

Spin discrimination of Higgs Portal DM at the Muon Collider

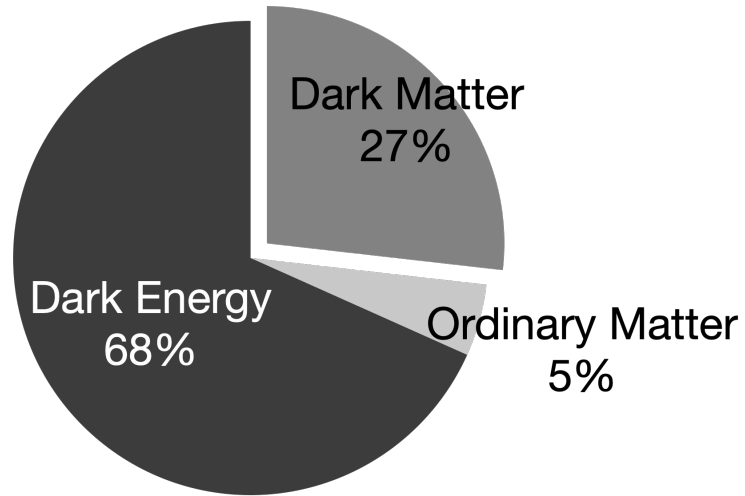
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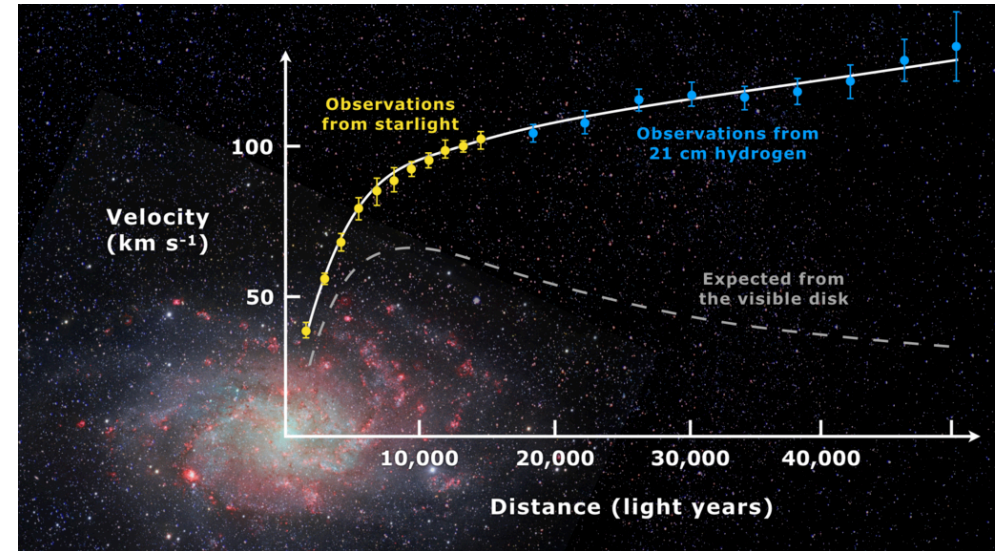
KIAS

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Introduction

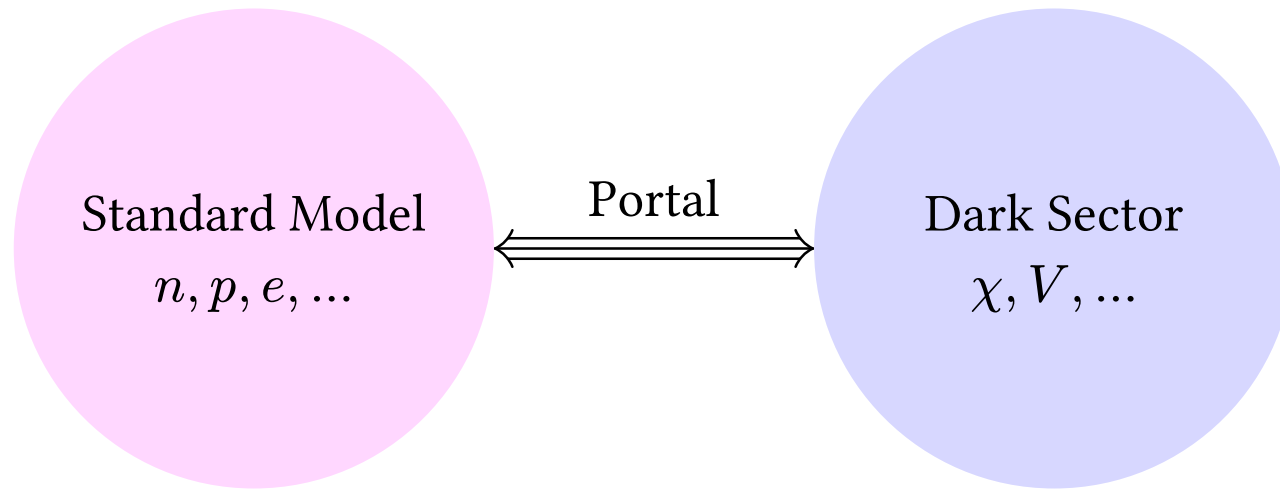


$$\Omega_{\text{DM}} h^2 = 0.1188 \pm 0.0022$$



Higgs Portal DM

- Portal model



- The SM particles can interact with the dark matter through the mediator particles.
- In the Higgs Portal model case, the mediator particles are the Higgs bosons.
- DM spin: 0 (SDM), 1/2 (FDM), 1 (VDM)

