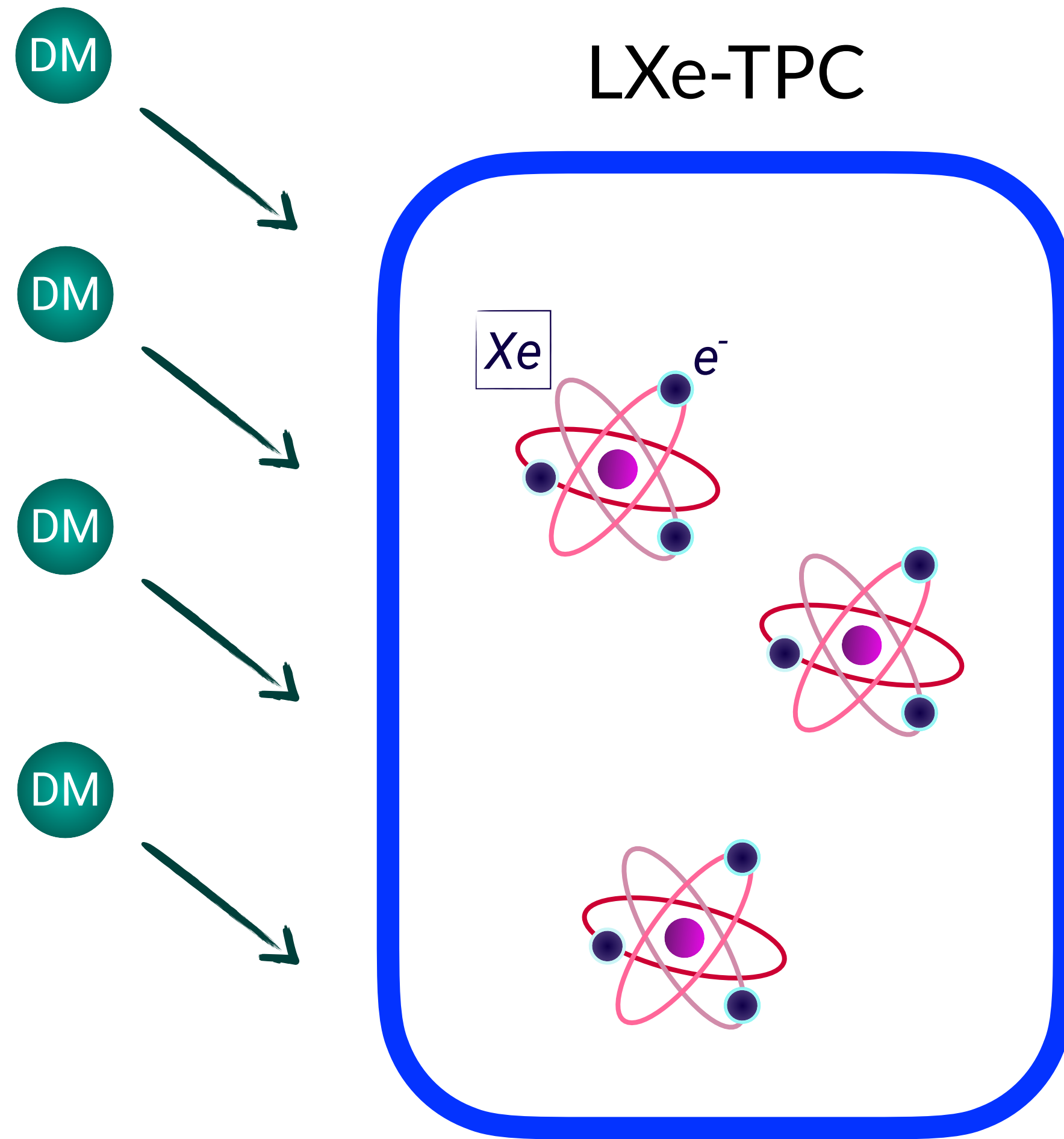


Future DM experiments (non-XLZD)

Christopher M^cCabe

Motivation: what does 'non-XLZD' mean?

My simple minded view of XLZD



Detector technology with:

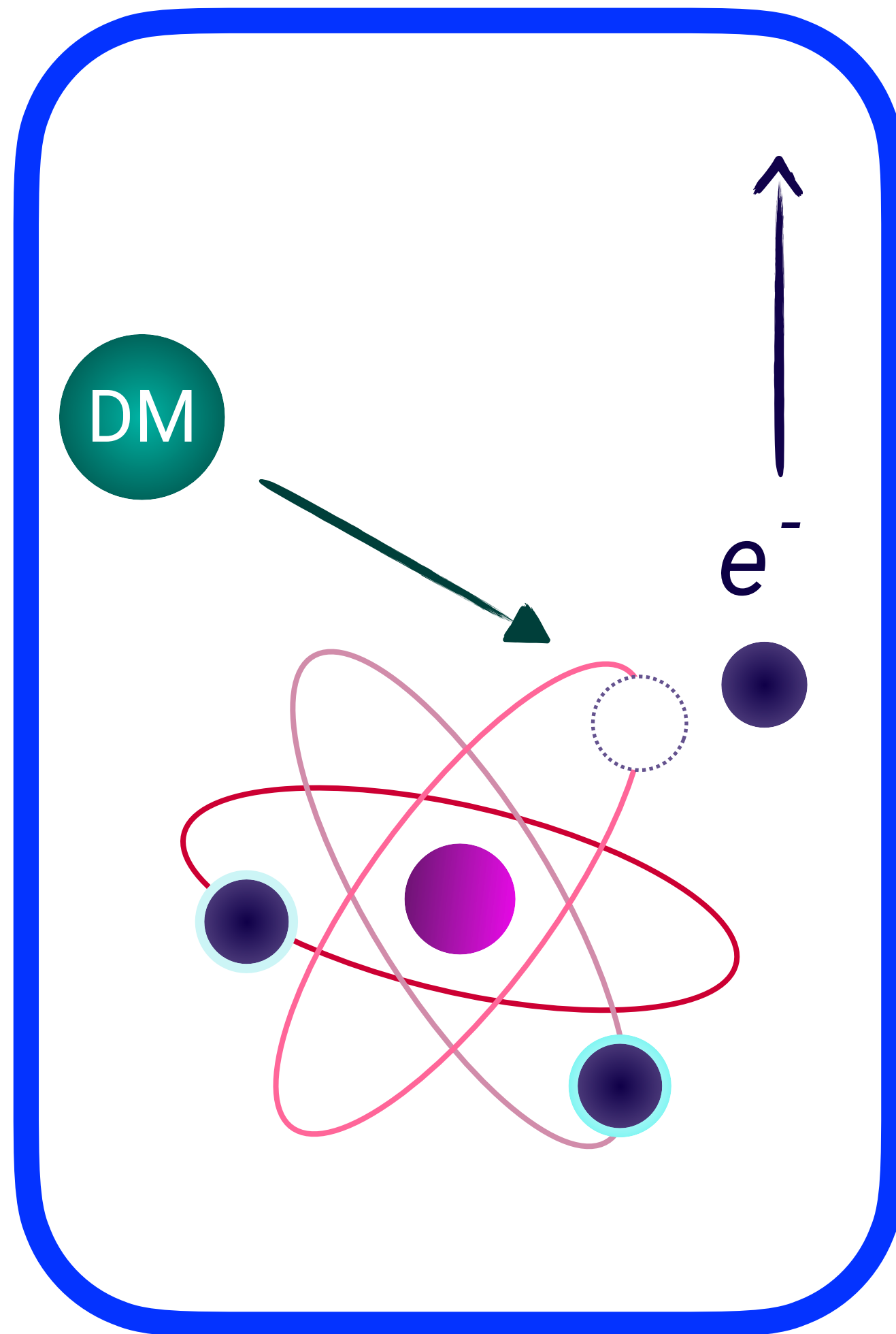
- Single-electron threshold
- Metre-scale
- Operate for a few years

These define boundaries of:

- Lightest dark matter mass
- heaviest dark matter mass

Lightest dark matter mass

LXe-TPC

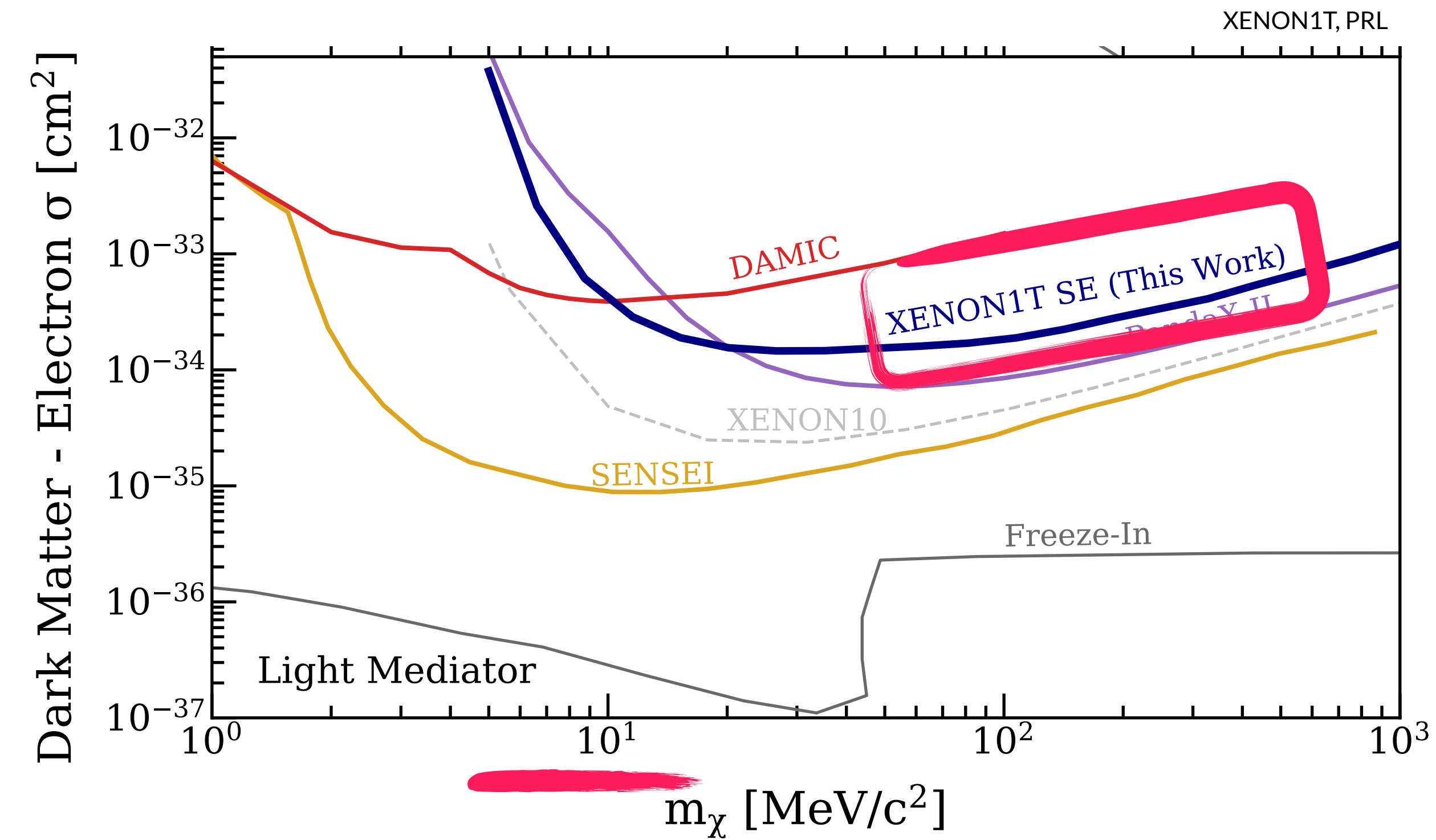


Signal generation if dark matter can deposit an energy above the electron binding energy

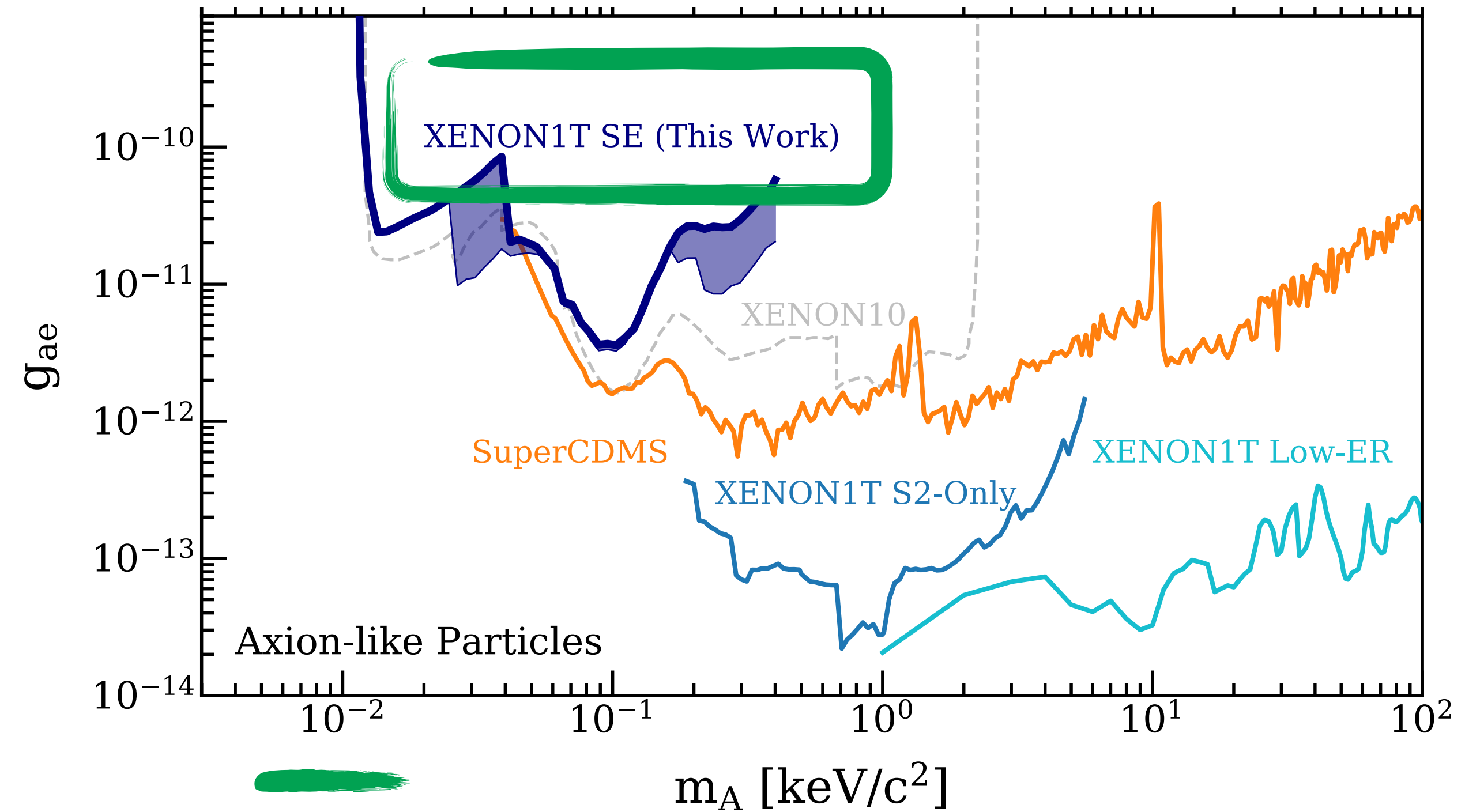
- Scattering: $E_{\text{DM}} \sim m_{\text{DM}} v_{\text{DM}}^2$
 $\sim 10 \text{ eV} \left(\frac{m_{\text{DM}}}{10 \text{ MeV}} \right)$
- Absorption: $E_{\text{DM}} \sim m_{\text{DM}}$
 $\sim 10 \text{ eV} \left(\frac{m_{\text{DM}}}{10 \text{ eV}} \right)$

Lightest dark matter mass

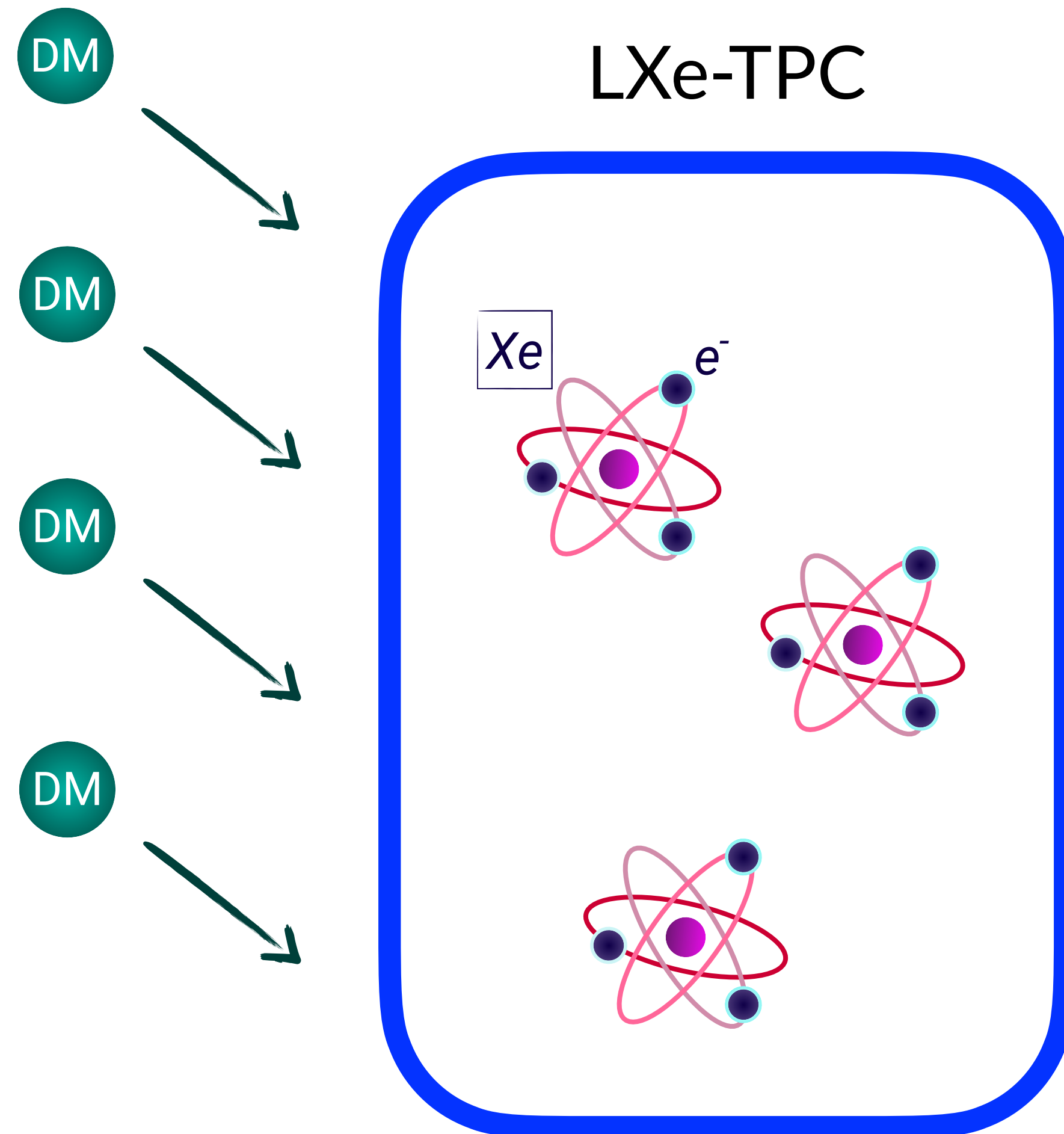
Scattering



Absorption



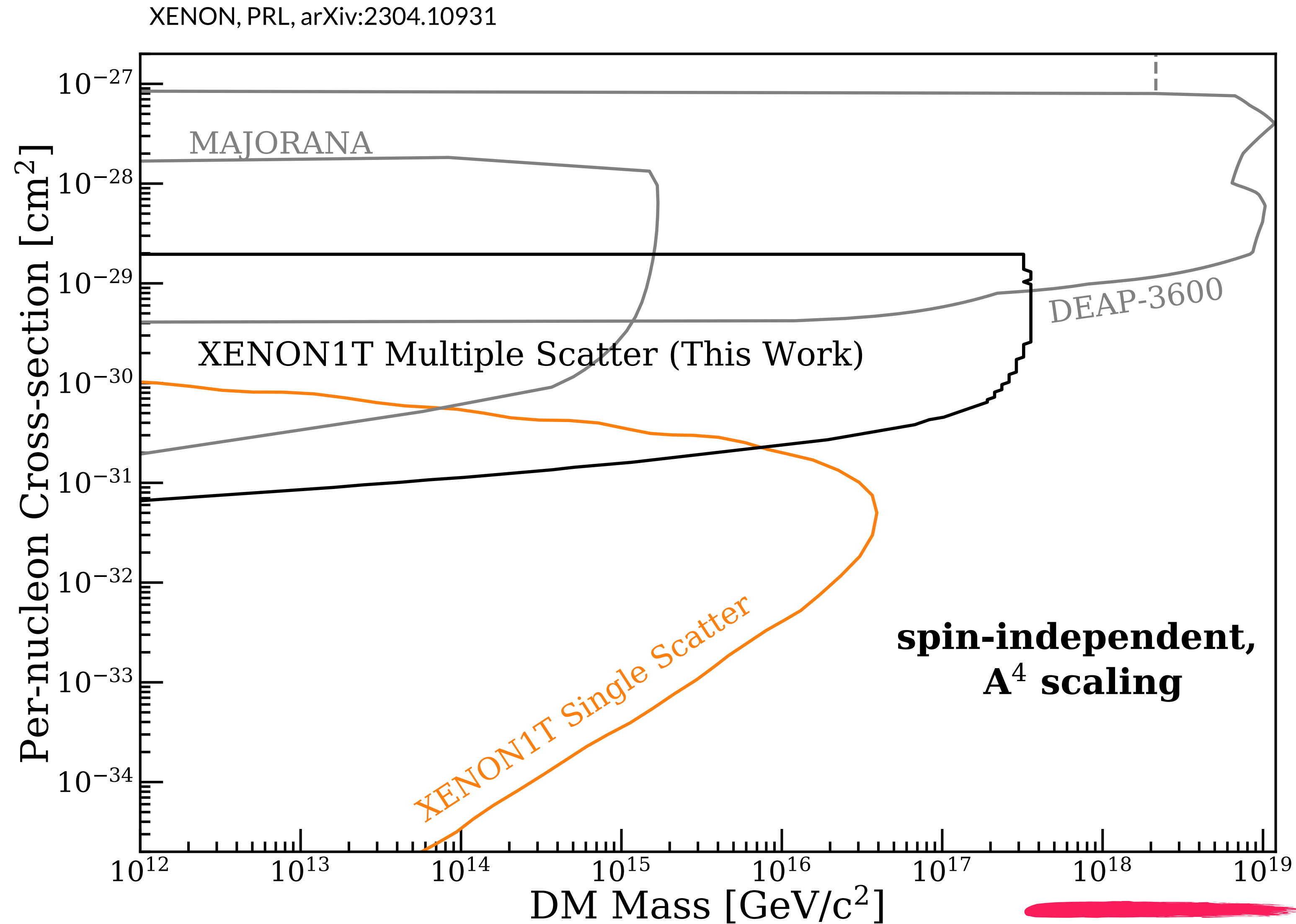
Heaviest dark matter mass



Signals generation if dark matter passes through the detector

$$\Phi_{\text{DM}} \sim 0.3 \left(\frac{10^{19} \text{ GeV}}{m_{\text{DM}}} \right) \text{m}^{-2} \text{yr}^{-1}$$

