

Phenomenology Beyond Colliders

Clare Burrage

University of Nottingham



University of
Nottingham
UK | CHINA | MALAYSIA

The Hunt for Light Scalar Fields at the Precision Frontier

Clare Burrage

University of Nottingham

Outline:

What can we do in space?

What can we do with atoms?

What can we do with photons?



University of
Nottingham

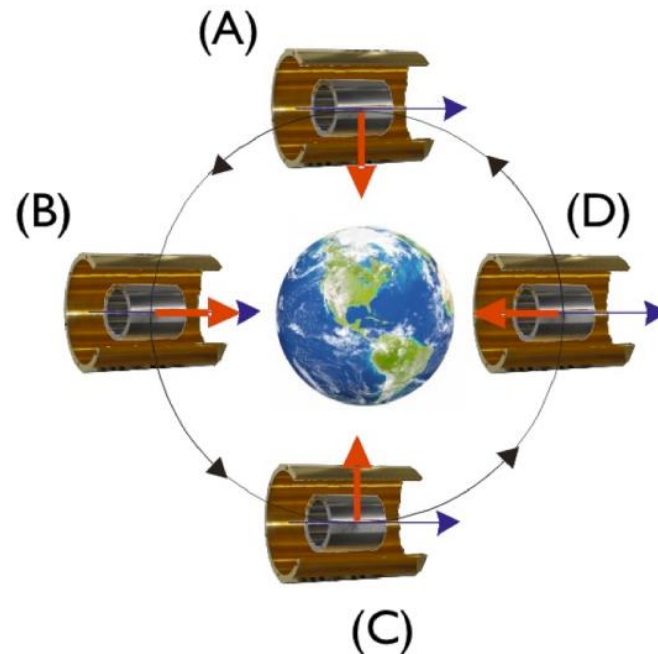
UK | CHINA | MALAYSIA

A deep space image showing a vast field of galaxies and stars against a black background. The galaxies are of various shapes and sizes, some appearing as bright, glowing clouds, others as more distant, faint points of light. The stars are scattered throughout the field, some appearing as bright, multi-pointed stars, others as smaller, dimmer dots. The overall scene is a dense, colorful mosaic of cosmic objects.

WHAT CAN WE DO IN SPACE?

MICROSCOPE

Testing the Weak Equivalence Principle
Concentric cylinders of platinum-rhodium alloy &
titanium-aluminium-vanadium alloy