Higgs Portal DM

- SDM: $Z_2 : S \rightarrow -S, \langle S \rangle = 0$

$$\mathcal{L}_{\text{SDM}} = \frac{1}{2} \partial_\mu S \partial^\mu S - \frac{1}{2} m_0^2 S^2 - \lambda_{HS} H^\dagger H S^2 - \frac{\lambda_S}{4!} S^4$$

$$\mathcal{L}_{\text{SDM}} \supset -h \left(\sum_f \frac{m_f}{v_h} \bar{f} f - \frac{2m_W^2}{v} W_\mu^+ W^{-\mu} - \frac{m_Z^2}{v} Z_\mu Z^\mu \right) - \lambda_{HS} v h S^2.$$

- FDM: $Z_2 : \chi \rightarrow -\chi, \langle S \rangle = v_s$

$$\mathcal{L}_{\text{FDM}} = \bar{\chi} (i \not{\partial} - m_\chi - g_\chi S) \chi + \frac{1}{2} \partial_\mu S \partial^\mu S - \frac{1}{2} m_0^2 S^2 - \lambda_{HS} H^\dagger H S^2 - \mu_{HS} S H^\dagger H - \mu_0^3 S - \frac{\mu_S}{3!} S^3 - \frac{\lambda_S}{4!} S^4$$

$$\mathcal{L}_{\text{FDM}} \supset -(H_1 \cos \alpha + H_2 \sin \alpha) \left(\sum_f \frac{m_f}{v_h} \bar{f} f - \frac{2m_W^2}{v_h} W_\mu^+ W^{-\mu} - \frac{m_Z^2}{v_h} Z_\mu Z^\mu \right) + g_\chi (H_1 \sin \alpha - H_2 \cos \alpha) \bar{\chi} \chi$$

$$\begin{pmatrix} h \\ S \end{pmatrix} = \begin{pmatrix} \cos \alpha & \sin \alpha \\ -\sin \alpha & \cos \alpha \end{pmatrix} \begin{pmatrix} H_1 \\ H_2 \end{pmatrix}$$

Higgs Portal DM

- VDM: Dark gauge $U(1)_X$, Dark Higgs field Φ , $D_\mu \Phi = (\partial_\mu - ig_V Q_\Phi V_\mu) \Phi$

$$\mathcal{L}_{\text{VDM}} = -\frac{1}{4}V_{\mu\nu}V^{\mu\nu} + D_\mu \Phi^\dagger D^\mu \Phi - \lambda_\Phi \left(\Phi^\dagger \Phi - \frac{v_\Phi^2}{2} \right)^2 - \lambda_{H\Phi} \left(H^\dagger H - \frac{v_H^2}{2} \right)^2 \left(\Phi^\dagger \Phi - \frac{v_\Phi^2}{2} \right)$$

$$\mathcal{L}_{\text{VDM}} \supset -(H_1 \cos \alpha + H_2 \sin \alpha) \left(\sum_f \frac{m_f}{v_h} \bar{f} f - \frac{2m_W^2}{v_h} W_\mu^+ W^{-\mu} - \frac{m_Z^2}{v_h} Z_\mu Z^\mu \right) - \frac{1}{2} g_V m_V (H_1 \sin \alpha - H_2 \cos \alpha) V_\mu V^\mu$$

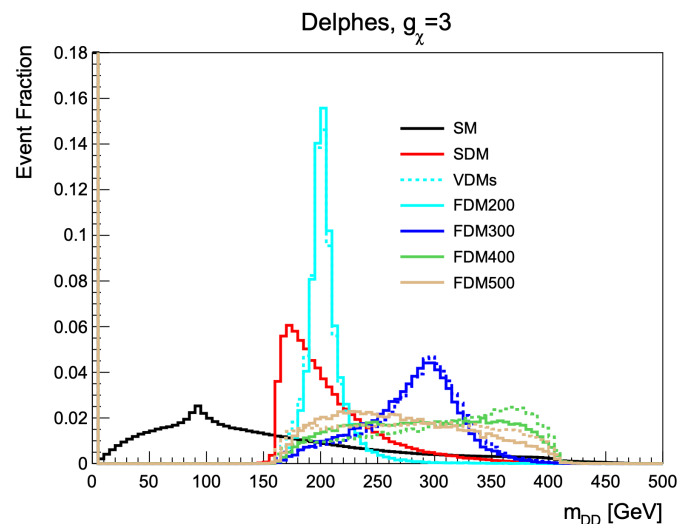
- $H_1 = h$, $M_h = 125$ GeV, $H_2 = H$, $M_H > 125$ GeV
- Can we distinguish the FDM and VDM at the collider?