Heaviest dark matter mass



$$M_{\text{Pl}} \approx 10^{19} \text{ GeV}$$

XLZD is a broadband, multi-purpose particle dark matter experiment sensitive to dark matter from the 10 eV to Planck scale

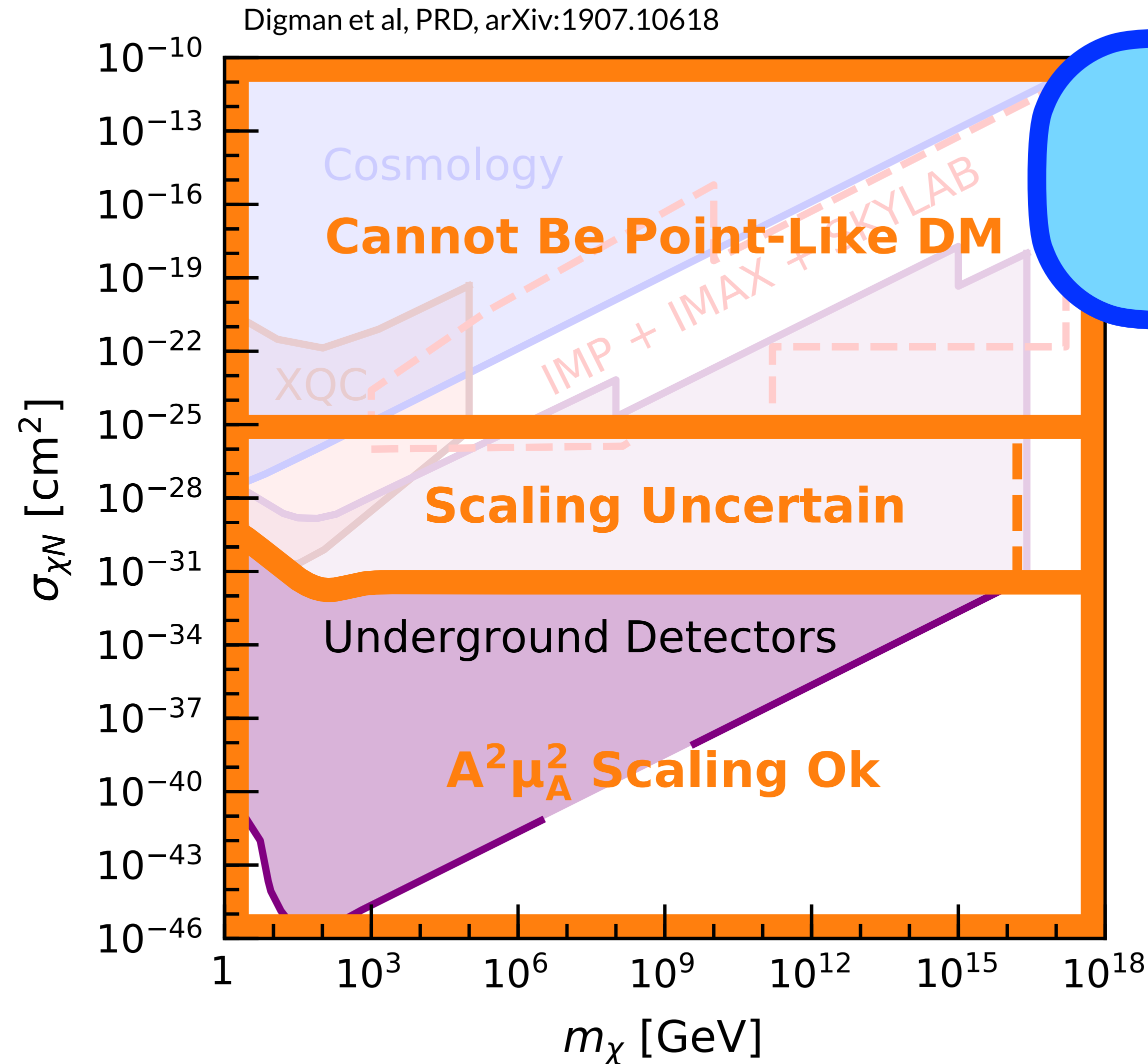
to me, non-XLZD means:

- 1. Beyond the Planck scale**
- 2. Below the eV scale**

1. Dark matter (particles) beyond the Planck scale



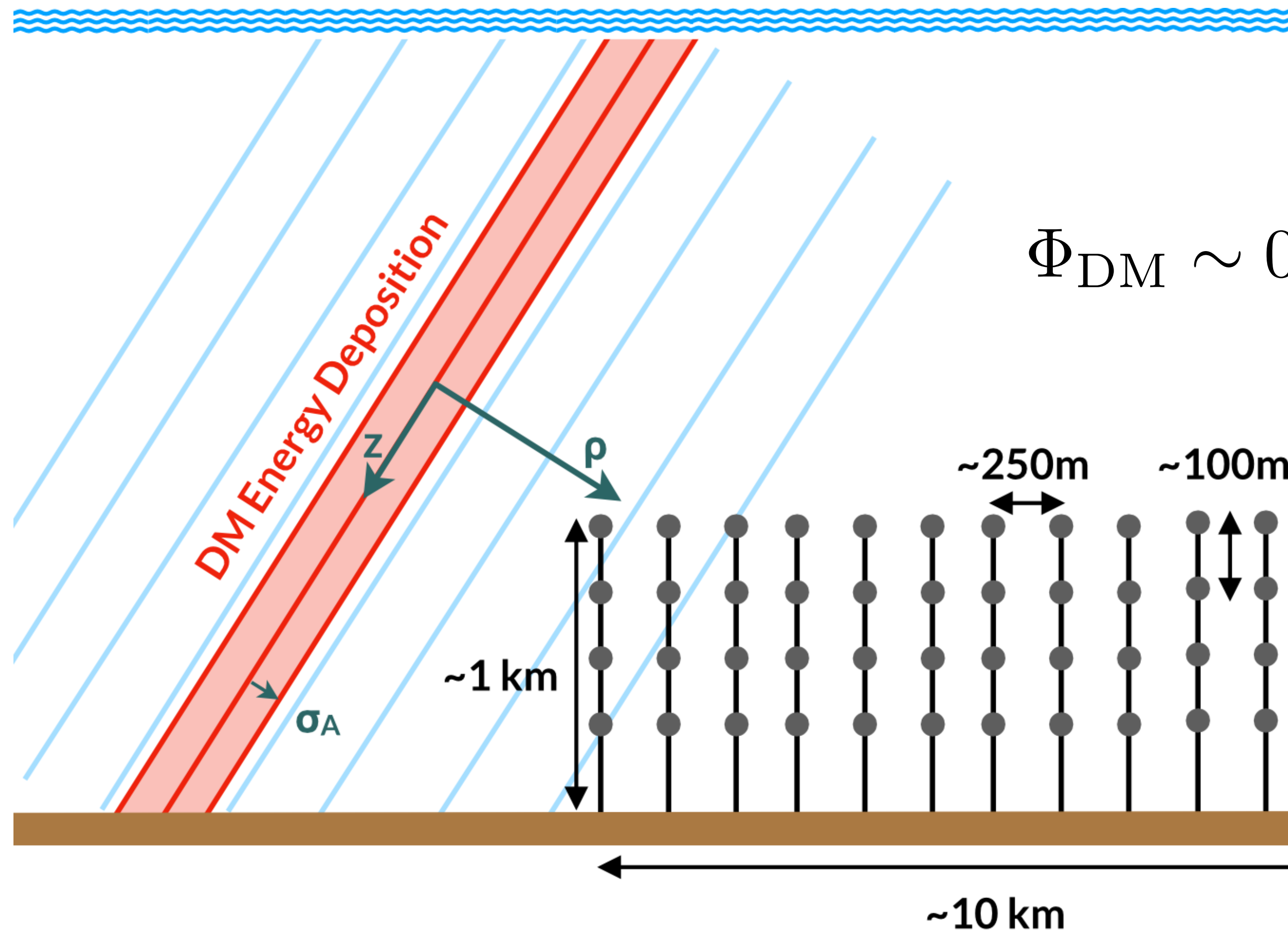
Theory landscape



When looking just beyond XLZD:

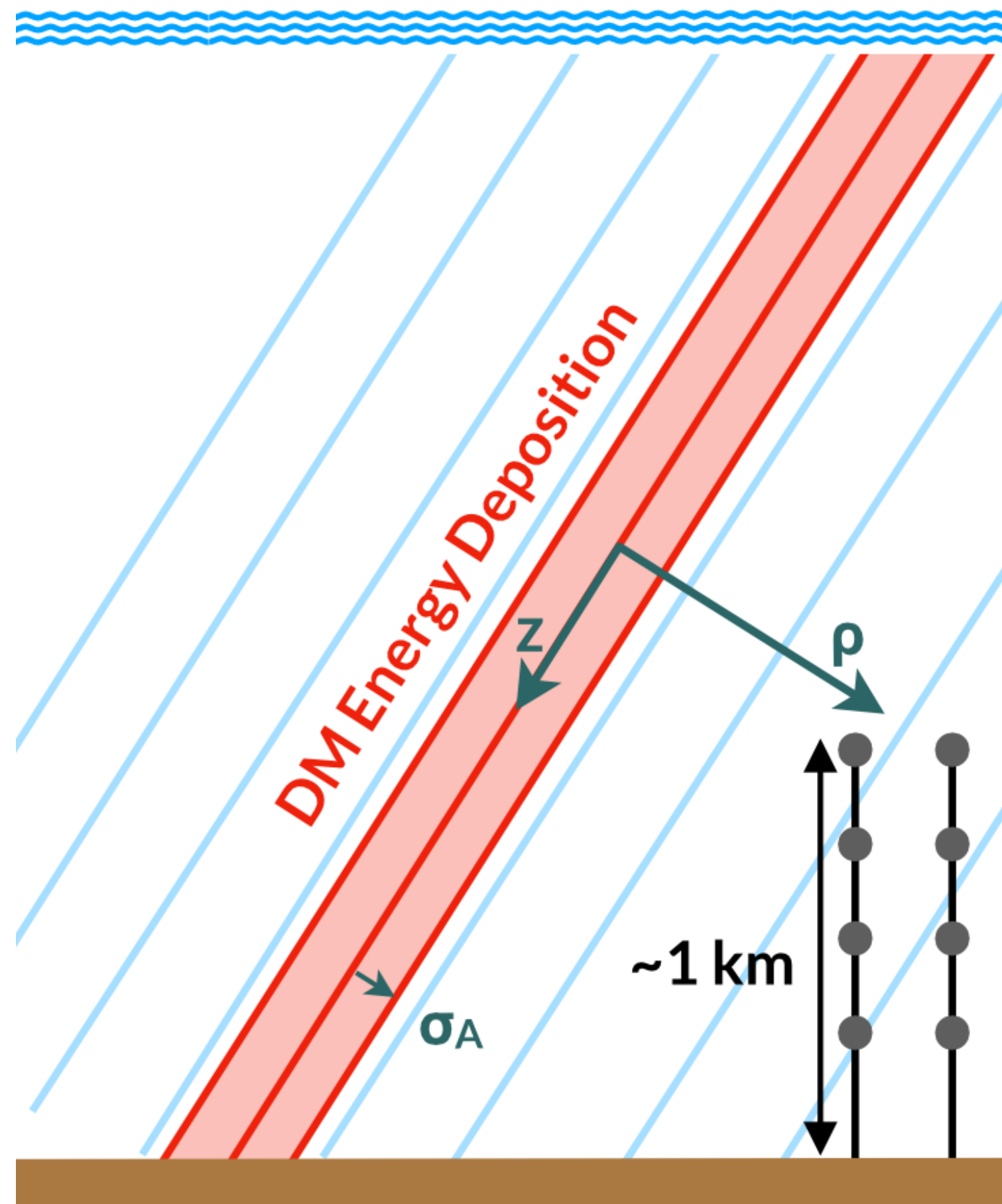
“For $\sigma > 10^{-25} \text{ cm}^2$, dark matter cannot be point-like. Contact interactions cannot achieve cross sections larger than the geometric cross section $\sigma \approx 4\pi r^2$. Dark matter with cross sections in this range must be composite.”

(Speculative) new proposal: large area instrument in seawater



$$\Phi_{\text{DM}} \sim 0.3 \left(\frac{10^{19} \text{ GeV}}{m_{\text{DM}}} \right) \text{ m}^{-2} \text{ yr}^{-1}$$

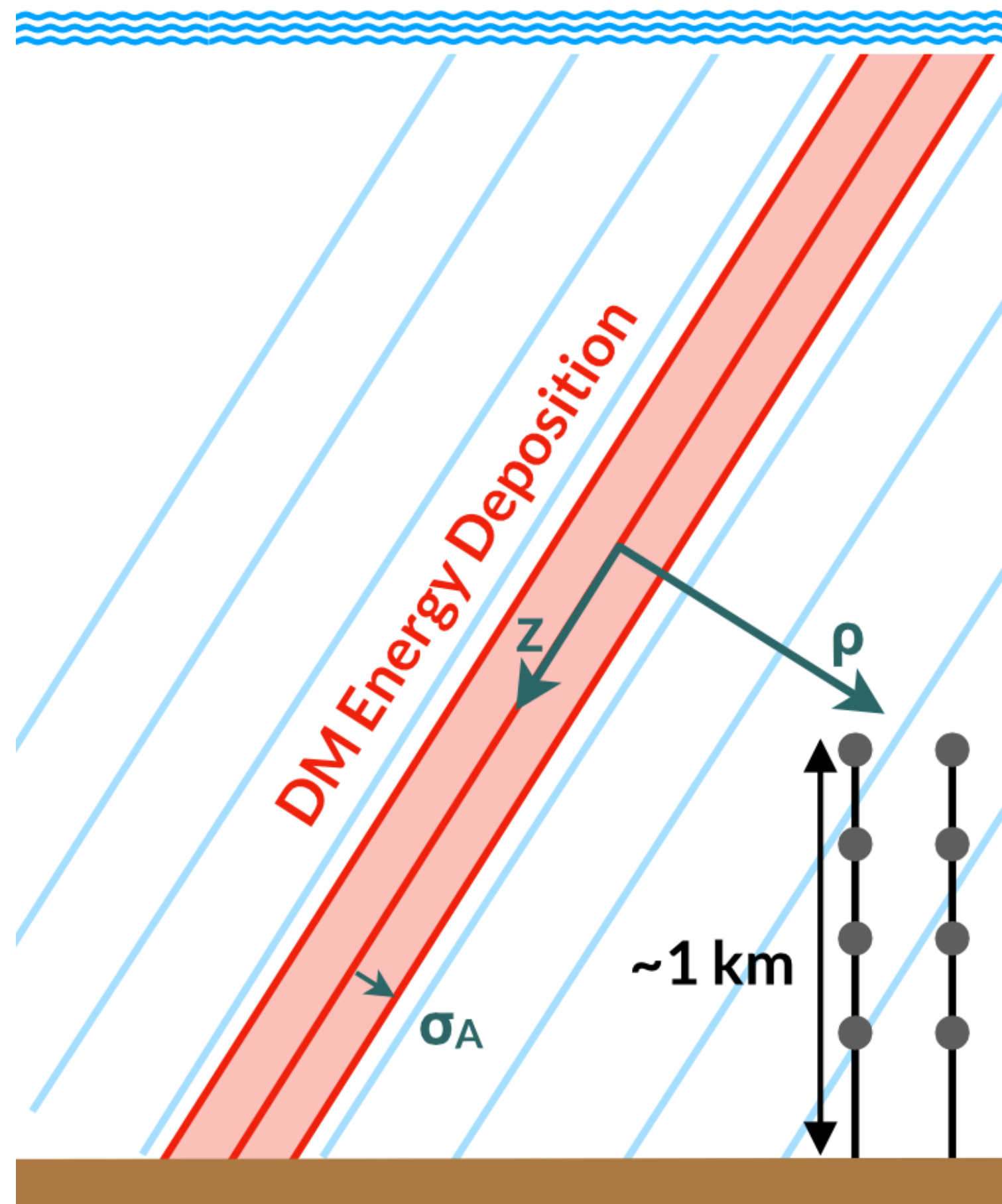
Track properties: very long and very thin



Dark matter scatters with every water nucleus in its path:

$$\lambda_\chi \simeq \underline{10^{-15} \text{ m}} \times (10^{-10} \text{ cm}^2) / \sigma_\chi$$

Track properties: very long and very thin



Dark matter scatters with every water nucleus in its path:

$$\lambda_\chi \simeq 10^{-15} \text{ m} \times (10^{-10} \text{ cm}^2) / \sigma_\chi$$

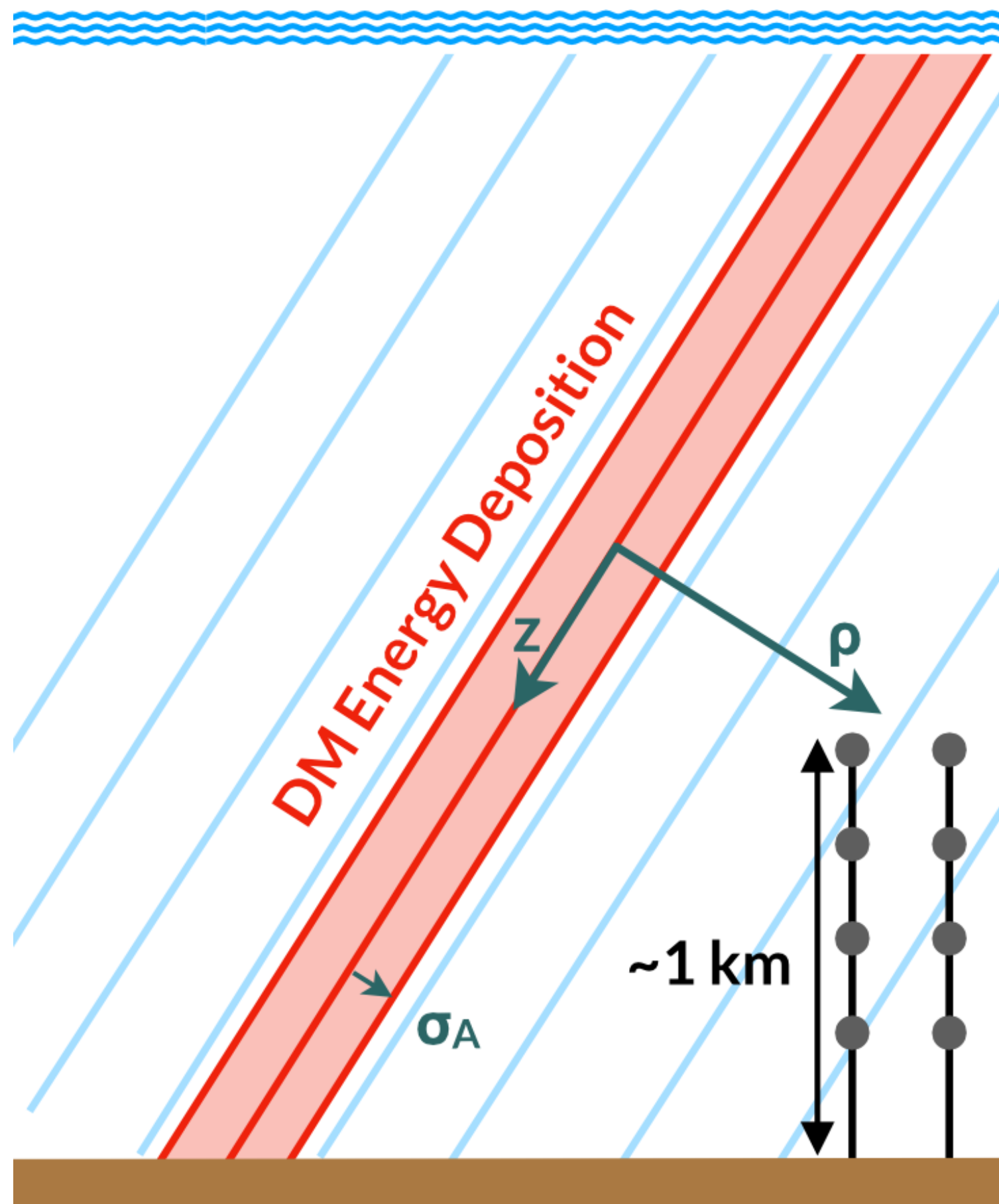
...but energy loss inefficient, so long tracks:

$$\frac{dE_\chi}{dz} = -\rho_{\text{sea}} \sigma_\chi v_\chi^2 \exp\left(-\frac{z}{\ell_{\text{sea}}}\right)$$

