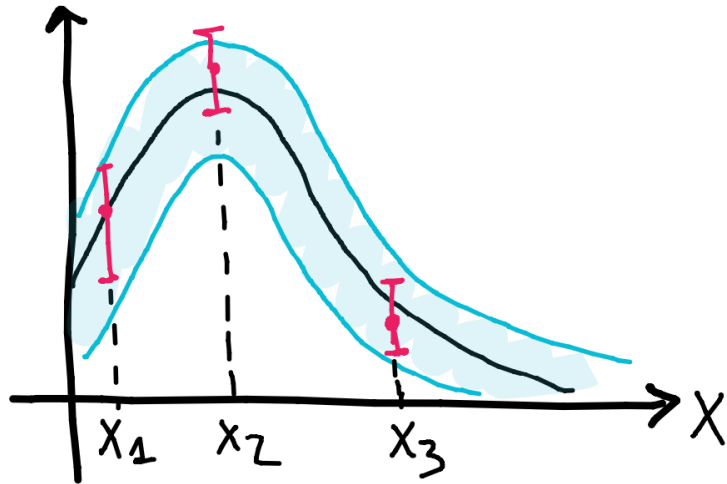
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Meaningful theory uncertainty:

- must reflect our degree of knowledge (or ignorance)
- provide correct **correlations** for different predictions
- have a statistical meaning needed for the interpretation of experimental measurements

Correlations and scale variations

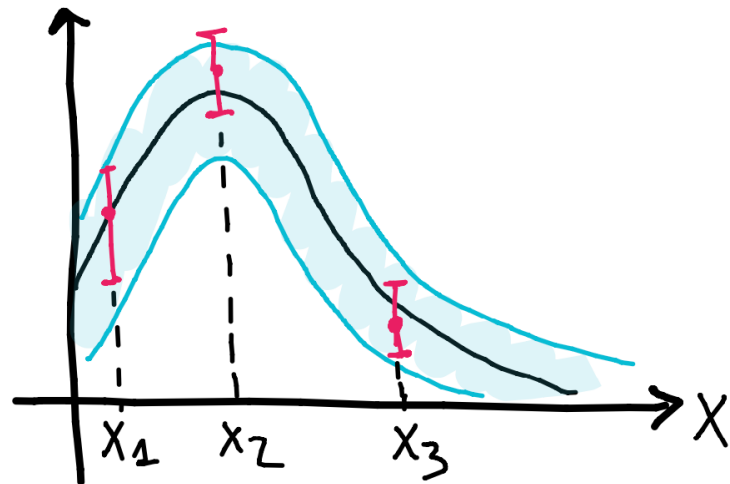
Taking a differential spectrum, each bin as separate predictions and separate measurements



- points close to each other are not intrinsically correlated, only their uncertainty is!

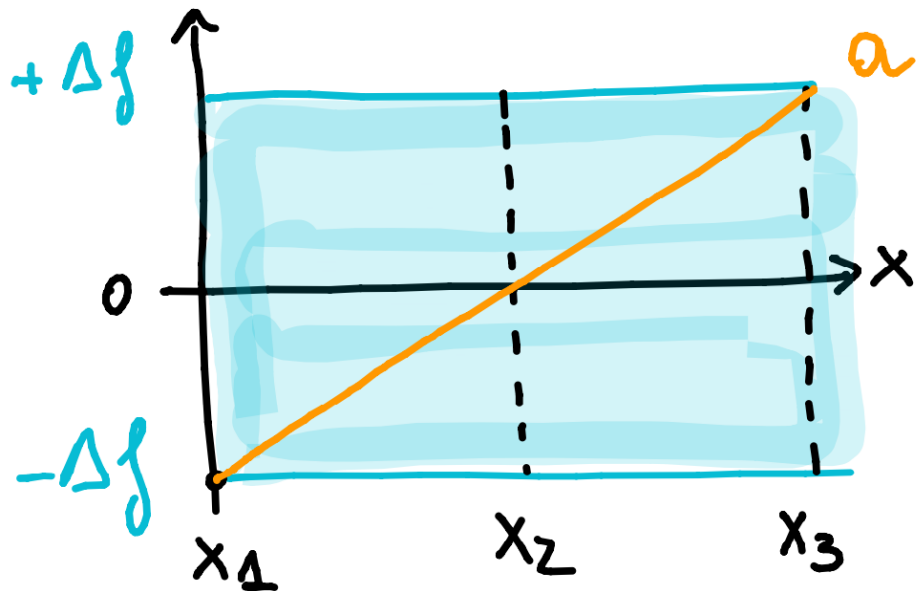
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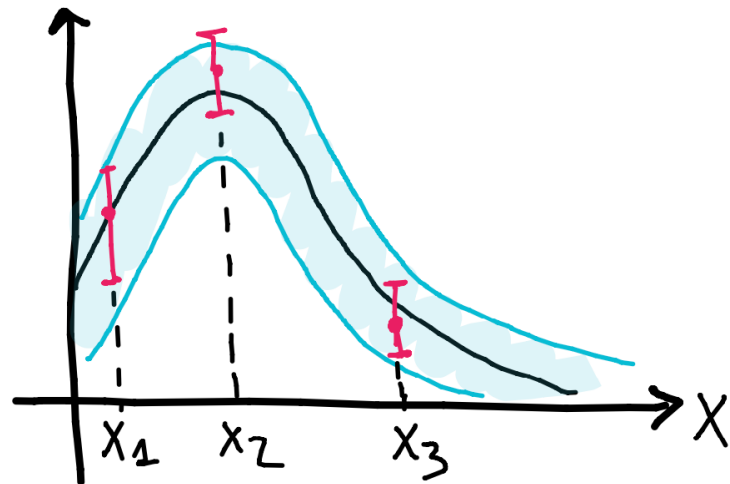
Let's be realistic: uncertainty band given by scale variations. What about its **shape**?



	ρ_{12}	ρ_{13}	ρ_{23}
a	0	-1	0

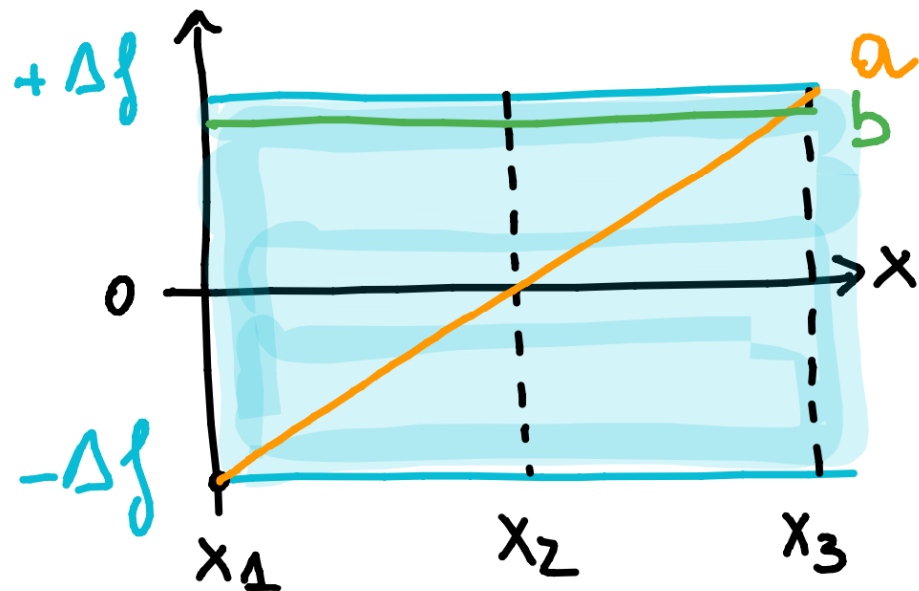
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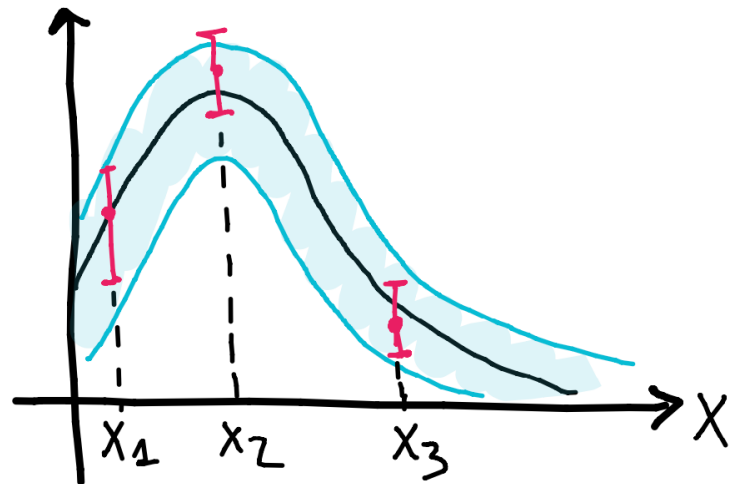
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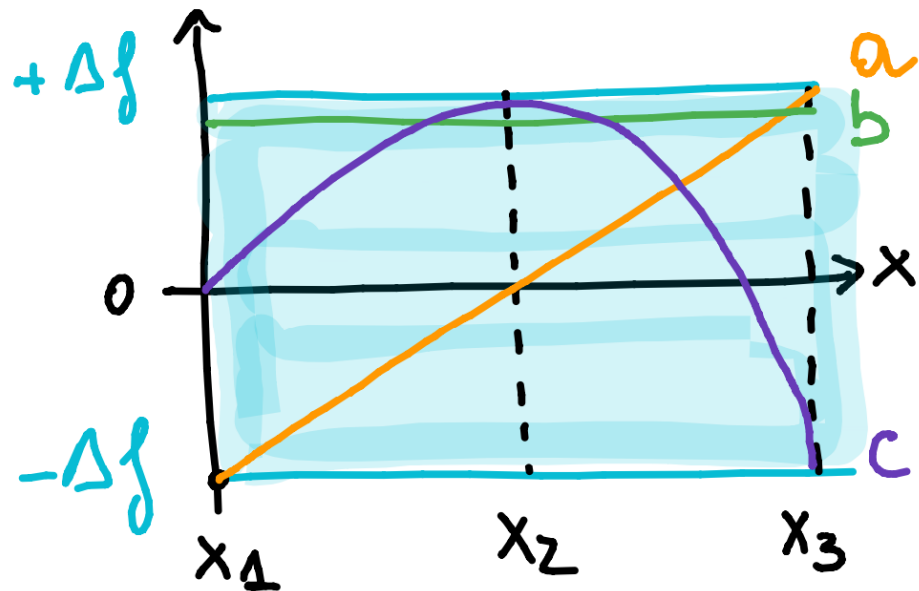
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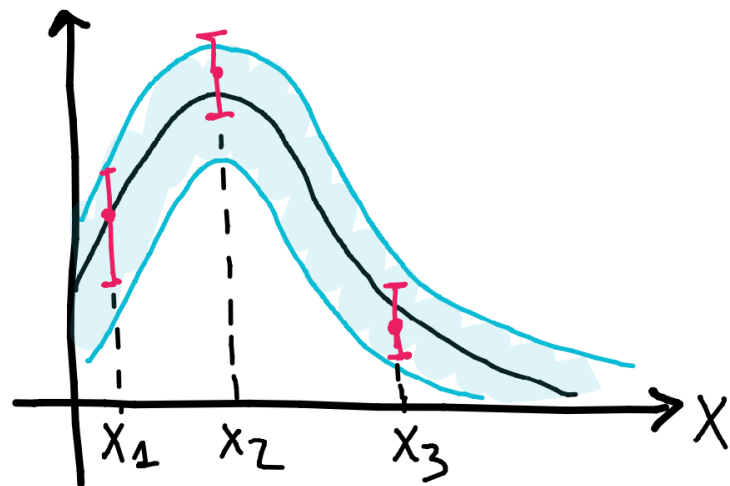


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every line (*a*, *b*, *c*) is a 100 % (anti-) correlated assumption

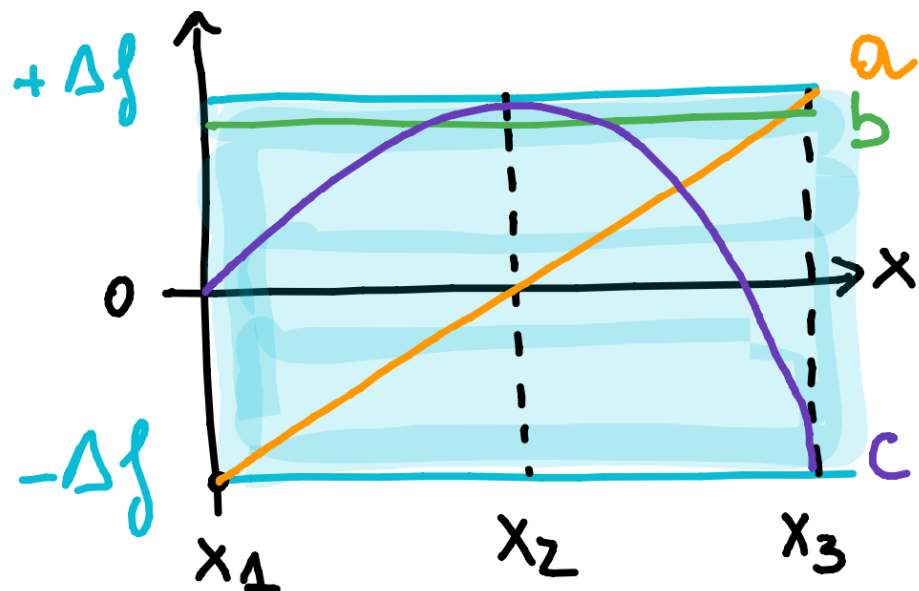
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➤ no idea about the correct shape of scale variations (and therefore correlation):

that's why we take **envelopes**!

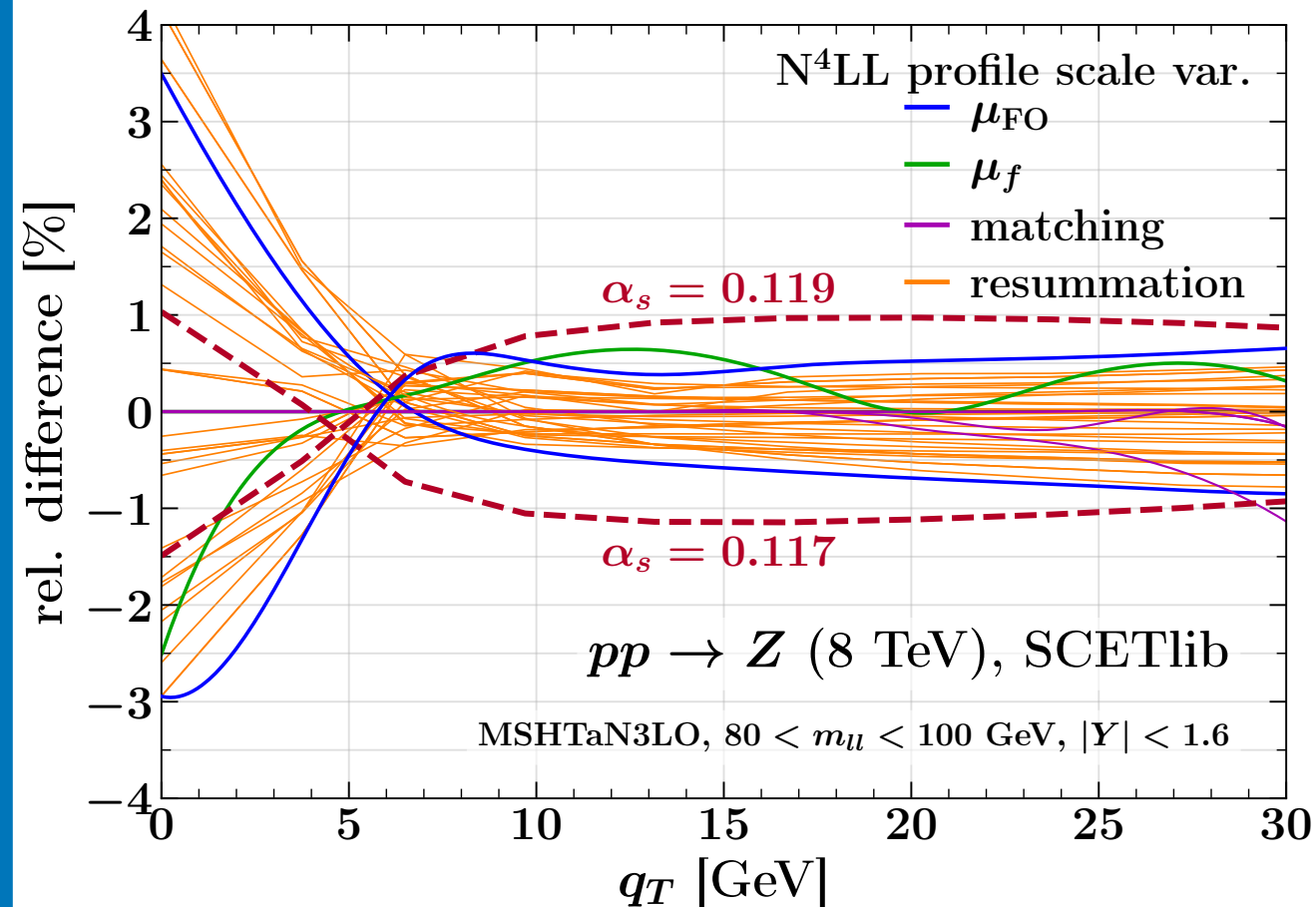
➤ to get correct correlation: breakdown into *independent uncertainty components* required

Extraction of $\Delta\alpha_s$ with scale variations

In the q_T spectrum each bin has its own theory prediction

» point-by-point correlation crucial for the determination of the α_s uncertainty

What are we used to do? Scale Variations!



