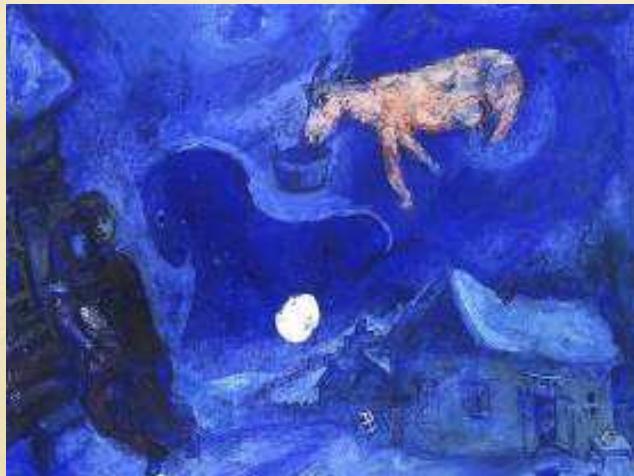


$$B \rightarrow V\gamma$$

## Beyond QCD Factorization



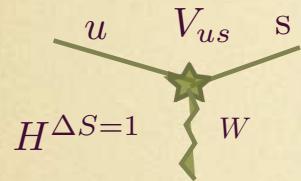
Roman Zwicky (Durham IPPP)

Flavour Physics at IPPP -- IPPP steering Committee

## Overview

- B-physics CKM-matrix generalities
- Exclusive decays  $H_{\text{eff}}$
- Hadronic matrix elements – Distribution Amplitudes
- Overview of our activities
- Focus  $B \rightarrow V\gamma$  beyond QCD factorization
- Particular focus on time dep. CP-asymmetry

# The Quark-Sector



$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \sim \begin{pmatrix} 1 & \delta & \delta^3 \\ \delta & 1 & \delta^2 \\ \delta^3 & \delta^2 & 1 \end{pmatrix} \quad \text{hierarchy : } \delta \sim 0.2$$

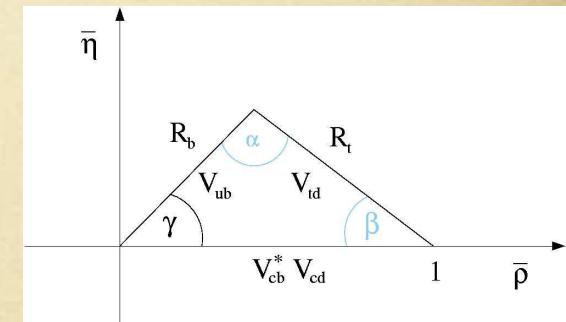
absolute values

- $V$  arbitrary **unitary** matrix  
9 parameters (5=6-1 absorbed in phases of quarks)  
 $\Rightarrow$  4 parameters  $\sim 3$  rotation & 1 complex phase
- $(CP)\mathcal{L}^{\text{weak}}(CP)^\dagger = \mathcal{L}^{\text{weak}} \Leftrightarrow V_{\alpha\beta} = V_{\alpha\beta}^* \text{ CP-violation}$
- unitarity  $VV^\dagger = 1$ , 9 conditions among  $B_d^0$  triangle

$$V_{du}V_{ub}^* + V_{dc}V_{cb}^* + V_{dt}V_{tb}^* = 0$$

known as the CKM-triangle ( $B_s^0$  and  $K^0$  triangles too flat)