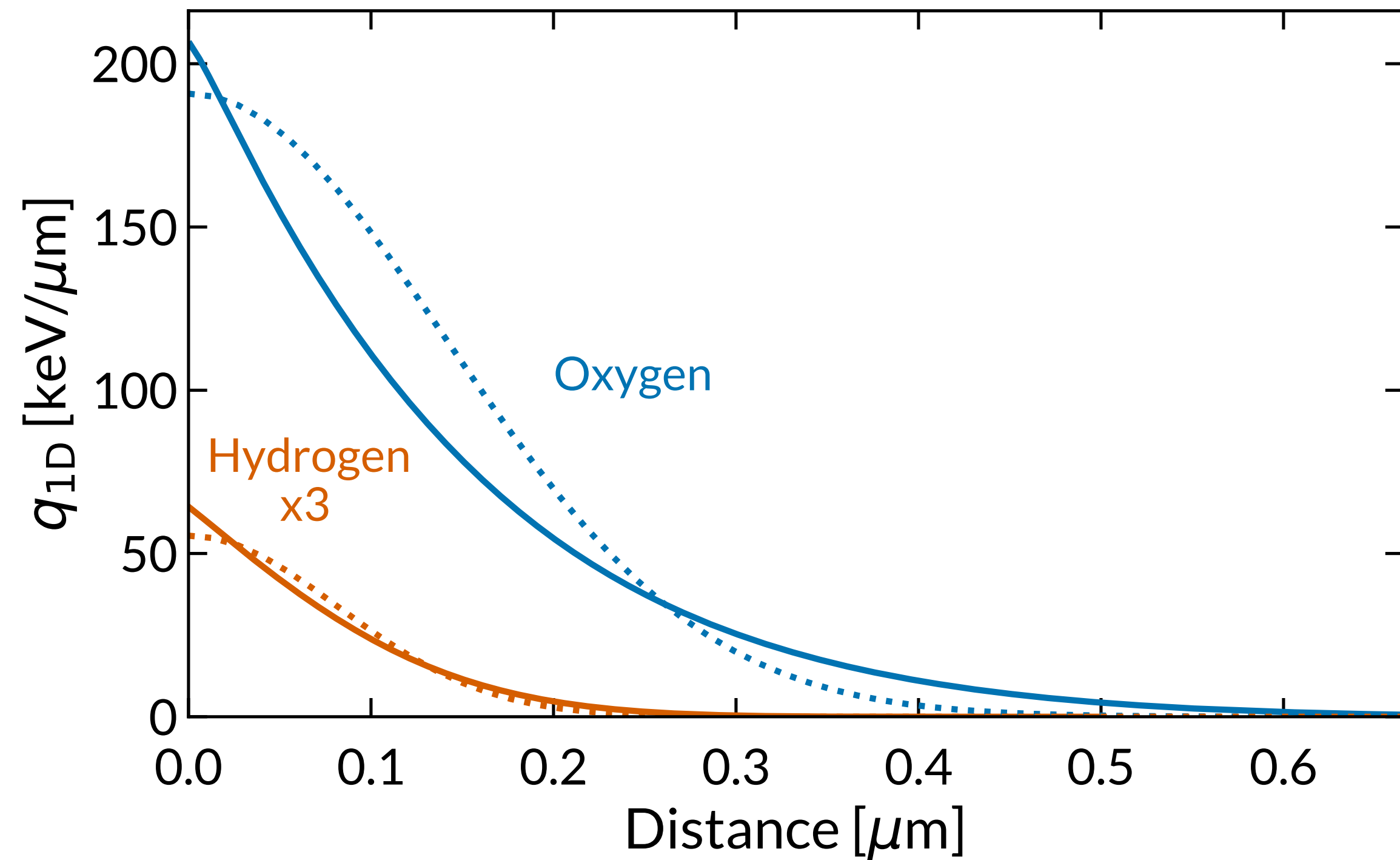
$$\ell_{\text{sea}} = \frac{m_\chi}{2\rho_{\text{sea}} \sigma_\chi}$$

$$\simeq 485 \text{ km} \times \left(\frac{m_\chi}{10^{-2} \text{ g}}\right) \left(\frac{10^{-10} \text{ cm}^2}{\sigma_\chi}\right)$$

Track properties: very long and very thin



Recoiling nuclei don't travel far:



Energy deposition in very long and thin cylinder

Energy deposition = sound waves

Thermo-acoustic model:

Rapid energy deposition creates localised heating, which instantaneously expands the seawater and produces a propagating acoustic pressure wave

$$\nabla^2 p - \frac{1}{c_s^2} \frac{\partial^2 p}{\partial t^2} = - \frac{\alpha}{c_p} \frac{\partial^2 q(\mathbf{r}, t)}{\partial t^2}$$

Speed of sound

Specific heat capacity

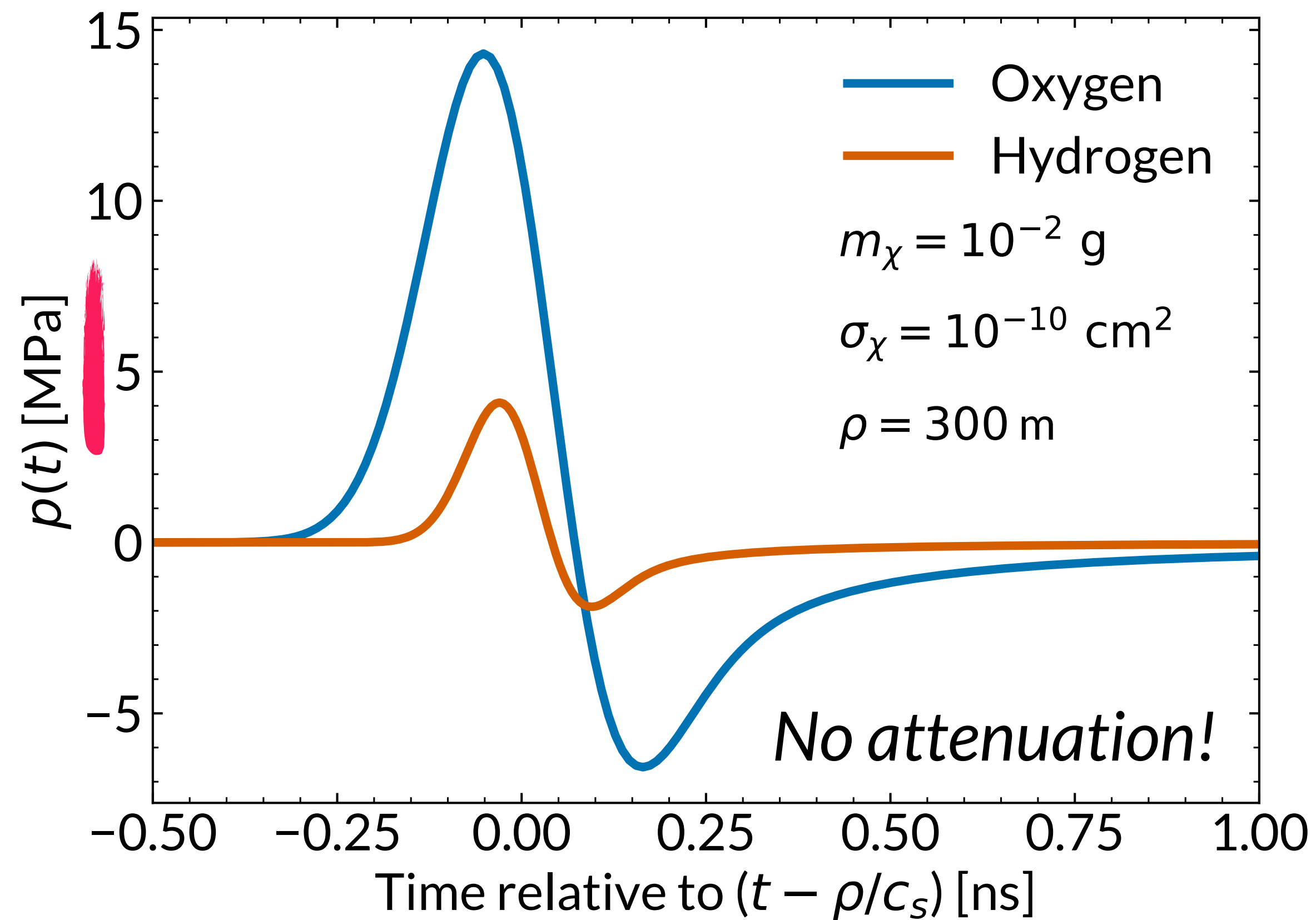
Thermal expansion coefficient

Energy deposition density

Energy deposition = sound waves

At first sight: extremely promising!

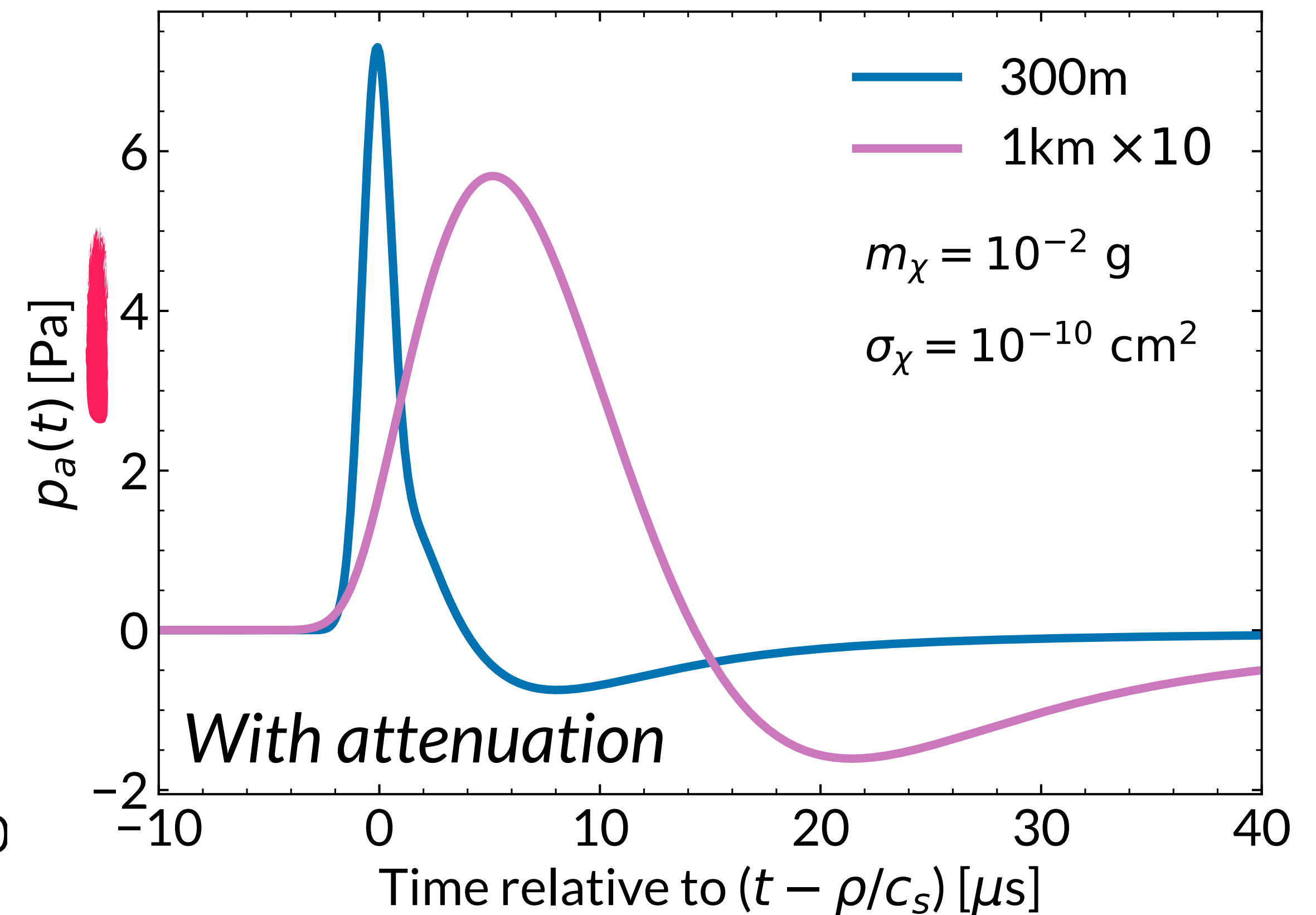
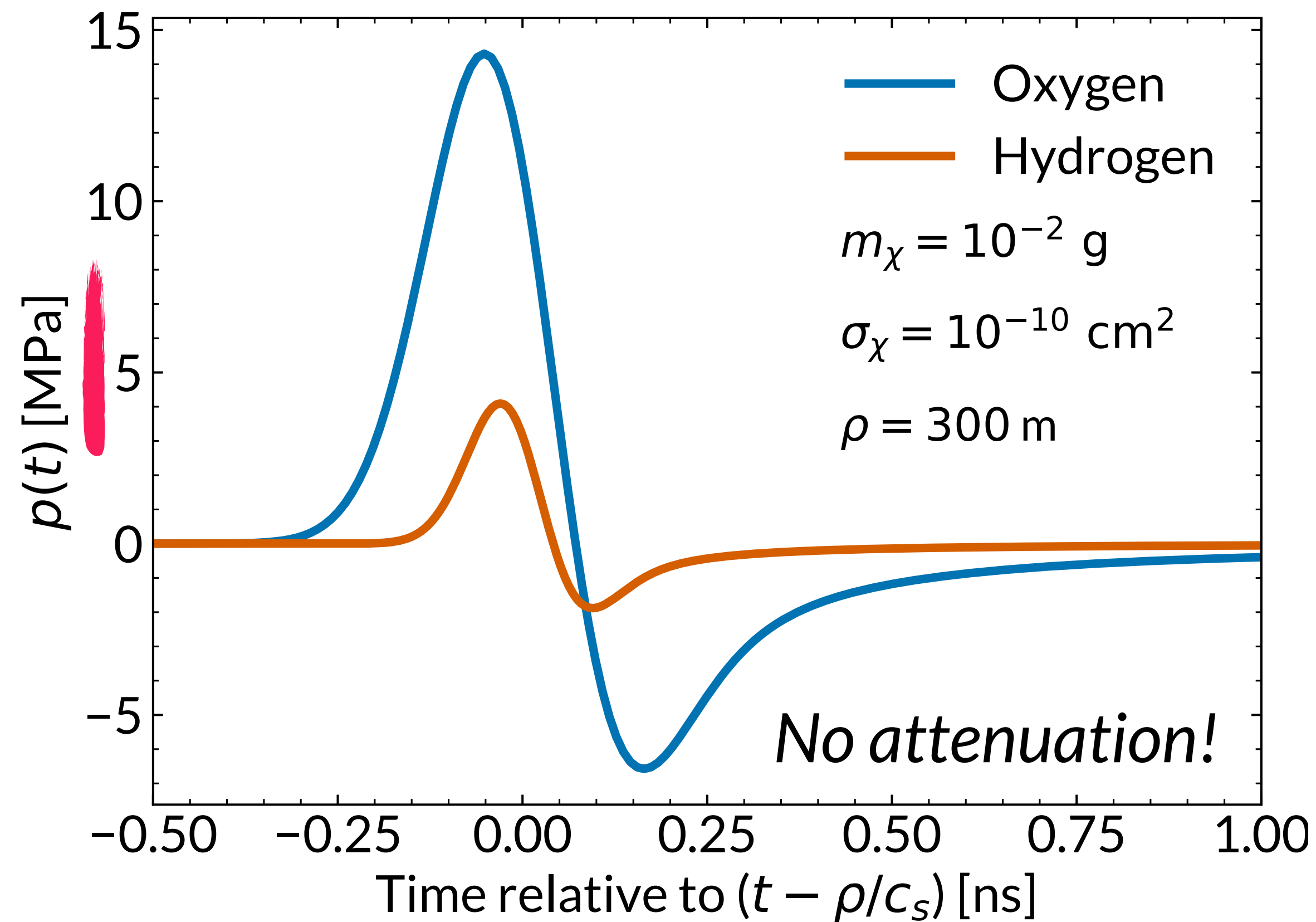
Hydrophone detection threshold ~ 5 mPa



Energy deposition = sound waves

At first sight: extremely promising!

Hydrophone detection threshold ~ 5 mPa



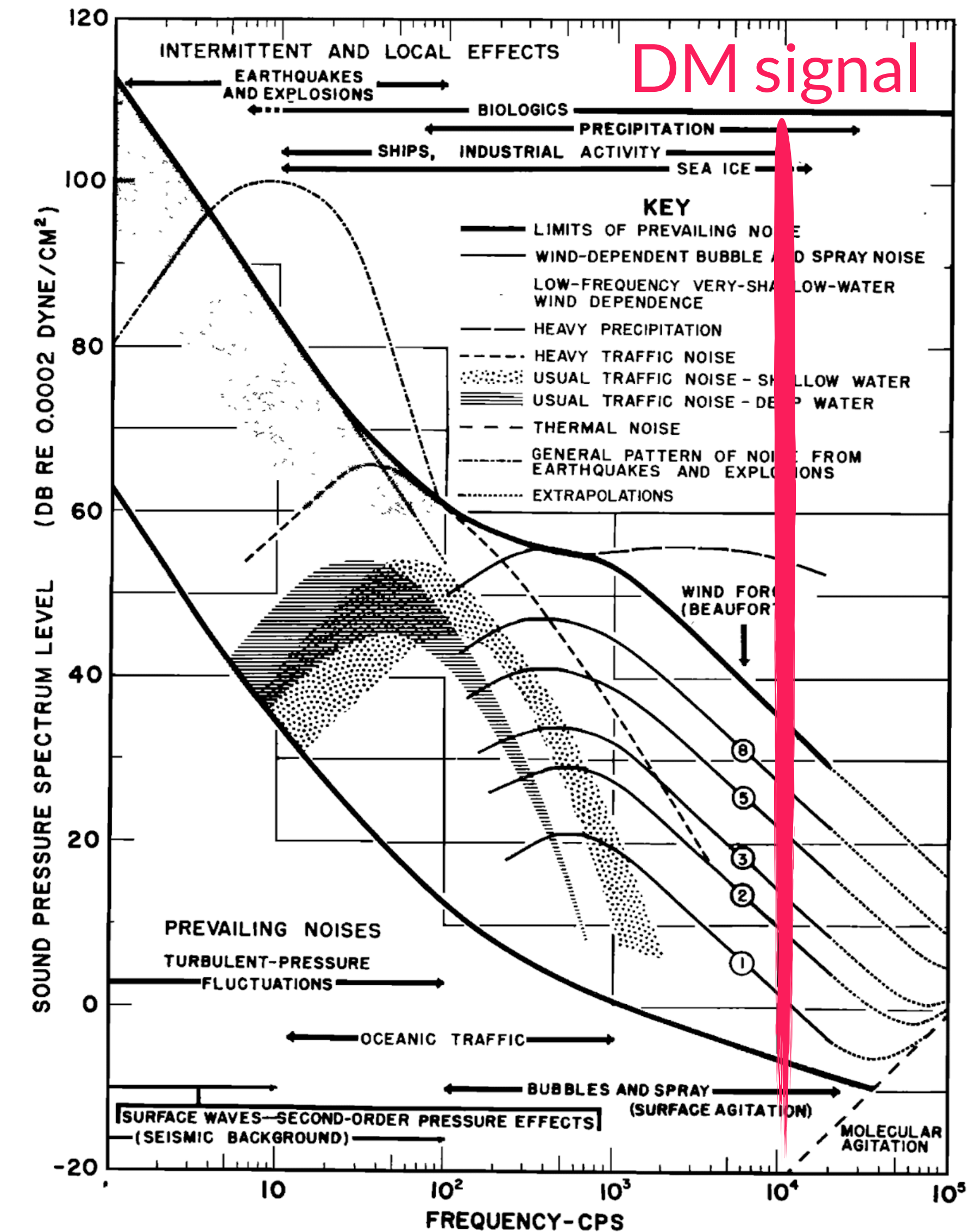
...still interesting even after attenuation included