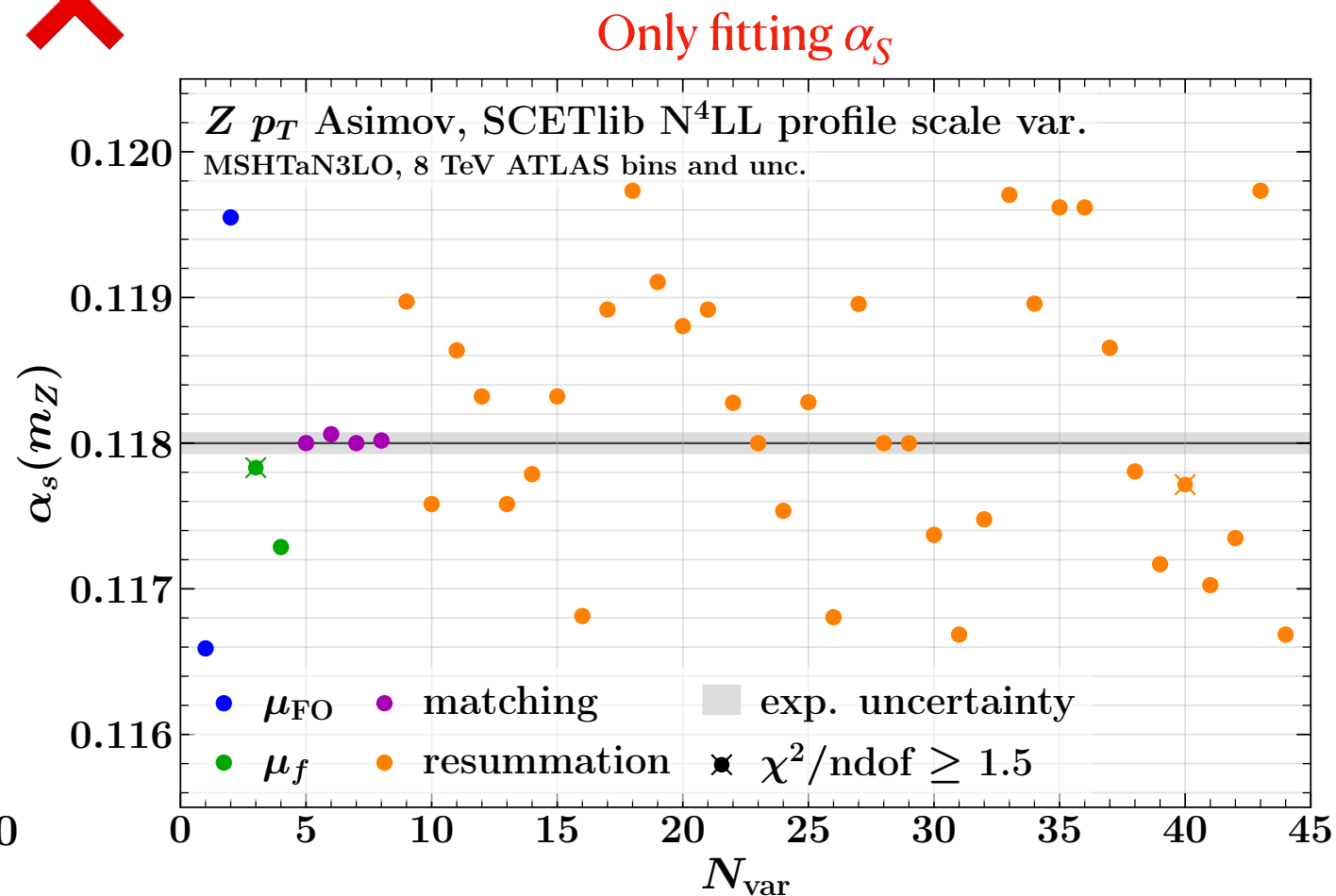
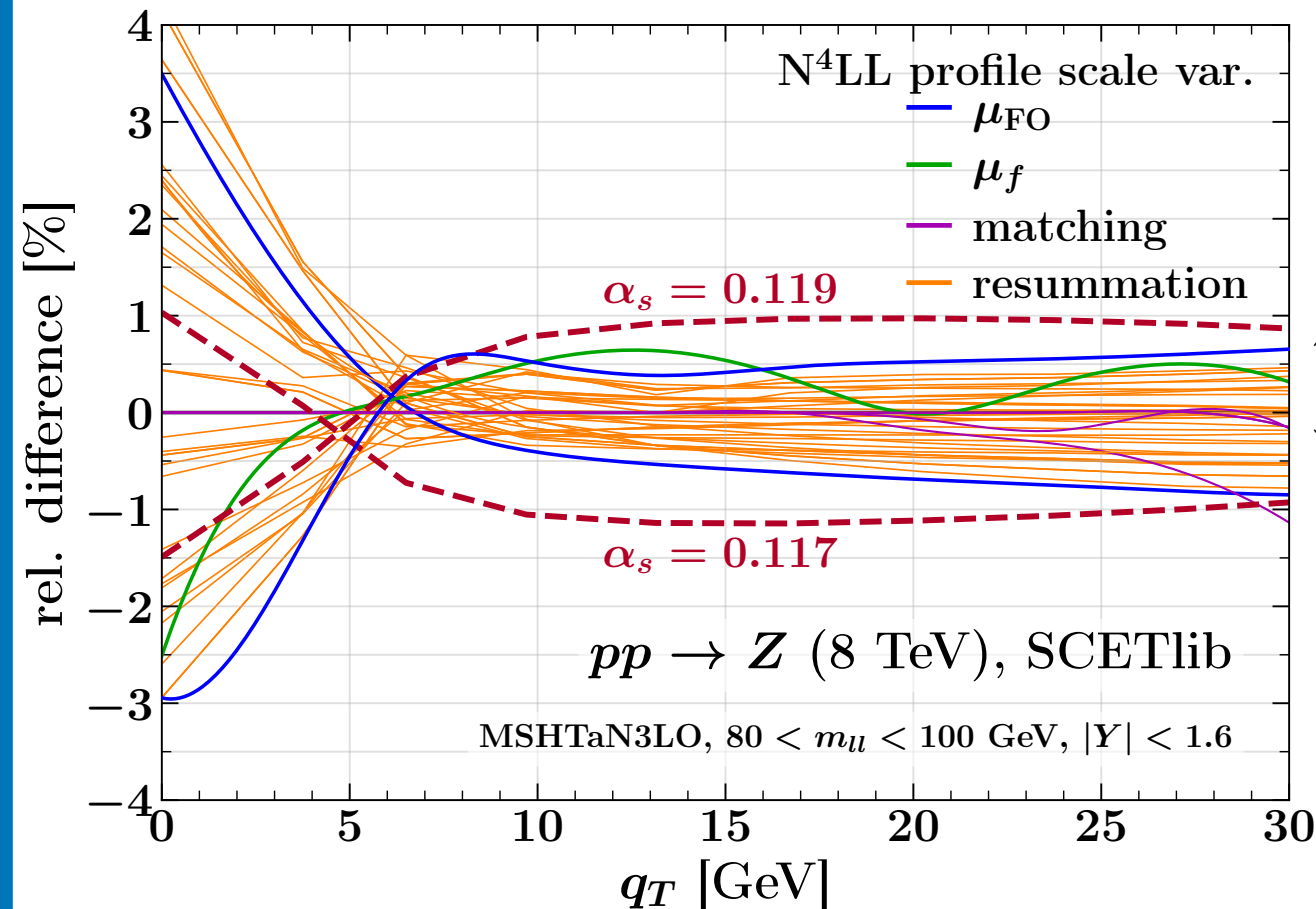
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Naive envelope: $\Delta_{scale} = 1.9$

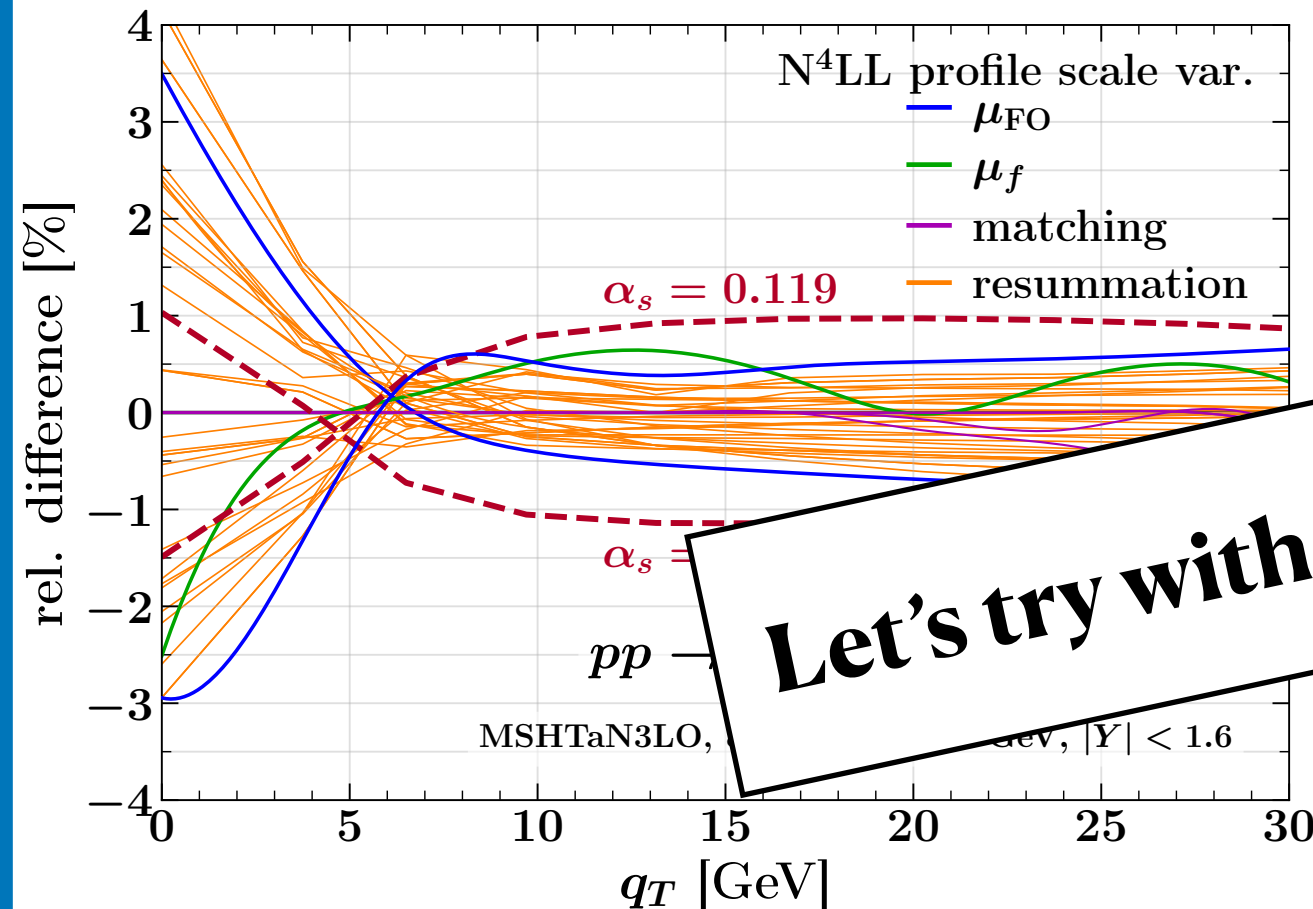
* uncertainties in units of 10^{-3}

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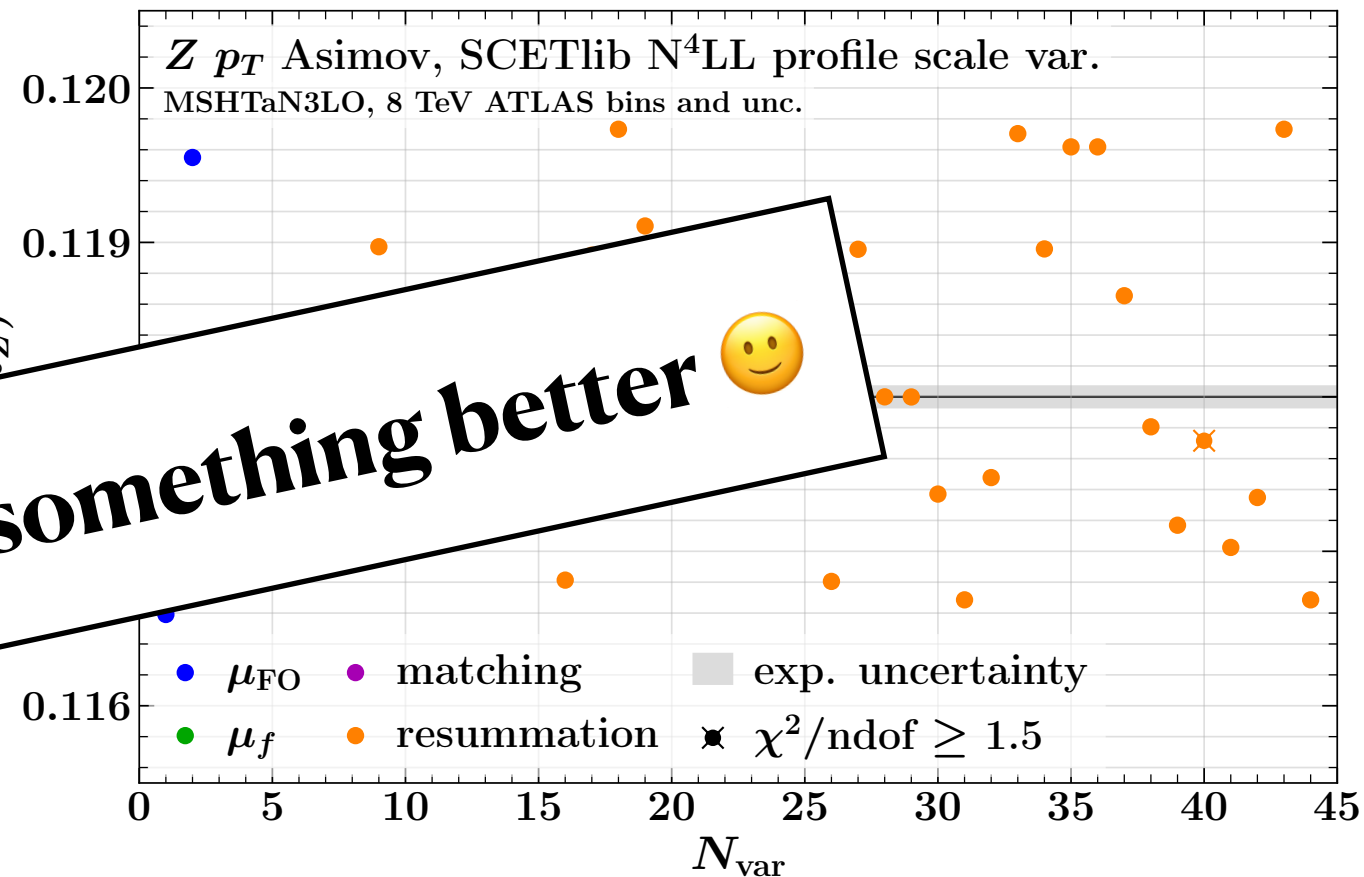
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What are we used to do? Scale Variations!



Only fitting α_s



Let's try with something better 😊

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Theory Nuisance Parameters (TNPs)

all details here [Tackmann '24!](#)

Consider the same series expansion:

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1 What is the **source of the uncertainty**?

$$\text{NNLO : } f(\alpha) = \hat{f}_0 + \alpha \hat{f}_1 + \alpha^2 \hat{f}_2 \pm \Delta f$$

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$$\text{NNLO} : f(\alpha) = \hat{f}_0 + \alpha \hat{f}_1 + \alpha^2 \hat{f}_2 \pm \Delta f$$

- 2 Parametrize and include the leading source of uncertainty:

$$\text{N}^{2+1}\text{LO} : f^{\text{pred}}(\alpha, \theta_3) = \hat{f}_0 + \alpha \hat{f}_1 + \alpha^2 \hat{f}_2 + \alpha^3 f_3(\theta_3)$$

using theory nuisance parameters θ_n ;

- θ_n have physical true value $\hat{\theta}_n$, such that $\hat{f}_n = f_n(\hat{\theta}_n)$
... and therefore encode correct theory correlations
- TNPs well-defined parameters with true but unknown value

Theory Nuisance Parameters (TNPs)

3 How to *define* these θ_n ?

all details here [Tackmann '24!](#)

- simplest case: $f_3(\theta_3) \equiv \theta_3$
- better: account for the internal structure of f_3
(given the process: partonic channels, color, ...)

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Consider the q_T spectrum, leading power q_T dependence is known to all orders:

$$q_T \frac{d\sigma}{dq_T} = \left[H \times B_a \otimes B_b \otimes S \right] (\alpha_S, L \equiv \ln q_T/m_Z) + \mathcal{O} \left(\frac{q_T^2}{m_Z^2} \right)$$

$F = \{H, B, S\}$ solution to RGE equations

$$F(\alpha_S, L) = \underbrace{F(\alpha_S)}_{\text{boundary conditions}} \exp \int_0^L dL' \left\{ \underbrace{\Gamma[\alpha_S(L')]}_{\text{anomalous dimensions}} L' + \underbrace{\gamma_F[\alpha_S(L')]}_{\text{anomalous dimensions}} \right\}$$

Theory Nuisance Parameters (TNPs)

3 Account for dependencies:

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➤ in which we need correlations

➤ those helping to obtain better theory constraints

$$F(\alpha_S) = 1 + \sum_{n=1} \left(\frac{\alpha_S}{4\pi} \right)^n F_n$$

$$\gamma(\alpha_S) = \sum_{n=0} \left(\frac{\alpha_S}{4\pi} \right)^{n+1} \gamma_n$$

$$F_n(\theta_n^f) = 4C_r(4C_A)^{n-1}(n-1)! \theta_n^f$$

$$\gamma_n(\theta_n^\gamma) = 4C_r(4C_A)^n \theta_n^\gamma$$

C_r leading color factor,
 C_A^{n-1} leading n -loop color factor

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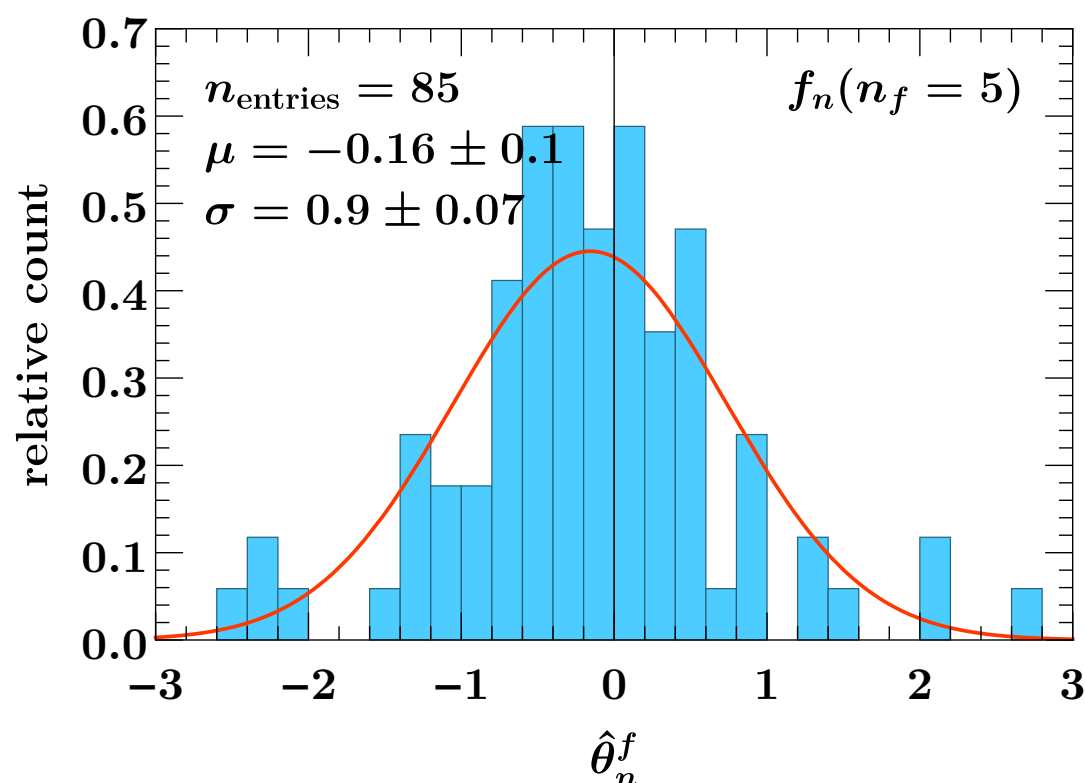
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4 How to *vary* θ_n ?