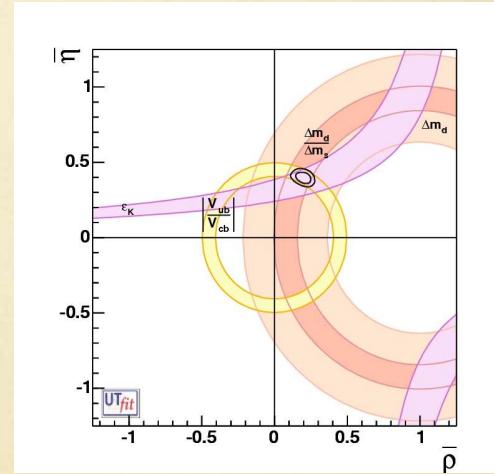
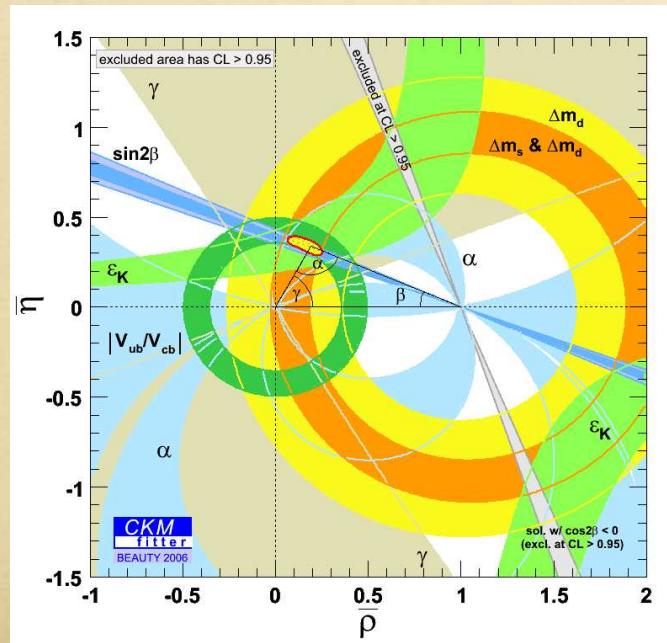
- CP violation  $\leftrightarrow$  non-zero area triangle
- only established source of CP-viol SM (MNS neutrinos ?)



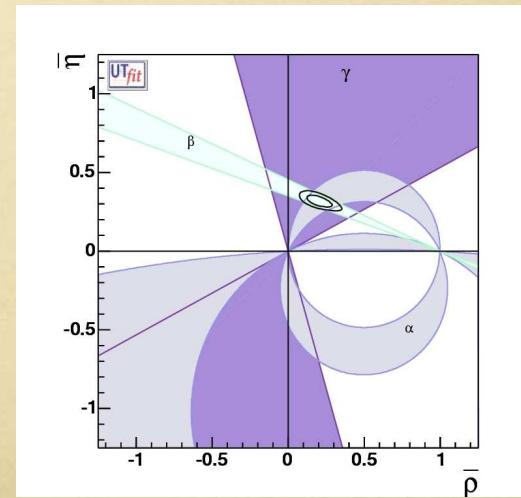
# Measuring CKM - B-Factories 1999-

## Results:

1. CKM mechanism works
  - triangle overconstrained
  - describes CP-viol. mesons quant.
2. Not enough Baryogenesis ( $\sim 10^{10}$ )  
 → T.Underwood's talk



From Sides (CP-cons.)



From Angles (CP-viol.)

## Challenge for Theory

Relate observables to fund. parameters (Hadronization)

$$\Gamma(B \rightarrow \pi e \nu), \Delta M_s \quad \leftrightarrow \quad m_b, V_{ub}, V_{ts}$$

A few simple examples

- $A_{CP}(B \rightarrow J/\Psi K_s) \sim \sin(2\beta) \sin(\Delta Mt)$  Bigi & Sanda 81
- $|V_{ub}|^{-2} \sim \int dq^2 |f_+(q^2)|^2$  one single form factor !

# Exclusive Decays

$$B \rightarrow X Y \leftrightarrow \langle XY | H_{\text{eff}} | B \rangle$$

	$H_{\text{eff}}$	remark	type
$B \rightarrow \pi e \nu$	$(\bar{b}u)_{V-A}$	tree – level	semileptonic
$\bar{B} \rightarrow K^* \gamma$	$(\bar{b}\sigma \cdot F s)_{V+A} + \dots$	FCNC	semileptonic
$B \rightarrow \pi \pi$	$(\bar{b}u)_{V-A}(\bar{d}u)_{V-A}$	tree – level	non – leptonic

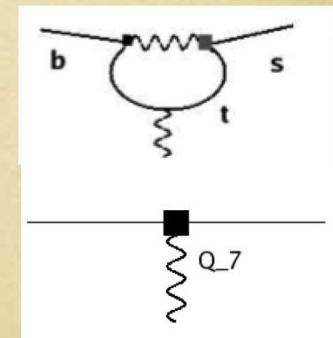


- $M_W$  to  $m_b$  described pert. QCD  $\Rightarrow$  concretely apply Wilsonian picture

$$H_{\text{eff}} = \sum \underbrace{C_i(\mu)}_{\text{UV}} \underbrace{O_i(\mu)}_{\text{IR}}$$