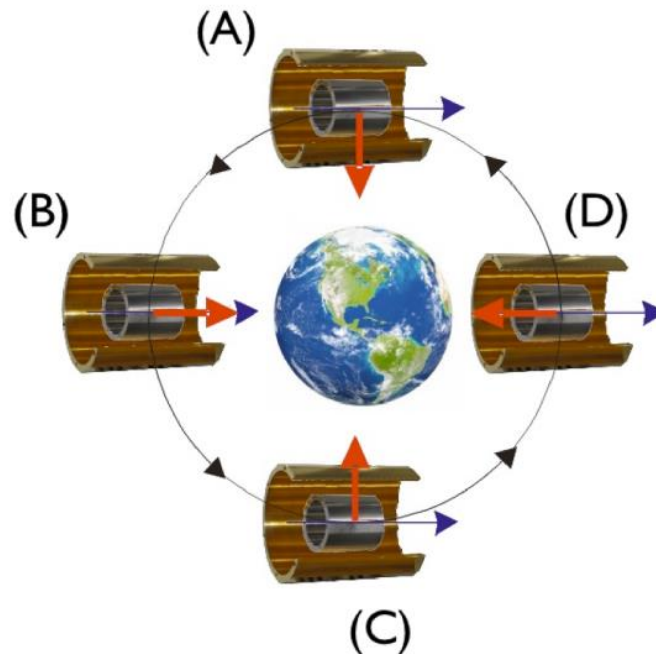


Orbited at a height of 710 km for ~135 days

MICROSCOPE

Eötvös parameter

$$\eta_{A,B} = 2 \frac{a_A - a_B}{a_A + a_B}$$

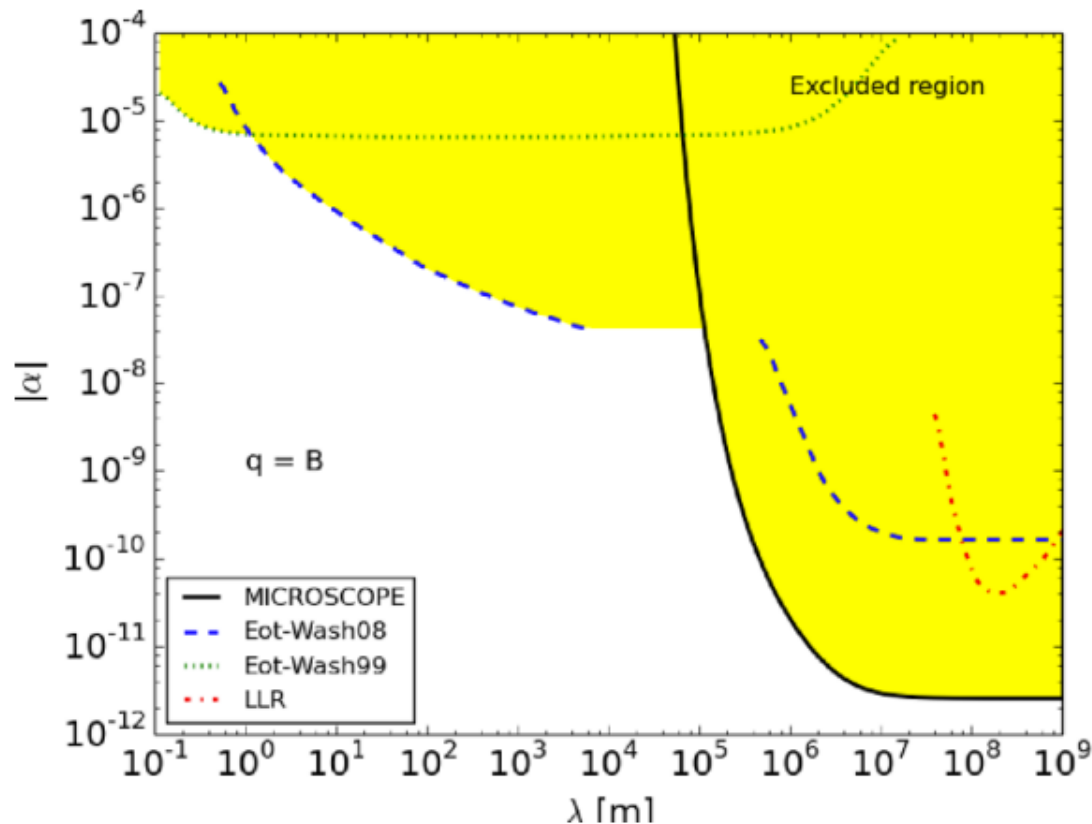


$$\eta(\text{Ti}; \text{Pt}) = [-1.5 \pm 2.3 \pm 1.5] \times 10^{-15}$$

Violations of WEP

New vector or scalar mediated Yukawa interactions

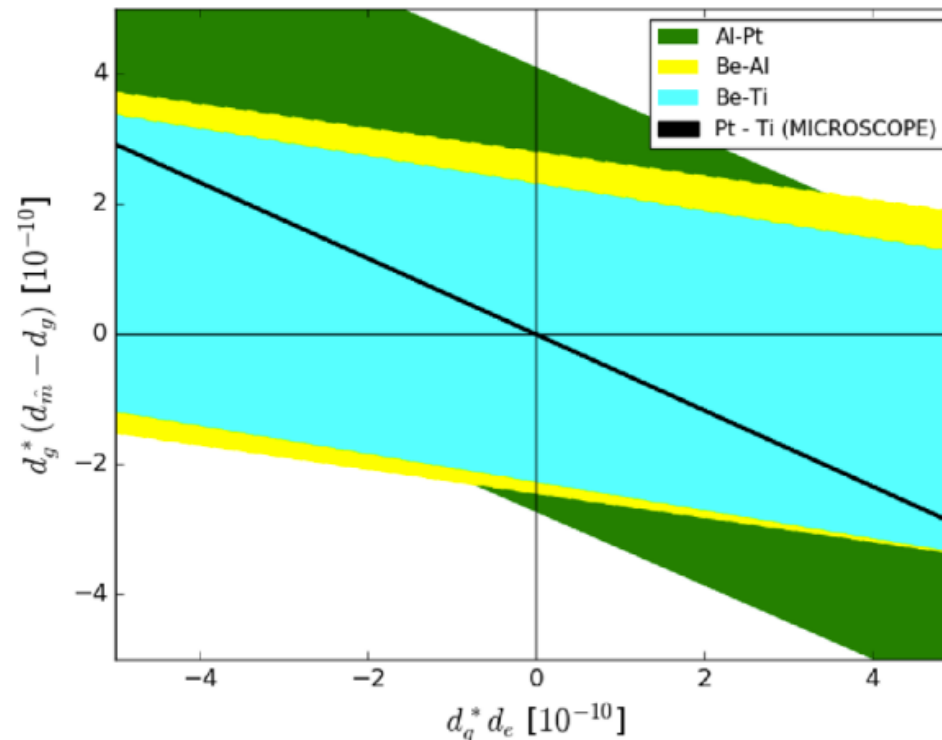
$$U(r) = -\frac{GM}{r} \left[1 + \alpha \exp\left(-\frac{r}{\lambda}\right) \right],$$



Violations of WEP

Constraints on a massless 'dilaton'

$$\mathcal{L}_{\text{int}} = \frac{\phi}{M_{\text{Pl}}^2} \left[\frac{d_e}{4e^2} F_{\mu\nu} F^{\mu\nu} - \frac{d_g \beta_3}{2g_3} F_{\mu\nu}^A F^{A\mu\nu} - \sum_{i=e,u,d} (d_{m_i} + \gamma_{m_i} d_g) m_i \bar{\psi}_i \psi_i \right]$$