With these normalizations,
 expected natural size $|\hat{\theta}_n| \lesssim 1$

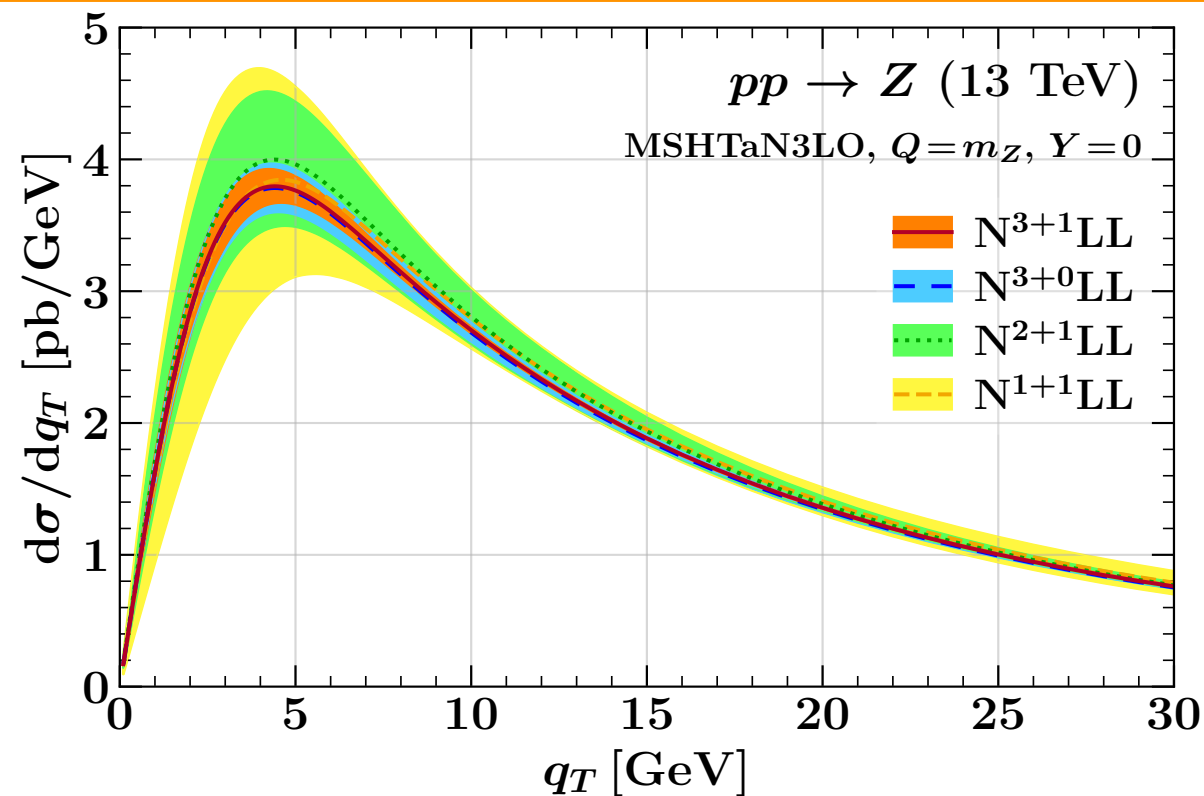
$$\theta_n = 0 \pm 1$$

look at other known n -loop coefficients from
 population sample [here](#),
 validated using known perturbative series

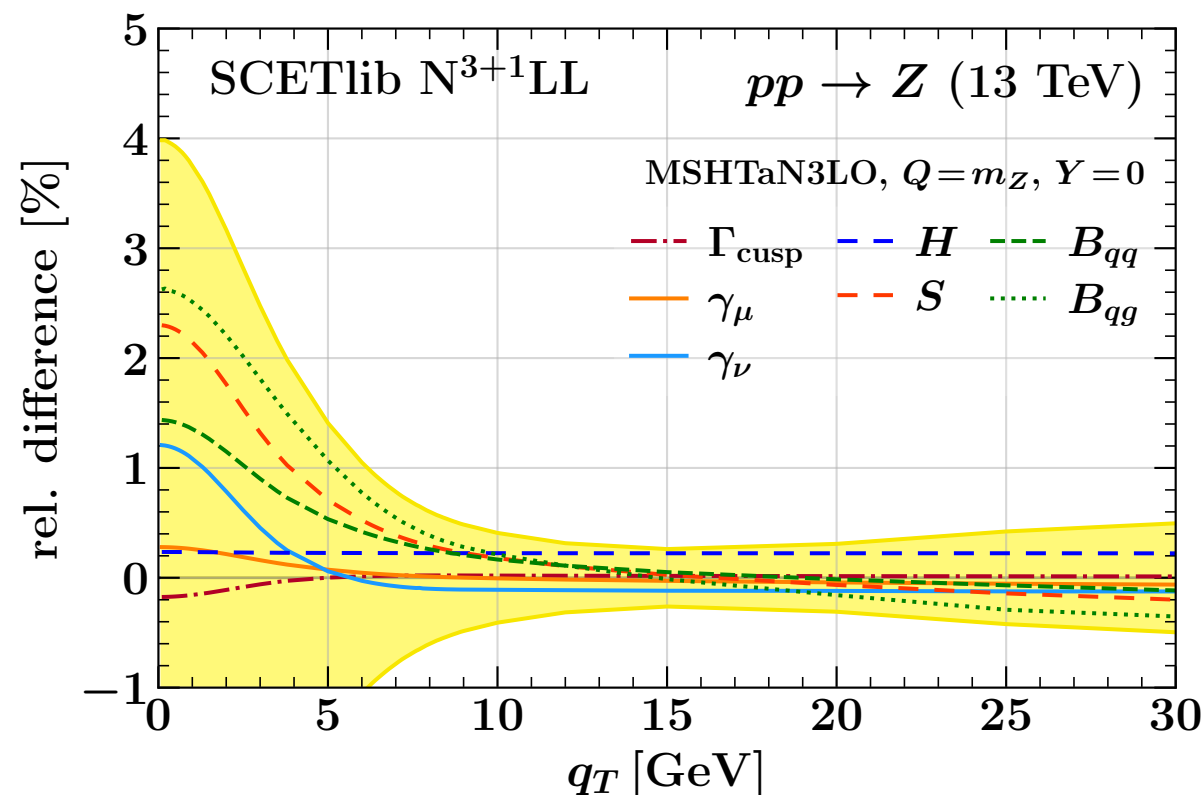
Uncertainty breakdown and correlations

all details here [Tackmann '24!](#)

Comparing different orders at 95% theory CL
($\Delta\theta_n = \pm 2$)



Breakdown of all the TNPs at $N^{3+1}LL$:



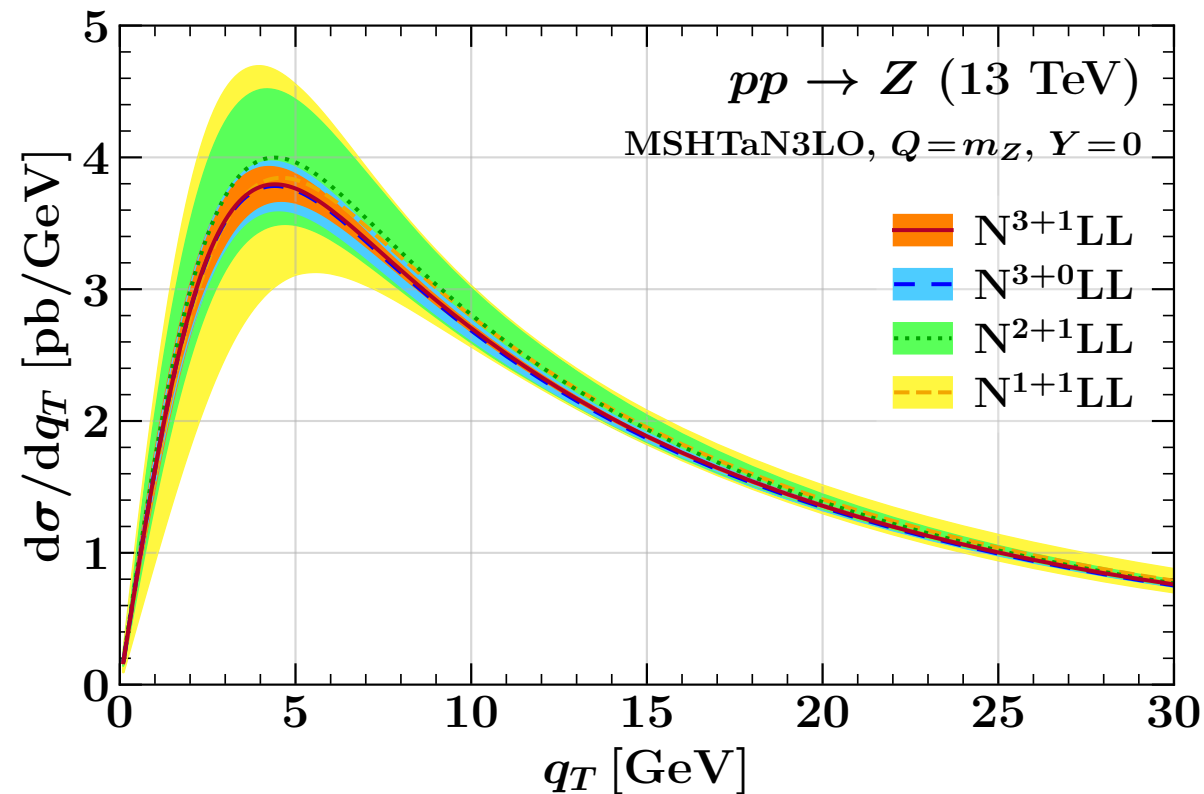
- varying each TNP by $\Delta\theta_n = \pm 1$ (68% CL)
- providing breakdown into independent sources of uncertainty
- encoding correct point-by-point correlations

* $B_{qj} : F_n(z, \theta_n) = 3/2 \theta_n \hat{F}_n(z)$, DGLAP splitting functions not varied

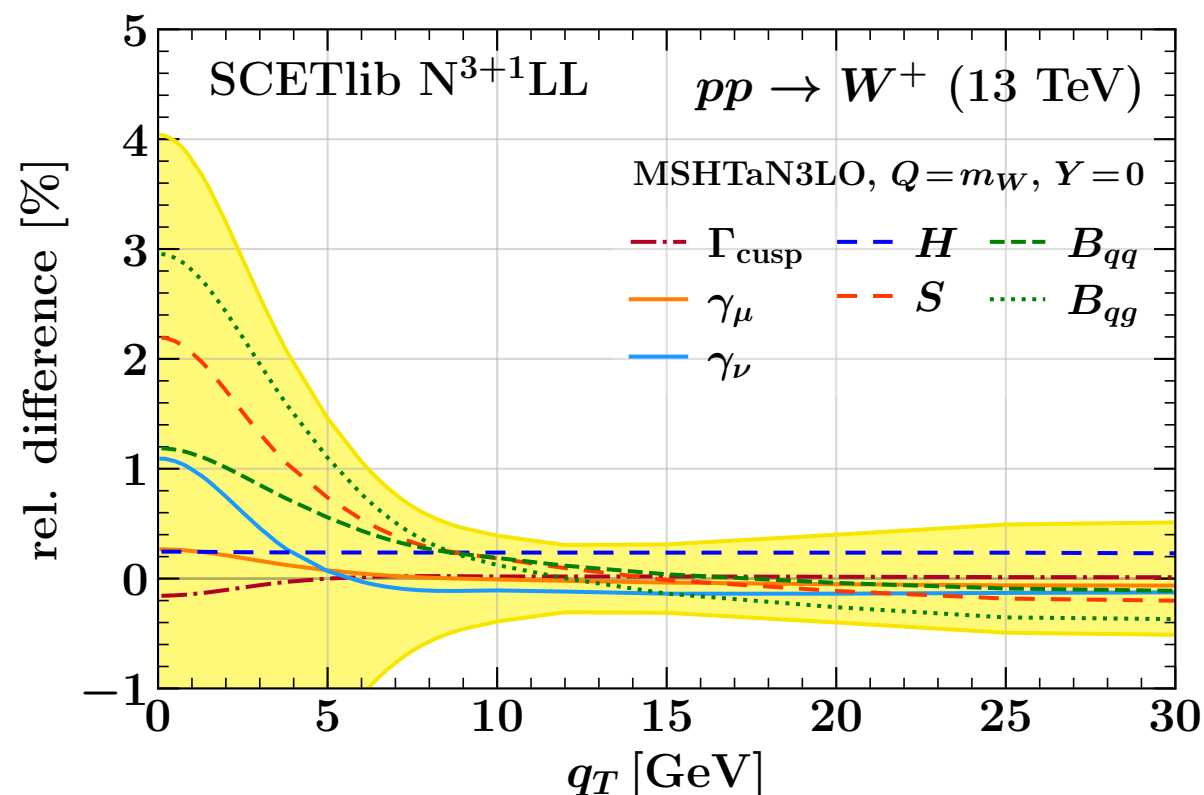
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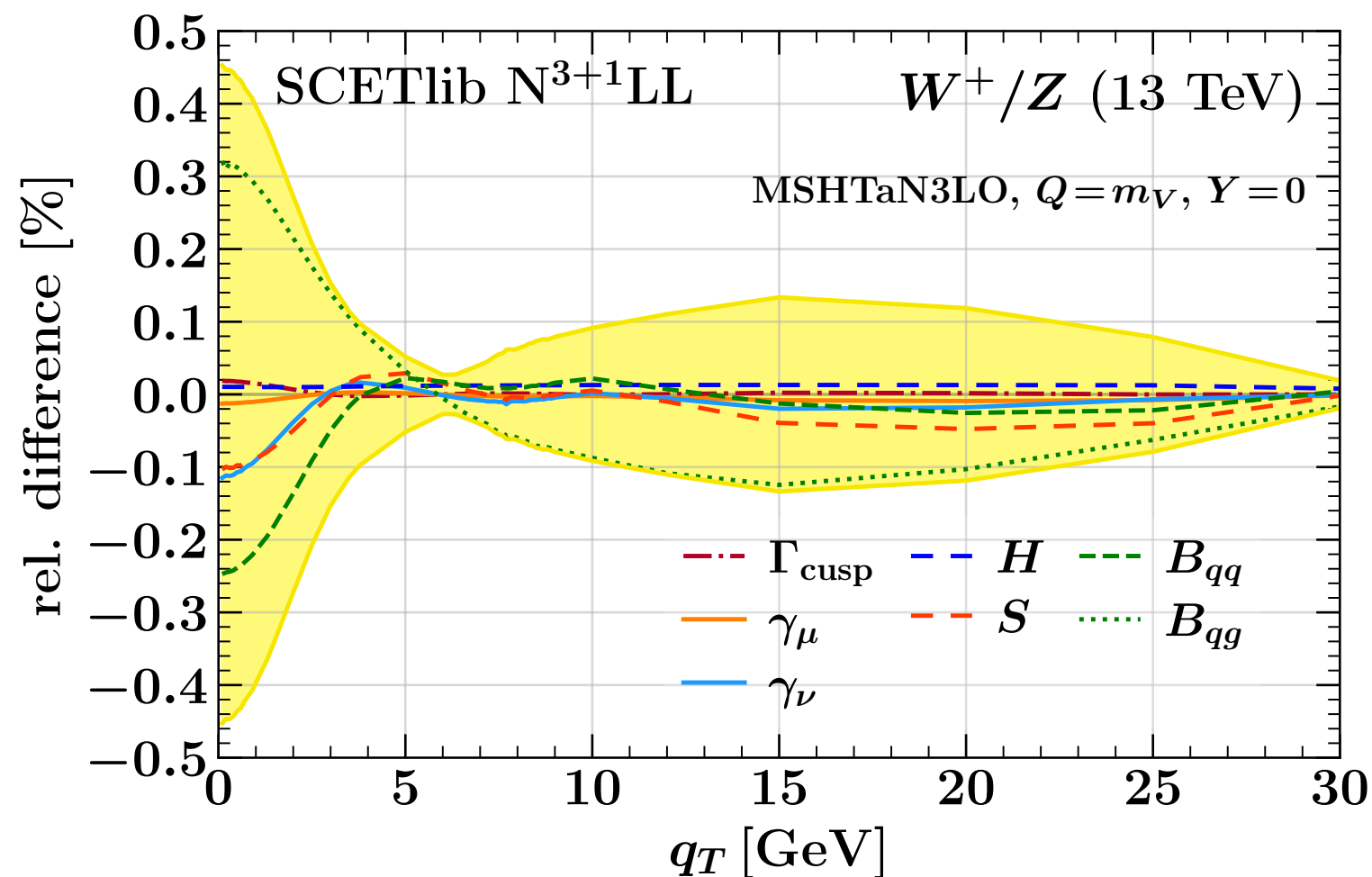
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Theory Nuisance Parameters (TNPs)

Relative impacts on W/Z :

all details here [Tackmann '24!](#)



- uncertainties very similar for Z and W processes: same TNPs for both
 - ➔ each TNP impacts are 100% correlated between the processes:
nice cancellation in the ratio!

*just for illustration: only leading **massless** contribution

Some applications

Asimov fits for $\alpha_s(m_Z)$ from $Z p_T$ spectrum

WIP Cridge, Marinelli, Tackmann

Asimov fits: standard procedure to estimate expected uncertainties in a fully controlled setting

- » study the *dominant* sources of uncertainty and their impact on the extracted α_s
not concerned with subleading effects: neglected both in pseudodata and theory model
 - still necessary for fitting real data