$Q_i$  operators (above),  $\mu \sim m_b$  factorisation scale ... steps:

1. matching scale  $M_W$  (initial condition)
2. RNG-running scale  $C_i \sim L^{\gamma_i} \dots \gamma_i$  anom. dim  $\rightarrow$  calc.
3. evaluate  $\langle XY | Q_i(\mu) B \rangle$      "domain of expertise"



# Matrix elements: semilept. form factors

- Parametrize matrix element grounds Lorentz-covariane

$$\langle \pi(p) | \bar{b} \gamma_\mu (1 - \gamma_5) q | B \rangle = 2 p_\mu f_+(q^2) \quad 0 < q^2 < (m_B - m_\pi)^2$$

Heavy-light (no sym. light-light or heavy-heavy)  $\Rightarrow$  non-pert object

**Methods:**

1. Light-Cone Sum Rules small  $q^2 \leftrightarrow E_\pi$  large, pion LC
  2. Lattice-QCD large  $q^2 \leftrightarrow E_\pi$  small, pion o/w not resolved by lattice
- are complementary
- LCSR are QCD sum-rules

$$\begin{aligned} \text{OPE} &\leftrightarrow \text{LC} - \text{OPE} \\ \text{Cond : } \langle \bar{q}q \rangle &\leftrightarrow \text{DistributionAmplitudes : } \phi_\pi(u) \end{aligned}$$

- Leading part the form factor in LCSR

$$f_+^{B \rightarrow \pi} \sim f_\pi (1 + O(1) a_2(\pi) \dots) + \dots$$

$a_2$  moment DA

# Meson Distribution Amplitudes (DA)

- relevant exclusive processes at large momentum transfer (LC)
- most important DA  $\leftrightarrow$  minimal number of partons

$$\langle \pi(p) | \bar{q}(x) \gamma_\mu \gamma_5 q(0) | 0 \rangle = i f_\pi p_\mu \int e^{i u p x} \phi_\pi(u) + O(x^2, m_\pi^2)$$

- called twist-2, twist = dim - spin = 3 -1  
correction: higher Fock states, dev. from LC
- variable  $u$  : col. momentum fraction of corr. quark in meson

- QCD e.o.m. different DA related

$$\partial_\mu \bar{q}(x) \gamma_\mu s(-x) = -i \int_{-1}^1 dv v \bar{q}(x) x_\alpha g G^{\alpha\mu}(vx) \gamma_\mu s(-x) - (m_s + m_q) \bar{q}(x) i(\gamma_5) s(-x)$$

Wilson lines omitted, take matrix elements

- what to do  $\phi_\pi(u)$  ?

A. models obeying theor. & exp. constraints

B. expand in Eigenfunction of evolution kernel

$$\phi_\pi(u, \mu) = 6u(1-u)(1 + \sum_{n \geq 1} a_n(\mu, \pi) C_n^{3/2}(2u-1))$$

- $a_n$  Gegenbauer moments
- $C_n$  Gegenbauer pol. rep. collinear  $SL(2, R) \in SO(4, 2)$  conformal
- asymptotic part  $6u(1-u)$  exact from perturbative QCD

- Why good expansion ?

1. convolution smooth kernel  $\int du T(u) C_n(u)$  decreasing ( $C_n$  n nodes)
2. det. of  $a_n$  indicate  $a_0 \equiv 1 > |a_1| > |a_2| \dots$   
"consistent" with  $\gamma_{n+1} > \gamma_n$  conformal hierarchy

- Determination ?

- a. fit to experiment .. (contam. other hadr. uncert.)
- b. calculate from local matrix element (Fourier inversion)
  - \* QCD sum rules
  - \* lattice QCD (new efforts)

## Overview results (selected)

1. calculation of all weak and penguin form factors  $B \rightarrow (P, V)$  from LCSR  
PRD 71, 14015 and PRD 71, 14029 2005 Ball RZ

$$|V_{ub}|_{\text{BZ}}^{\text{HFAG}} = 3.38 \pm 0.12^{+0.56}_{-0.37} 10^{-3}$$

Open parenthesis (NP ?)

$$|V_{ub}|_{\text{CKM-fit}} = 3.5 \pm 0.18 10^{-3} \quad |V_{ub}|_{\text{incl}}^{\text{DGE}} = 4.46 \pm 0.2 \pm 0.2 10^{-3}$$

2. Kaon DA  $K, K^{*,\parallel}, K^{*,\perp}$  emph. SU(3)-break.  
PLB 633,289 2006, JHEP 0602:034 2006 Ball RZ

$$a_1(K) = 0.06 \pm 0.03, \quad a_1^\parallel(K^*) = 0.03 \pm 0.02, \quad a_1^\perp(K^*) = 0.04 \pm 0.03.$$

Later QCDSF/HPQCD & QCDSF same values advert. smaller uncert.