$$\begin{pmatrix} h \\ S \end{pmatrix} = \begin{pmatrix} \cos \alpha & \sin \alpha \\ -\sin \alpha & \cos \alpha \end{pmatrix} \begin{pmatrix} H_1 \\ H_2 \end{pmatrix}$$

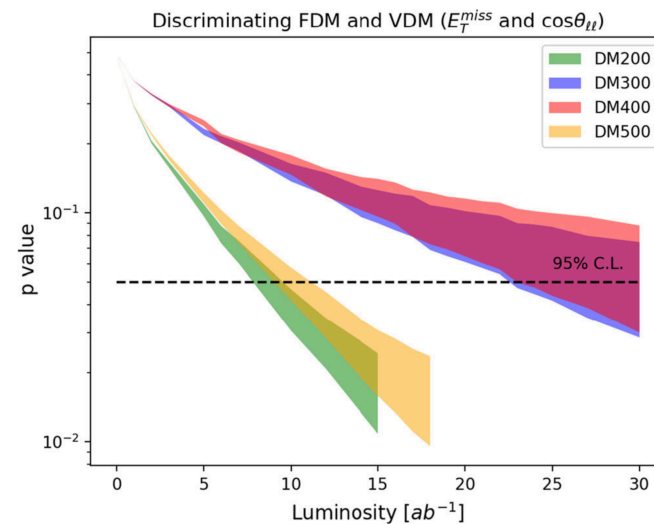
Previous study

- Previous studies of spin discrimination between FDM and VDM were conducted at the ILC and LHC.

ILC [1]



FCC [2]



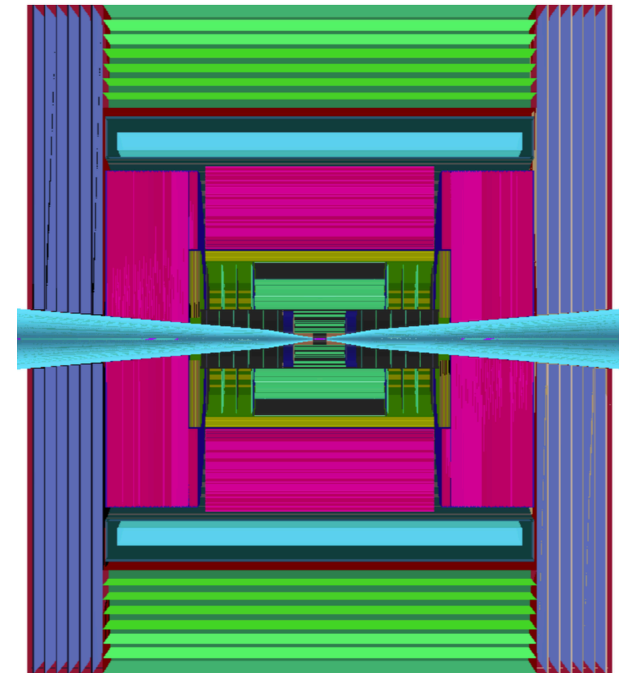
- SDM vs FDM discrimination is possible
- FDM vs VDM discrimination is not possible
- Both SDM vs FDM and FDM vs VDM discrimination are possible
- Only for $M_H < 500$ GeV

Why Muon Collider?

- e^+e^- collider: Low backgrounds, low CM energy.
- $pp(p\bar{p})$ collider: High CM energy, huge backgrounds.
- $\mu^+\mu^-$ collider: **Low backgrounds, high CM energy.**