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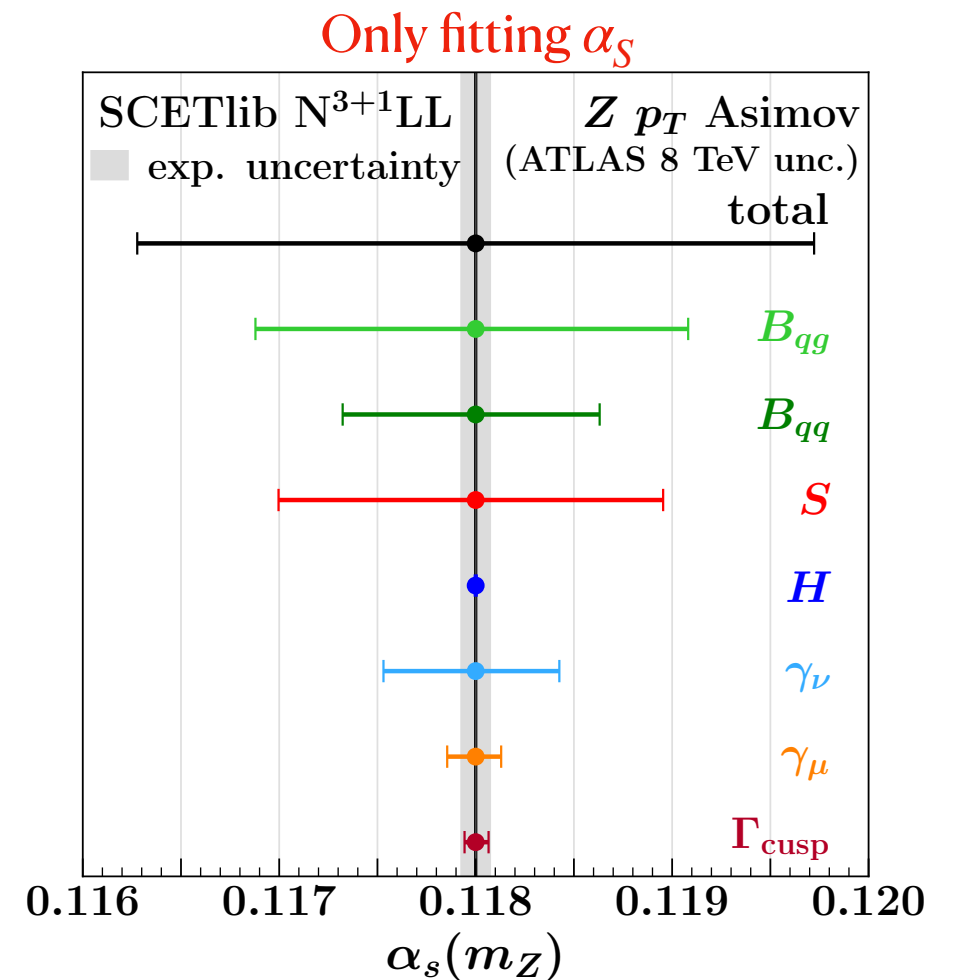
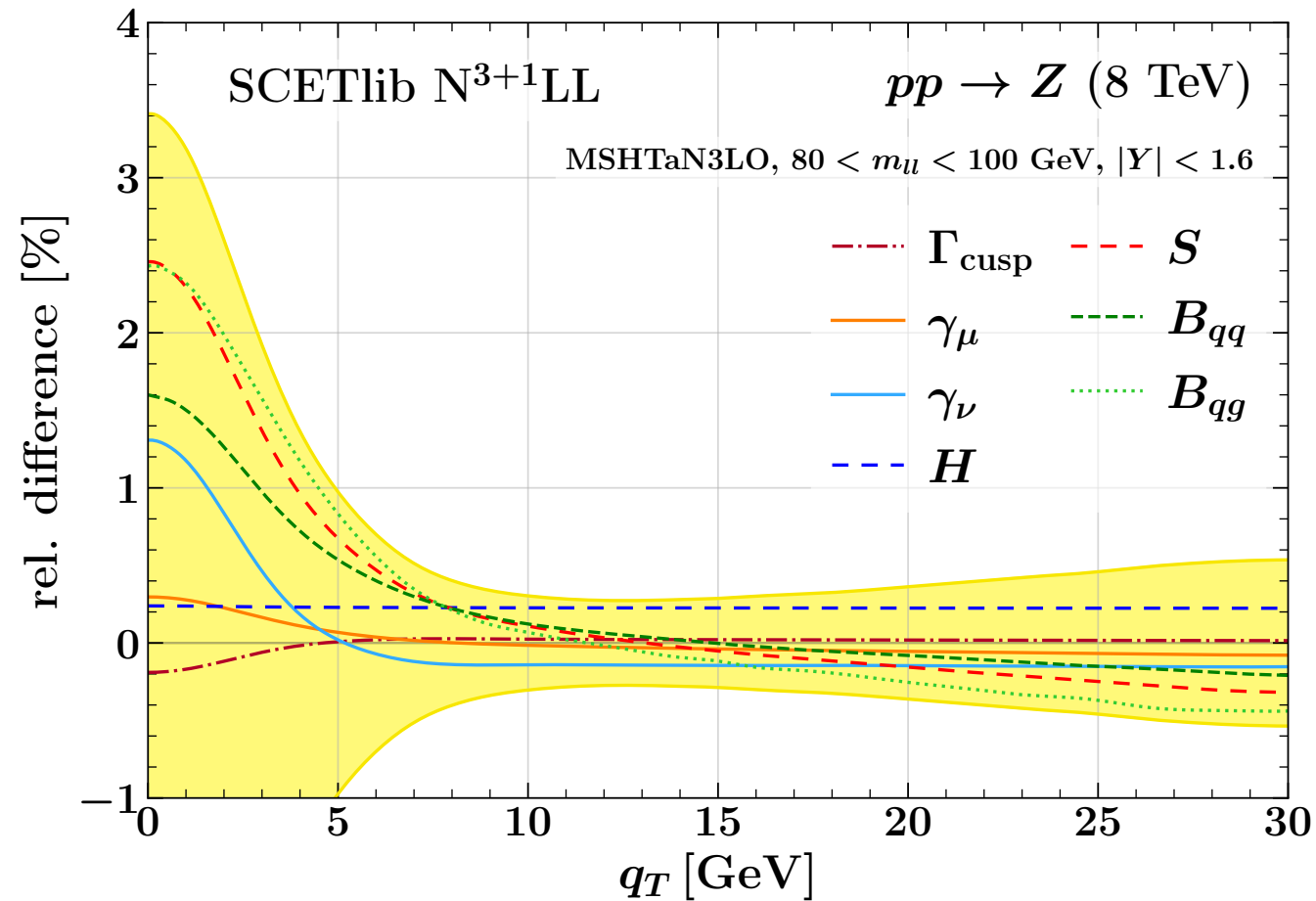
Our theory inputs:

- » SCETlib N^{3+1} LL and N^4 LL only resummed contribution

Our toy data:

- » Data defined as central theory prediction [$\alpha_s = 0.118$]
[fixed nonp. params, MSHT20aN3LO PDF set]
- » 72 data points in ATLAS binning,
9 q_T bins in $[0, 29]$ GeV for each 8 Y bin in $[0.0, 3.6]$
[integrated in q_T , Y and Q]
- » Using ATLAS exp. uncertainties and complete correlations for all 72 bins [arXiv:2309.12986](https://arxiv.org/abs/2309.12986)

Scanning TNPs



Repeat fit separately varying each TNP by $\Delta\theta_n = \pm 1$

- providing breakdown into independent sources of uncertainty
- encoding correct point-by-point correlations
- can now sum in quadrature $\Delta_{\text{total}} = 1.7$

still does not let the fit decide between moving α_s or theory

rather profiling

Perturbative uncertainty with profiling TNPs

Profiling: fitting α_s *together* with all TNPs (allows the fit to decide what to do)

- TNPs are proper parameters, included in the fit with Gaussian constraint $\theta_n = 0 \pm 1$
- allows data to constrain TNPs and thereby reduce theory uncertainty

Pseudodata: central $[\alpha_s = 0.118]$ N⁴LL prediction

- simulates the fit to real data, which contains the all-order result

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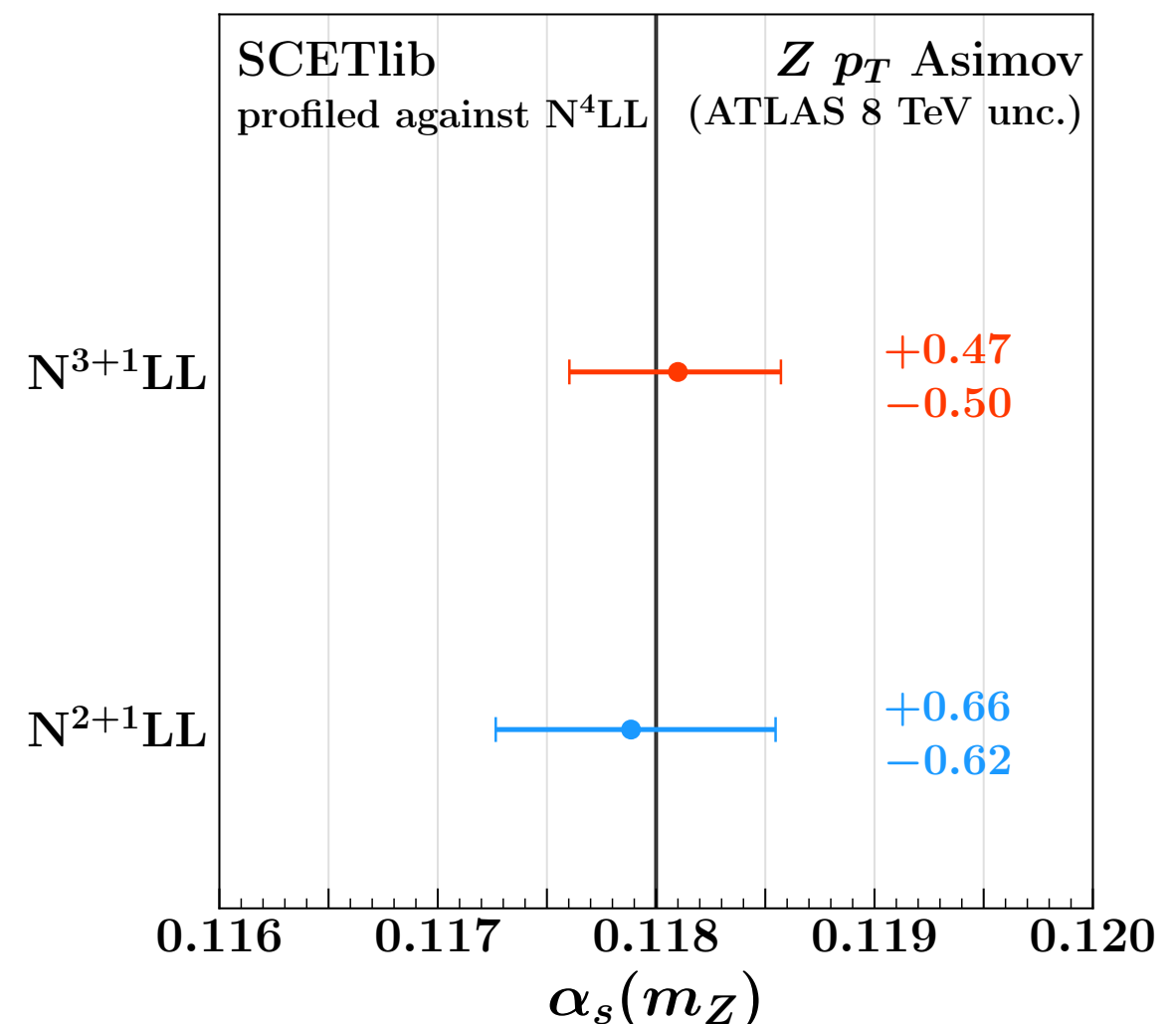
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1 N²⁺¹LL theory model

2 N³⁺¹LL theory model

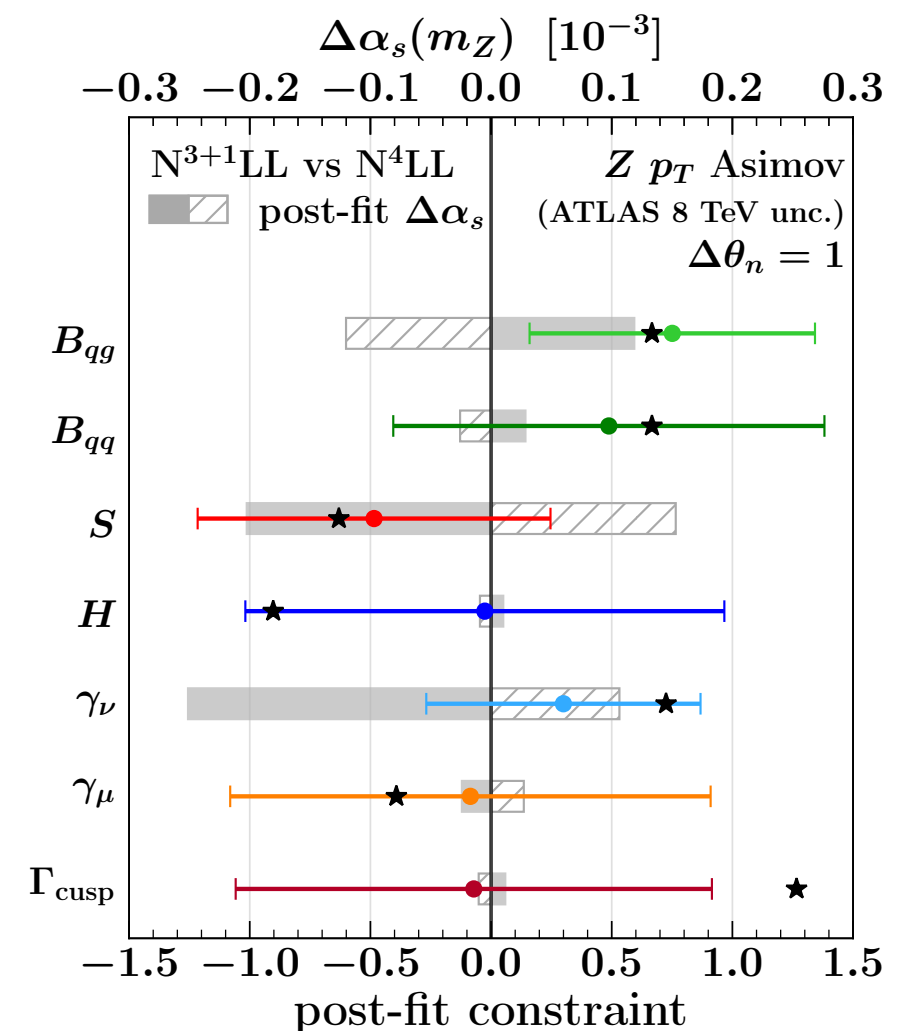
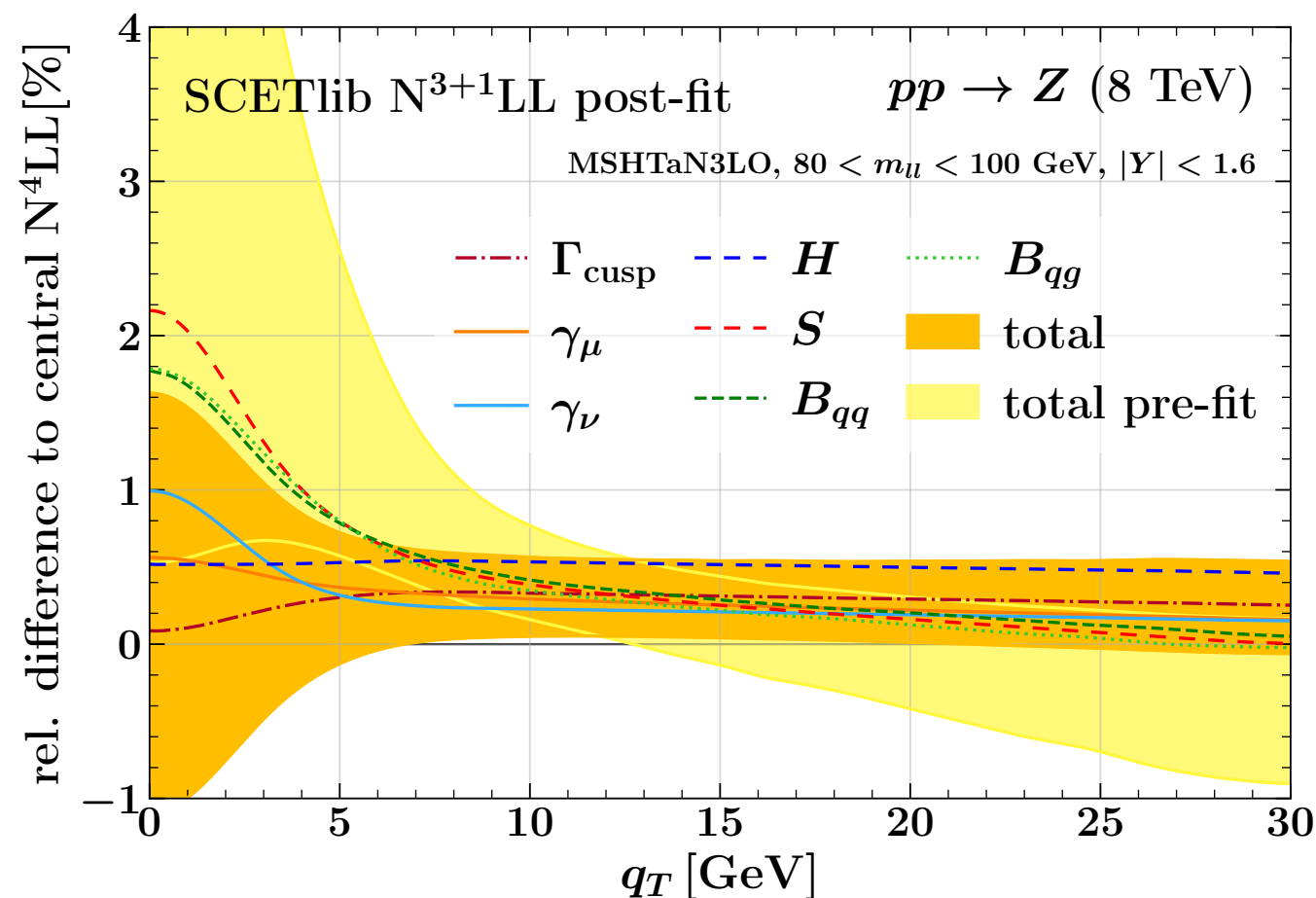
➔ look at the post-fit constraints on TNPs



* uncertainties in units of 10^{-3}

Post-fit constraints on TNPs

Profiling lower order against higher order: $N^{3+1}LL$



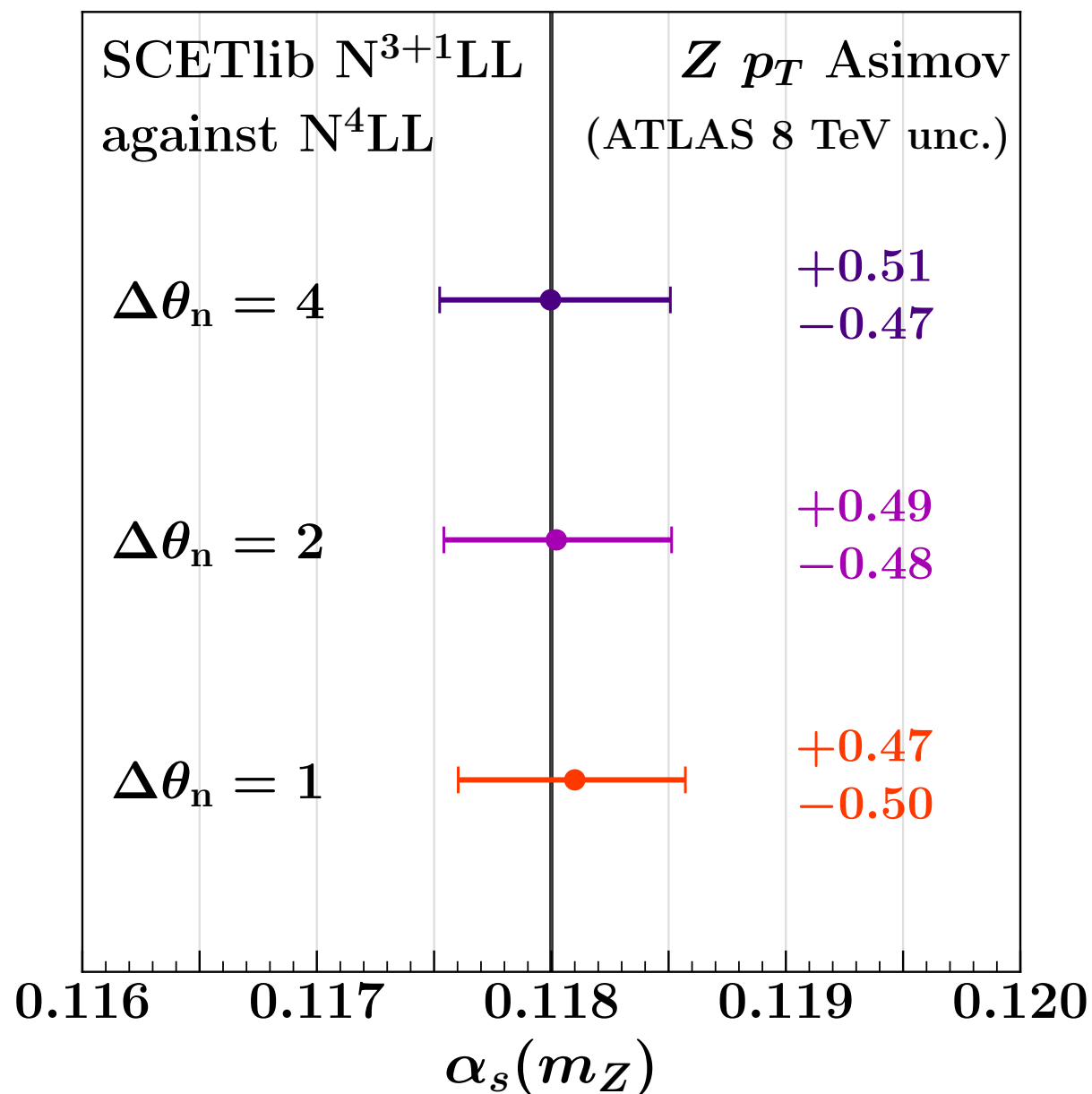
- $N^{3+1}LL$ pulled, toward correct true values [★]
- post-fit prediction for q_T spectrum driven by constraints from data
- grey → TNP down variation, dashed grey → TNP up variation
- γ_ν , S and B_{qg} have the largest remaining impact on $\alpha_s(m_Z)$ after profiling
- ➔ for the exact correlation between parameters, look at the post-fit covariance matrix!