Backgrounds

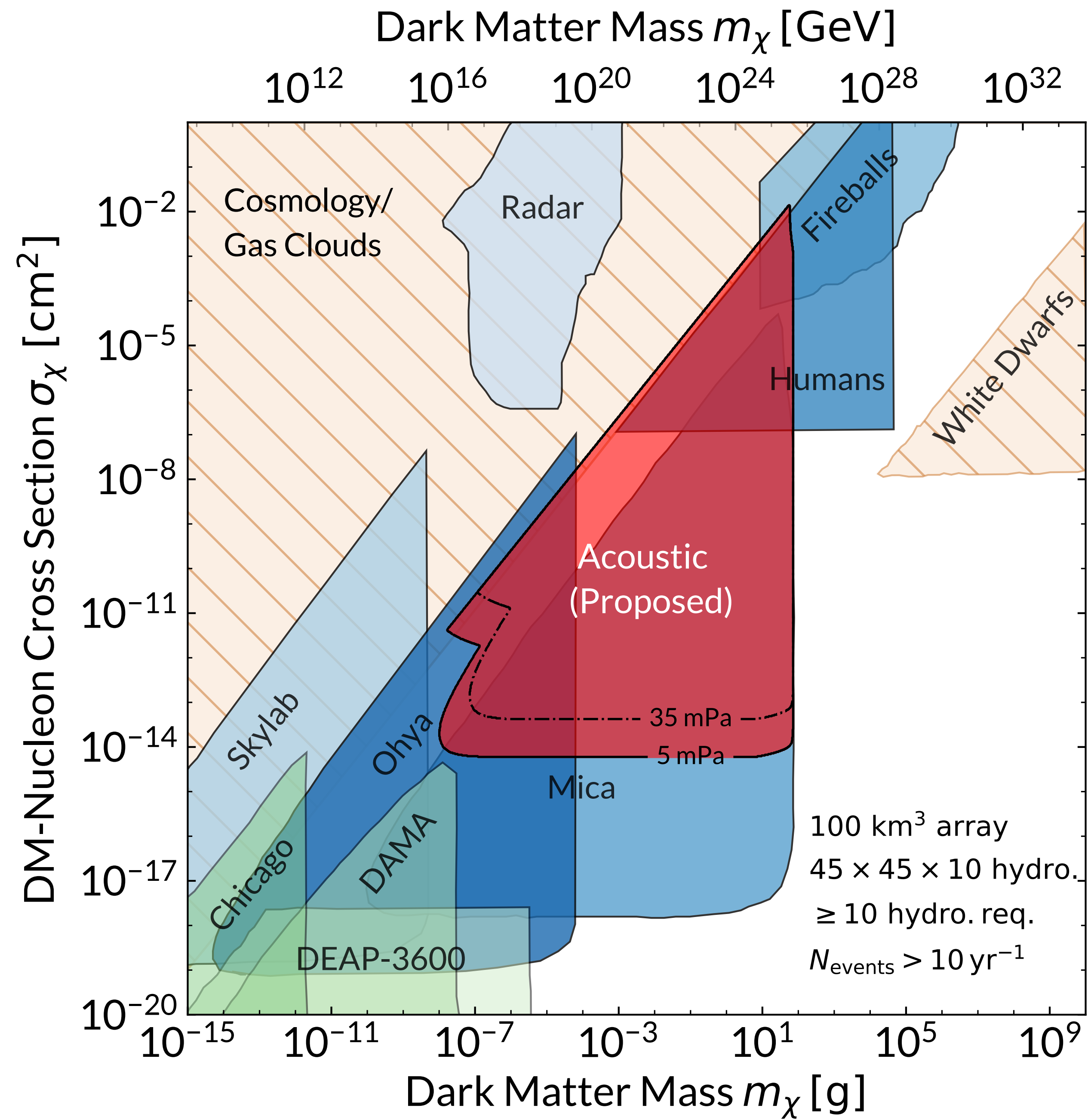
Backgrounds understood:

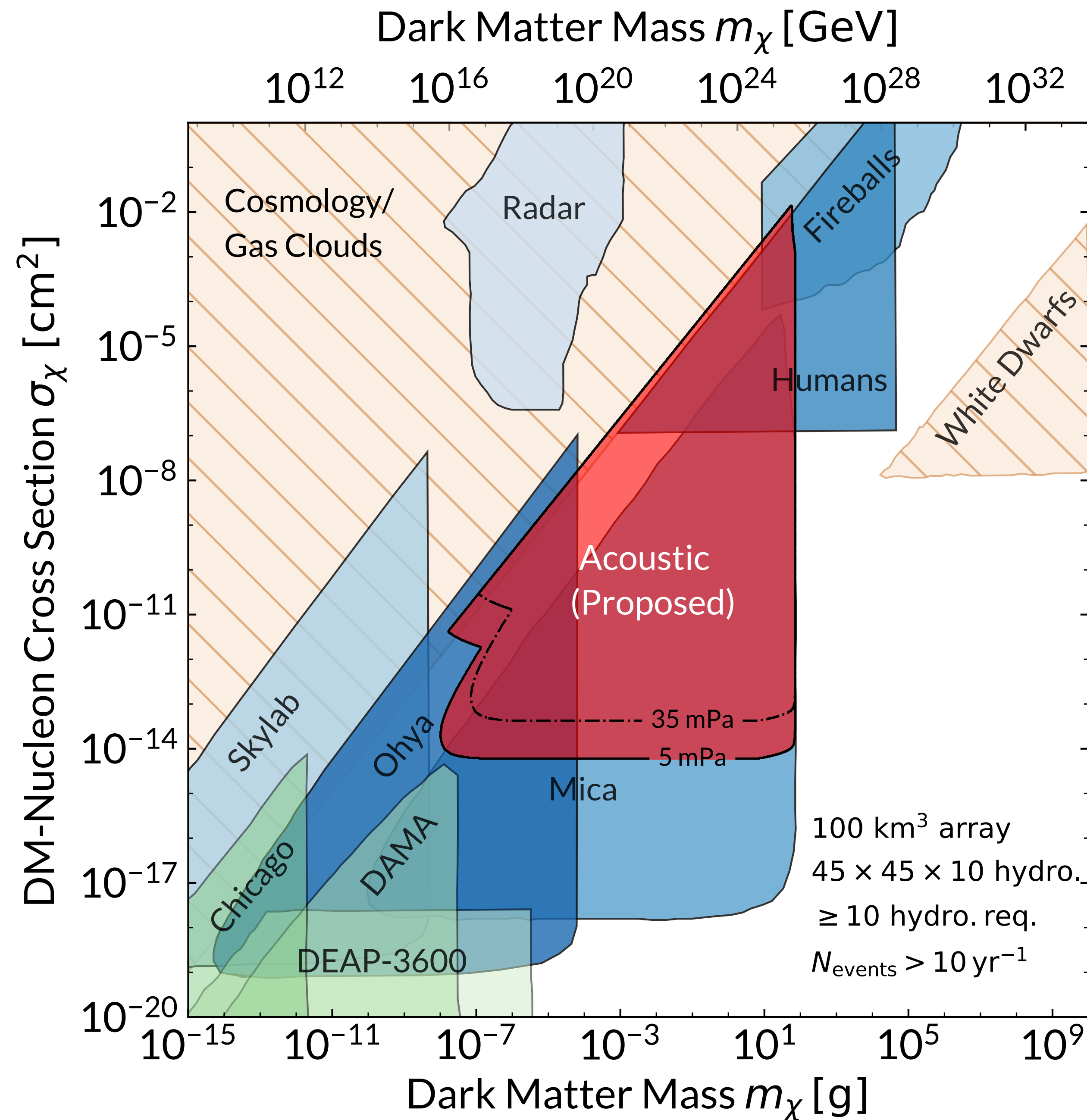
- Seawater waves
- Sea traffic (ships)
- Animals (dolphins, shrimps)

Can produce pulses similar to dark matter
...but none extend over km-long tracks



Wenz, J. Acoust. Soc. Am 1962





Challenge:

$$\Phi_{\text{DM}} \sim 0.3 \left(\frac{10^{19} \text{ GeV}}{m_{\text{DM}}} \right) \text{ m}^{-2} \text{ yr}^{-1}$$

Solution:

To go beyond XLZD, need large exposure through larger areas, larger measuring times, or both!

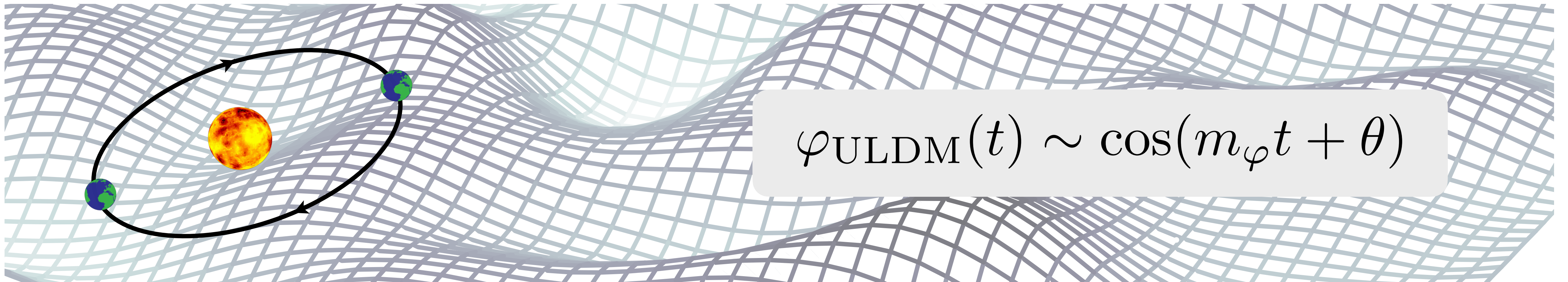
2. Dark matter below the eV scale



Ultra-light dark matter

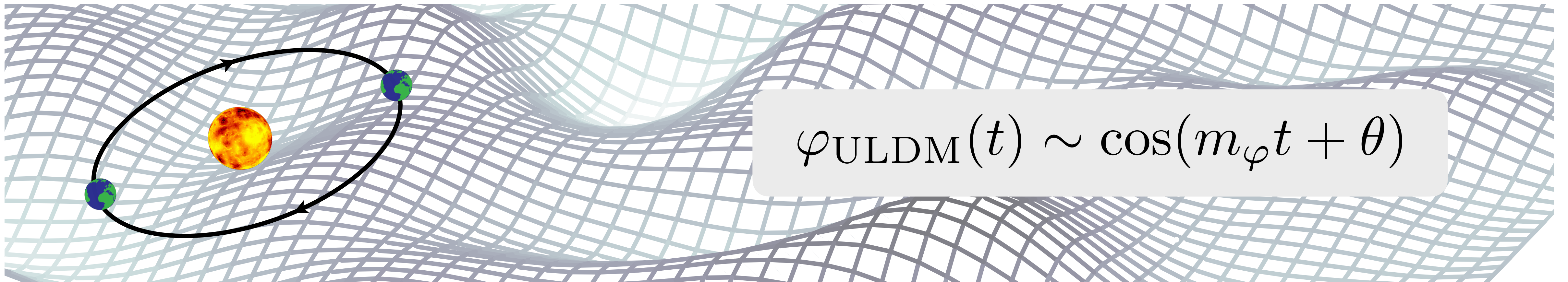
DM lighter than ~few eV behaves as a classical wave

Angular frequency set by the ULDM mass: $\omega \simeq m_\varphi (1 + \mathcal{O}(v^2))$



e.g., Foster et al arXiv: 1711.10489
Derevianko arXiv:1605.09717

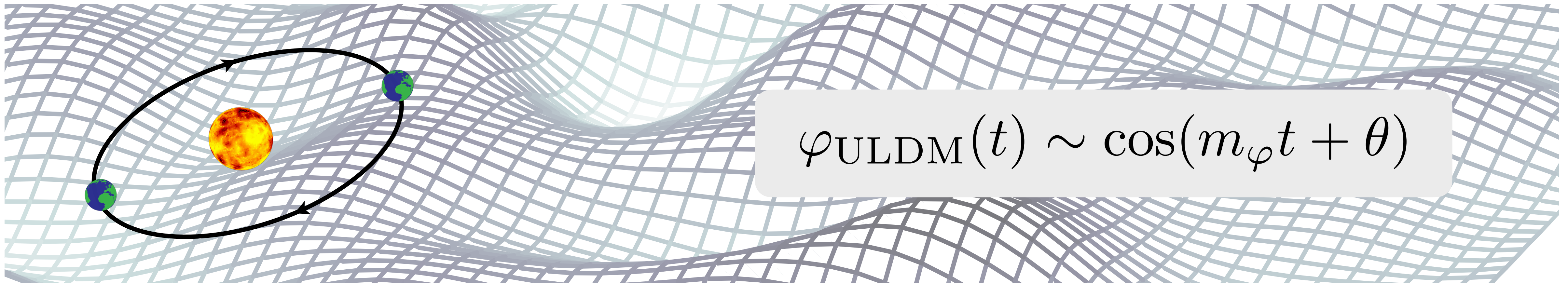
Ultra-light scalar dark matter candidates



$$\mathcal{L} \supset \sqrt{4\pi G_N \phi} \left[\overbrace{d_{m_e} m_e \bar{e} e}^{\text{red bar}} - \underbrace{\frac{d_e}{4} F_{\mu\nu} F^{\mu\nu}}_{\text{green bar}} \right] \begin{matrix} \xrightarrow{\text{red arrow}} m_e(t, \mathbf{x}) = m_e \left[1 + d_{m_e} \sqrt{4\pi G_N \phi(t, \mathbf{x})} \right] \\ \xrightarrow{\text{green arrow}} \alpha(t, \mathbf{x}) = \alpha \left[1 + d_e \sqrt{4\pi G_N \phi(t, \mathbf{x})} \right] \end{matrix}$$

See e.g., Geraci et al, arXiv:1605.04048
and Arvanitaki et al, arXiv:1606.04541

Ultra-light scalar dark matter candidates



$$\mathcal{L} \supset \sqrt{4\pi G_N \phi} \left[\overbrace{d_{m_e} m_e \bar{e} e}^{\text{red bar}} - \underbrace{\frac{d_e}{4} F_{\mu\nu} F^{\mu\nu}}_{\text{green bar}} \right] \begin{matrix} \xrightarrow{\text{red arrow}} m_e(t, \mathbf{x}) = m_e \left[1 + d_{m_e} \sqrt{4\pi G_N \phi(t, \mathbf{x})} \right] \\ \xrightarrow{\text{green arrow}} \alpha(t, \mathbf{x}) = \alpha \left[1 + d_e \sqrt{4\pi G_N \phi(t, \mathbf{x})} \right] \end{matrix}$$

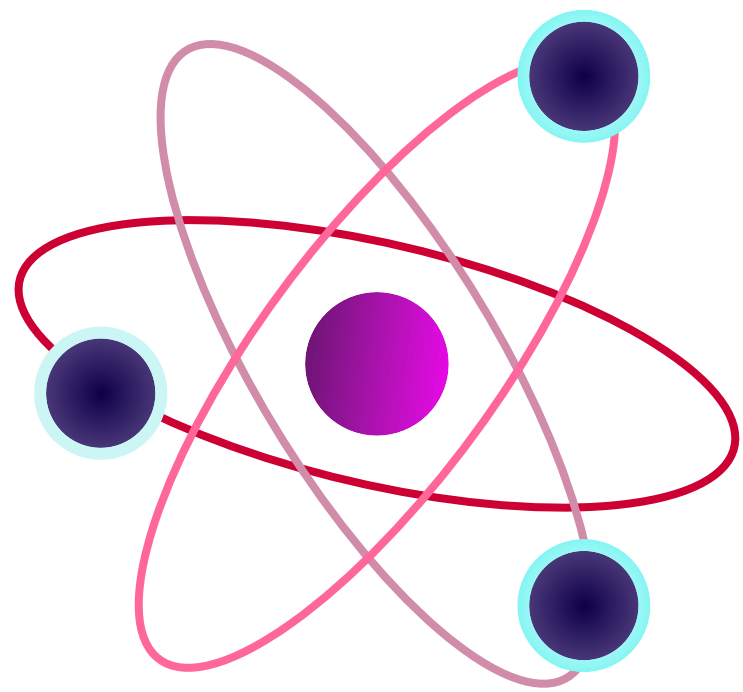
Impact of scalars is to make ‘fundamental constants’ time-dependent
with oscillation frequency set by the ULDM mass

See e.g., Geraci et al, arXiv:1605.04048
and Arvanitaki et al, arXiv:1606.04541

Ultra-light scalar dark matter induces changes in atoms

$$m_e(t, \mathbf{x}) = m_e \left[1 + d_{m_e} \sqrt{4\pi G_N} \phi(t, \mathbf{x}) \right]$$

$$\alpha(t, \mathbf{x}) = \alpha \left[1 + d_e \sqrt{4\pi G_N} \phi(t, \mathbf{x}) \right]$$



Typical atomic size ~ Bohr radius $a_0 \sim \frac{1}{m_e \alpha}$

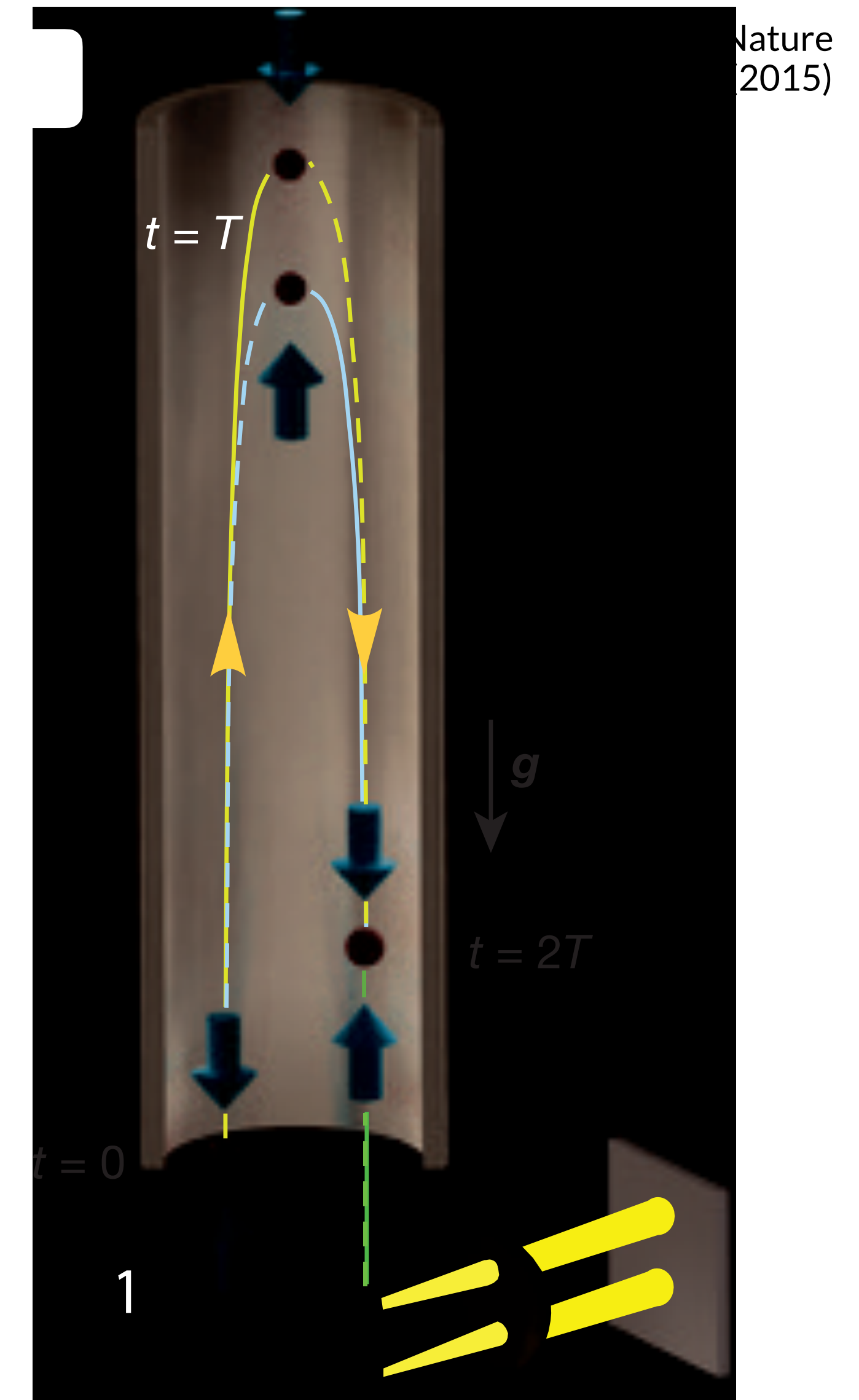
Typical atomic energy ~ Rydberg energy $R_\infty \sim m_e \alpha^2$