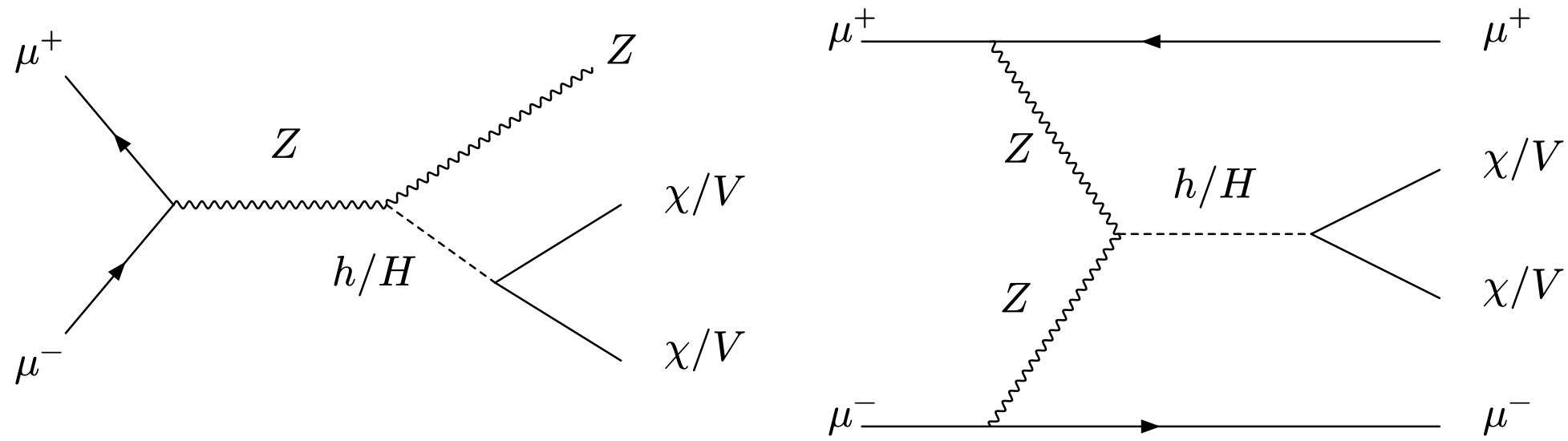
→ We can study TeV scale H_2 .

- CM energy: $\sqrt{s} = 3, 10$ TeV, $\mathcal{L}_{\text{tot}} : 1, 10 \text{ ab}^{-1}$
- Angular coverage:
 - Central: $|\eta| < 2.5$
 - Forward: $|\eta(\mu)| < 8$



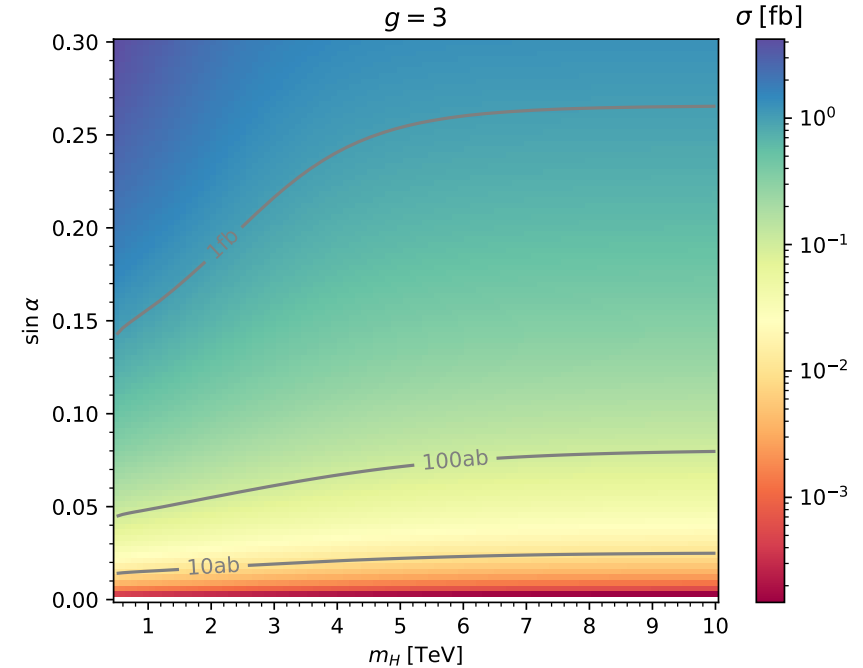
Muon collider detector concept [3]

Production of FDM and VDM at MuC



Production of FDM and VDM at MuC

- FDM production cross section at $\sqrt{s} = 10$ TeV MuC.
- $\mu^+\mu^- \rightarrow \mu^+\mu^-\chi\bar{\chi}$, $\eta(\mu) < 8$
- $M_\chi = 150$ GeV, $g_\chi = 3$
- Increasing $\sin \alpha$ gives an increasing σ .
- Increasing M_H gives a decreasing σ .
 - For large M_H , σ becomes constant.
- SM cross section: $\sigma(\mu\mu\nu\bar{\nu}) \sim 2$ pb
- g_V : $\sigma_{\text{VDM}}(g_V) = \sigma_{\text{FDM}}(g_\chi = 1, 3)$