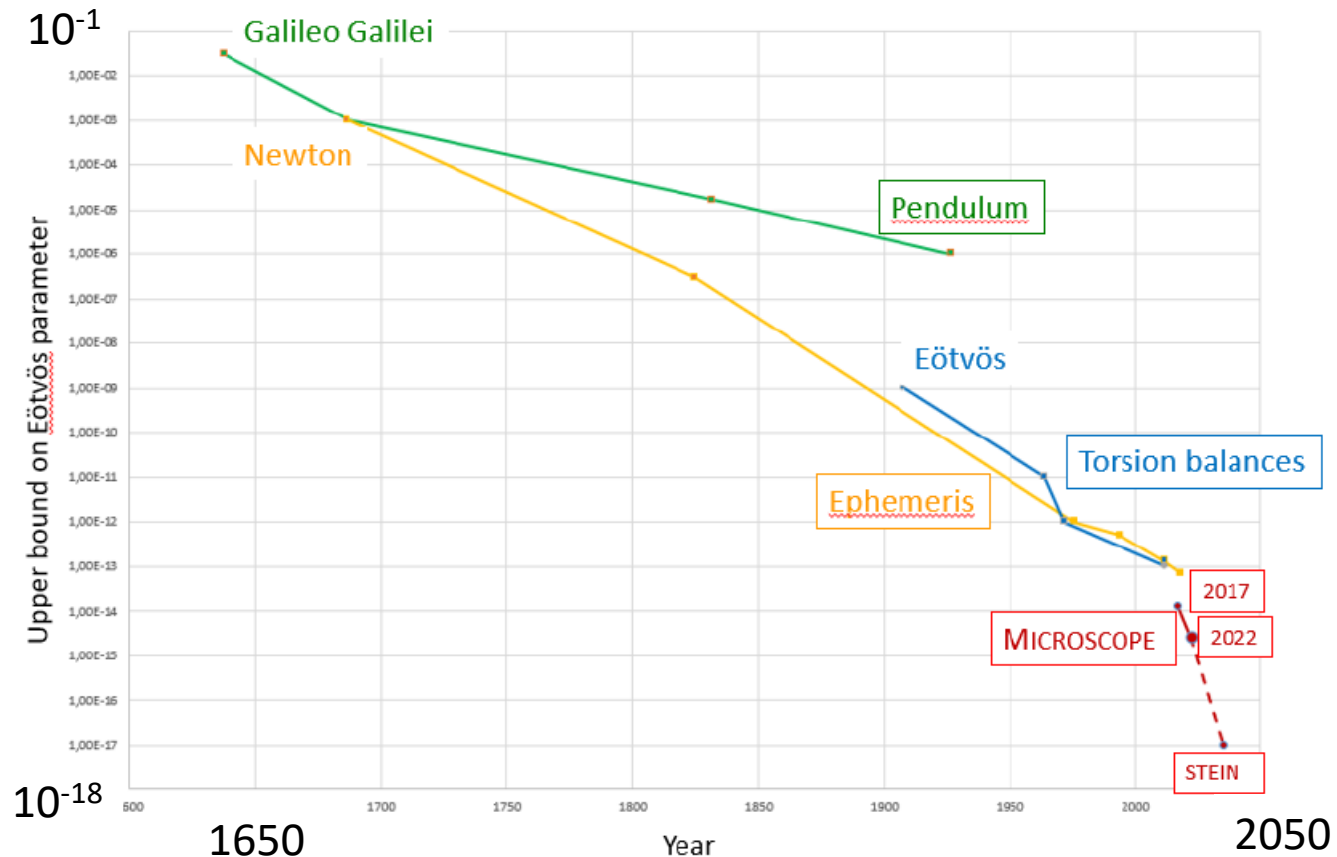


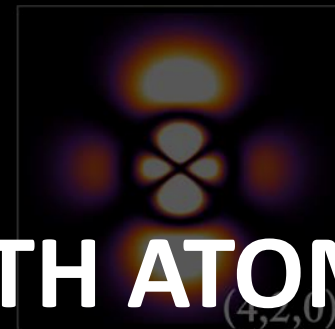
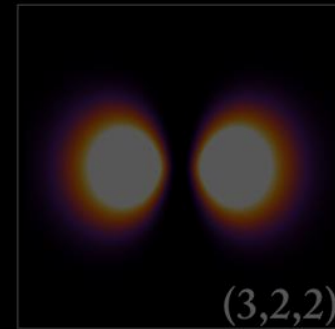
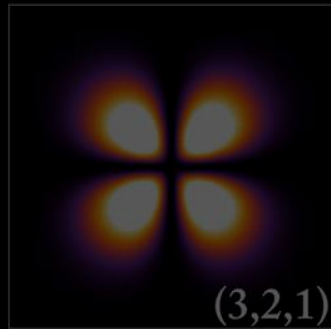
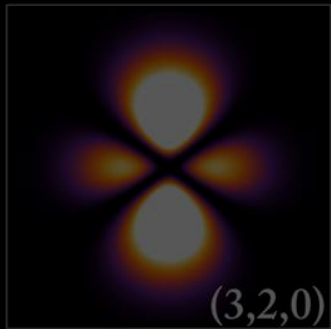
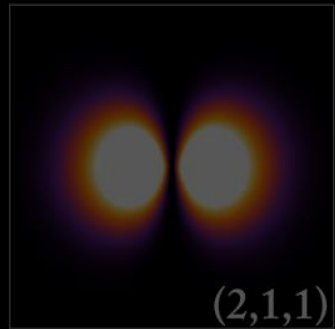
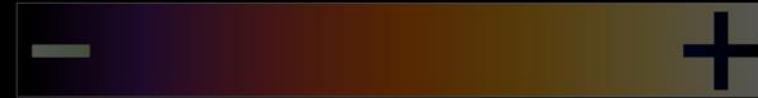
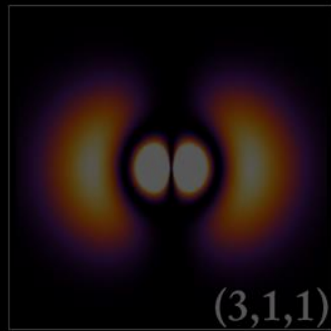
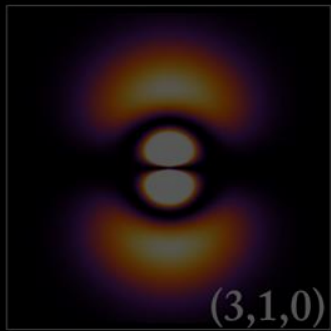
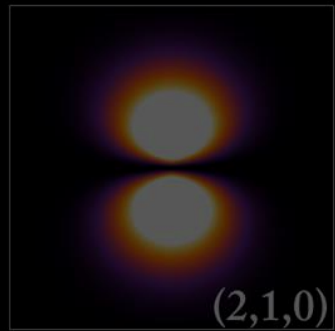
Berge. Rep. Prog. Phys. 86 (2023) 066901

Damour, Donoghue. PRD (2010) 8 084033, CQG (2010) 20 202001,

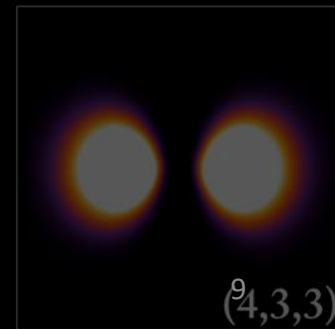
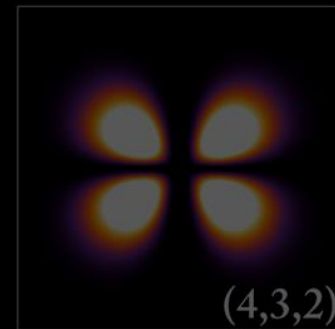
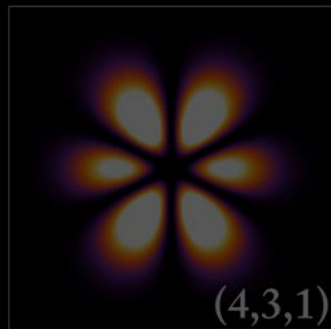
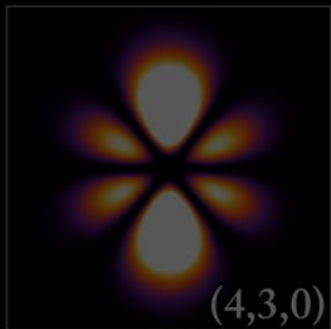
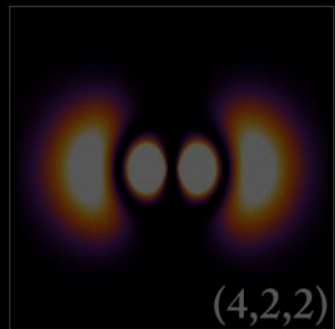
Future Searches

Satellite Test of Einstein's gravitation theory (STEIN) Step 2 proposal for an ESA F-Class mission



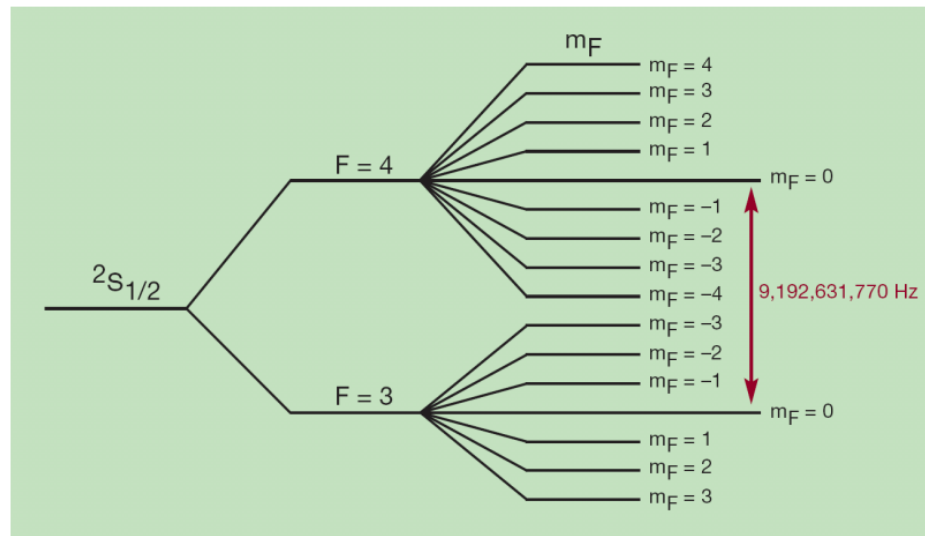


WHAT CAN WE DO WITH ATOMS?