Recent developments:
use very precise quantum sensors to search for these changes

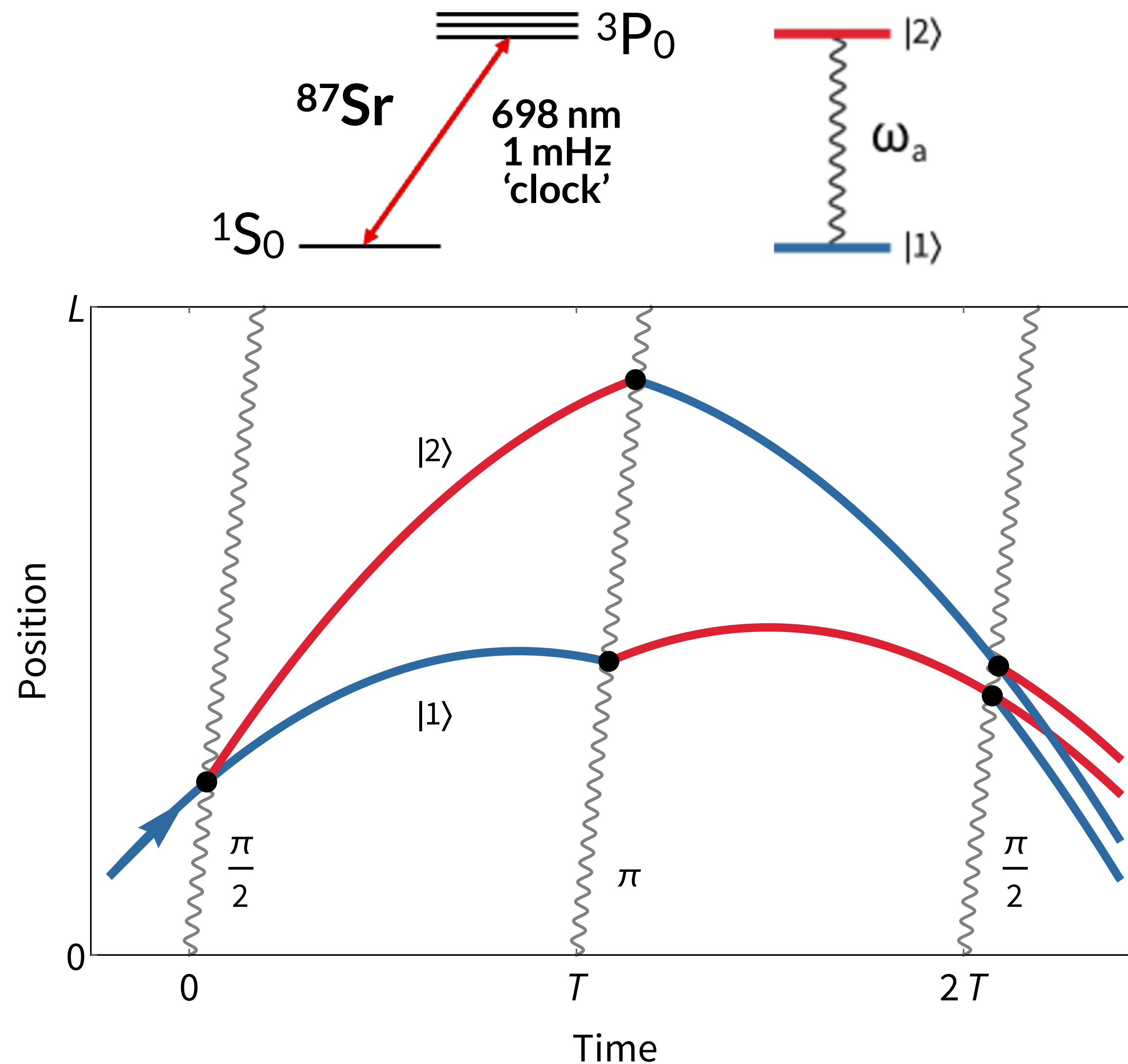
See e.g., Geraci et al, arXiv:1605.04048
and Arvanitaki et al, arXiv:1606.04541

Light-pulse atom interferometry (physical-space)

- Launch ultra-cold cloud of atoms into an atomic fountain
- Sequence of optical pulses manipulate the atoms
- Quantum superposition over macroscopic distances (>50cm achieved)
- Interfere using a final optical pulse when they spatially overlap
- Image the two interferometer output ports



Light-pulse atom interferometry (space-time)



Two-level system separated by optical frequency difference ω_a

Initial pulse: 'beamsplitter'

Middle pulse: 'mirror'

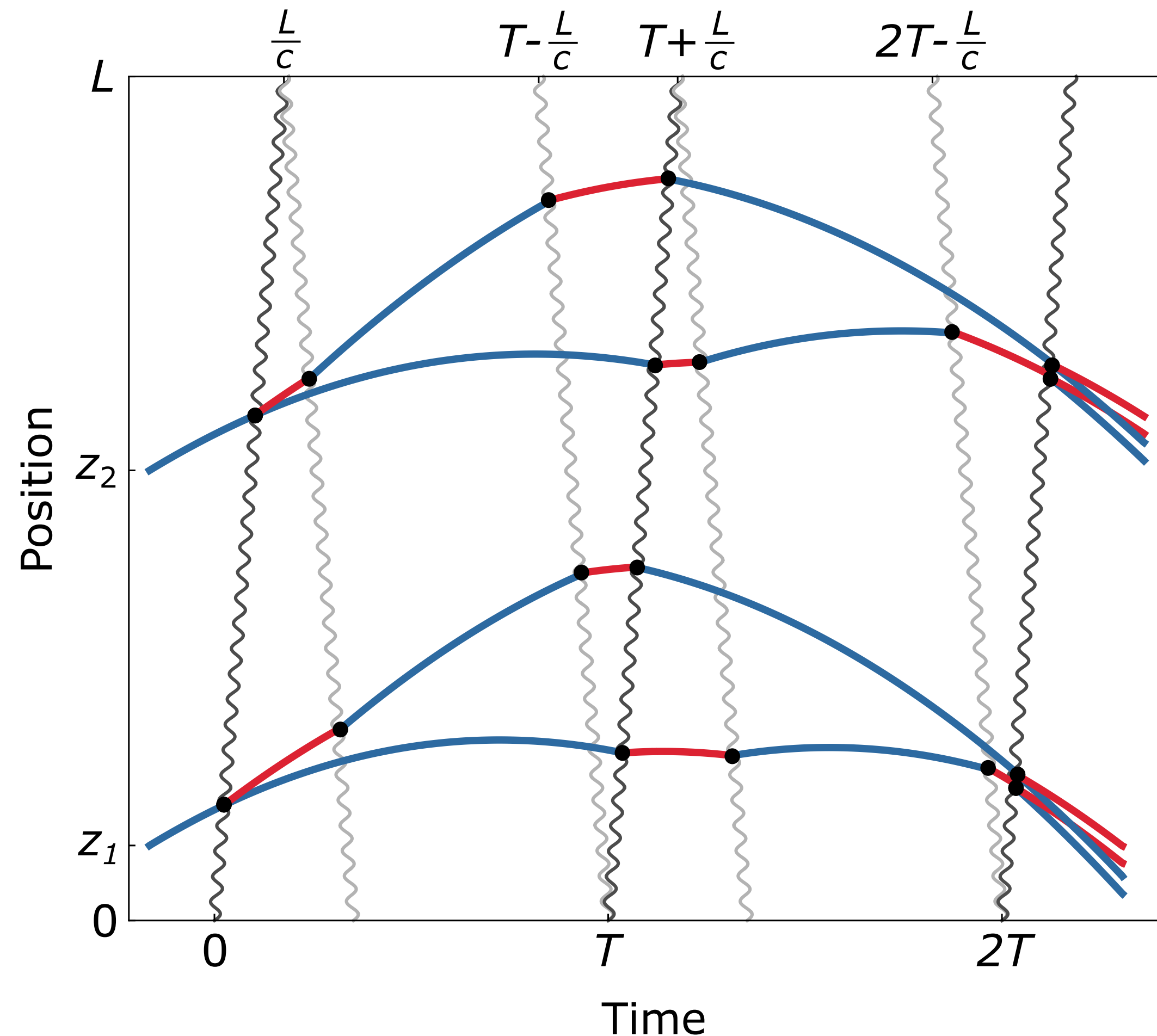
Final pulse: 'beamsplitter (interfere)'

Atom evolves extra clock phase:

$$\frac{1}{\sqrt{2}} |1\rangle + \frac{1}{\sqrt{2}} |2\rangle e^{-i\omega_a T}$$

Phase sensitive to changes in timings, atomic structure, and local accelerations

A bit more involved... operate as a gradiometer

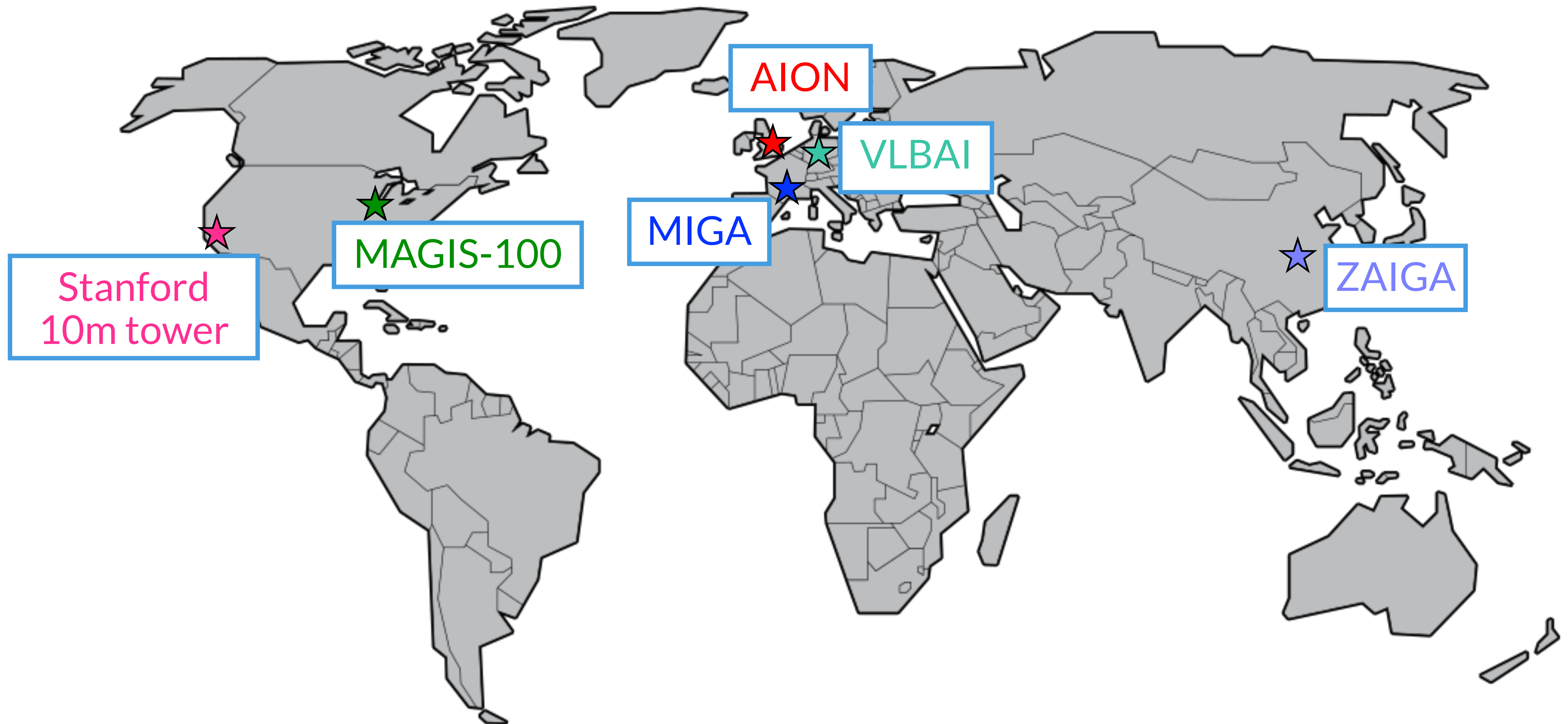


Run **two atom interferometers** simultaneously with the same laser ('gradiometer')

State-of-the-art single photon **strontium** atom interferometry with large momentum transfer (LMT) techniques

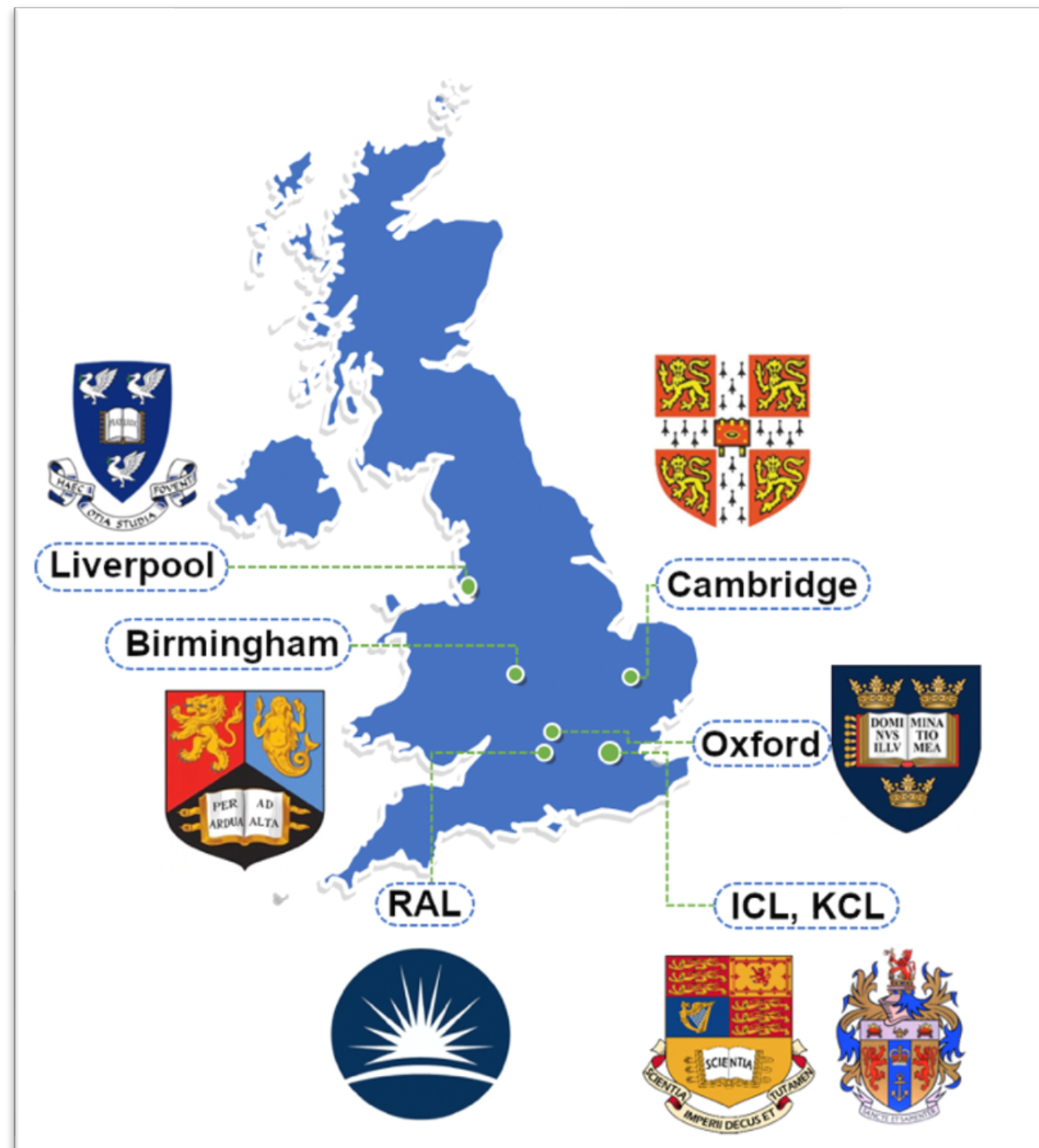
Cancels one of the leading noise sources (laser phase noise)

New atom interferometers across the world coming online

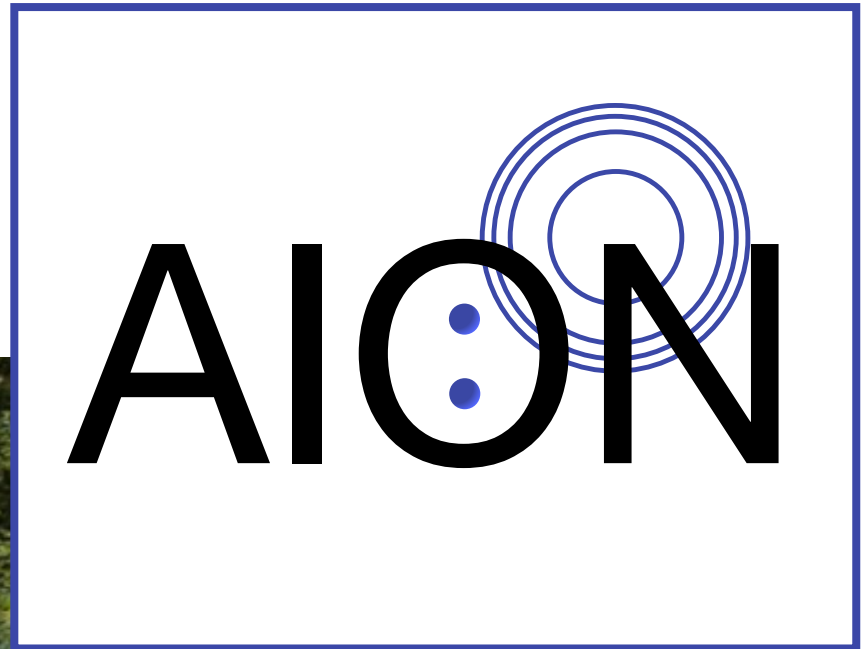


MAGIS-100, [arXiv:2104.02835](https://arxiv.org/abs/2104.02835); MIGA, [arXiv:1703.02490](https://arxiv.org/abs/1703.02490); AION, [arXiv:1911.11755](https://arxiv.org/abs/1911.11755); VLBAI, [arXiv:2003.04875](https://arxiv.org/abs/2003.04875); ZAIGA, [arXiv:1903.09288](https://arxiv.org/abs/1903.09288)