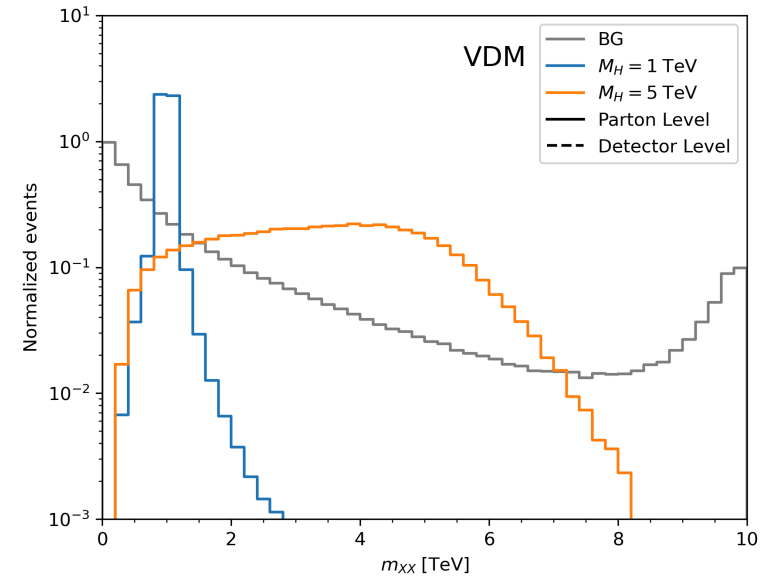
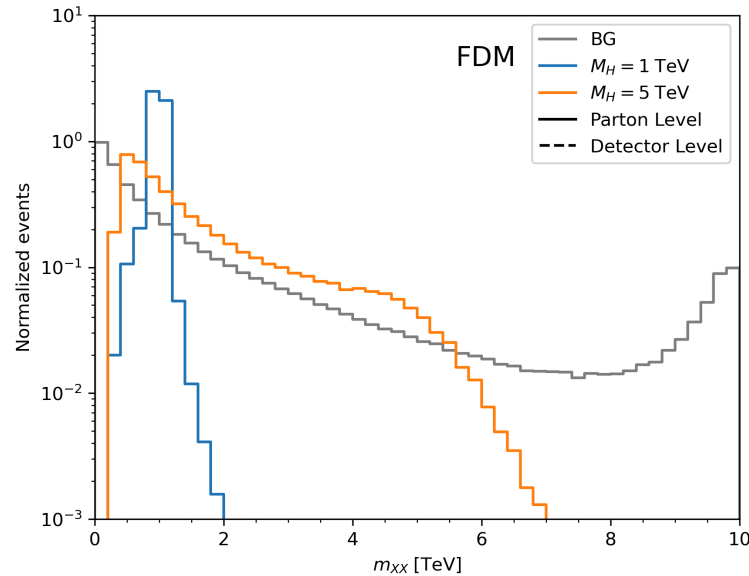


Production of FDM and VDM at MuC

- Recoil Mass

$$m_{XX}^2 = s + m_{\mu\mu}^2 - 2E_{\mu\mu}\sqrt{s}$$

- m_{XX} : Invariant mass of the DM pair (or neutrino pair)
- $m_{\mu\mu}, E_{\mu\mu}$: Invariant mass and energy of the muon pair

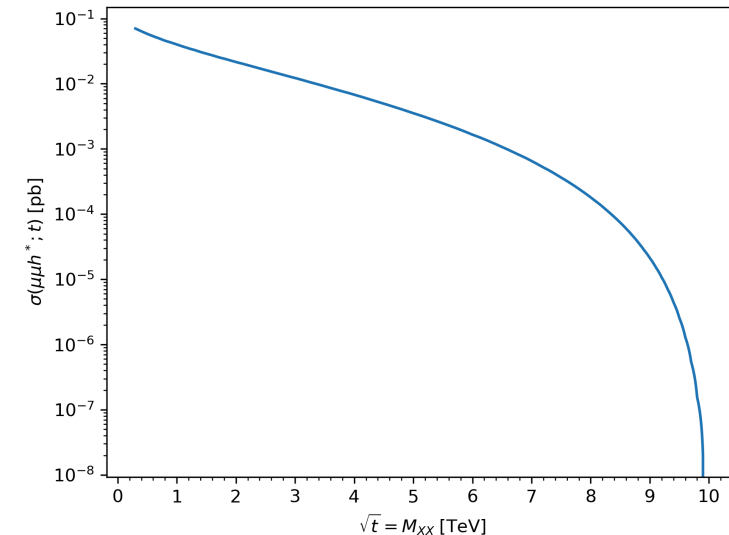
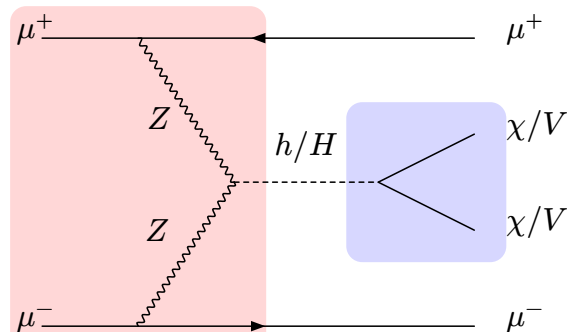


Production of FDM and VDM at MuC

The DM production cross section at MuC is given by:

$$\sigma(\mu\mu \rightarrow \mu\mu XX) = \int_{4m_X^2}^s dt \quad \sigma(\mu\mu h^*; t) \quad \left| \frac{1}{t - m_h^2 + im_h\Gamma_h} - \frac{1}{t - m_H^2 + im_H\Gamma_H} \right|^2 \times \frac{2t}{\pi} \quad \Gamma_{h^* \rightarrow XX}(t) \quad s_\alpha c_\alpha$$