Different theory constraints on TNPs

What happens by changing the prior theory constraint?

Using now $\theta_n = 0 \pm \Delta\theta_n$ with $\Delta\theta_n = 1, 2, 4$

Fit N^{3+1} LL against N^4 LL data



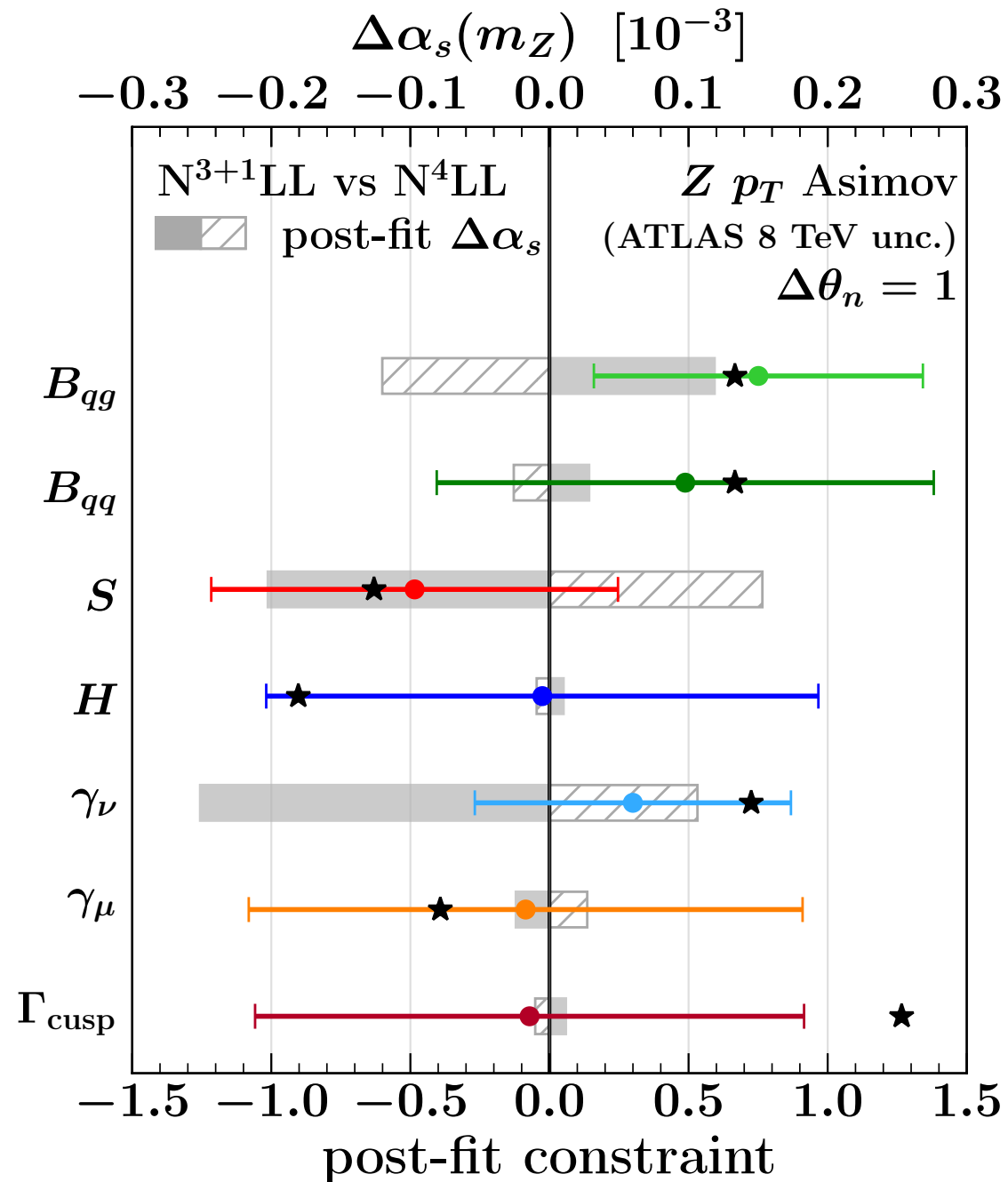
- profiling substantially reduces the dependence on theory constraint (with scanning, α_s unc. directly depends on choice of $\Delta\theta_n$)
- the effect relative to the theory constraint strongly depends on the power of the experimental constraint
- only slight difference in the uncertainties when relaxing the TNP constraint

Different theory constraints on TNPs

Using now $\theta_n = 0 \pm \Delta\theta_n$ with $\Delta\theta_n = 1$

Fit N^{3+1} LL against N^4 LL data

1 $\Delta\theta_n = 1$ start seeing the exp. constraint

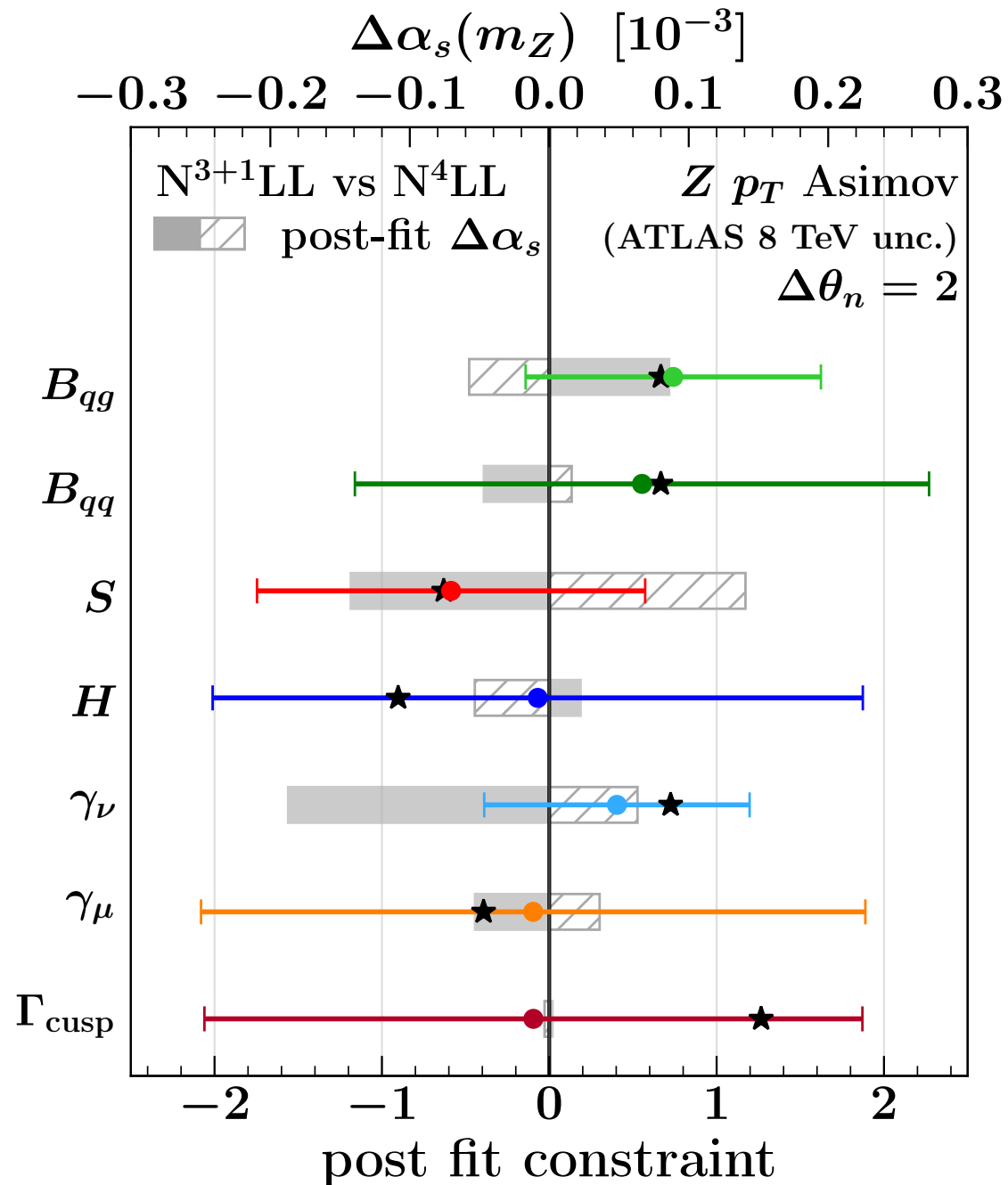


[don't be fooled by the different x-range!]

Different theory constraints on TNPs

Using now $\theta_n = 0 \pm \Delta\theta_n$ with $\Delta\theta_n = 2$

Fit N^{3+1} LL against N^4 LL data



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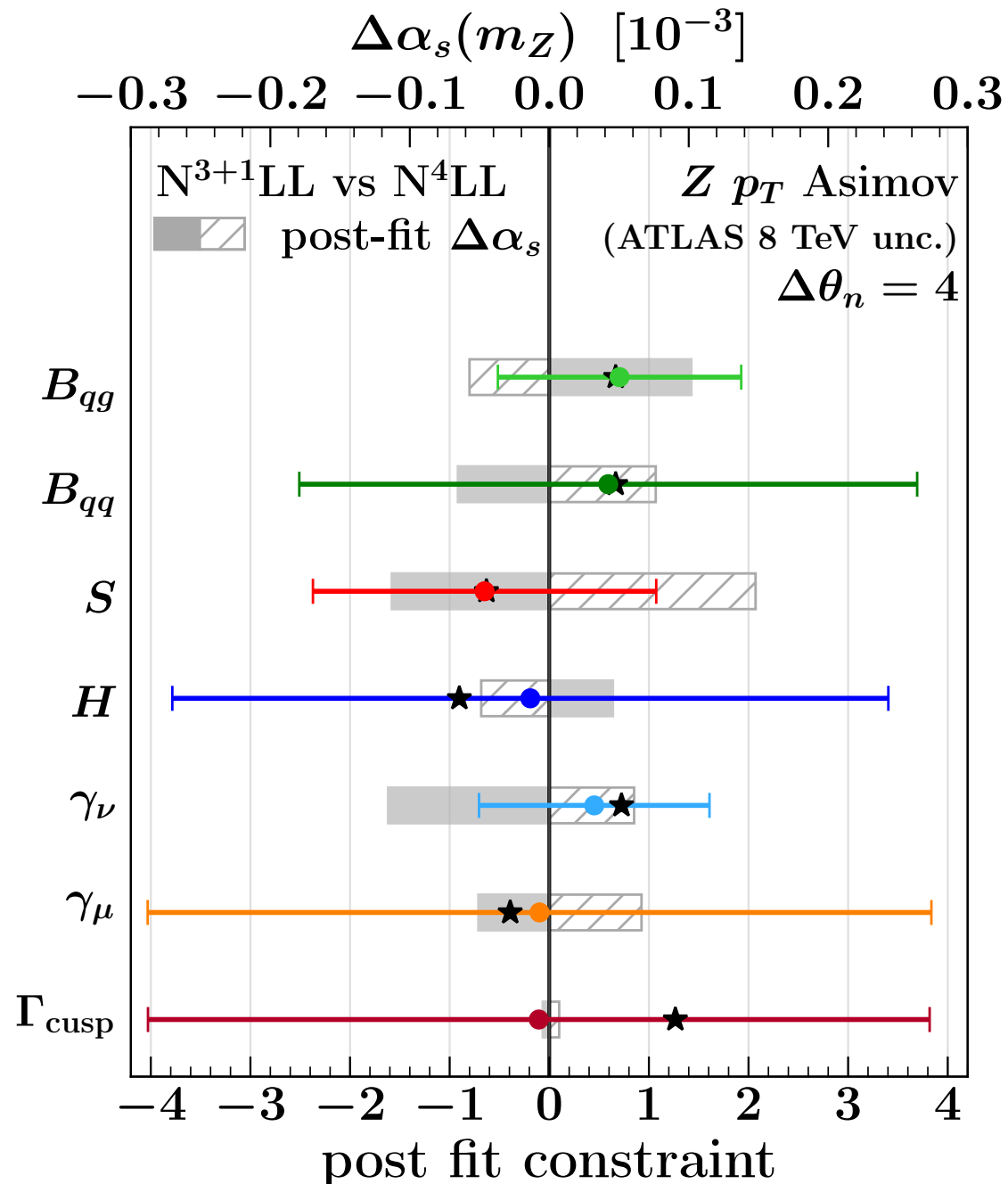
2 $\Delta\theta_n = 2$ it's basically a factor
2 w.r.t $\Delta\theta_n = 1$

[don't be fooled by the different x-range!]

Different theory constraints on TNPs

Using now $\theta_n = 0 \pm \Delta\theta_n$ with $\Delta\theta_n = 4$

Fit N^{3+1} LL against N^4 LL data



- 1 $\Delta\theta_n = 1$ start seeing the exp. constraint
- 2 $\Delta\theta_n = 2$ it's basically a factor 2 w.r.t $\Delta\theta_n = 1$
- 3 $\Delta\theta_n = 4$ data can constrain TNPs more
- 4 remaining impact on $\alpha_s(m_Z)$ after profiling proportional to $\Delta\theta_n$

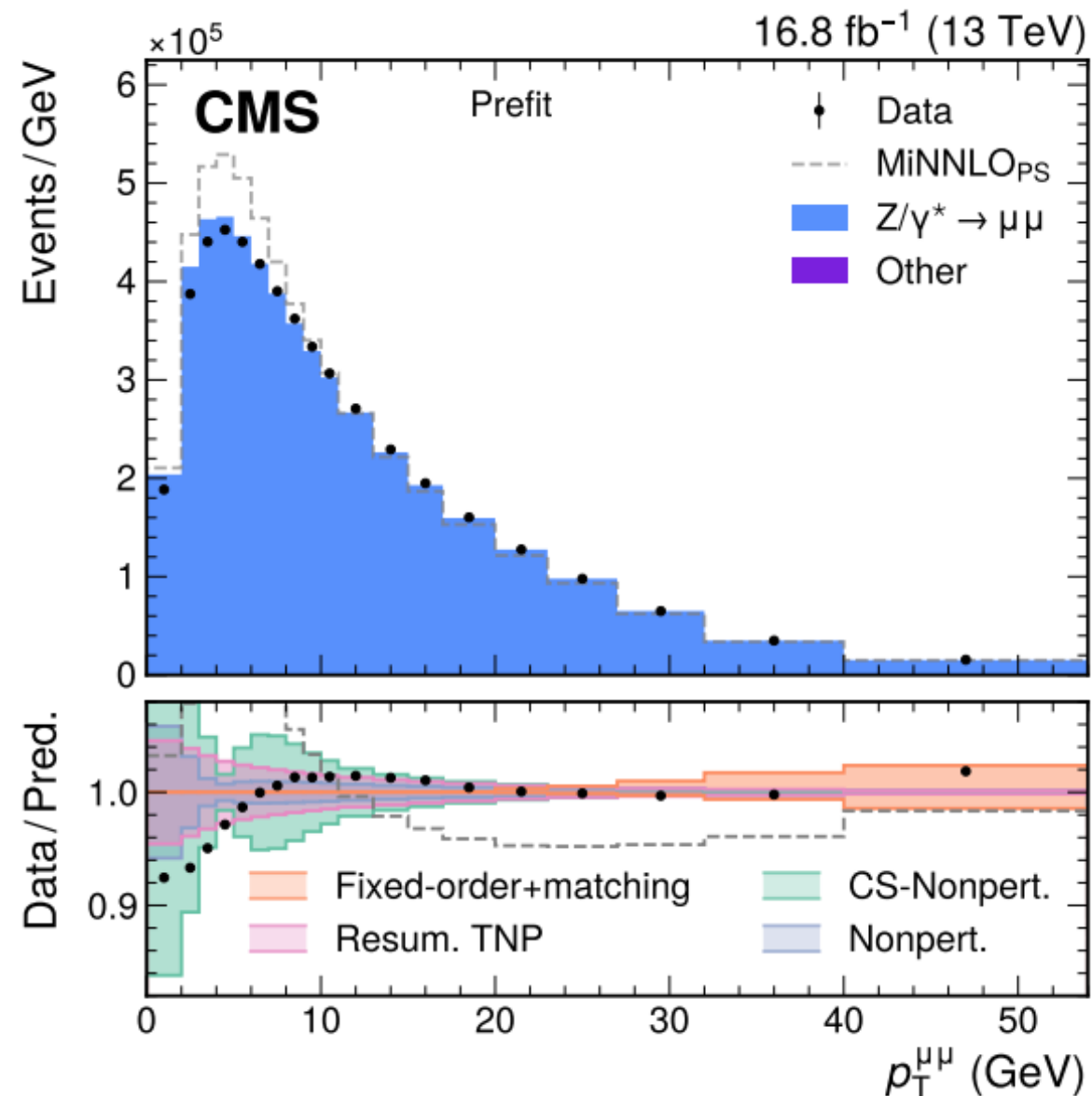
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CMS W mass measurement

See Chiara's talk!