AION: Atom Interferometer Observatory and Network



7 institutes in the UK



Autumn 2021

Collaboration ~50 people
Cold atom: fundamental physics ratio is ~50:50

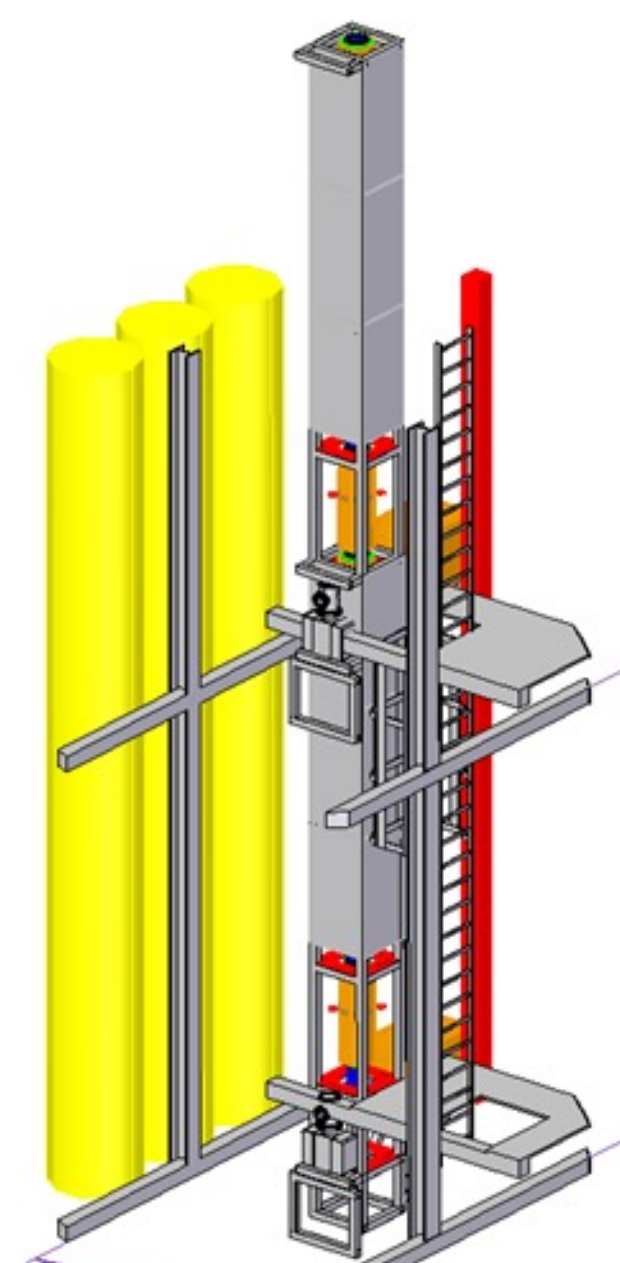
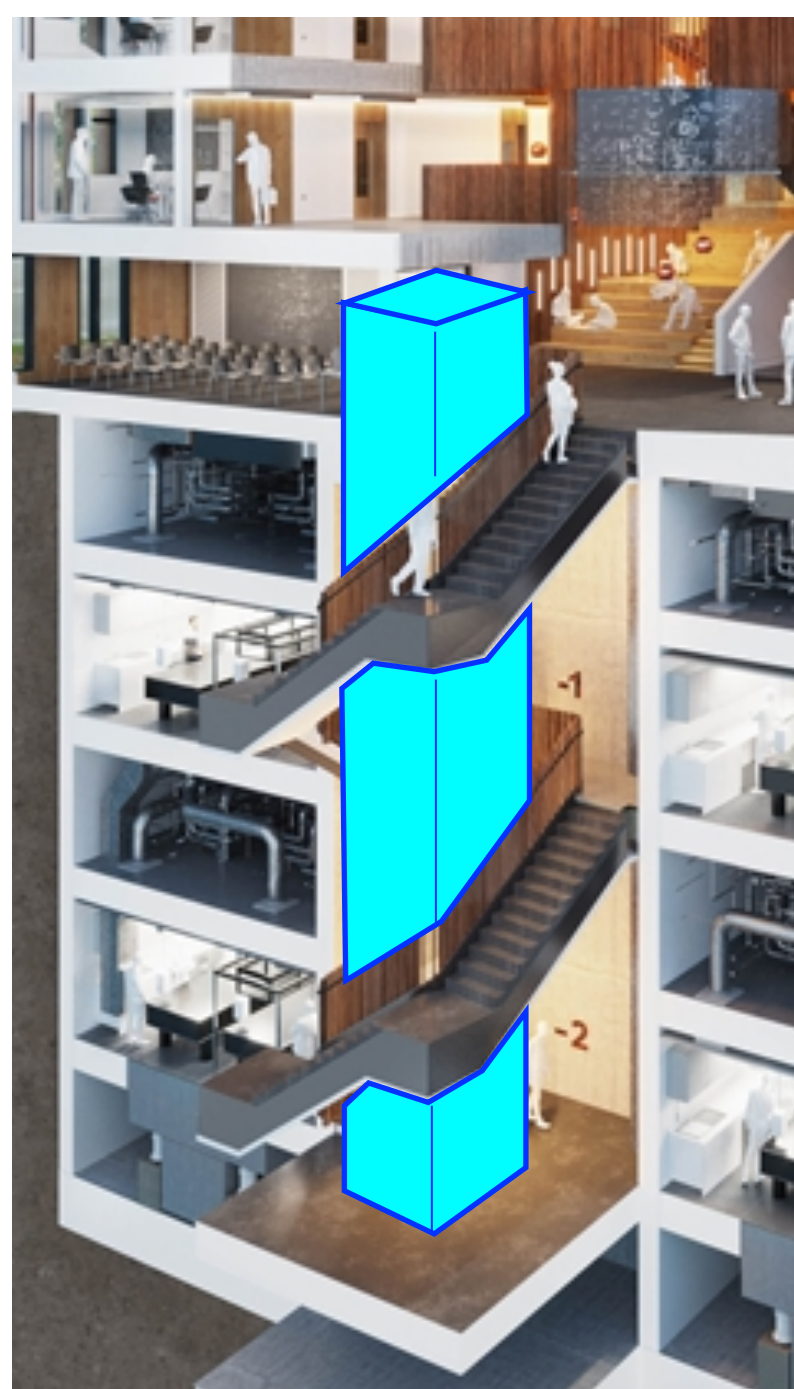
AION & TVLBAI: Towards very-long baselines

AION-10
2025+ ~10m
instrument in
Oxford

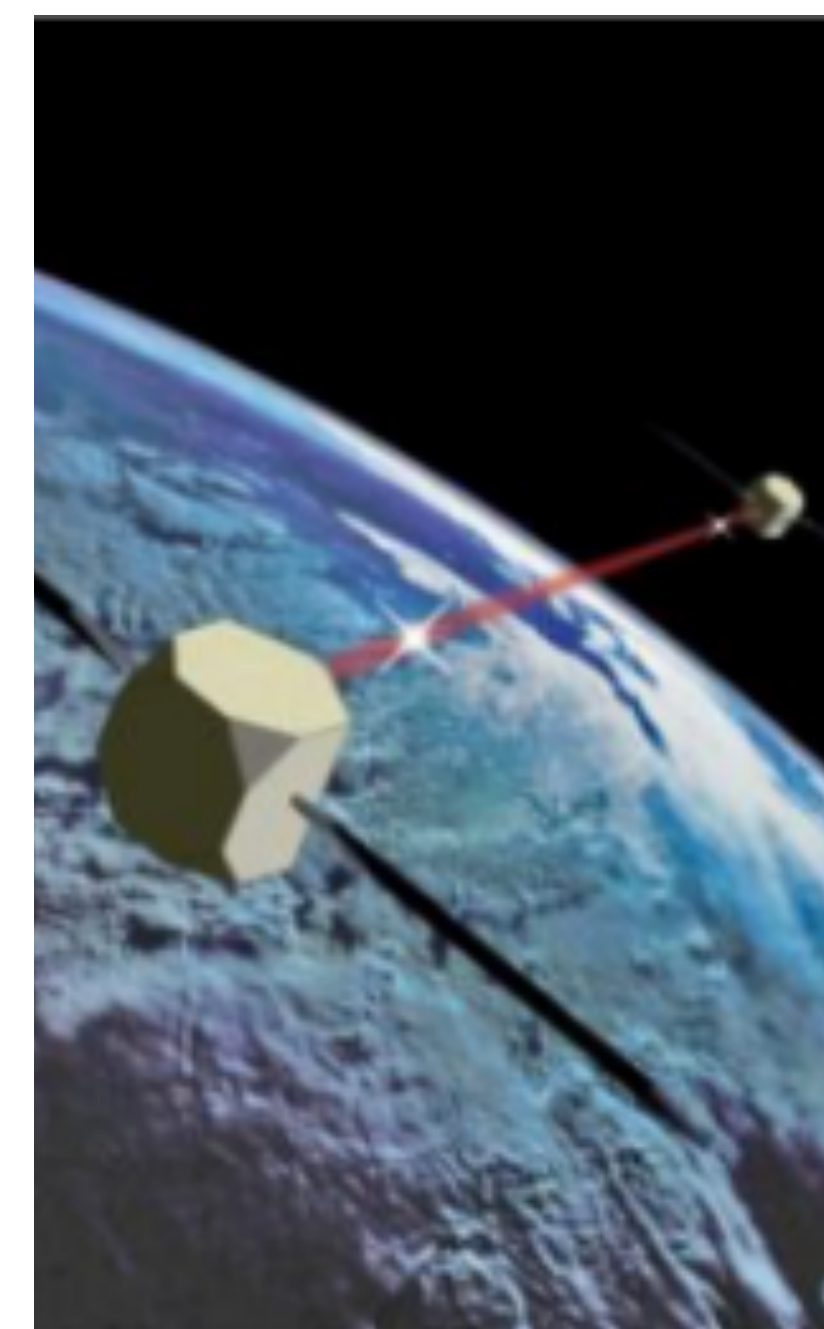
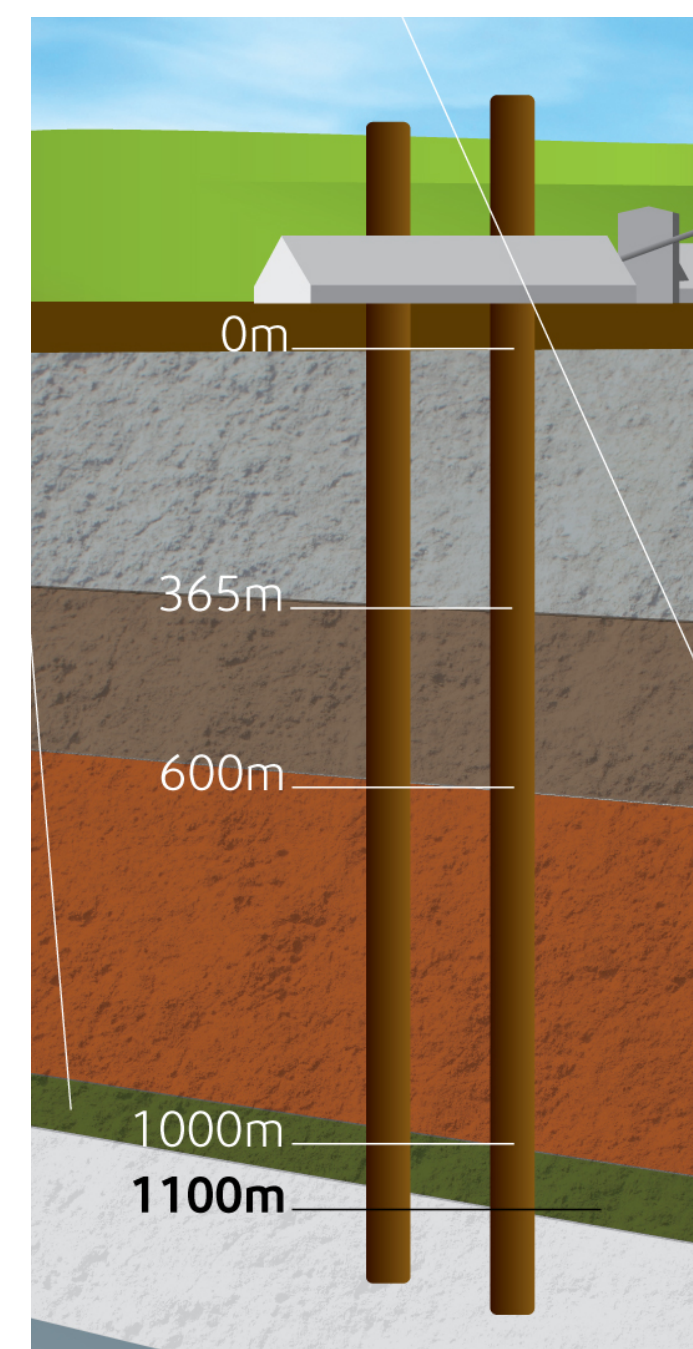
AION-100
2030s ~100m
instrument at
Boulby/CERN/...?

km-instrument
2040s major
international
project(s)

Space-instrument
2050s
detectors with
~ 10^7 km baseline



Boulby SHAFT 3



AION-10 technical design: [arXiv:2508.03491](https://arxiv.org/abs/2508.03491); CERN studies: [arXiv:2304.00614](https://arxiv.org/abs/2304.00614) & [arXiv:2509.11867](https://arxiv.org/abs/2509.11867); AEDGE, [arXiv:1908.00802](https://arxiv.org/abs/1908.00802); Cold atoms in Space, [arXiv:2201.07789](https://arxiv.org/abs/2201.07789); TVLBAI [arXiv:2503.21366](https://arxiv.org/abs/2503.21366)

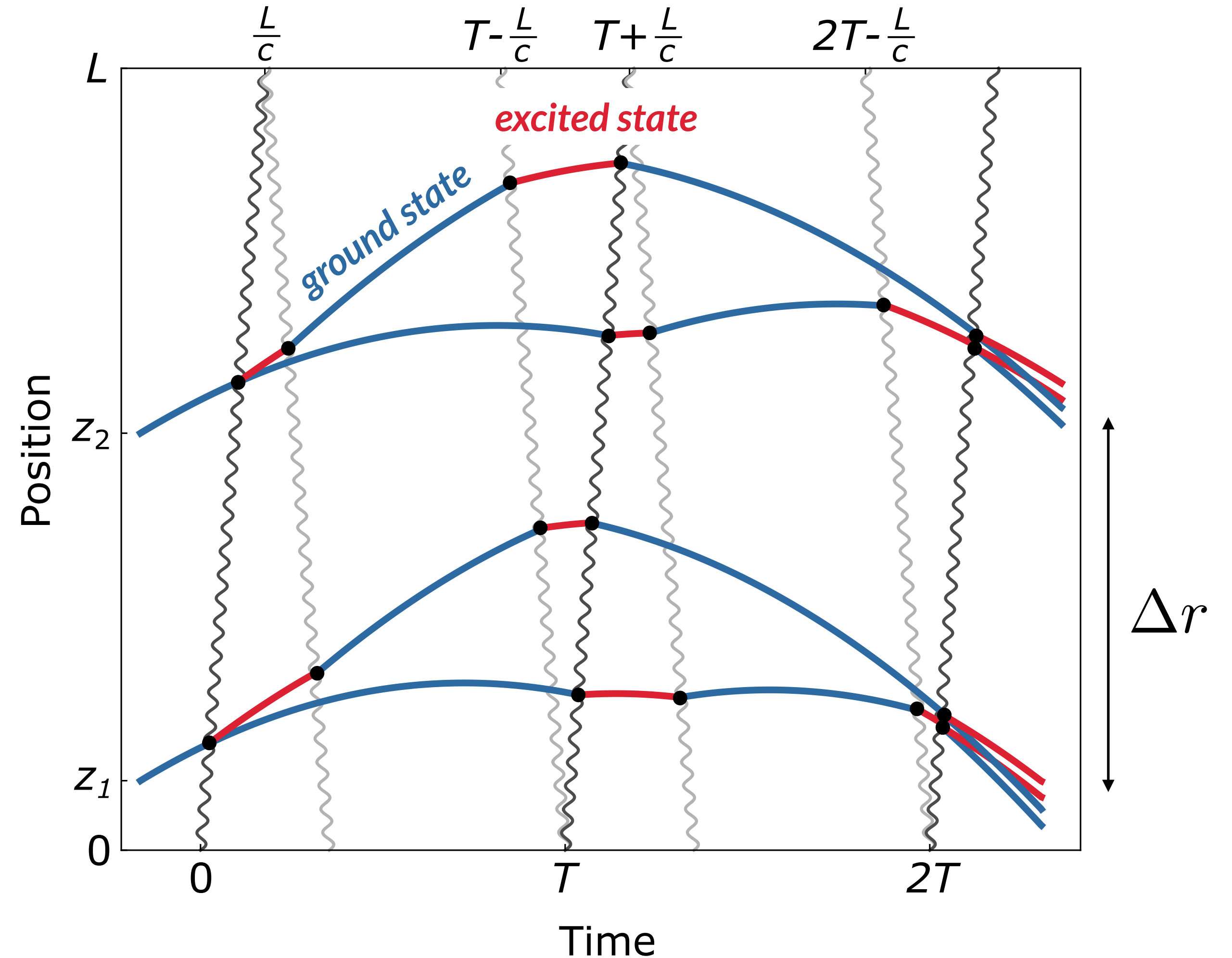
Scalar ULDM signal

Phase is accumulated by the **excited state** relative to the **ground state** along all paths:

$$\Phi_{t_1}^{t_2}(\mathbf{r}) = \int_{t_1}^{t_2} \Delta\omega_a(t, \mathbf{r}) dt$$

$$\Delta\omega_A(t) \sim [d_{m_e} + \xi_A d_e] \cos(m_\phi t + \theta)$$

t_1, t_2 = time in excited state



Many parameters to tune to reach sensitivity

$$d_{m_e}^{\text{best}} \sim \left(\frac{1}{T}\right)^{5/4} \frac{1}{C n \Delta r} \left(\frac{\Delta t}{N_a}\right)^{1/2} \left(\frac{1}{T_{\text{int}}}\right)^{1/4}$$

Handles to optimise (in order of priority):

$T \sim 1$ s (interrogation time)

$C \sim 0.1 - 1$ (contrast)

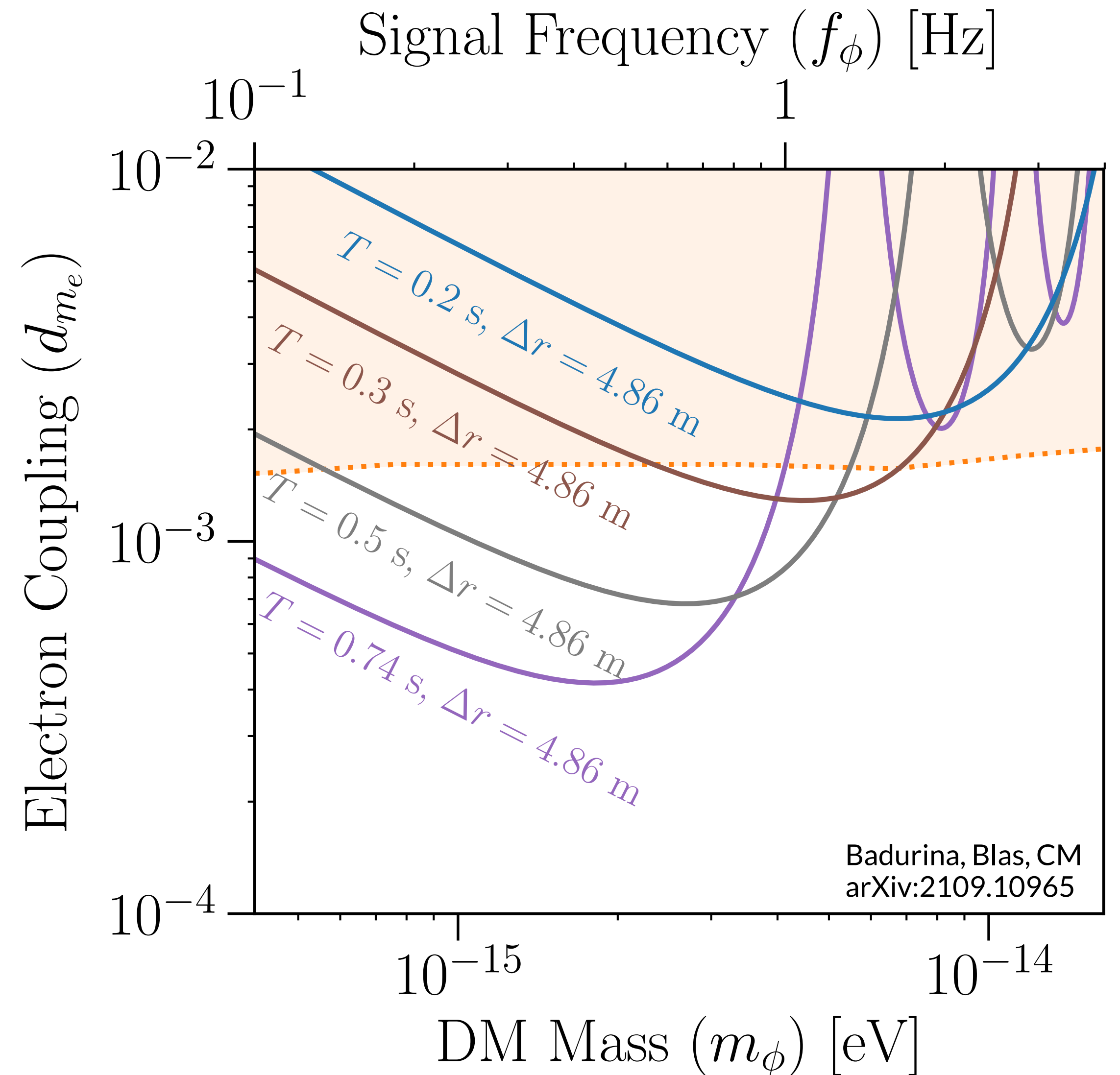
$n \sim 1000$ (LMT)