3. apply these results  $B \rightarrow V\gamma$   
JHEP 0604:046,2006, PLB B642:478 RZ Ball  
all new hep-ph/0612081 Ball Jones RZ topic rest

## $B \rightarrow V\gamma$ beyond QCD Factorisation

- whole zoo:  $B \rightarrow (\rho^{(0,+)}, \omega, K^{*,(0,+)})$  and  $B_s \rightarrow (\bar{K}^*, \phi)$
- FCNC new physics hidden loops ?
- $\mathcal{B}(b \rightarrow X_s \gamma)_{\text{incl}}$  no hints NP
- history

1993  $B \rightarrow K^* \gamma$  CLEO

2005  $B \rightarrow \rho \gamma$  Belle (BaBar 2006)

200x  $B_s \rightarrow (\bar{K}^*, \phi) \gamma$  LHCb

$$\mathcal{H}_{\text{eff}}^{(d,s)} = \frac{G_F}{\sqrt{2}} \sum_{U=u,c} \underbrace{\lambda_U^{(d,s)}}_{\text{CKM}} \left[ C_1 Q_1 + C_2 Q_2 + \sum_{i=3,\dots,8} C_i Q_i \right] + \text{BSM}$$

Effective Hamiltonian (OPE-description), CKM-factors  $\lambda_t^s = V_{bt} V_{ts}^*$

$$Q_1^U = (\bar{s}_i U_j)_{V-A} (\bar{U}_j b_i)_{V-A}$$

$$Q_3 = (\bar{s}b)_{V-A} \sum_q (\bar{q}q)_{V-A}$$

$$Q_5 = (\bar{s}b)_{V-A} \sum_q (\bar{q}q)_{V+A}$$

$$Q_7 = \frac{e}{8\pi^2} m_b \bar{s}_i \sigma^{\mu\nu} (1 + \gamma_5) b_i F_{\mu\nu}$$

$$Q_2^U = (\bar{s}U)_{V-A} (\bar{U}b)_{V-A}$$

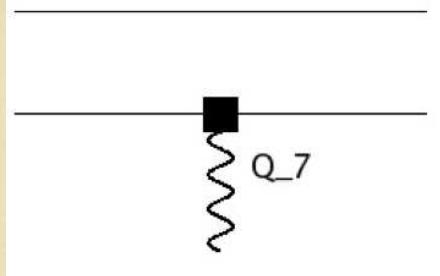
$$Q_4 = (\bar{s}_i b_j)_{V-A} \sum_q (\bar{q}_j q_i)_{V-A}$$

$$Q_6 = (\bar{s}_i b_j)_{V-A} \sum_q (\bar{q}_j q_i)_{V+A}$$

$$Q_8 = \frac{g}{8\pi^2} m_b \bar{s}_i \sigma^{\mu\nu} (1 + \gamma_5) T_{ij}^a b_j G_{\mu\nu}^a$$

- $C_i$  Wilson Coeff. NLO  $\alpha_s$  (NNLO first estimate) from “inclusive people”
- $C_7$  leading  $1/m_b$  and  $\alpha_s$   $b \rightarrow (d, s)\gamma$ , add  $O(m_s)$ -part later
- $C_2$   $1/m_b$  (power)-correction in  $b \rightarrow (d, s)\gamma$ , but  $C_2 \sim C_7$
- $Q_2^U$ -flavour sensitive, discr. final state