Atomic Clocks

Driven oscillations in a 'two state' system



1 second = 9,192,631,770 transitions between
hyperfine levels of the ground state of caesium 133
(at 0K)

Atomic Clocks

Typical accuracy $\Delta f / f \sim 10^{-18}$

Compare two different clocks:

Different transitions have different sensitivities to different fundamental constants



Ultra-Light Scalar Dark Matter

Light and feebly interacting bosons can be dark matter

A scalar field oscillating in a quadratic potential

$$\phi_{\text{homog}}(t) = \phi_{\infty} \cos(mt + \delta)$$

Gives zero pressure and dark matter energy density

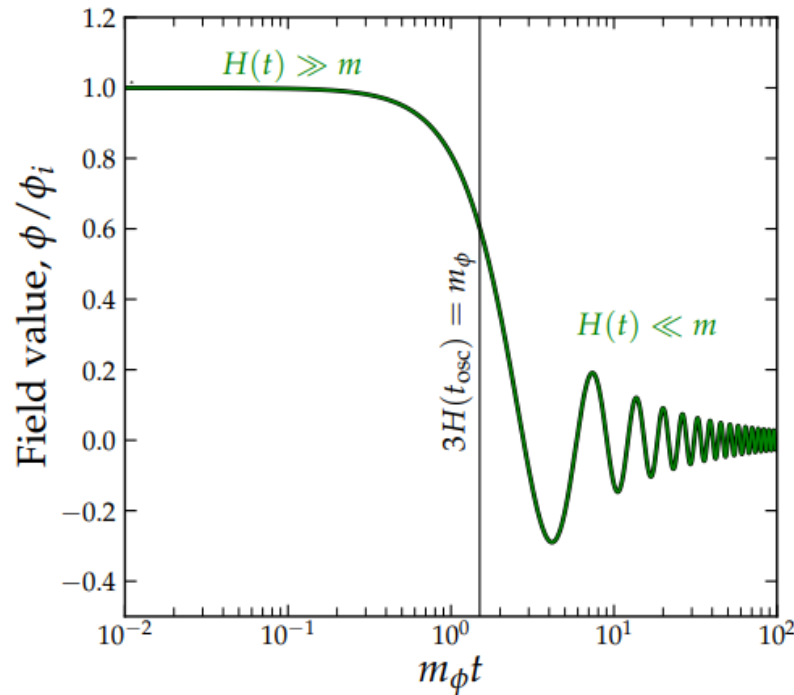
$$\rho_{\text{DM}} = \frac{1}{2} m^2 \phi_{\infty}^2$$

Preskill, Wise, Wilczek. Phys. Lett.B 120 (1983) 127–132. Abbott, Sikivie. Phys. Lett.B 120 (1983) 133–136. Dine, Fischler, Phys. Lett. B 120 (1983) 137–141.

For a review see e.g. Ferreira. Astron. Astrophys. Rev. 29 (2021) 1, 7

Ultra-Light Scalar Dark Matter

Light and feebly interacting bosons can be dark matter



High occupation numbers mean they can be described as classical fields exhibiting a wave-like behaviour

Atomic Clocks

Interactions with matter

$$\mathcal{L} = (\kappa\phi)^n \left(d_\gamma^{(n)} \frac{1}{4} F_{\mu\nu} F^{\mu\nu} - d_{m_e}^{(n)} m_e \bar{\psi}_e \psi_e \right)$$

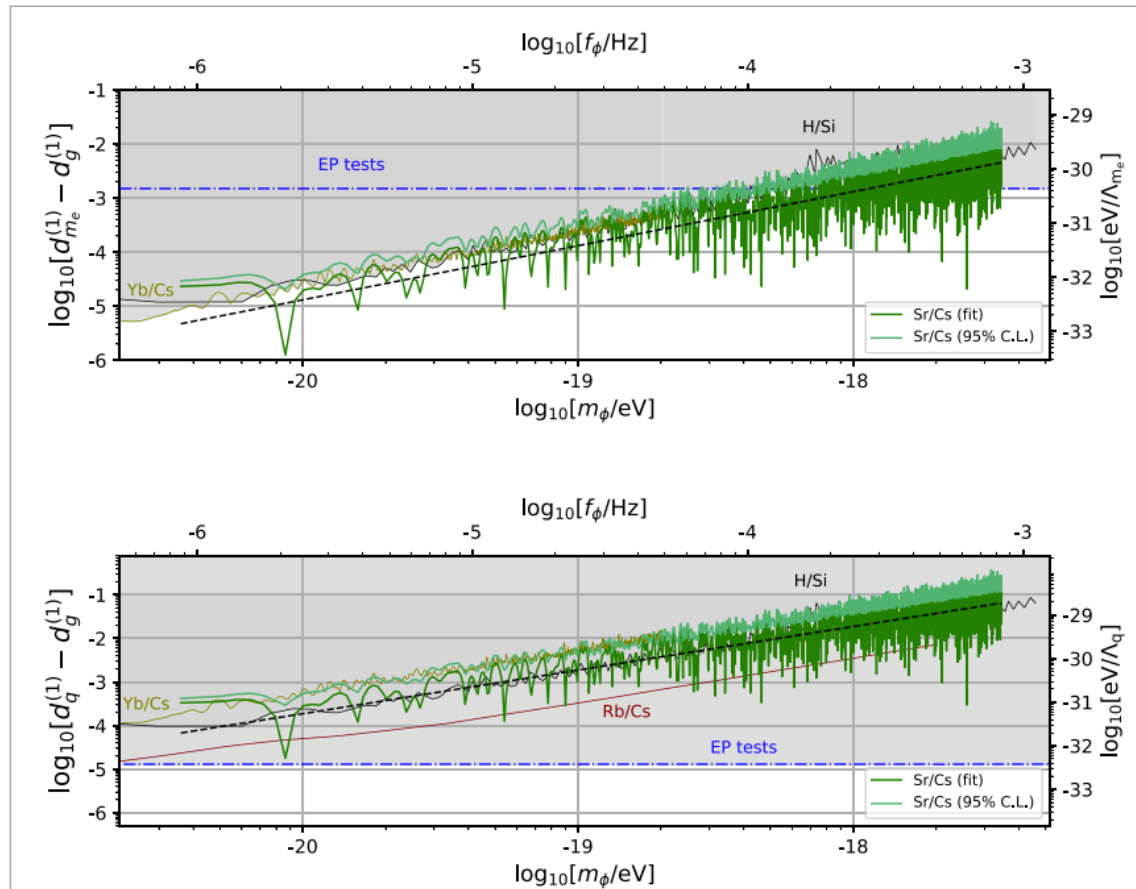
Lead to variations in fundamental constants

$$\frac{\delta\alpha}{\alpha} \equiv d_\gamma^{(n)} (\kappa\phi)^n \quad \frac{\delta m_f}{m_f} \equiv d_{m_f}^{(n)} (\kappa\phi)^n \quad \frac{\delta\Lambda_{\text{QCD}}}{\Lambda_{\text{QCD}}} \equiv d_g^{(n)} (\kappa\phi)^n$$