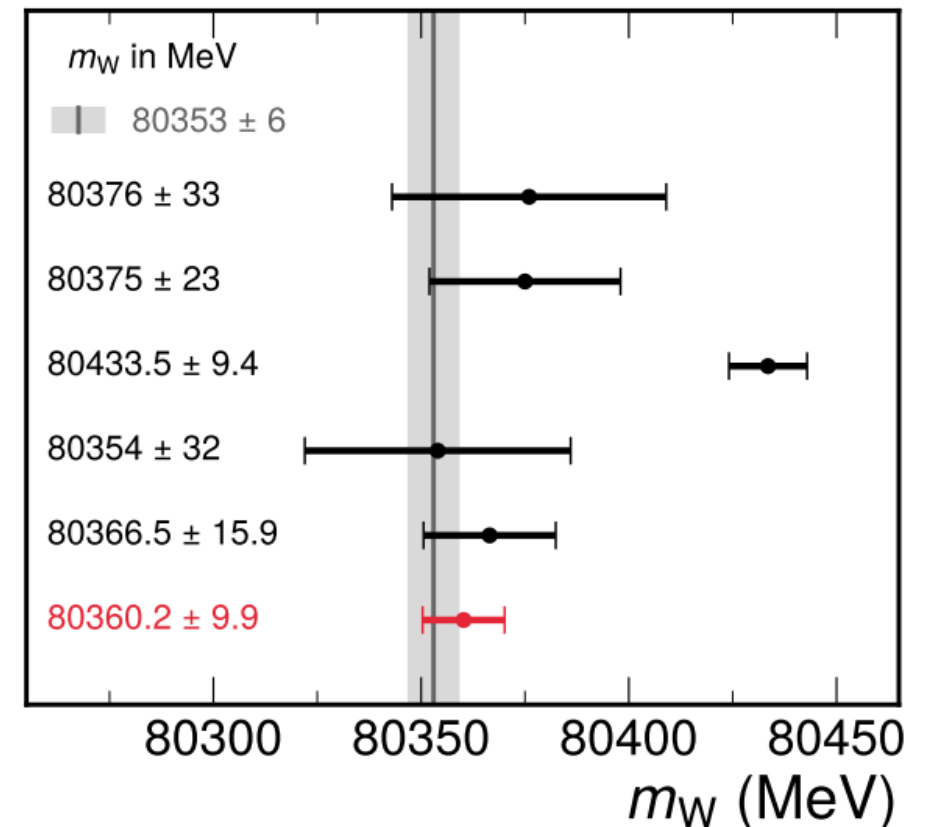
Recent CMS W mass measurement [arXiv:2412.13872](https://arxiv.org/abs/2412.13872)

Theory input: $N^{3+0}LL+NNLO$
(SCETlib and DYturbo)



Electroweak fit
PRD 110 (2024) 030001
LEP combination
Phys. Rep. 532 (2013) 119
D0
PRL 108 (2012) 151804
CDF
Science 376 (2022) 6589
LHCb
JHEP 01 (2022) 036
ATLAS
arXiv:2403.15085
CMS
This work

CMS



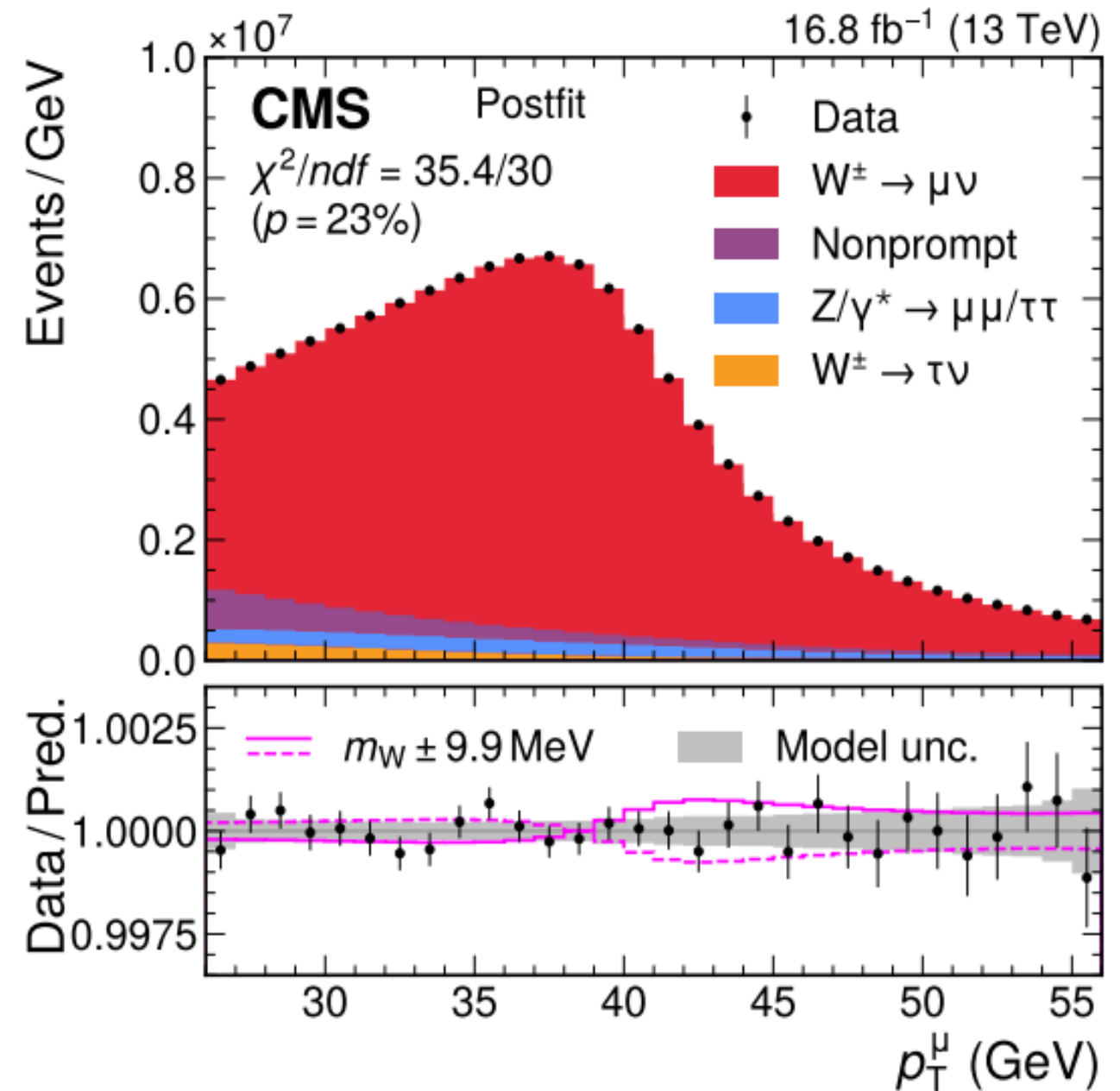
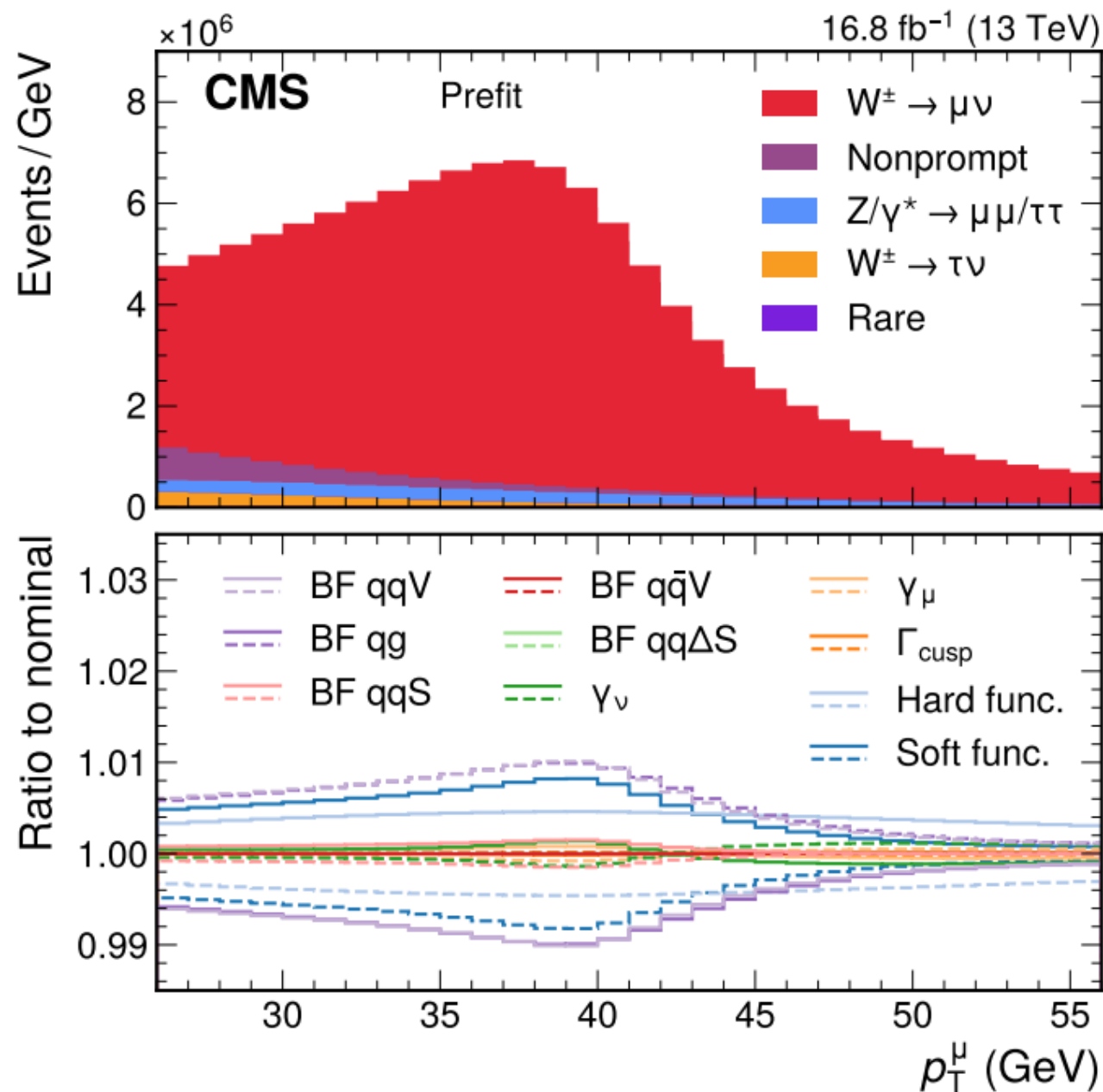
- p_T^W modeling fundamental: uncertainties in the low p_T region affect the shape as m_W variation
- theory correlations are crucial: uncertainty propagated from p_T^W to p_T^μ to m_W !

CMS W mass measurement

See Chiara's talk!

Perturbative uncertainties in the resummed prediction: $N^{3+0}LL$ SCETlib*

➤ contribution of all theoretical and experimental uncert. before and after profiling



* for details about this order look [here](#)

Summary

p_T Z/W crucial benchmark observables for LHC precision physics program
→ that's why we need meaningful theory uncertainties!

Correlations are fundamental for interpretation of precision measurements:
having meaningful theory uncert. is as important as meaningful exp. uncert.!

1 Theory Nuisance Parameters perfect candidate

- » include correct point-by-point correlations across the q_T spectrum, different processes ,...
- » can be constrained by data reducing theory uncertainty
- » first applications work as advertised

2 Perturbative uncertainty with TNPs

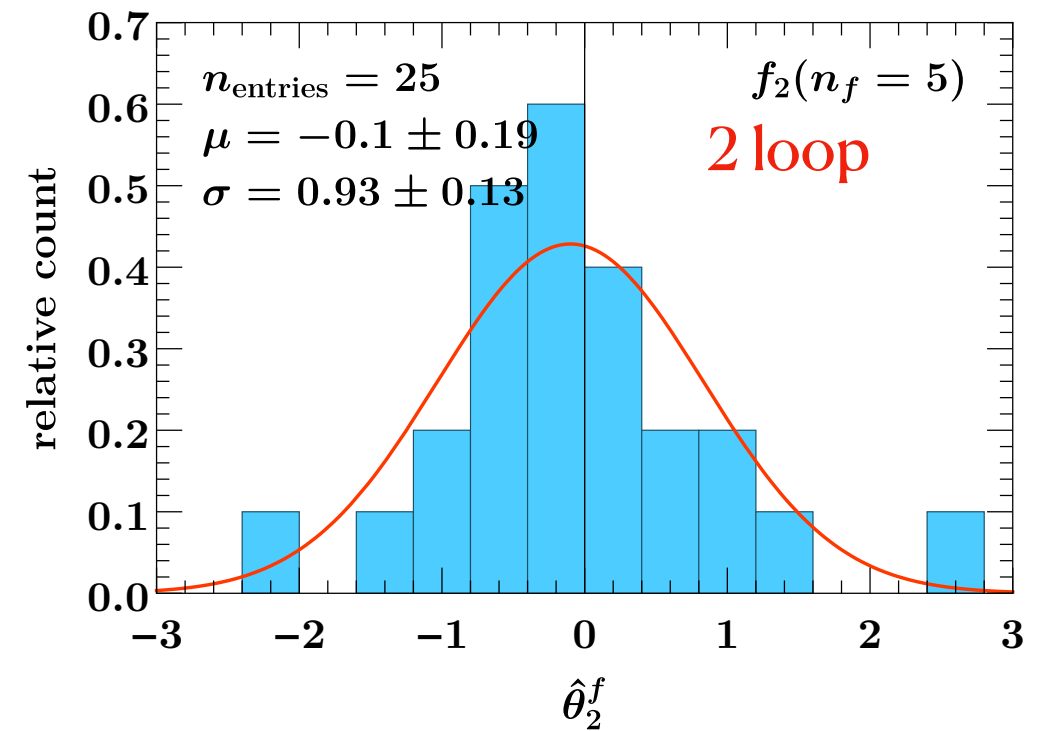
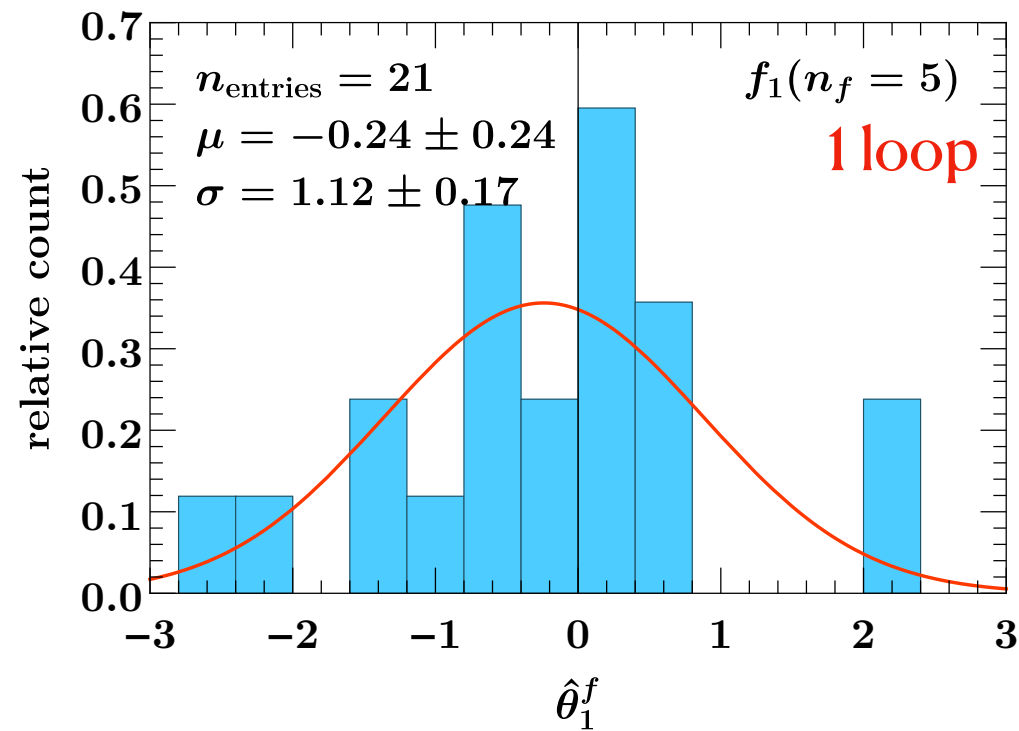
- » perturbative uncertainty can be correctly profiled
- » TNPs not “easy and cheap” as scale variation, but worth it!

THANK YOU!

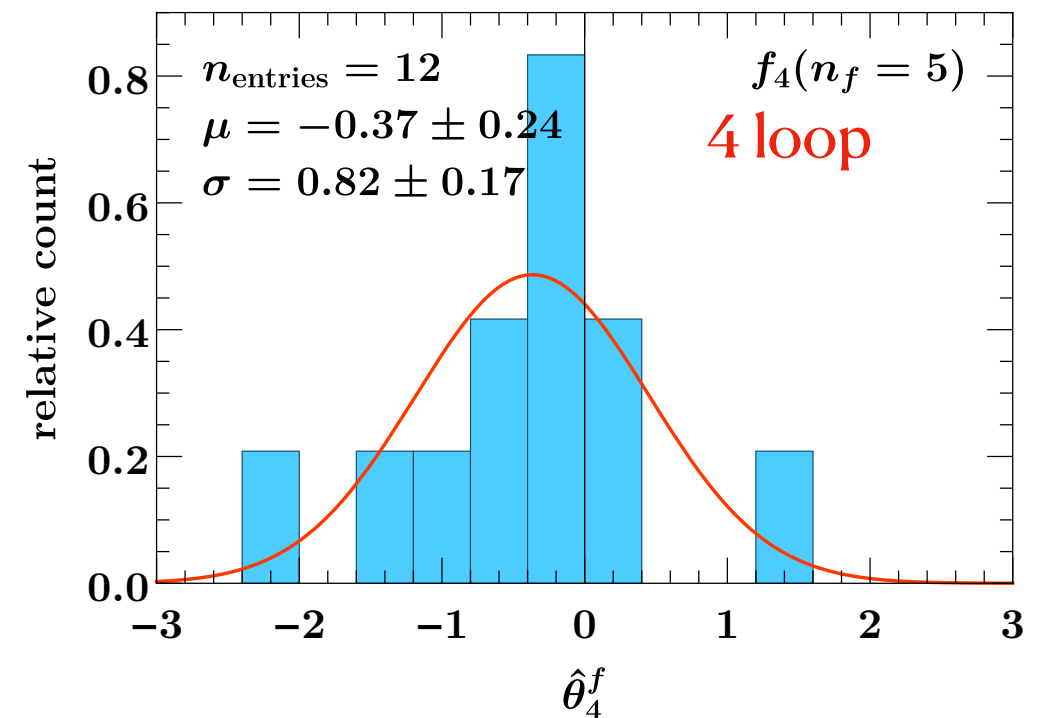
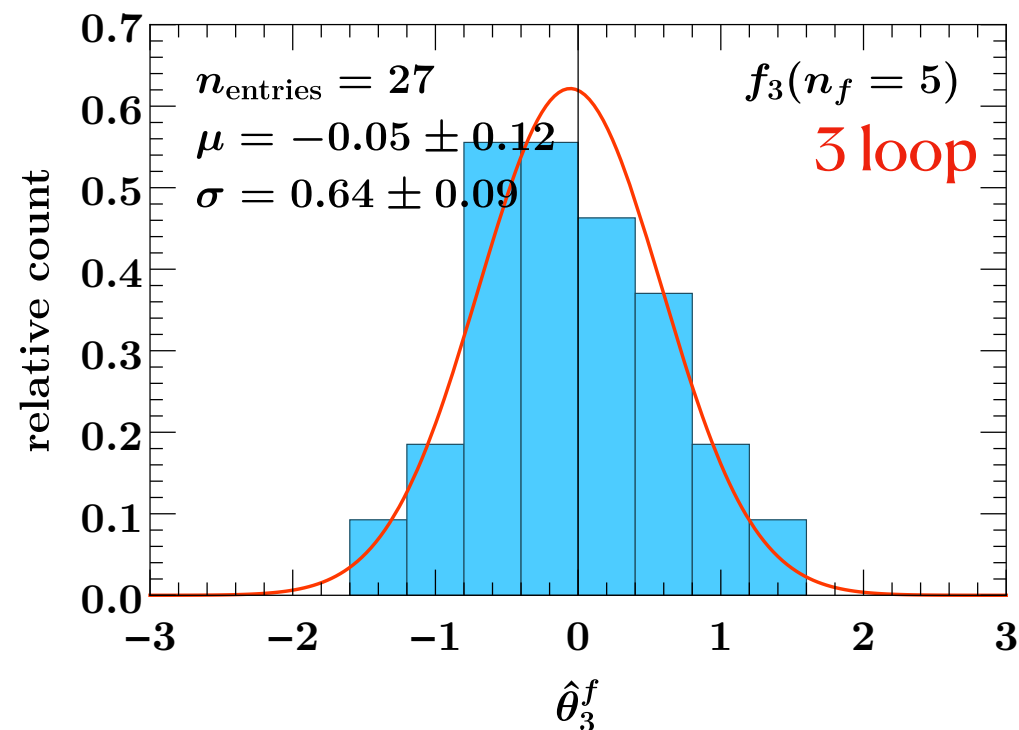
Backup slides