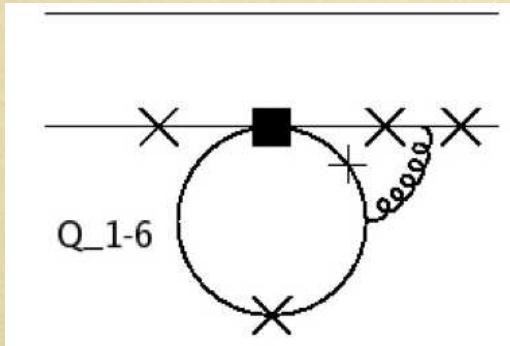
- At  $O(\alpha_s^0)$ , LO  $1/m_b$  only electric penguin



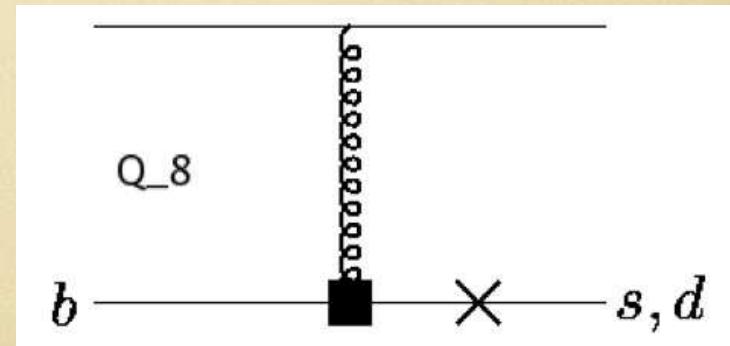
$$\langle V\gamma | \underbrace{\bar{b}\sigma \cdot F(1+\gamma_5)q}_{Q_7} | B \rangle = kin.T_1(0) \quad \text{FF from LCSR}$$

- At  $O(\alpha_s)$ , LO  $1/m_b$  **QCD factorization**  
 ( Bosch et al, Beneke et al, Ali et al 01):

$$\begin{aligned} \langle V\gamma | Q_i | B \rangle &= T_i^I F(B \rightarrow V_\perp) + \int_0^1 d\xi du \phi_B(\xi) \phi_{V_\perp}(u) T_i^{II}(\xi, u) + O(\Lambda/m_b) \\ &= \# \langle V\gamma | Q_7 | B \rangle + O(\Lambda/m_b) \end{aligned}$$



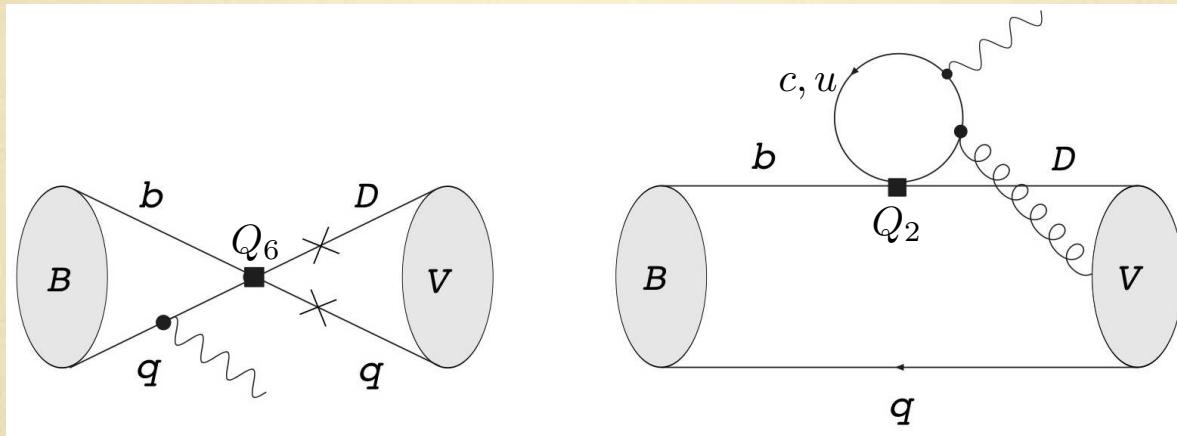
Ex. hard-vertex,  $T^I$



Ex. hard-spectator,  $T^{II}$

# Beyond QCD Factorization

power suppressed non-factorizable contribution



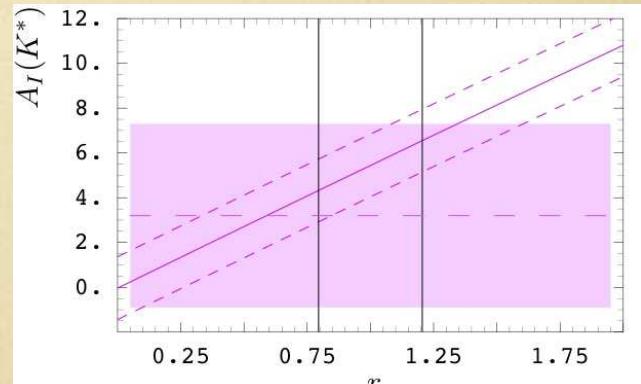
annihilation (drives isospin break.)