Atomic Clocks – Constraints on ULDM

$$\phi_{\text{homog}}(t) = \phi_{\infty} \cos(mt + \delta)$$



Sherrill et al, New J. Phys. 25 093012 (2023)

For dark energy and modified gravity constraints and forecasts see:

Elder et al. Phys.Rev.D 112 (2025) 4, 044053

Ultra-Light Scalar Dark Matter

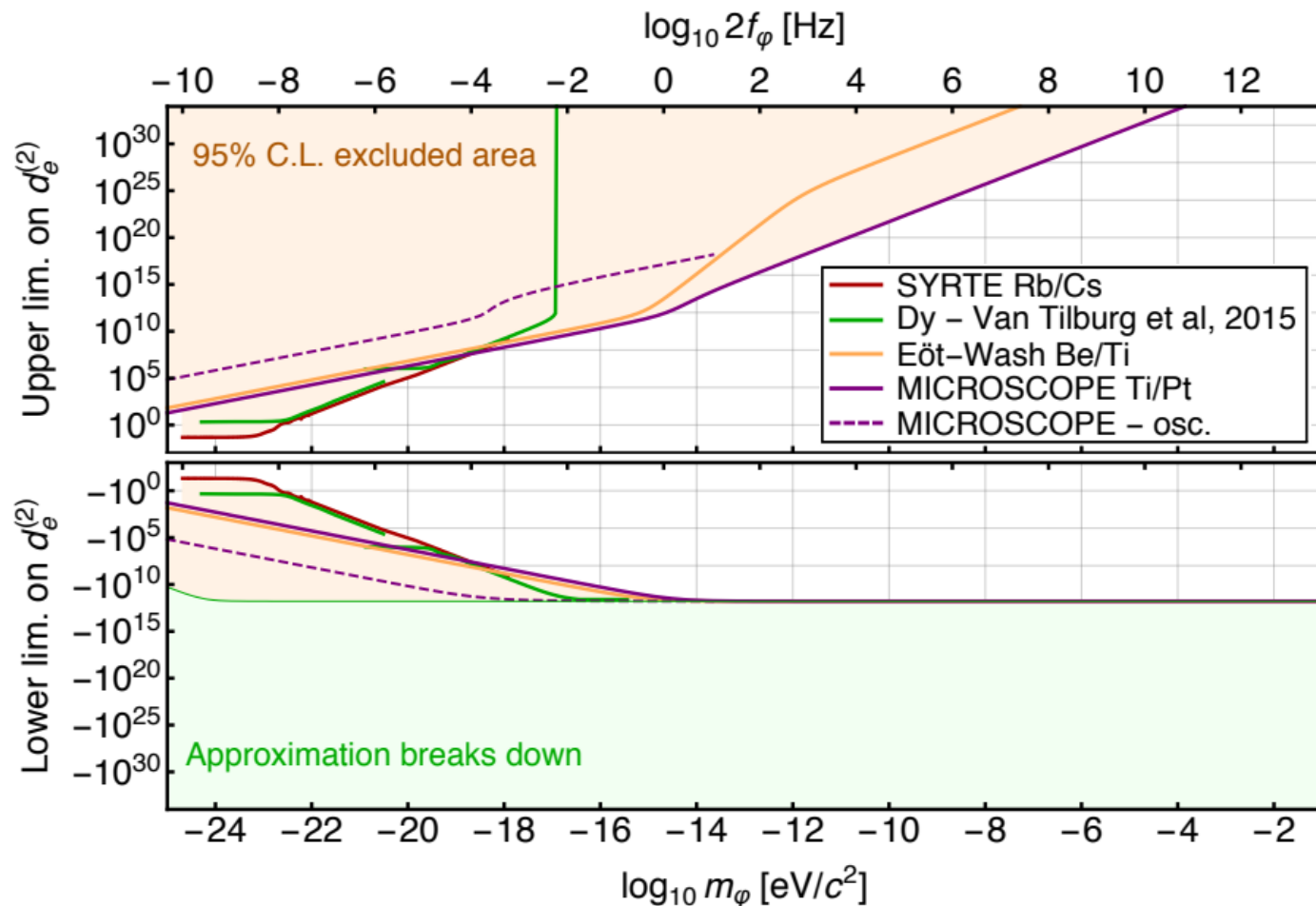
Quadratic interactions from
e.g. symmetric Higgs portals, axion models

$$\mathcal{L}_{int} = \frac{(\kappa\phi)^2}{2} \left[\frac{d_e}{4} F_{\mu\nu} F^{\mu\nu} - \frac{d_g \beta_3}{2g_3} F_{\mu\nu}^A F^{A\mu\nu} - \sum_{i=u,d,e} (d_{m_i} + \gamma_{m_i} d_g) m_i \bar{\psi}_i \psi_i \right]$$

Leading to: Varying fundamental ‘constants’
Macroscopic couplings to non-relativistic energy
density

$$S = \int d^4x \left(\frac{1}{2} (\partial\phi)^2 - \frac{1}{2} m^2 \phi^2 - \frac{1}{2} \rho \alpha \phi^2 \right)$$