- t : Squared recoil mass m_{XX}^2
- $\sigma(\mu\mu h^*; t)$: production cross section with mediator scalar mass $m_\phi = \sqrt{t}$
- $\Gamma_{h^* \rightarrow XX}(t)$: Decay width of the mediator scalar to the DM pair. $XX = \chi\bar{\chi}, VV$
- $\Gamma(H \rightarrow \bar{\chi}\chi) = g_\chi^2 \frac{M_H}{8\pi} (1 - \tau_\chi)^{\frac{3}{2}}$
- $\Gamma(H \rightarrow VV) = \frac{g_V m_V}{2048\pi} \frac{M_H^3}{m_V^2} (1 - \tau_V)^{\frac{1}{2}} (4 - 4\tau_V + 3\tau_V^2)$

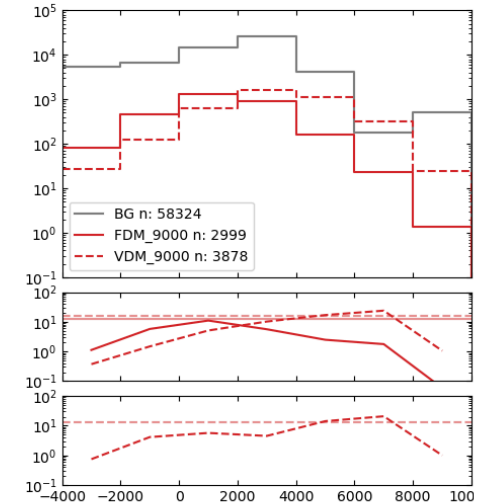
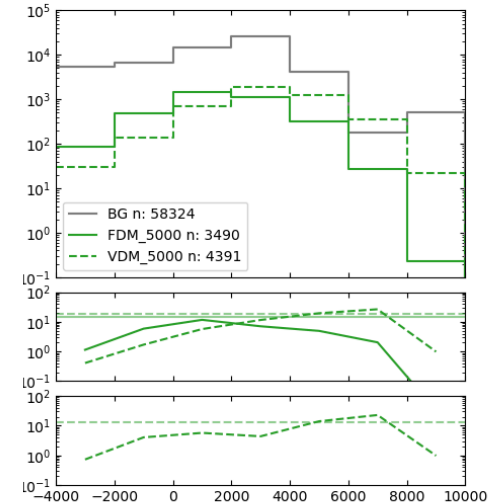
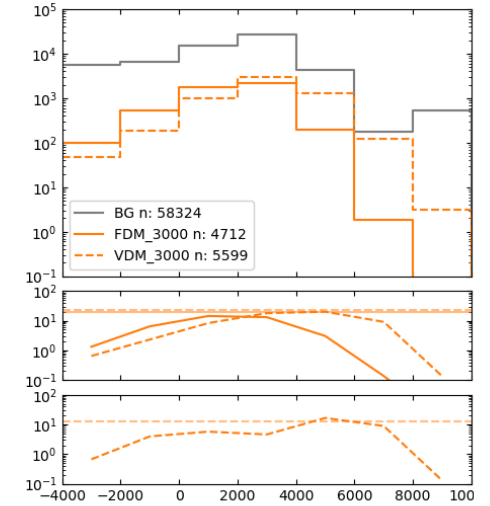
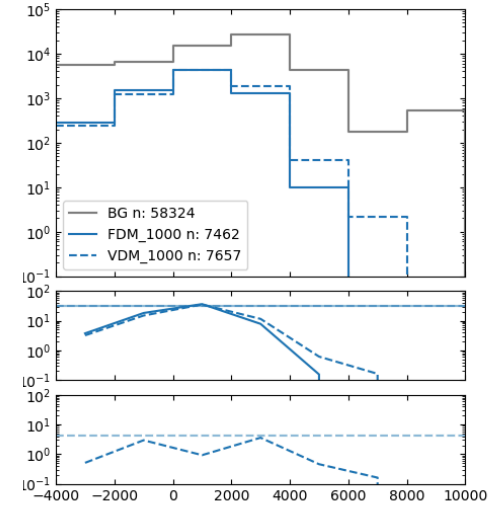