$c, u$ -loops (Time-dep CP asymmetry)

- $\mathcal{B}(B \rightarrow V\gamma) \sim 20(30)\%$
- Isospin asymmetries
  - $(\rho^0, \rho^+)$  depend angle  $\gamma$
  - $(K^{*0}, K^{*+})$  .. interesting
- $R \equiv \frac{B(B \rightarrow (\rho, \omega)\gamma)}{B(B \rightarrow K^*\gamma)} \sim \left| \frac{V_{td}}{V_{ts}} \right|^2$  red. uncert.

extract  $|V_{td}/V_{ts}|_{\text{BaBar}}^{BJZ} = 0.199(22)(14)$

compare  $|V_{td}/V_{ts}|_{\text{CDF}}^{\Delta M_s} = 0.206(8)$



$$r = C_6/C_6^{\text{SM}}$$

# Time dependent CP asymmetry

- Key idea  $\bar{B}^0$  decay ( $V-A$ )-int.produce predominantly left handed photons  
Beyond SM anything possible .. enhanced  $m_{\text{NP}}/m_b$

•

$$A_{CP} = \frac{\Gamma(\bar{B}^0(t) \rightarrow \bar{K}^{*0}\gamma) - \Gamma(B^0(t) \rightarrow K^{*0}\gamma)}{\Gamma(\bar{B}^0(t) \rightarrow \bar{K}^{*0}\gamma) + \Gamma(B^0(t) \rightarrow K^{*0}\gamma)} = S \sin(\Delta m_B t) - C \cos(\Delta m_B t),$$

- Gronau,Attwood & Soni 97 since LO operator

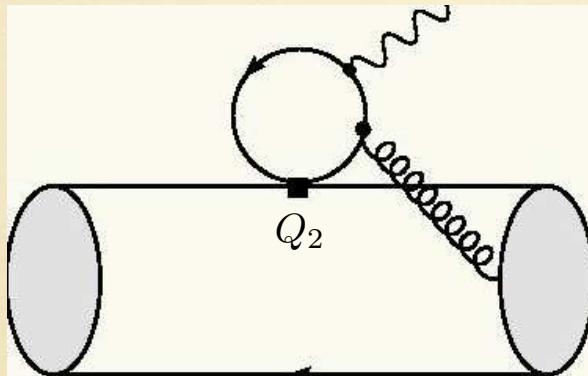
$$Q_7 = \frac{e}{8\pi^2} [m_b \bar{s} \sigma_{\mu\nu} (1 + \gamma_5) b + \cancel{m_s} \bar{s} \sigma_{\mu\nu} (1 - \gamma_5) b] F^{\mu\nu} \equiv Q_7^L + \frac{\cancel{m_s}}{m_b} Q_7^R$$

Time dependent CP asymmetry ( $\sim L \cdot R/L^2$ )

$$S_{K^*\gamma}^{\text{SM}, s_R} = -\sin(2\beta) \frac{m_s}{m_b} (2 + O(\alpha_s)) \sim (2-3)\%$$

- $O(\alpha_s)$ -correction calculable QCD-F
- $\mathcal{B} \sim (L^2 + R^2)/L^2$  right-handed "not sizable"  $O(m_s^2/m_b^2)$

- Grinstein et al 04 further contribution from soft gluon emission, c-loop from  $Q_2^c$  (QCD V-interaction)



- SCET based analysis resorts to dimensional estimate of matrix element and obtain

$$|S_{K^*\gamma}^{\text{SM,soft gluons}}| = 2 \sin(2\beta) \left| \frac{C_2}{3C_7} \right| \frac{\Lambda_{\text{QCD}}}{m_b} \approx 0.06.$$

reach conclusion:  $S_{K^*\gamma}^{\text{SM}} \sim 10\%$  with large uncertainties