Ultra-Light Scalar Dark Matter - Constraints



Hees, Minazzoli, Savalle, Stadnik, Wolf. Phys. Rev. D 98 no. 6, (2018) 064051

See also: Banerjee, Perez, Safronova, Savoray. JHEP 10 (2023) 042

Bartnick, Springmann, Stelzl, Weiler. arxiv:2509.25305

Bauer, Chakraborti, Rostagni, JHEP 05 (2025) 023

Ultra-Light Scalar Dark Matter

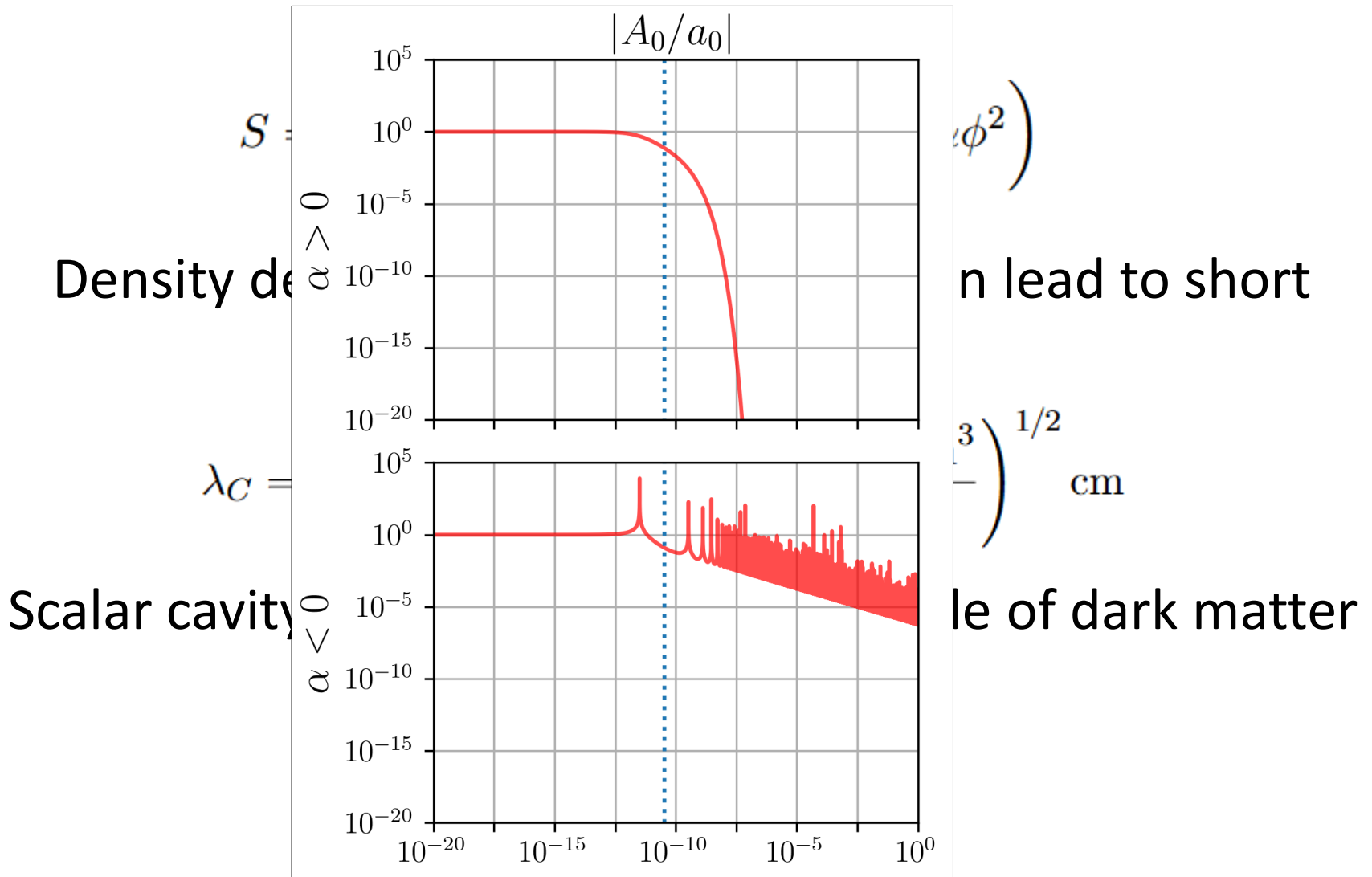
$$S = \int d^4x \left(\frac{1}{2}(\partial\phi)^2 - \frac{1}{2}m^2\phi^2 - \frac{1}{2}\rho\alpha\phi^2 \right)$$

Density dependent effective mass, can lead to short
Compton wavelengths

$$\lambda_C = \frac{h}{m_{\text{eff}}c} \approx 10^{-5} \left(\frac{\text{GeV}^{-2}}{\alpha} \right)^{1/2} \left(\frac{\text{g/cm}^3}{\rho} \right)^{1/2} \text{ cm}$$

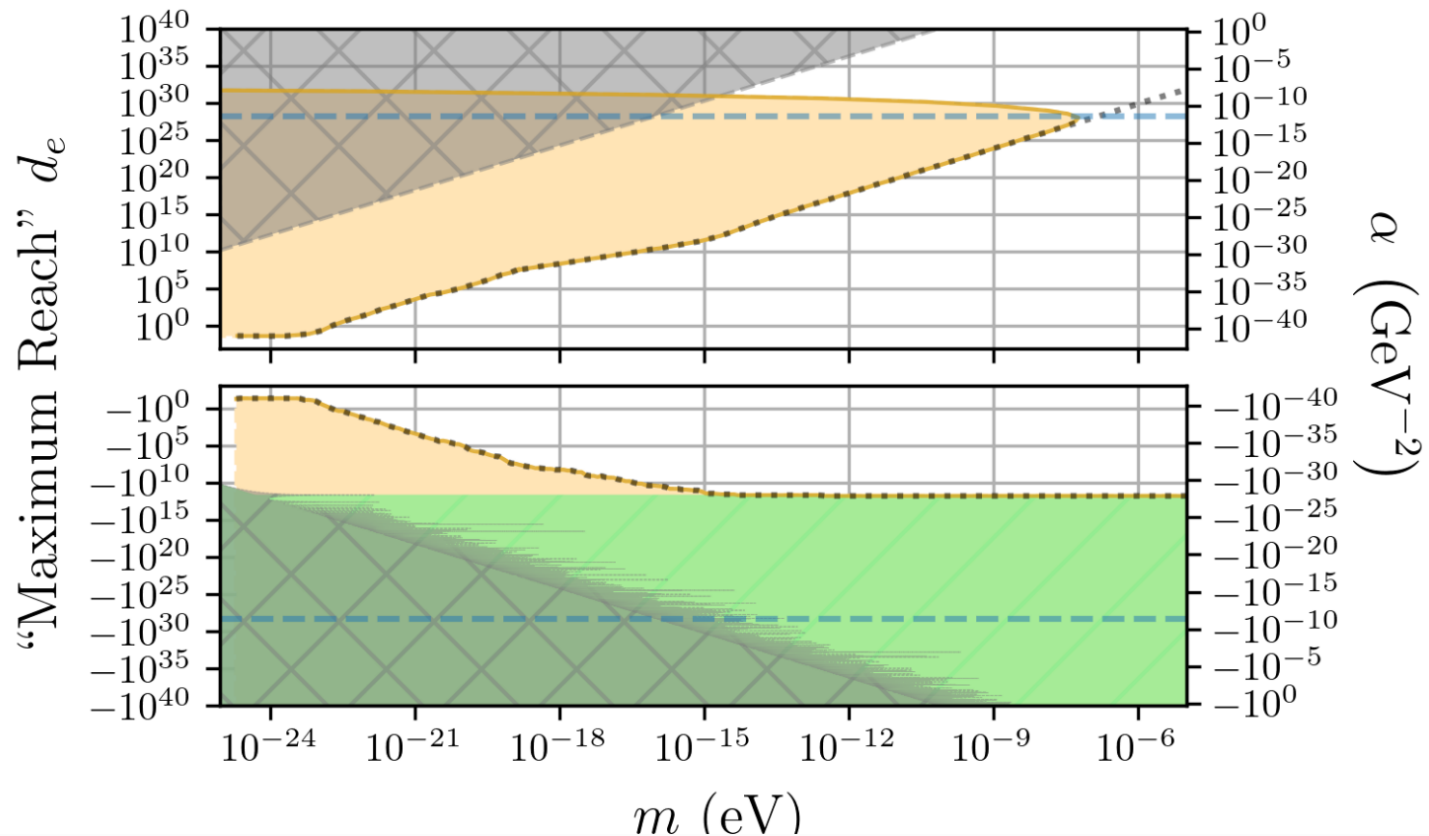
Scalar cavity can suppress the amplitude of dark matter
oscillations