ML approach

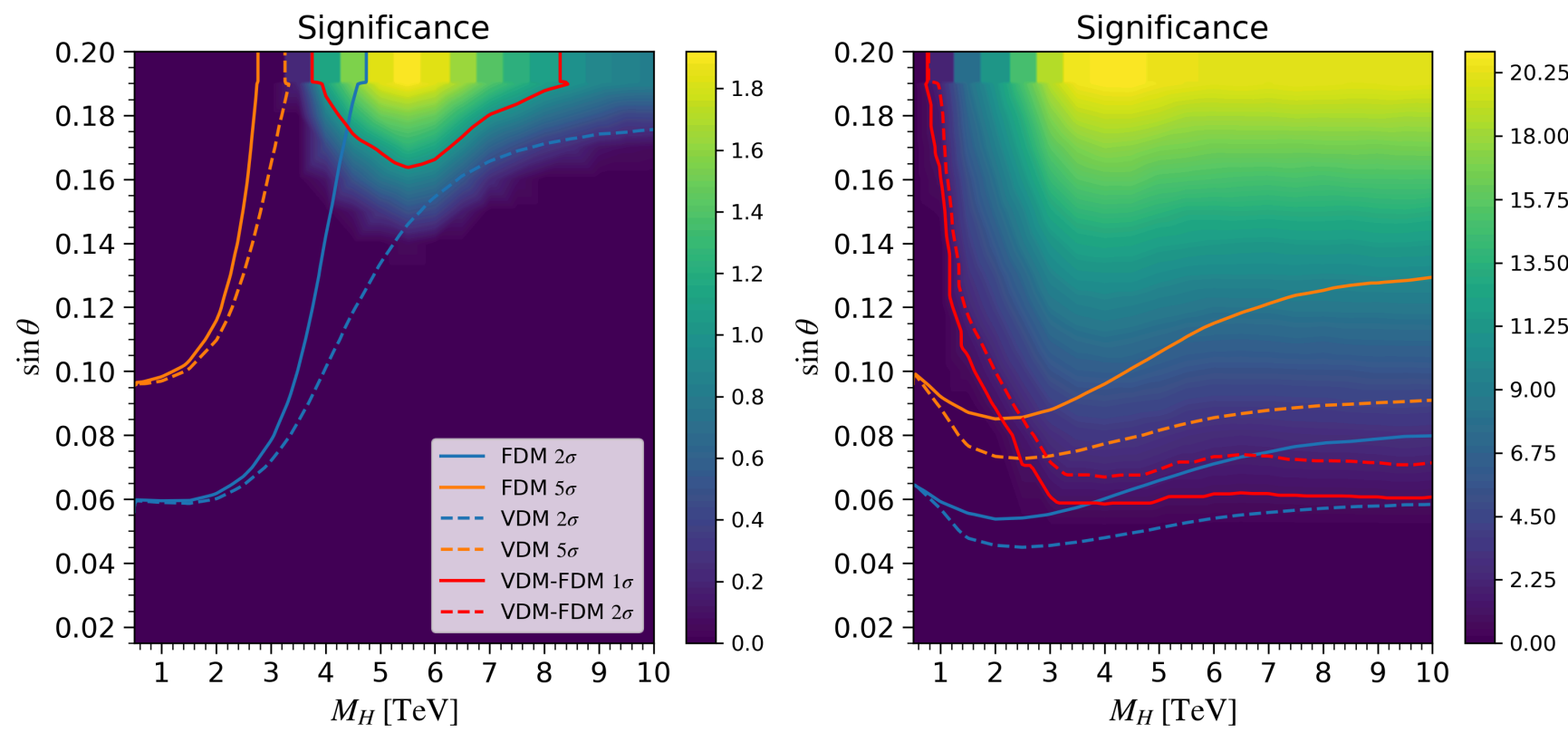
- For a given $m_{XX} (= \sqrt{t})$, the distributions of all kinematic variables $(P^T(\mu), \eta(\mu), m_{\mu\mu}, \dots)$ are independent of the choice of BSM parameters (spin, M_H , $\sin \alpha$, g , $\dots$).
- If the NN can discriminate signal and background for the full range of t , this network can be used for any parameter choice.
- Training
 - Kinematic Variables:
 - $E(\mu), \eta(\mu), \phi(\mu)$ of $\mu_0, \mu_1, \mu\mu, \Delta R(\mu_0, \mu_1), \Delta\phi(\mu_0, \mu_1), \Delta\eta(\mu_0, \mu_1)$
 - $p_T^{\text{miss}}, \phi^{\text{miss}}, M(\mu^+\mu^-), m_{XX}$
 - 100k events each of signal and background, with flat m_{XX} distribution.
 - 70%, 15%, 15% for training, validation, and test

Results

- Network
 - FC Layer [128, 64, 32]
 - 1 Output
 - Layer Norm, ReLU, Dropout(0.25)
- Signal efficiency and background refaction
 - $\varepsilon_{\text{Signal}} = 0.6$
 - $1/\varepsilon_{\text{Background}} = 483$
- Spin discrimination
 - $$\delta\chi^2 = \sum_i^{\text{nbin}} \left(\frac{N_i^{\text{VDM}} - N_i^{\text{FDM}}}{\sqrt{N_i^{\text{VDM} + \text{SM}}}} \right)^2$$



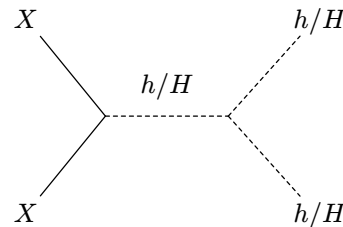
Results



Dark Matter Phenomenology

- Relic Density Constraints

- ▶ $\Omega h^2 \propto \frac{1}{\langle \sigma v \rangle}$
 - ▶ $\lambda_{hhH}, \lambda_{hHH}, \lambda_{HHH}$



- Direct Detection

- ▶ To avoid DD constraints, small σ_{DN} is needed
 - ▶ Small g_X , small $|\sin \alpha|$, large m_H .