$\Delta r \sim \text{AI separation}$

$\Delta t \sim \text{sampling time}$

$N_a \sim \text{atoms in cloud}$

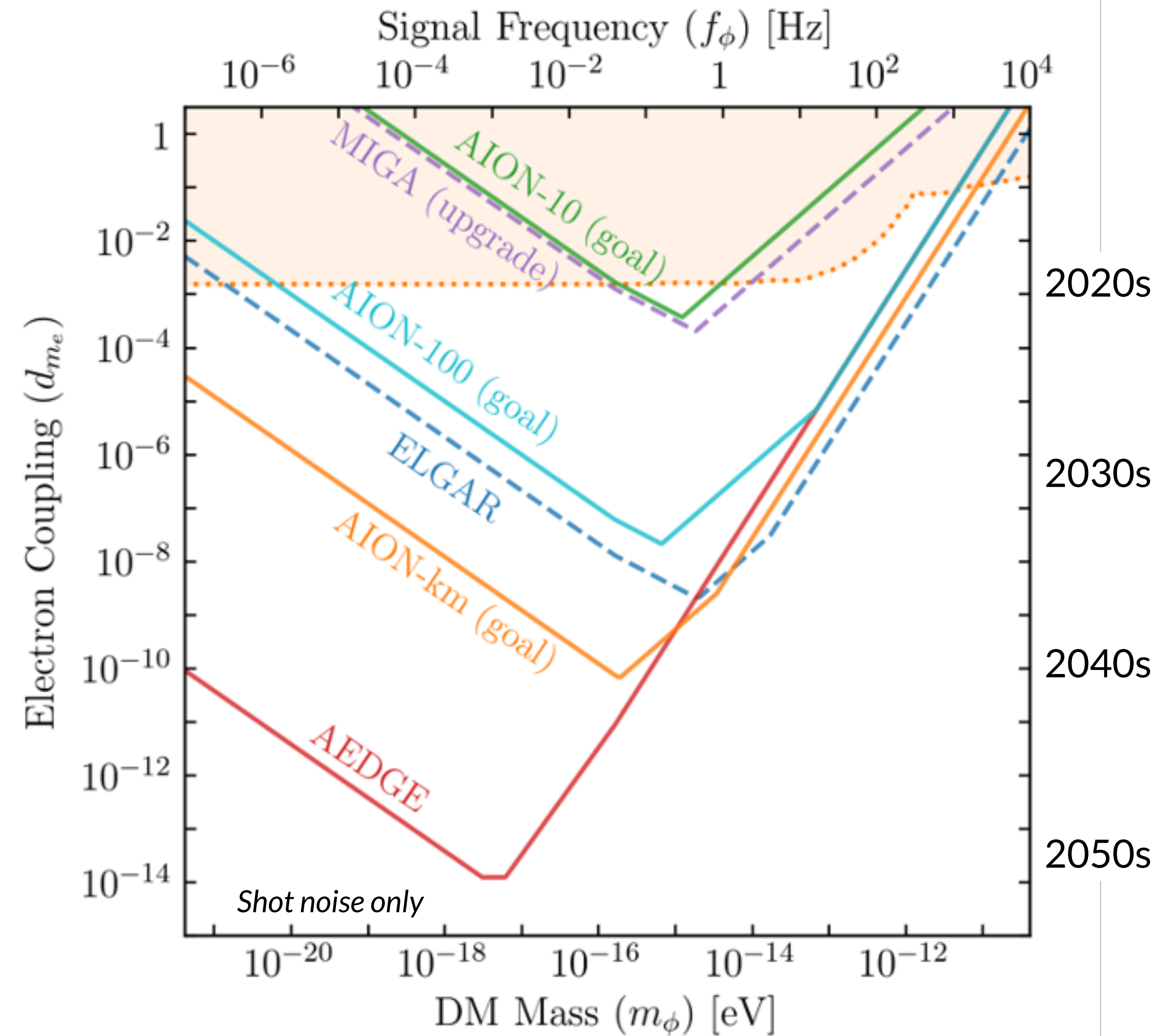
$T_{\text{int}} \sim 10^7$ s (integration time)



Near- and long-term prospects (Scalar ULDM)

Sensitivity Scenario	L [m]	T_{int} [sec]	$\delta\phi_{noise}$ [$1/\sqrt{\text{Hz}}$]	LMT [number n]
AION-10 (initial)	10	1.4	10^{-3}	100
AION-10 (goal)	10	1.4	10^{-4}	1000
AION-100 (initial)	100	1.4	10^{-4}	1000
AION-100 (goal)	100	1.4	10^{-5}	40000
AION-km	2000	5	0.3×10^{-5}	40000

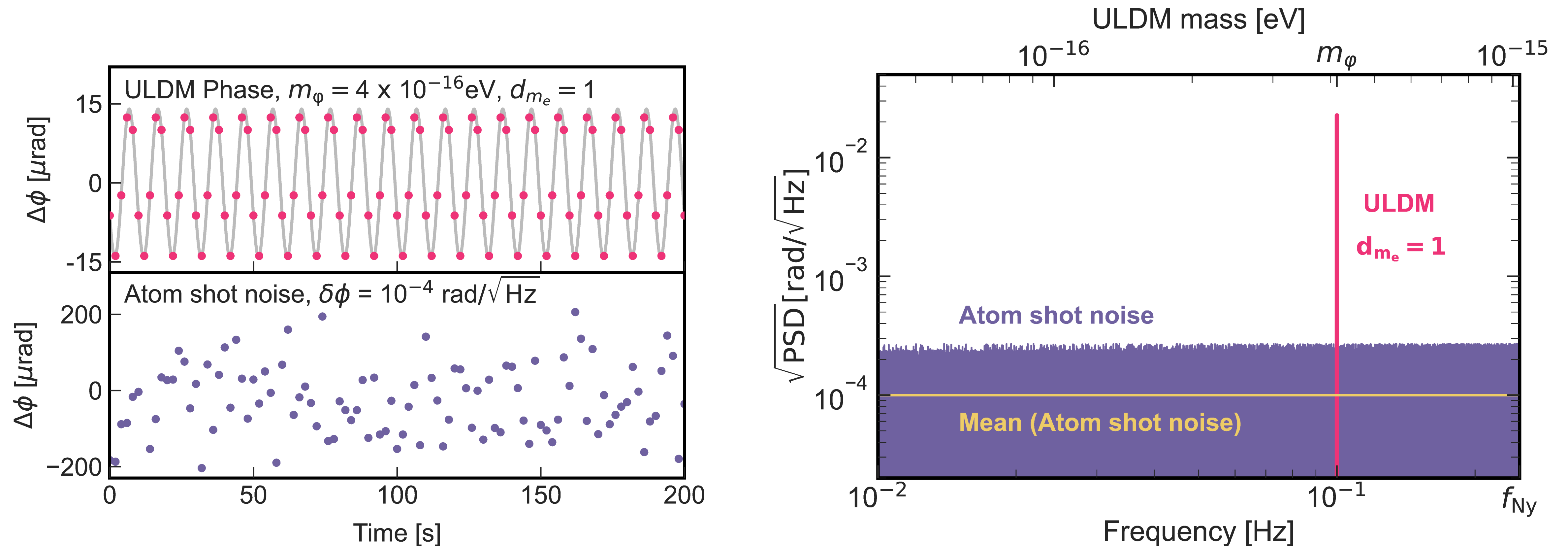
Badurina, CM, et al, arXiv:1911.11755, 2108.02468



Searching for a signal (ideal setup)

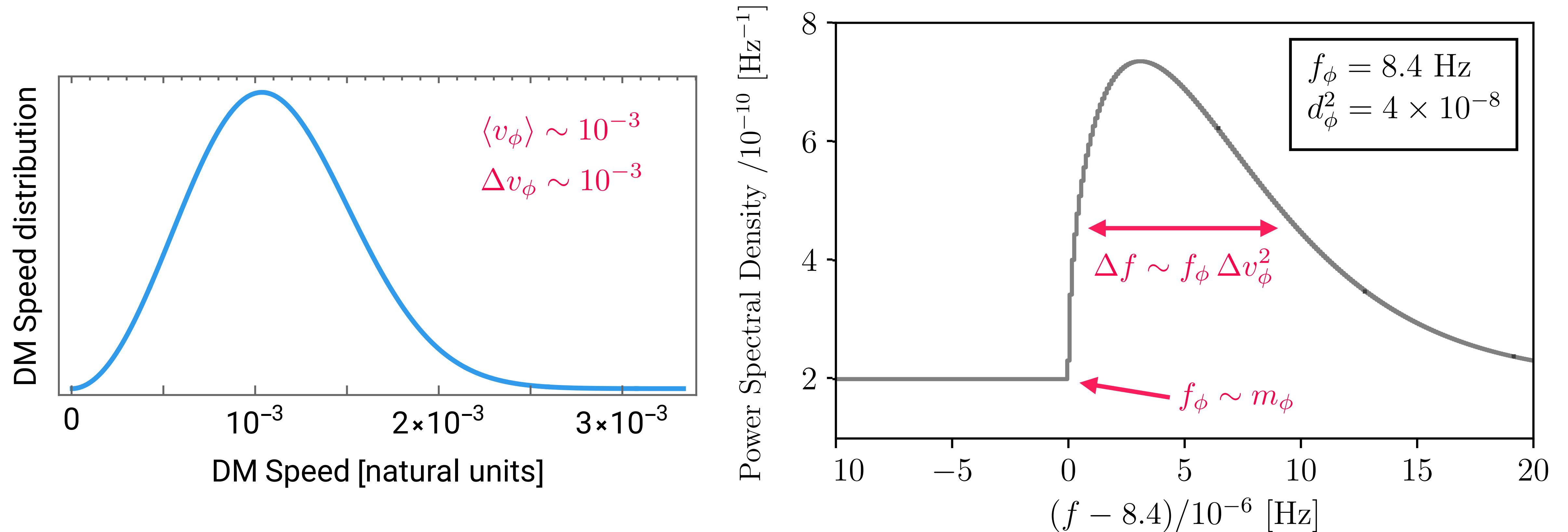
ULDM signal has frequency set by the ULDM mass

Most natural to search for ULDM in frequency space (power spectral density)

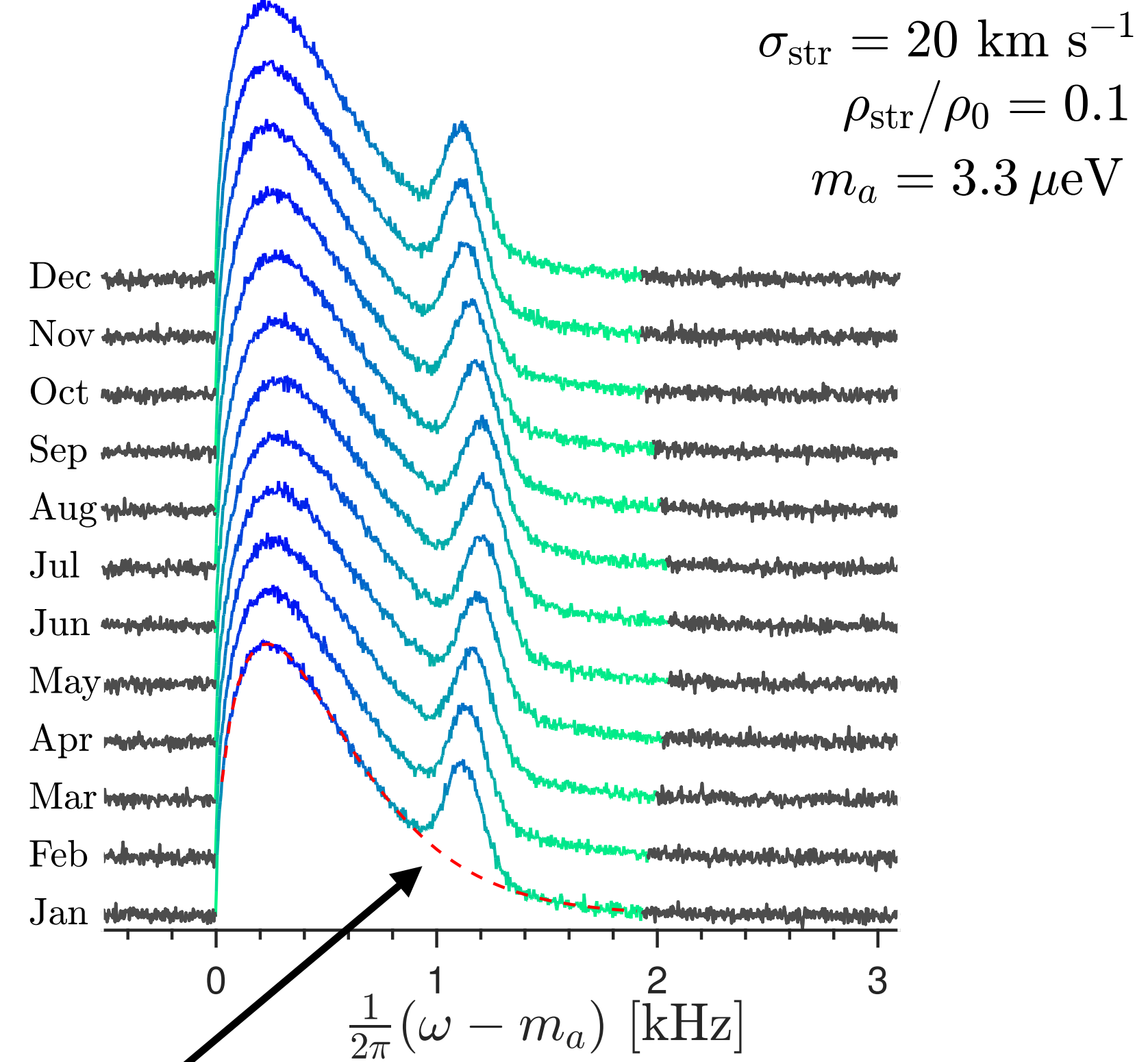
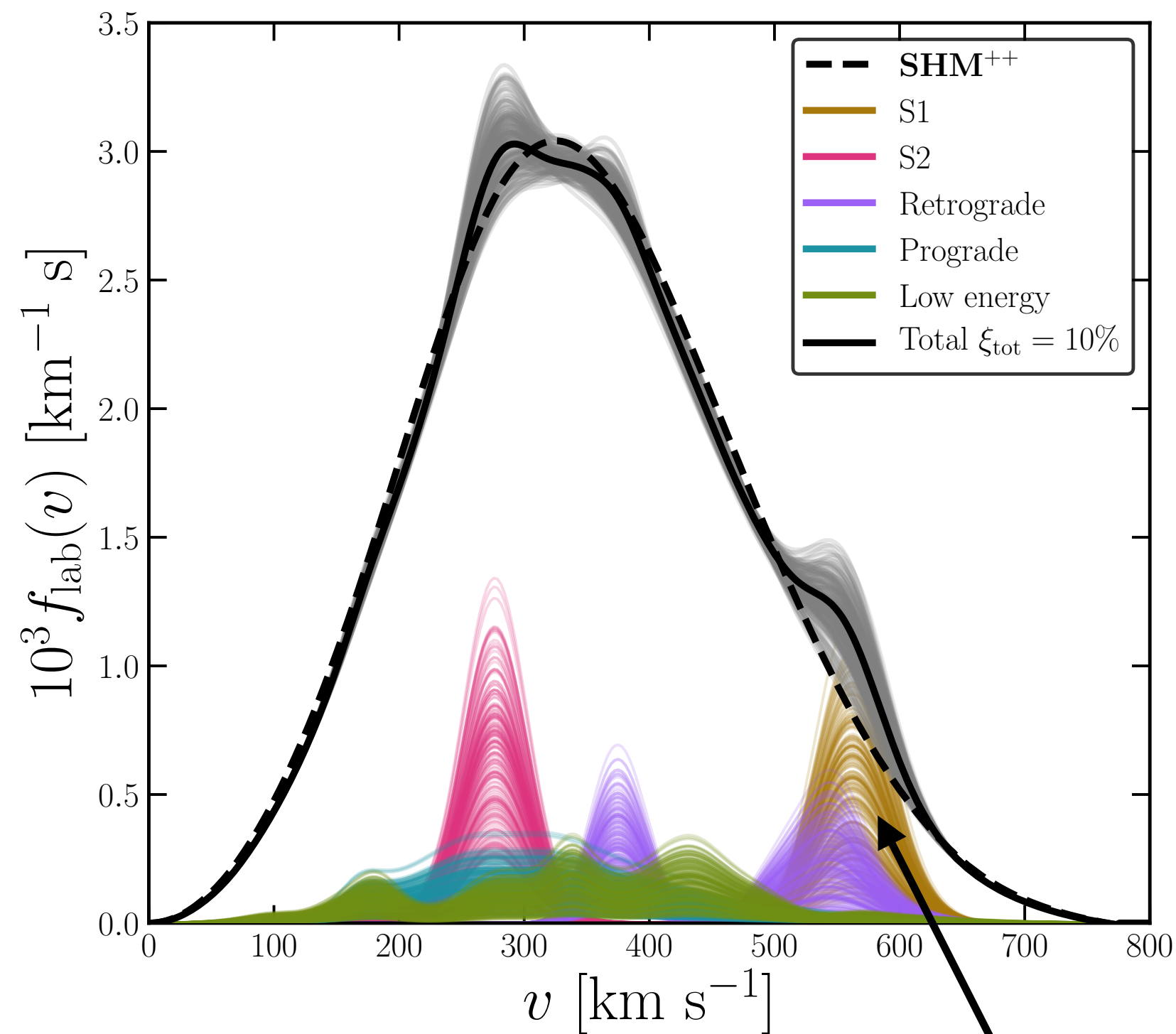


Zooming in on ULDM ‘spike’

The ‘spike’ has a characteristic distribution set by ULDM properties in the Solar System

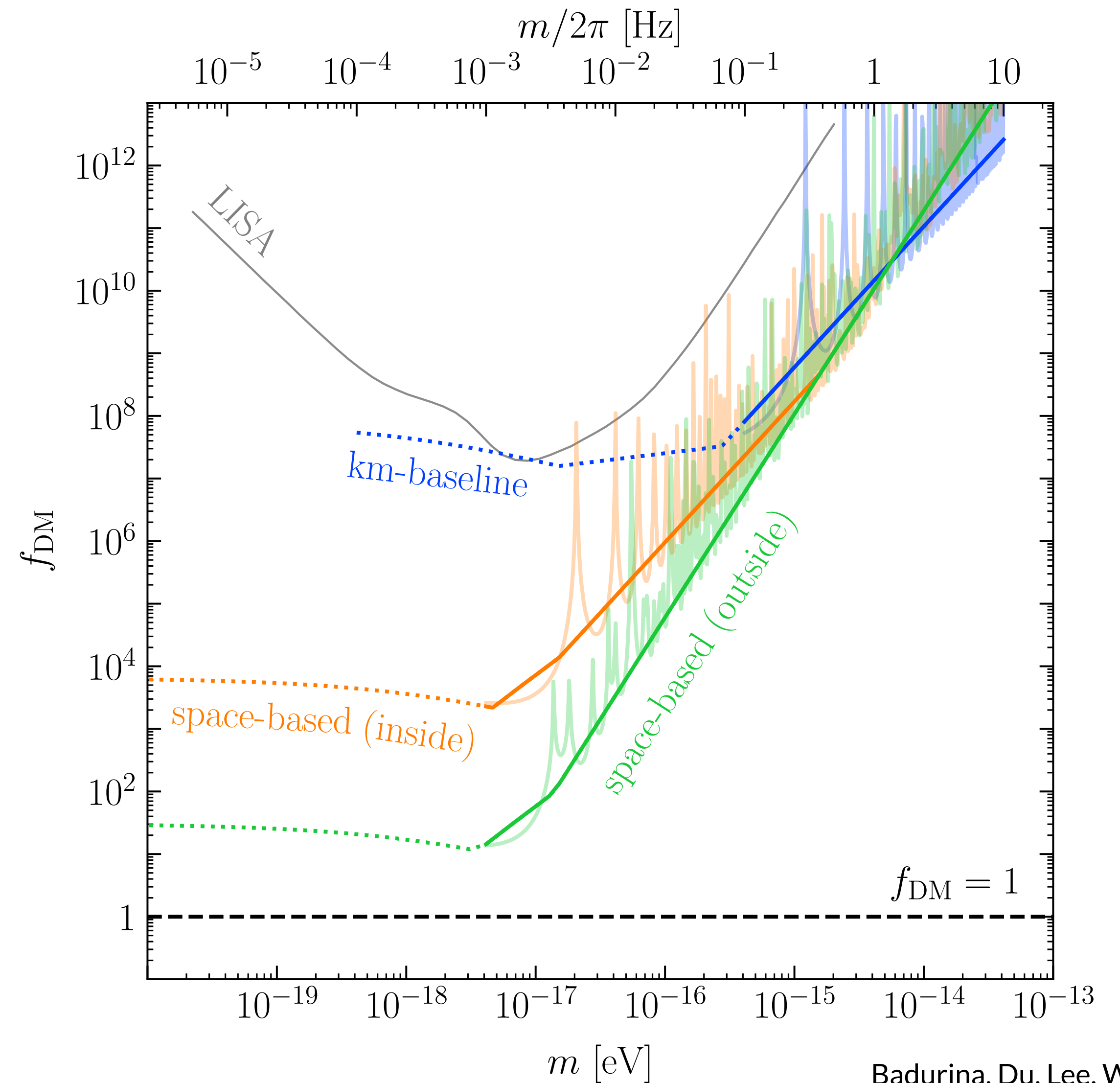
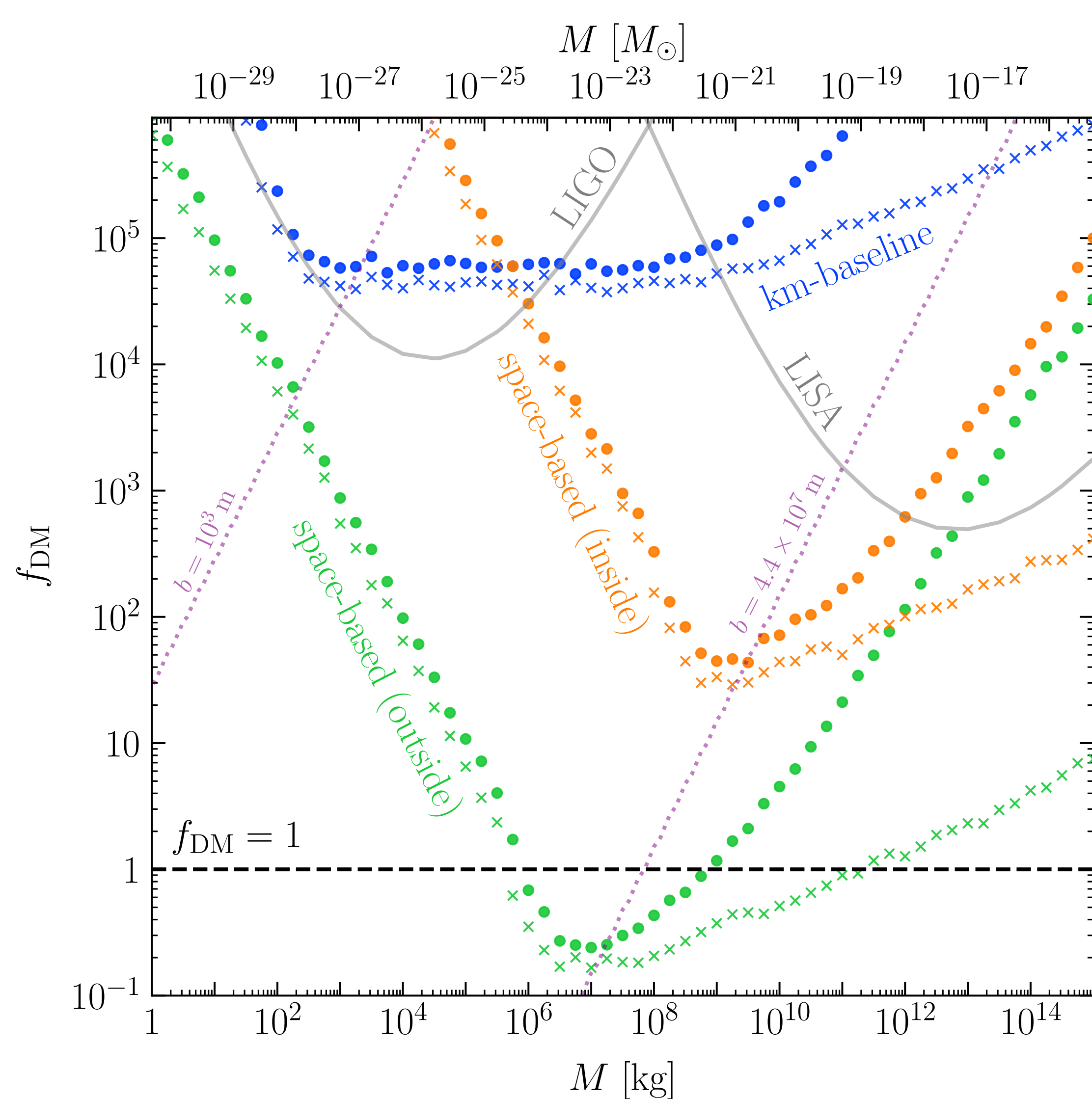