Ultra-Light Scalar Dark Matter



Effects of a Spherical Cavity

Naive rescaling weakens constraints at strong coupling



Also new opportunities for space-based searches

Networks of atomic clocks

Search for spatial gradients from dark matter velocity dispersion

$$\phi(t, \mathbf{x}) = \phi_0 \cos(\omega t - \mathbf{k} \cdot \mathbf{x})$$

Induced frequency shift

$$\begin{aligned} \frac{\delta(\nu^A / \nu^B)}{\nu^A / \nu^B} &= \kappa \phi_0 d_{\text{eff}} [\cos(\omega t) - \cos(\omega t - \delta)] \\ &\approx \kappa \phi_0 d_{\text{eff}} \sin(\omega t) \delta, \end{aligned}$$

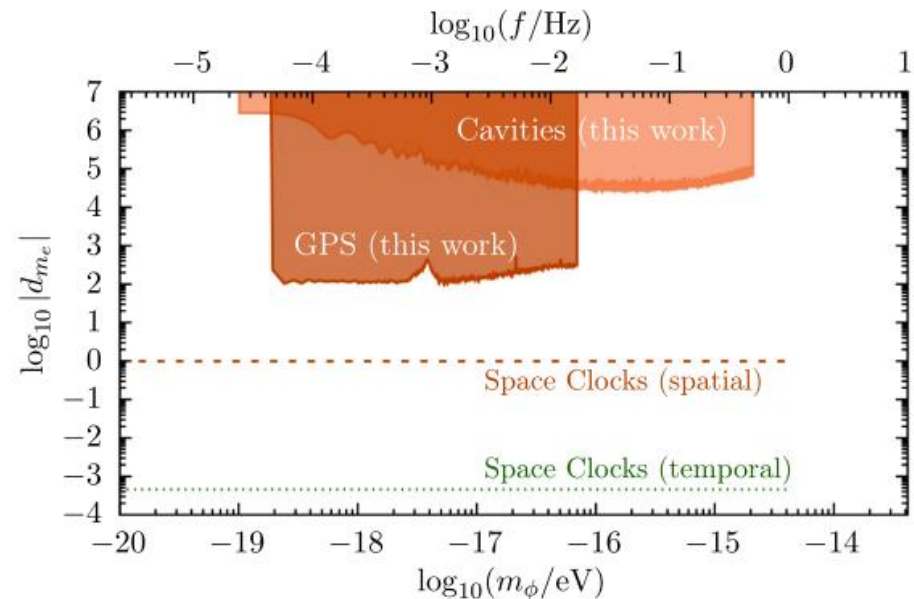
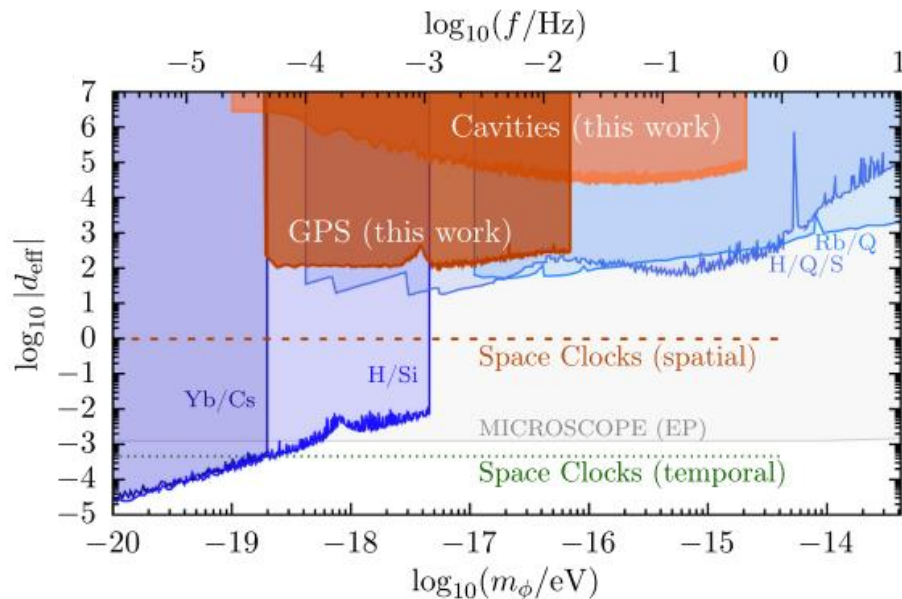
$$\delta = \omega \Delta t - \mathbf{k} \cdot \Delta \mathbf{x},$$

Δt the time delay due to signal propagation

$\Delta \mathbf{x}$ the distance between the oscillators

Networks of atomic clocks

Constraints from optical cavities, 2220 km separation,
Atomic clocks on GPS satellites, 30 satellites orbiting at
a radius of 26,000 km



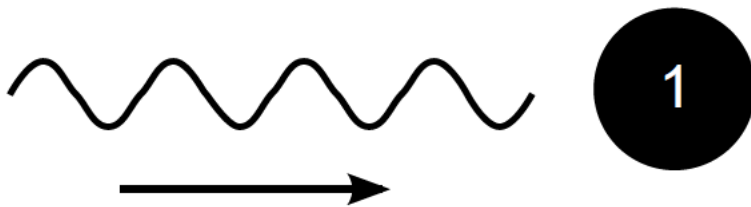
Filzinger et al. Phys Rev Lett. 134 031001 (2025)

Schioppo et al. Nat. Commun. 13, 212 (2022)