# Quantitative Analysis

- idea: on-shell photon, c-quark heavy perform a local OPE

$$\begin{aligned} Q_F &= ie^*\mu \int d^4x e^{iqx} T\{[\bar{c}(x)\gamma_\mu c(x)] Q_2^c(0)\} \\ &= -\frac{1}{48\pi^2 m_c^2} (D^\rho F^{\alpha\beta}) [\bar{s}\gamma_\rho(1-\gamma_5) g\tilde{G}_{\alpha\beta}^a \frac{\lambda^a}{2} b] + \dots \end{aligned}$$

- Remains calculate matrix element  $\langle K^*(p, \eta)\gamma(q, e)|Q_F|B\rangle = L \text{ kin}_1 + \tilde{L} \text{ kin}_2$
- use LCSR (generous uncertainty, 3-particle DA,  $1/m_c$ -corrections)

$$S_{K^*\gamma}^{\text{SM, soft gluons}} = -2 \sin(2\beta) \left( -\frac{C_2}{C_7} \frac{L - \tilde{L}}{36m_b m_c^2 T_1^{B \rightarrow K^*}(0)} \right) = 0.5 \pm 1\%$$

$$\Rightarrow S^{\text{SM}} = -2.2 \pm 1.2_{-\text{1}_{\alpha_s}}^{+0} \%$$

$$S_{\text{BaBar}} = -21 \pm 40 \text{ (stat)} \pm 5 \text{ (syst)} \% \quad \text{BaBar} \quad (232 \cdot 10^6 B\bar{B} \text{ pairs}),$$

$$S_{\text{Belle}} = -32_{-33}^{+36} \text{ (stat)} \pm 5 \text{ (syst)} \% \quad \text{Belle} \quad (535 \cdot 10^6 B\bar{B} \text{ pairs}),$$

and  $S_{\text{HFAG}} = -28 \pm 26\%$  waiting for BaBar update !

## Other transitions:

- Results in all transitions,  $S(\phi\gamma)$  LHCb !!

$$S(\rho\gamma) = (\underbrace{0.01}_{m_D/m_b} + \underbrace{0.02}_{\text{LD WA}} + \underbrace{0.20}_{\text{soft } g} \pm 1.6)\% = (0.2 \pm 1.6)\%,$$

$$S(\omega\gamma) = (0.01 - 0.08 + 0.22 \pm 1.7)\% = (0.1 \pm 1.7)\%,$$

$$S(K^*\gamma) = -(2.9 - 0 + 0.6 \pm 1.6)\% = -(2.3 \pm 1.6)\%,$$

$$S(\bar{K}^*\gamma) = (0.12 + 0.03 + 0.11 \pm 1.3)\% = (0.3 \pm 1.3)\%,$$

$$S(\phi\gamma) = -(0.5 - 0 - 5.3 \pm 8.2) \times 10^{-2}\% = -(0.1 \pm 0.1)\%.$$

- $b \rightarrow d$  transitions CKM-hierarchy demands u-quark loops !

Problem because photon momentum  $q^2 = 0$

- Method devised to calculate it: key ingredients:

1. calculate off-shell photon  $q^2$  deep euclidian
2. use dispersion relation (make sure no subtraction terms)
3. estimate the non-perturbative  $s \in [0, \text{cont}]$  part of the spectral function with LCSR and  $s \in [\text{cont}, \infty[$  by analytic continuation above threshold

# Epilogue

Efforts in

1. form factors from LCSR
2. distribution amplitudes (DA)
3. new methods (light-quark loops)

led to

- a. CKM-determination (tree-level)  $|V_{ub}|$
- b. "CKM-determination" (loop-induced)  $|V_{td}/V_{ts}|$
- c. plenty observables  $B \rightarrow V\gamma$  for NP  
most exciting  $S_{V\gamma}$  time-dep. CP asymmetry
- d. The future:  $B \rightarrow K^*l^+l^-$  (LHCb), even more obsevables, continue programme, hadr. param. and new methods valuable

In particular NP associated prod  $\gamma_R$  in  $B \rightarrow V\gamma$

$$\mathcal{B}(B \rightarrow V\gamma) \sim 1 + (R/L)^2 \quad S_{V\gamma} \sim R/L$$

Look at all possible observables ! Do not get discouraged by  $\mathcal{B}(b- > X_s)_{\text{incl}}$  !