TNPs for Boundary Conditions

$$F_n(\theta_n) = 4C_r(4C_A)^{n-1}(n-1)!\theta_n^f$$

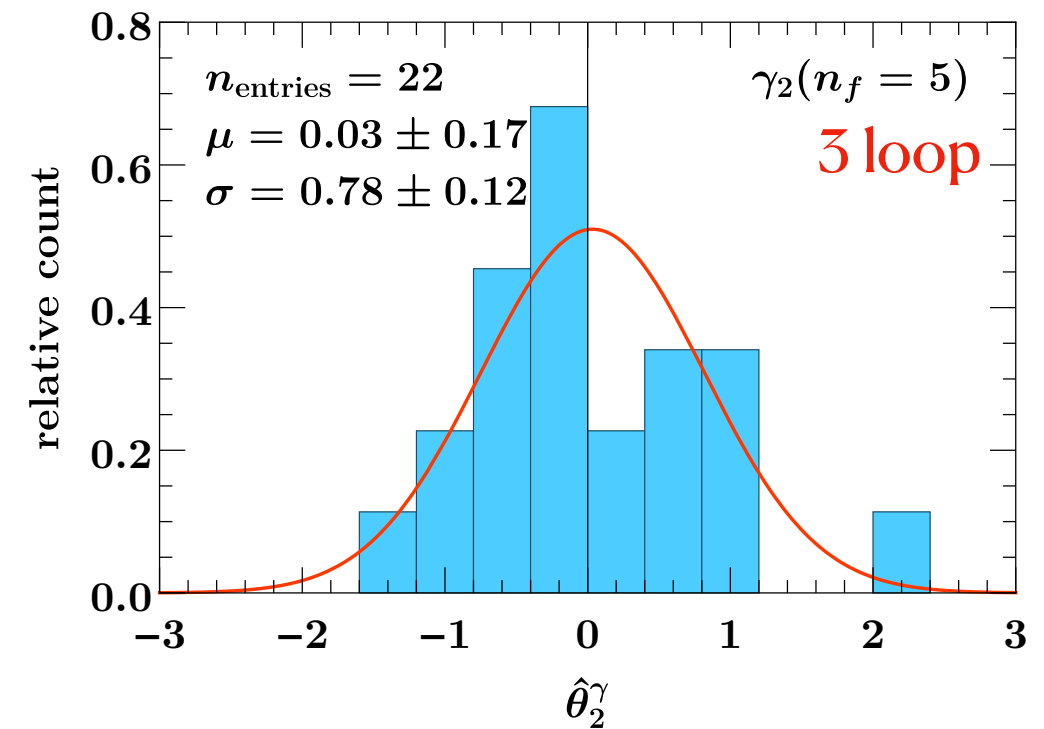
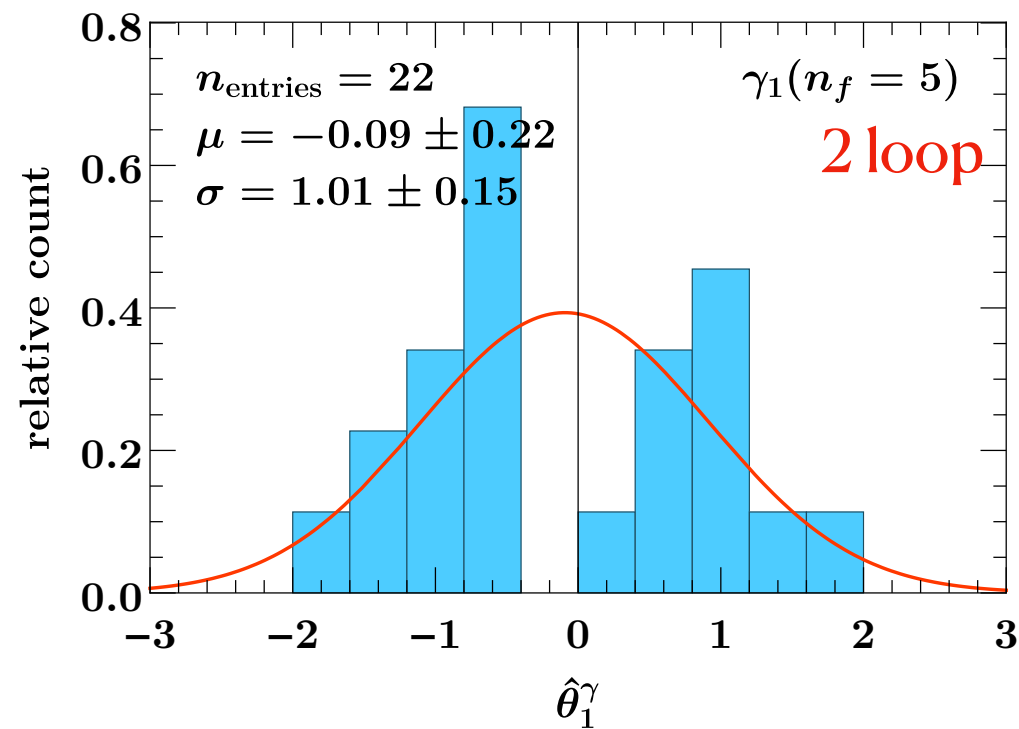


Good fit to a Gaussian with $\theta_n \approx 0$ and $\Delta\theta_n \approx 1$

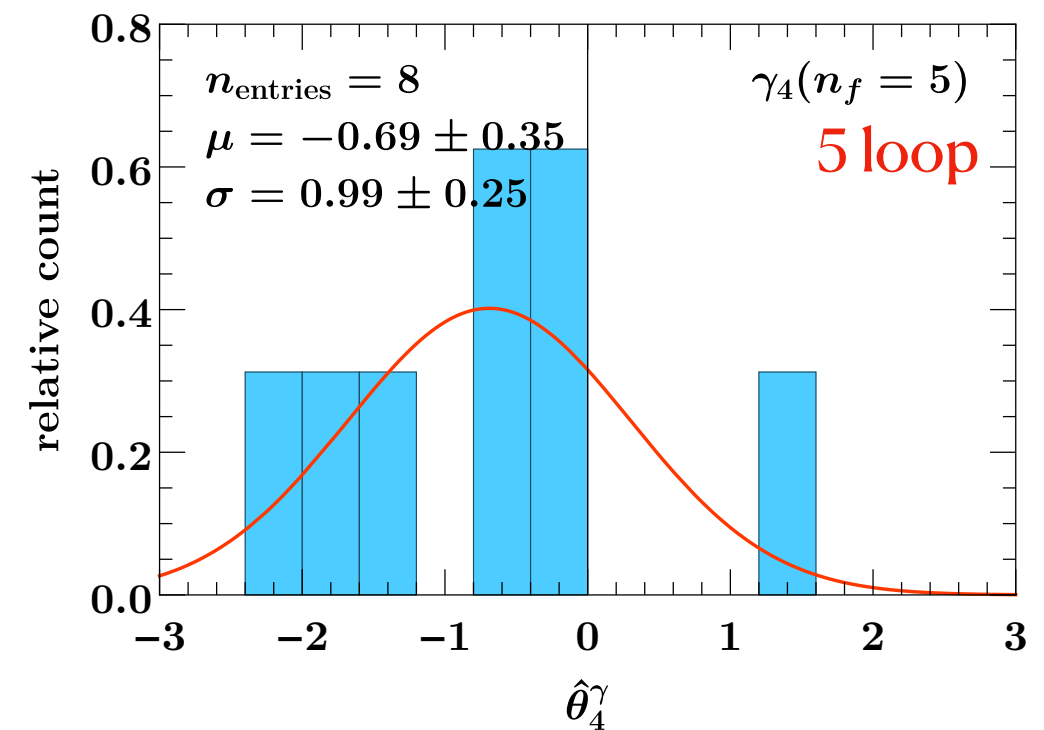
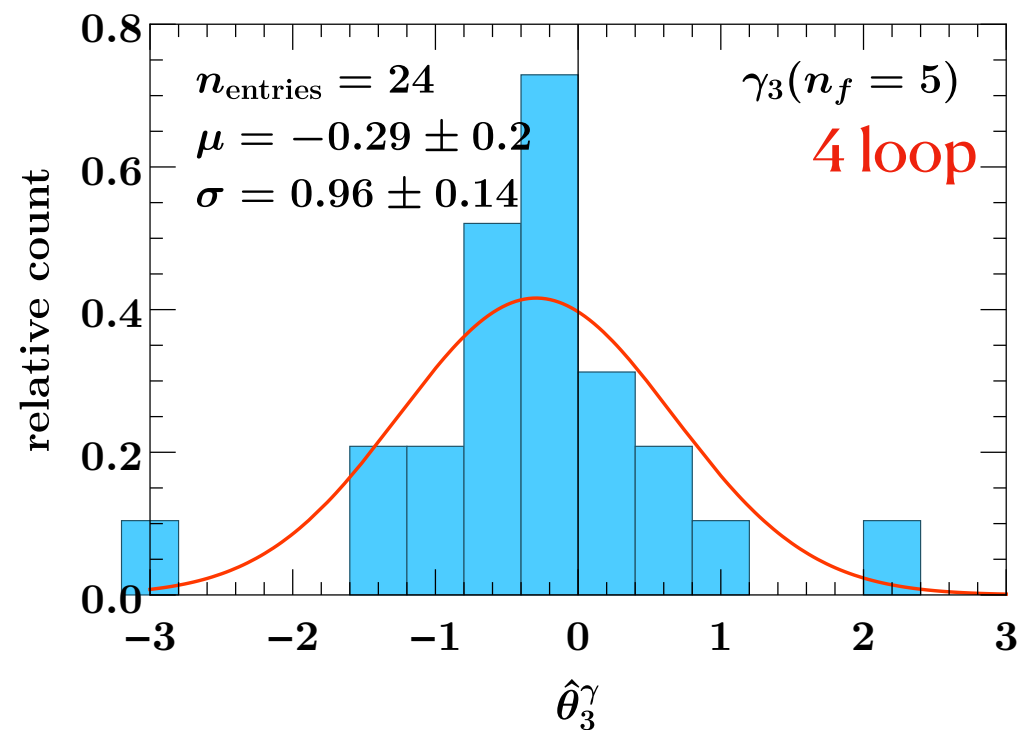


TNPs for Anomalous Dimensions

$$\gamma_n(\theta_n) = 4C_r(4C_A)^n \theta_n^\gamma$$



Good fit to a Gaussian with $\theta_n \approx 0$ and $\Delta\theta_n \approx 1$

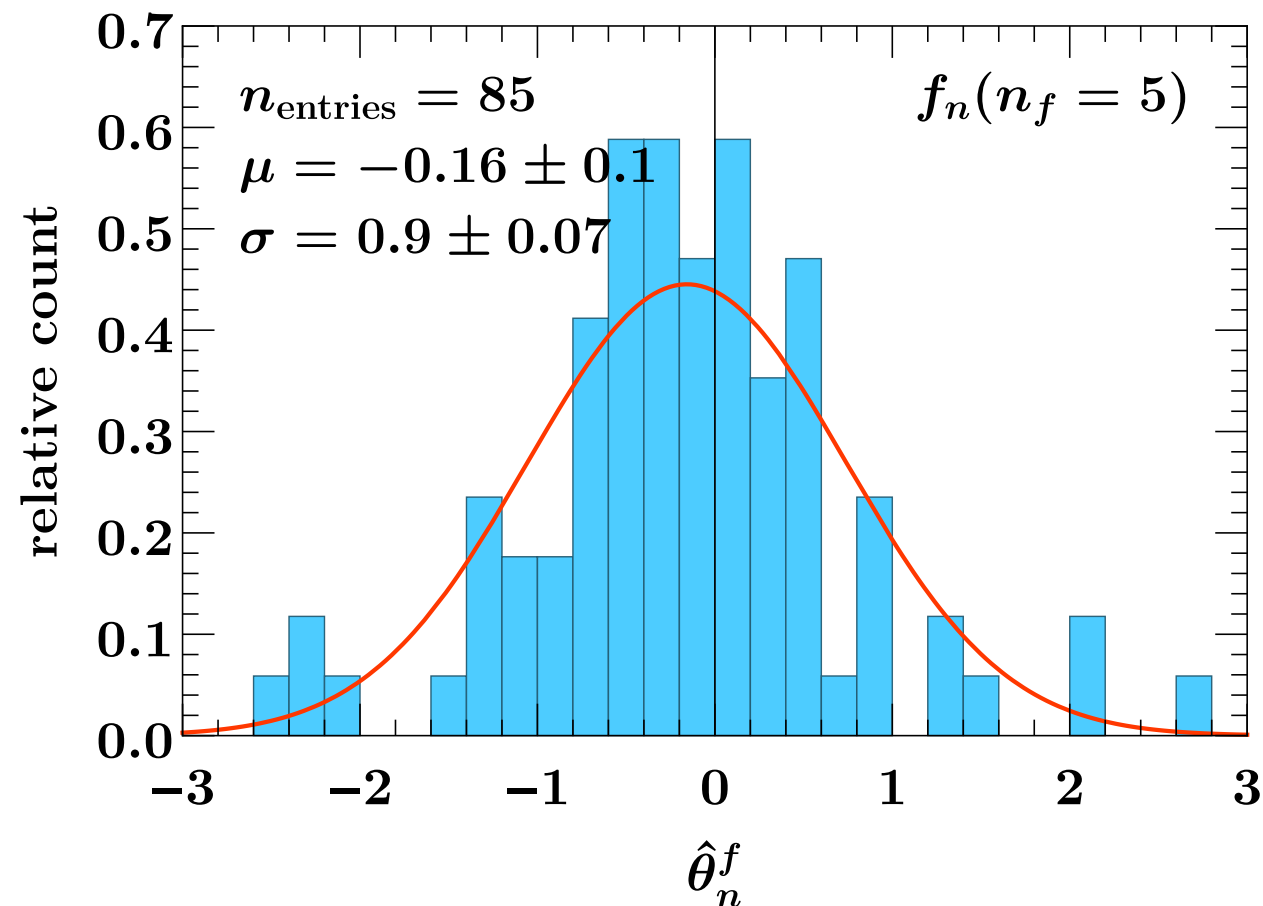


TNPs for BC and ADm

Considering all together from 1 to 4 loop:

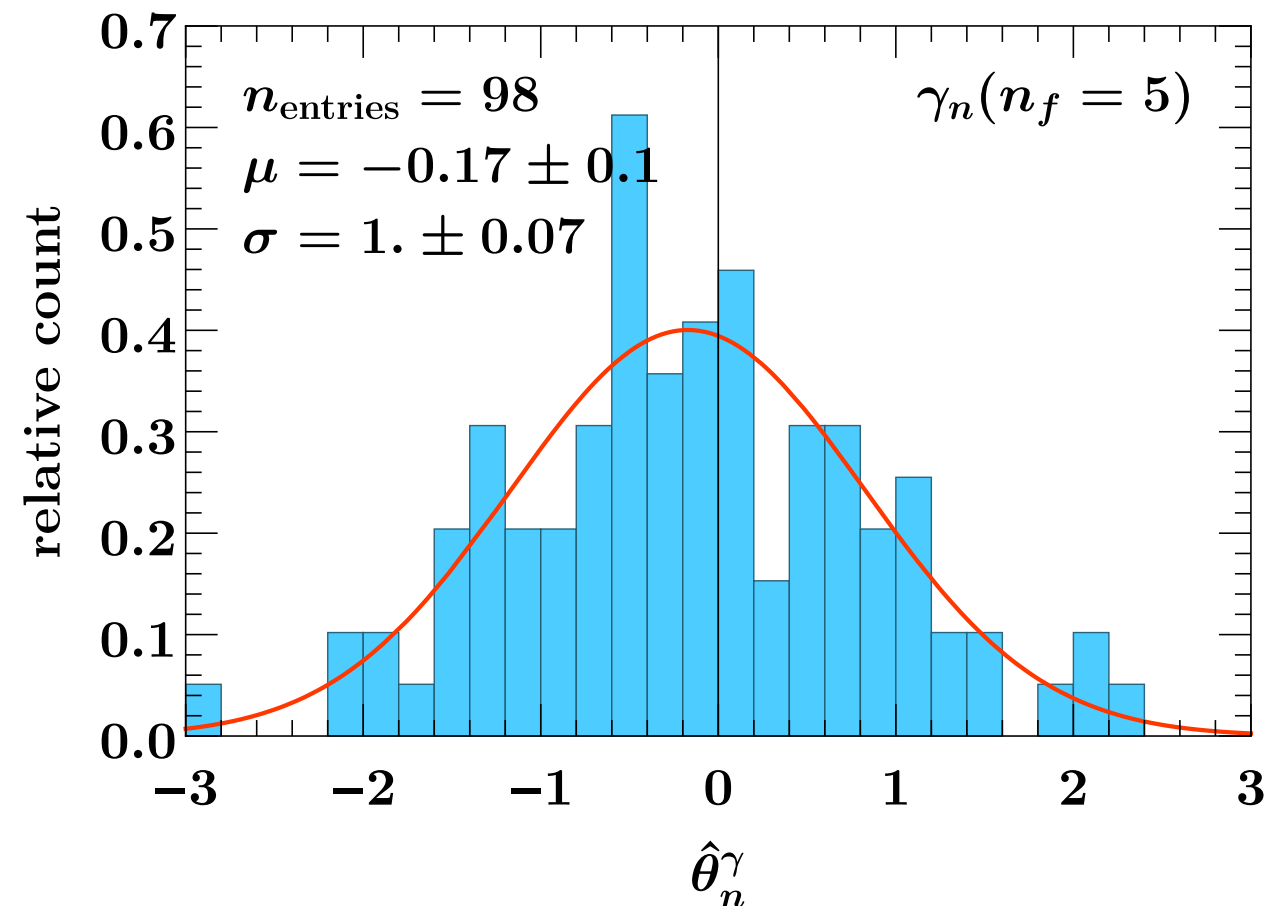
Boundary Conditions

$$F_n(\theta_n) = 4C_r(4C_A)^{n-1}(n-1)! \theta_n^f$$



Anomalous Dimensions

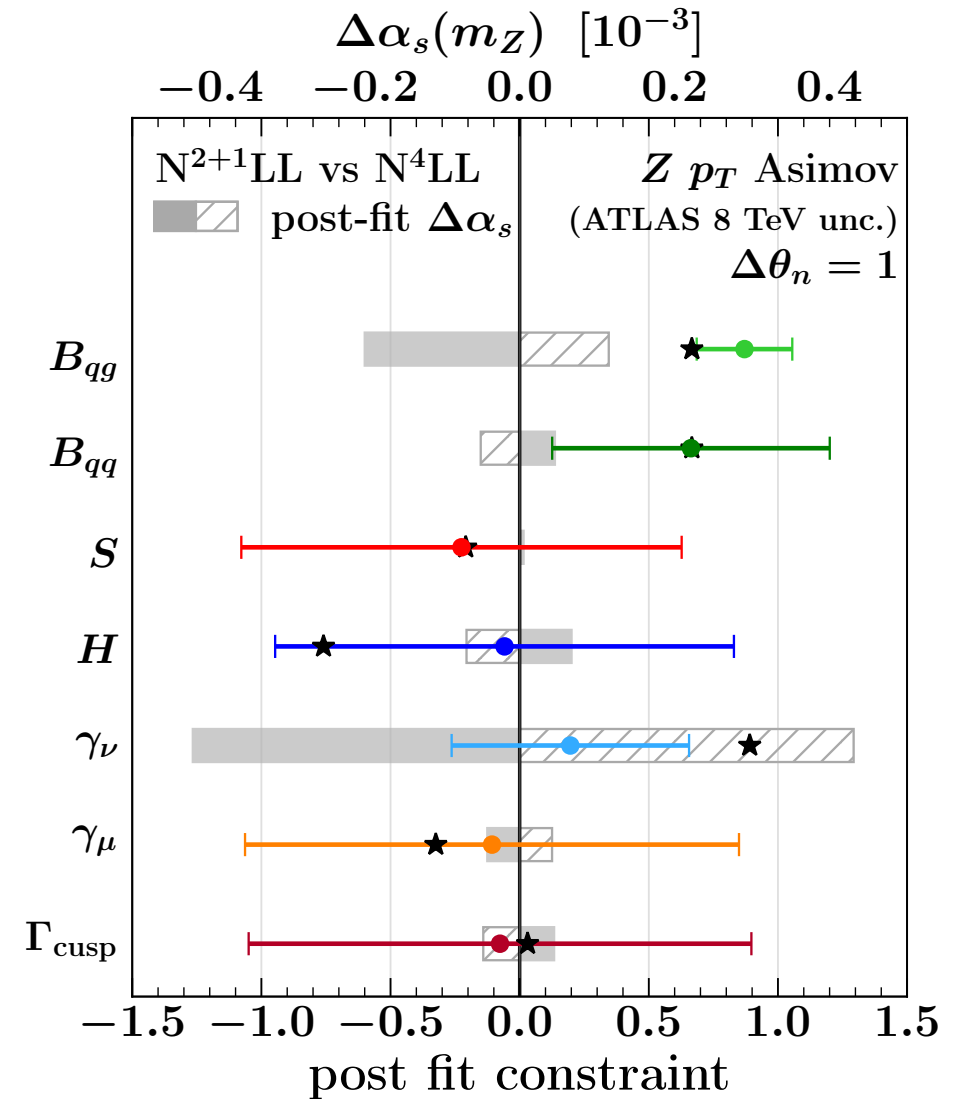
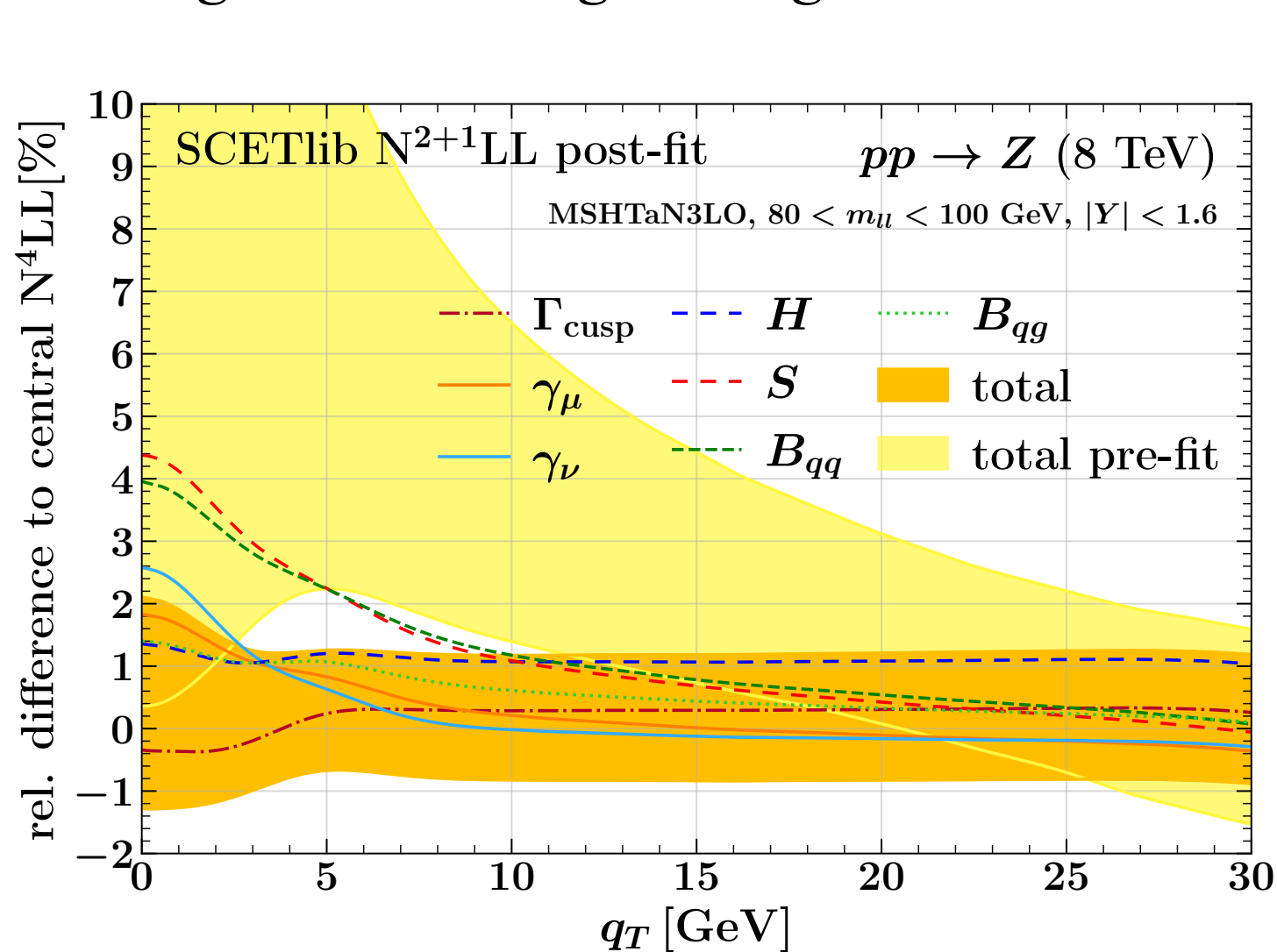
$$\gamma_n(\theta_n) = 4C_r(4C_A)^n \theta_n^\gamma$$



Very good fit to a Gaussian with $\theta_n \approx 0$ and $\Delta\theta_n \approx 1$ ✓

Post-fit constraints on $N^{2+1}LL$

Profiling lower order against higher order: $N^{2+1}LL$

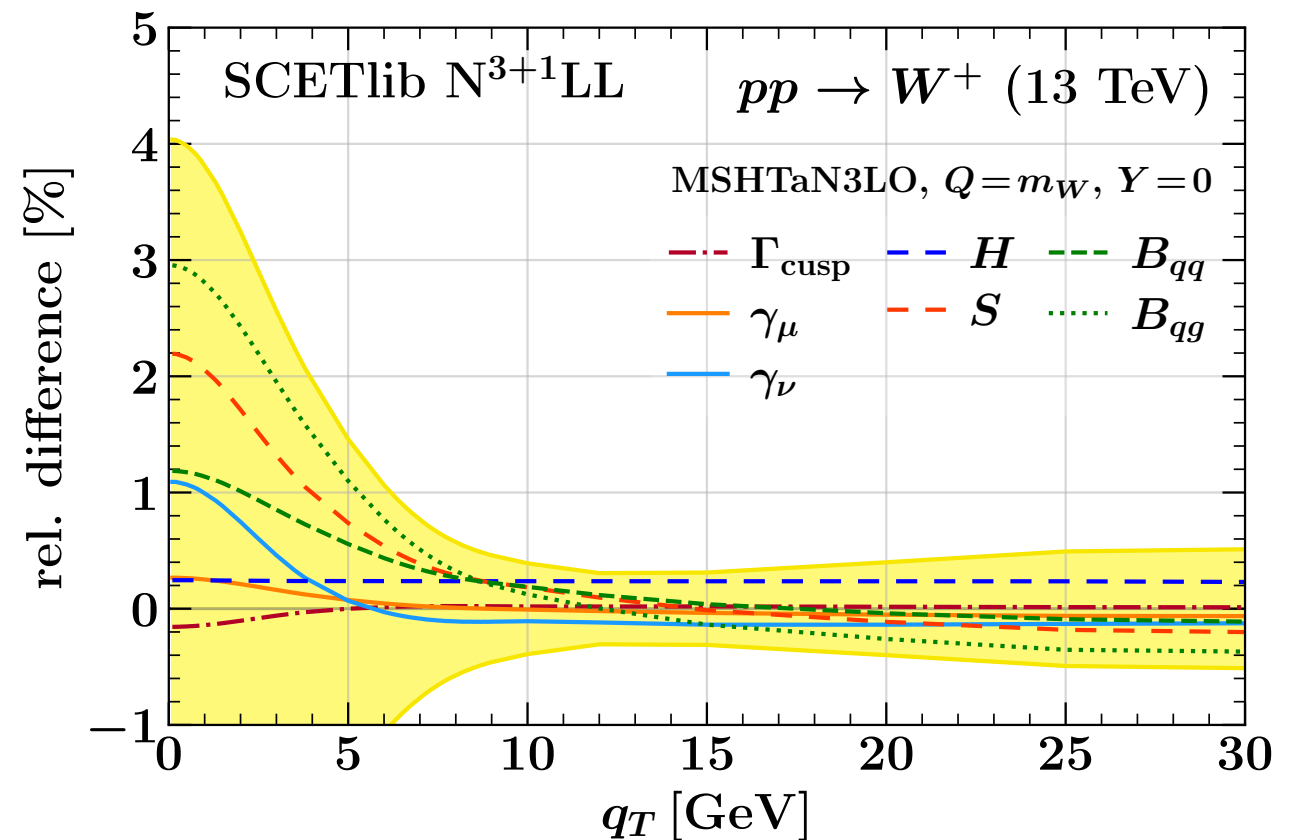
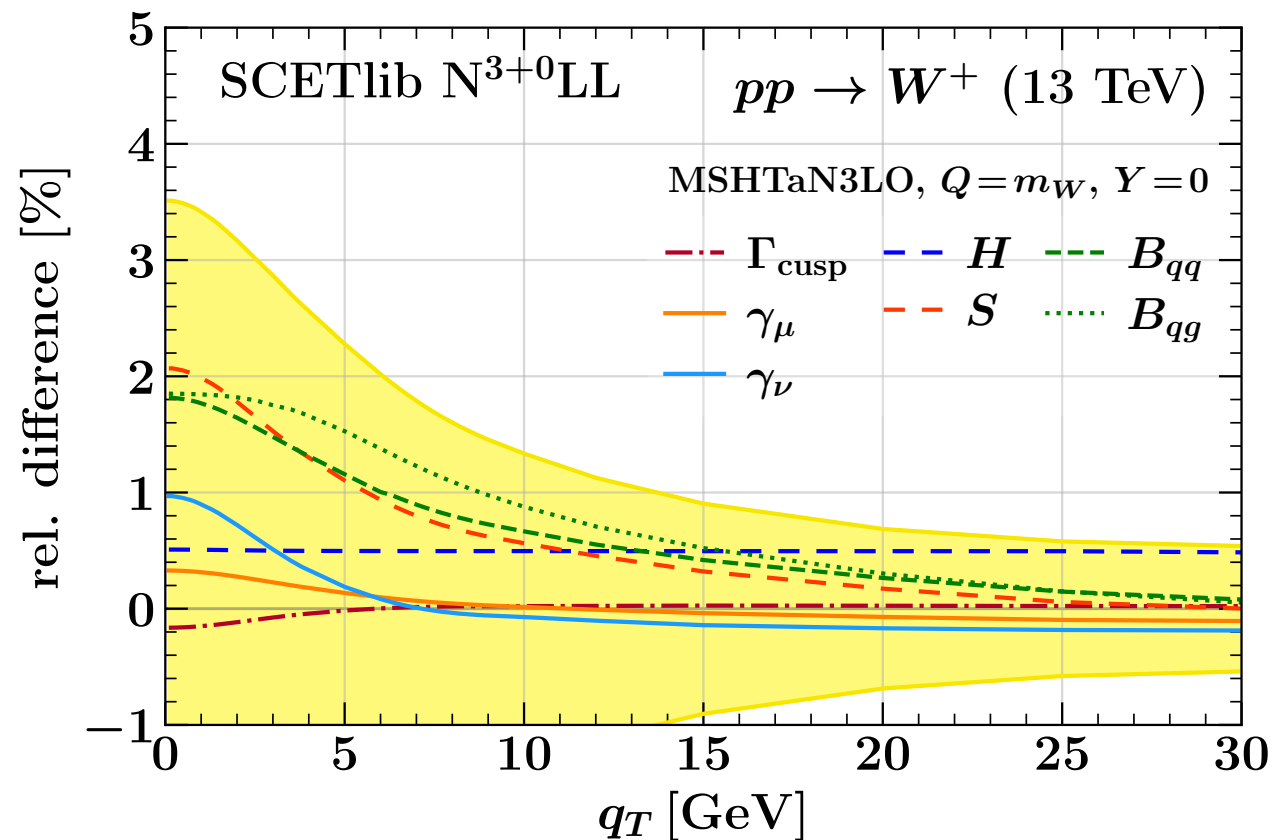


- $N^{2+1}LL$ strongly pulled, toward correct true values [★]
- post-fit prediction for q_T spectrum driven by constraints from data

TNPs at N^{3+0} LL

N^{3+0} LL is an approximation of N^{3+1} LL:

consider the N^3 LL structure but absorb the N^{3+1} LL TNPs uncert. term into the N^3 LL structure



- The impact of the parameters is only approximately correct!
 - Limited effect on the overall size of theory uncert., but may have bigger effect on theory correlations
- ➔ if possible prefer the N^{m+1} LL prescription!

SCETlib nonperturbative models

➤ Collins-Soper (CS) kernel

$$\tilde{\gamma}_\nu(b_T) = \tilde{\gamma}_\nu^{\text{pert}}\left(b_6^*(b_T)\right) + \tilde{\gamma}_\nu^{\text{nonp}}(b_T)$$

$$\tilde{\gamma}_\nu^{\text{nonp}}(b_T) = -\lambda_\infty f_\nu\left(\frac{\lambda_2}{\lambda_\infty}b_T^2 + \frac{\lambda_4}{\lambda_\infty}b_T^4\right)$$

➤ Transverse Momentum Distributions (TMDs)

$$\tilde{f}(b_T) = \tilde{f}_{\text{pert}}(b_T) \tilde{f}_{\text{nonp}}(b_T)$$

$$\ln\left(\tilde{f}_{\text{nonp}}(b_T)\right) = -\Lambda_\infty b_T f\left(\frac{\Lambda_2}{\Lambda_\infty}b_T + \frac{\Lambda_4}{\Lambda_\infty}b_T^3\right)$$

λ_2, λ_4 and Λ_2, Λ_4 quadratic/quartic small b_T coefficients

$\lambda_\infty, \Lambda_\infty$ determine $b_T \rightarrow \infty$ behavior

What about $f_\nu(x)$ and $f(x)$? [Collins and Rogers '14]

$$\begin{aligned} \tilde{\gamma}_\nu^{\text{nonp}}(b_T \rightarrow 0) &\sim b_T^2, & \tilde{\gamma}_\nu^{\text{nonp}}(b_T \rightarrow \infty) &\sim \text{const} & f_\nu\left(\frac{\lambda_2}{\lambda_\infty}b_T^2 + \frac{\lambda_4}{\lambda_\infty}b_T^4\right) &= \tanh\left(\frac{\lambda_2}{\lambda_\infty}b_T^2 + \frac{\lambda_4}{\lambda_\infty}b_T^4\right) \\ \log\left(\tilde{f}_{\text{nonp}}(b_T \rightarrow 0)\right) &\sim b_T^2, & \log\left(\tilde{f}_{\text{nonp}}(b_T \rightarrow \infty)\right) &\sim b_T & f\left(\frac{\Lambda_2}{\Lambda_\infty}b_T + \frac{\Lambda_4}{\Lambda_\infty}b_T^3\right) &= 2 \tanh\left(\frac{\Lambda_2}{\Lambda_\infty}b_T + \frac{\Lambda_4}{\Lambda_\infty}b_T^3\right) \end{aligned}$$

$\alpha_s(m_Z)$ determination from ATLAS [arXiv:2309.12986]

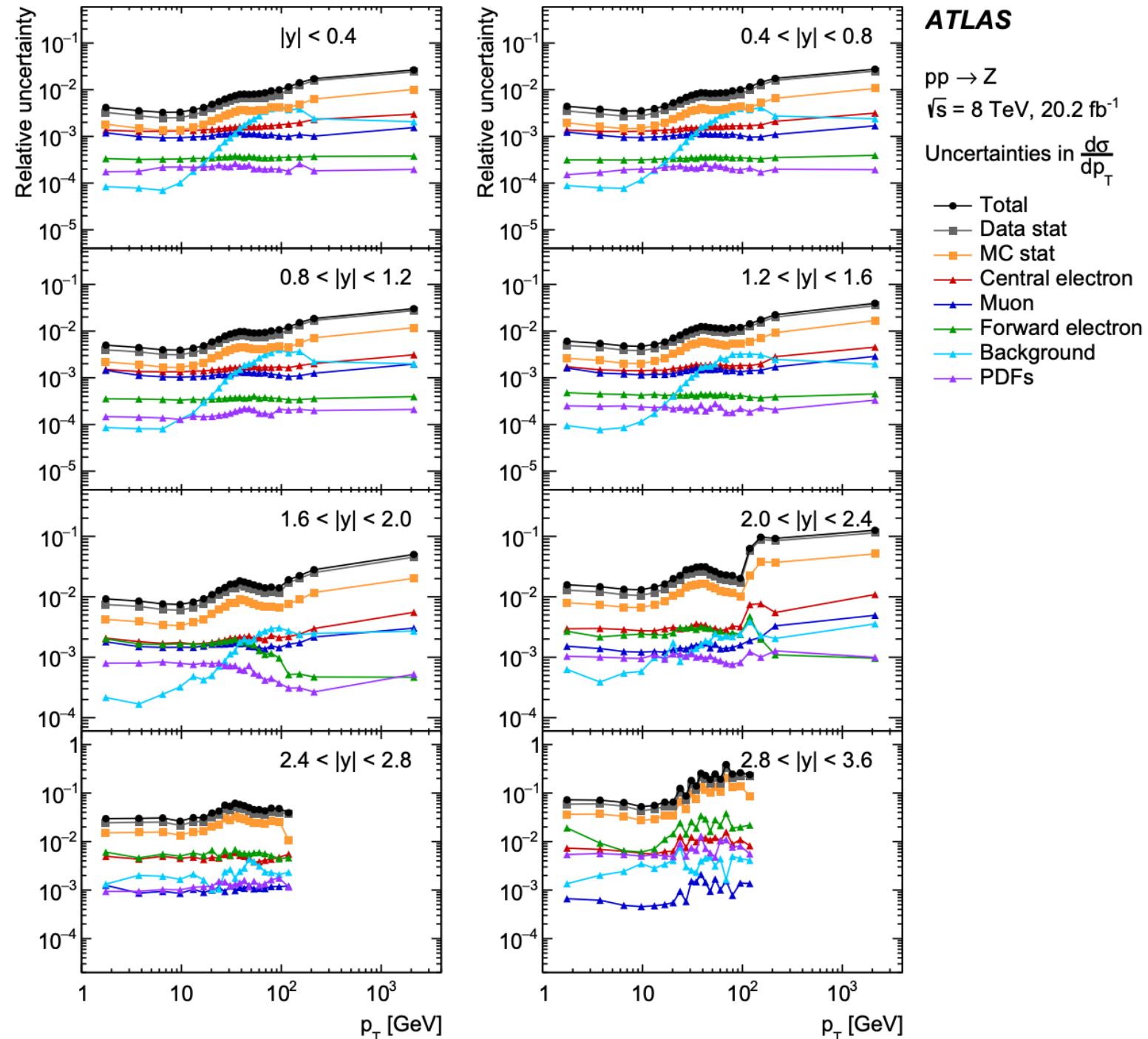
Fit of strong coupling constant using p_T in $[0, 29]$ GeV

based on $N^3\text{LO}+N^4\text{LLa}$ from DYTurbo

$$\alpha_s(m_Z) = 0.1183 \pm 0.0009$$

Breakdown of uncertainties on α_s
in units of 10^{-3} :

Experimental uncertainty	± 0.44	
PDF uncertainty	± 0.51	
Scale variation uncertainties	± 0.42	
Matching to fixed order	0	-0.08
Non-perturbative model	+0.12	-0.20
Flavour model	+0.40	-0.29
QED ISR	± 0.14	
$N^4\text{LL}$ approximation	± 0.04	
Total	+0.91	-0.88



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