Zooming in on ULDM ‘spike’



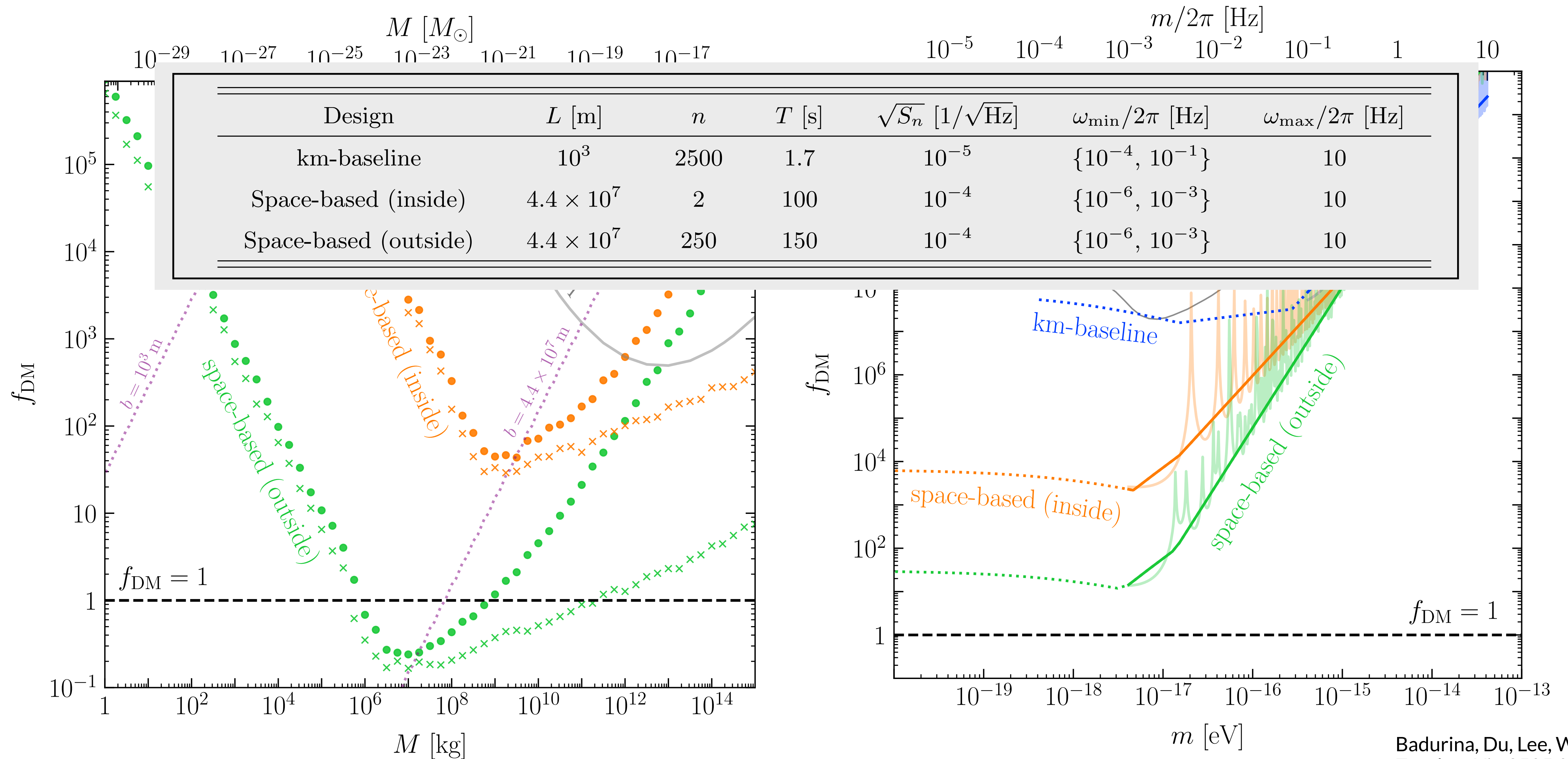
Features in the speed distribution would show up
in the measured signal

Very long-term: dark matter gravitational interactions



Badurina, Du, Lee, Wang,
Zurek, arXiv:2505.00781

Very long-term: dark matter gravitational interactions

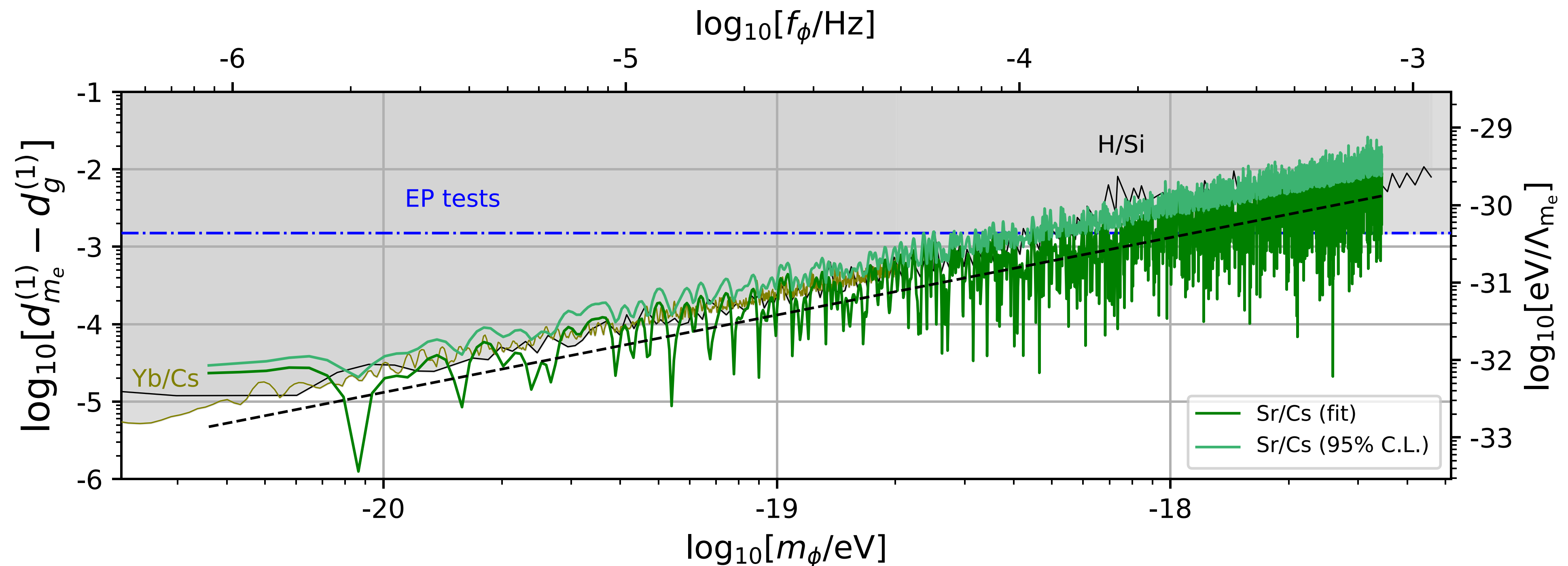


Badurina, Du, Lee, Wang,
Zurek, arXiv:2505.00781

Atomic clocks are complementary (lower mass)

Atomic clocks also sensitive to changes in energy levels

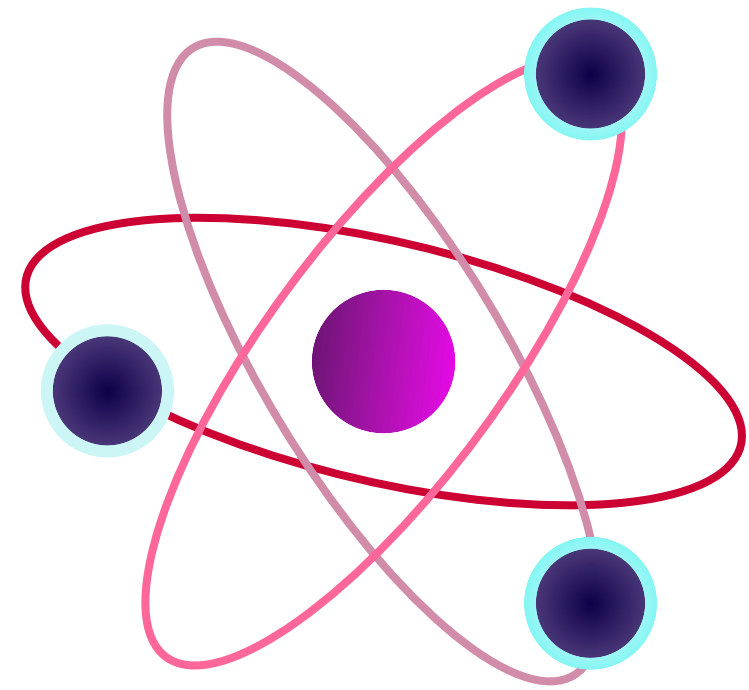
Always need to compare two clocks (use one as a reference)



[atom interferometers probe $\sim 10^{-16}$ to 10^{-13} eV]

QSNET
New J. Phys. 25 093012 (2023)

Laser interferometers are complementary (higher mass)

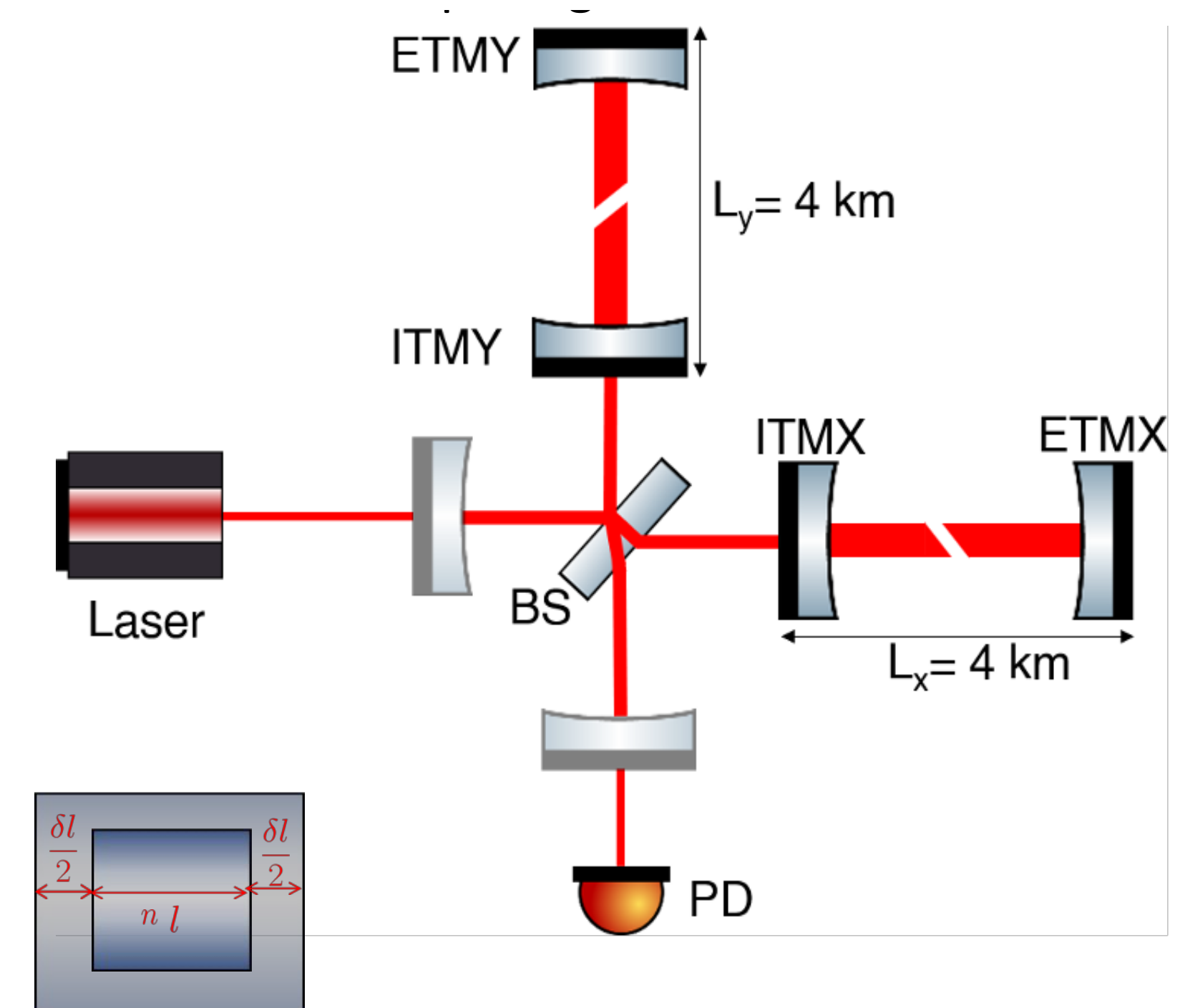


Typical atomic size ~ Bohr radius $a_0 \sim \frac{1}{m_e \alpha}$

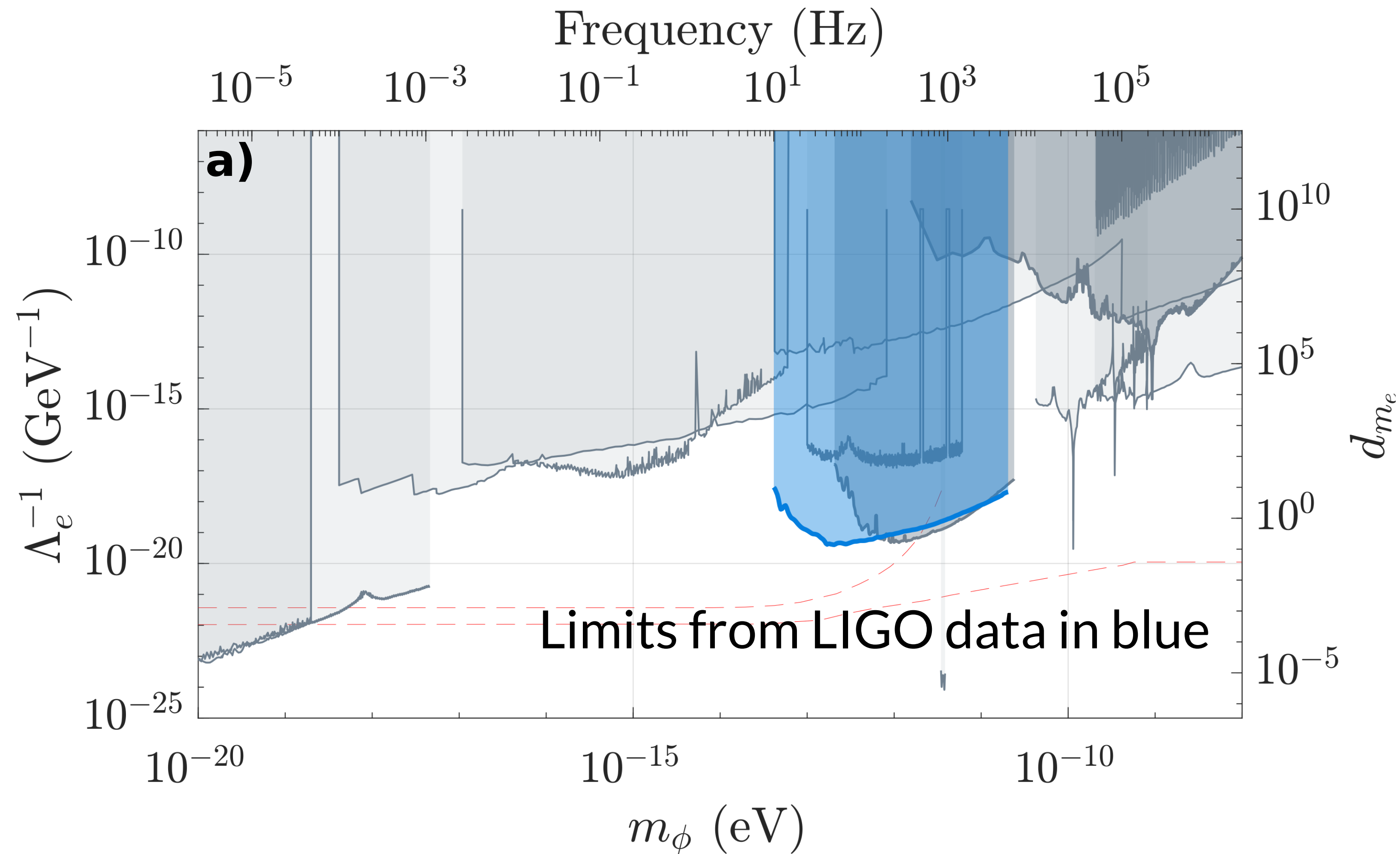
Scalar ULDM leads to oscillations in the beamsplitter and the test masses

$$\frac{\delta l}{l} = - \left(\frac{\delta \alpha}{\alpha} + \frac{\delta m_e}{m_e} \right),$$

$$\frac{\delta n}{n} = -5 \cdot 10^{-3} \left(2 \frac{\delta \alpha}{\alpha} + \frac{\delta m_e}{m_e} \right)$$



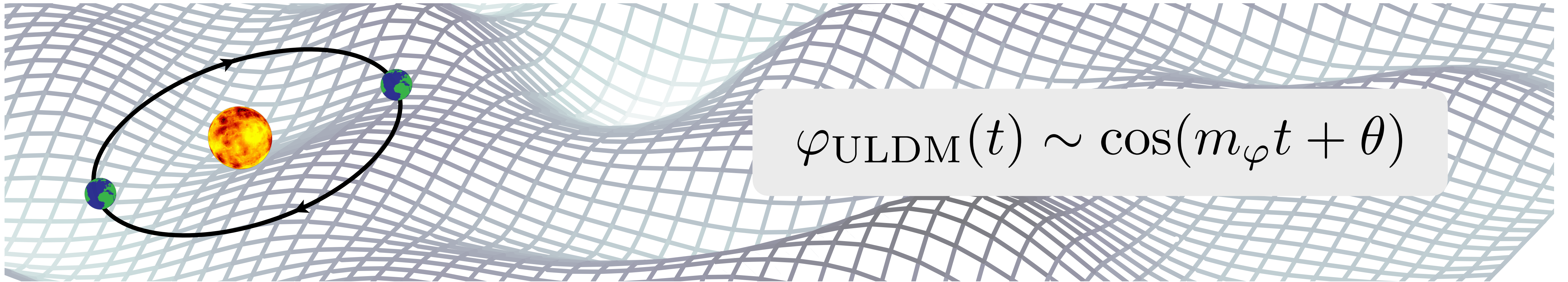
Laser interferometers are complementary (higher mass)



[atom interferometers probe $\sim 10^{-16}$ to 10^{-13} eV]

Göttl et al, PRL
arXiv:2401.18076

Much more could be said...



ULDM can give rise to:

1. Changes in fundamental constants (scalar ULDM)
2. Accelerations on test masses (vector ULDM)
3. Precession of spins (pseudoscalar ULDM)
4. Mimic GW signals (spin-2 ULDM)

Lots of activity in this area, including many groups in the UK
(apologies to all of those I wasn't to mention)

Summary

XLZD is an incredible detector with potential to detect many types of particle dark matter candidate

And there is a very active and comprehensive programme of activity to search for dark matter across