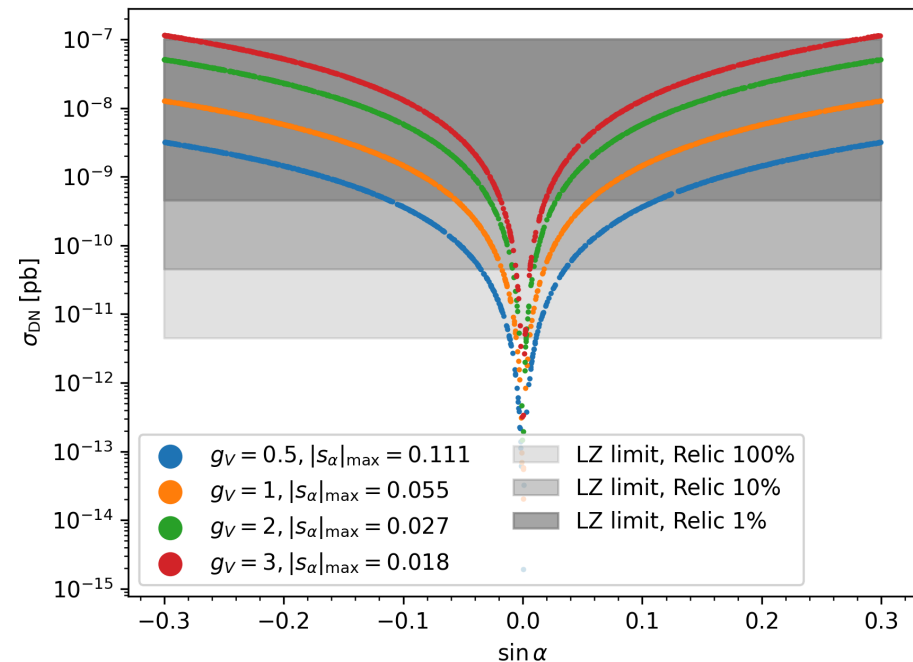
- Other Explanations

- ▶ Uncertainty in local DM density
 - ▶ Excited Dark Matter [4]

$$- \mathcal{L} \supset \sum \bar{\chi}_i (i \not{D} - m_i - y_i S) \chi_i - [\bar{\chi}_1 (y_S + i y_p \Gamma_5) S \chi_2 + \text{h.c}]$$

$$- \frac{1}{4} V_{\mu\nu} V^{\mu\nu} + D_\mu S D^\mu S - \frac{1}{2} m_S^2 S^2 - V(H, S)$$

$$- \chi_2 \rightarrow \chi_1 \gamma_D$$



Conclusion

- We study the spin discrimination of Higgs Portal DM at the Muon Collider.
- At the muon collider, the recoil mass m_{XX} is a useful variable to measure dark matter properties.
- Due to the different growth of FDM and VDM with M_H , the FDM and VDM show different distributions in m_{XX} .
- For $g_\chi = 1$, the separation of FDM and VDM is limited and it is challenging to discriminate the spin of DM. For $g_\chi = 3$ with $\sin \alpha \gtrsim 0.5$, we can observe the signal above 2σ . For most of this region, we can discriminate the spin of DM.
- The parameter space is excluded by direct detection constraints. However, this may be evaded by introducing Excited Dark Matter.

Thank you for your attention!

References

- [1] T. Kamon, P. Ko, and J. Li, “Characterizing Higgs portal dark matter models at the ILC,” *Eur. Phys. J. C*, vol. 77, no. 9, p. 652, 2017, doi: 10.1140/epjc/s10052-017-5240-8.
- [2] B. Dutta, T. Kamon, P. Ko, and J. Li, “Prospects for discovery and spin discrimination of dark matter in Higgs portal DM models and their extensions at 100 TeV pp collider”, *Eur. Phys. J. C*, vol. 78, no. 7, p. 595, 2018, doi: 10.1140/epjc/s10052-018-6071-y.
- [3] C. Accettura and others, “Interim report for the International Muon Collider Collaboration (IMCC),” *CERN Yellow Rep. Monogr.*, p. 176, 2024, doi: 10.23731/CYRM-2024-002.
- [4] J. L. W.-I. P. S. Baek P. Ko, “(Paper in preparation).”