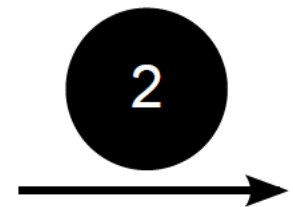
Atom Interferometry

An interferometer where the wave is made of atoms

Atoms can be moved around by absorption of laser photons

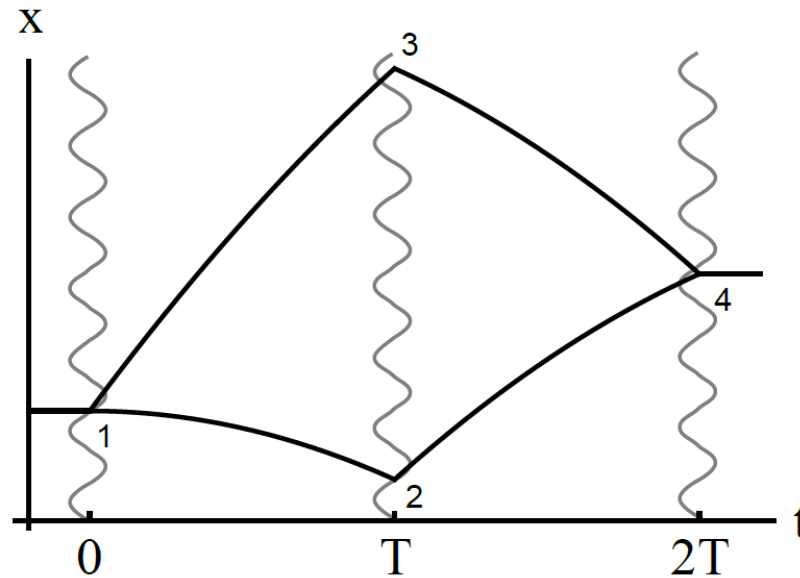


Photon Momentum = k
Atom in ground state



Atom in excited state
with velocity = V

Atom interferometry

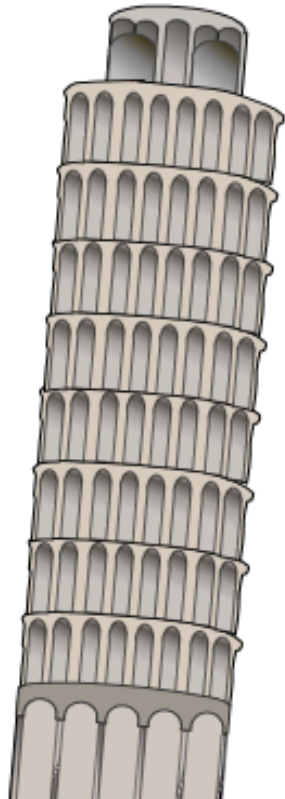


Phase difference accumulates from:
the propagation phase, the laser interaction
phase, and the final wavepacket separation phase

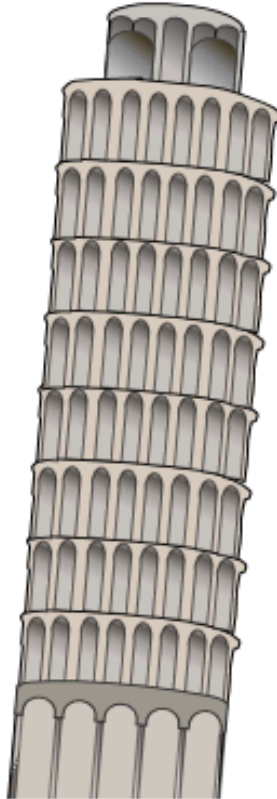
Determines probability atoms exit in ground or excited
state

As a test of the equivalence principle

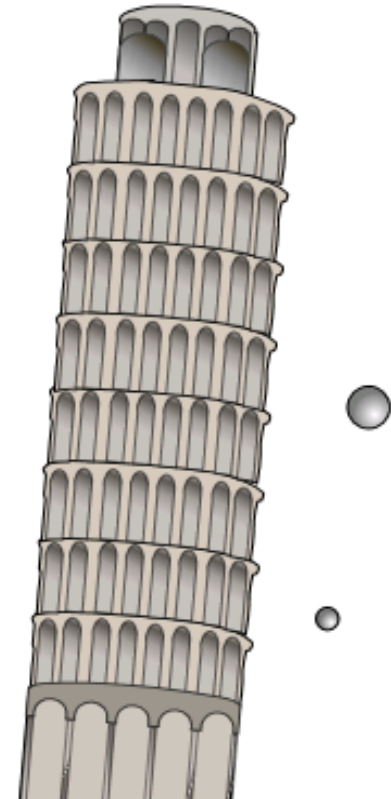
Do objects with different composition fall at the same rate?



Old idea



Galileo



New physics?

The Chameleon



A scalar field with canonical kinetic terms, non-linear potential, and direct coupling to matter

$$S_\phi = \int d^4x \sqrt{-g} \left(-\frac{1}{2}(\partial\phi)^2 - V(\phi) - A(\phi)\rho_m \right)$$

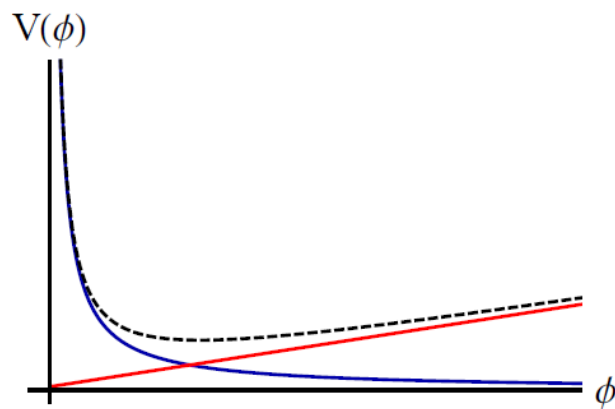
$$V(\phi) = \frac{\Lambda^5}{\phi}, \quad A(\phi) = \frac{\phi}{M} ,$$

Varying Mass

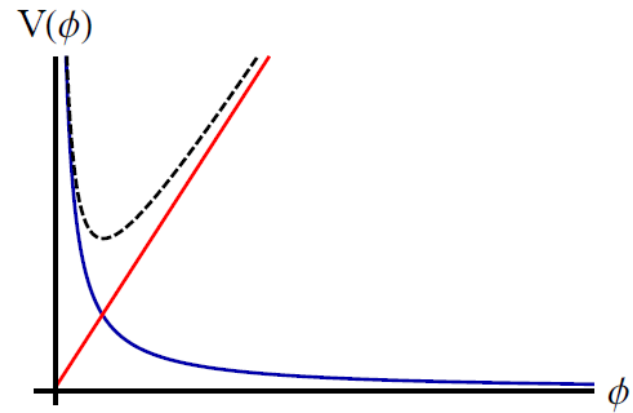
Dynamics governed by an effective potential

$$V_{\text{eff}} = \frac{\Lambda^5}{\phi} + \frac{\phi}{M} \rho$$

Non-linearities in the potential mean that the mass of the field depends on the local energy density



Low density



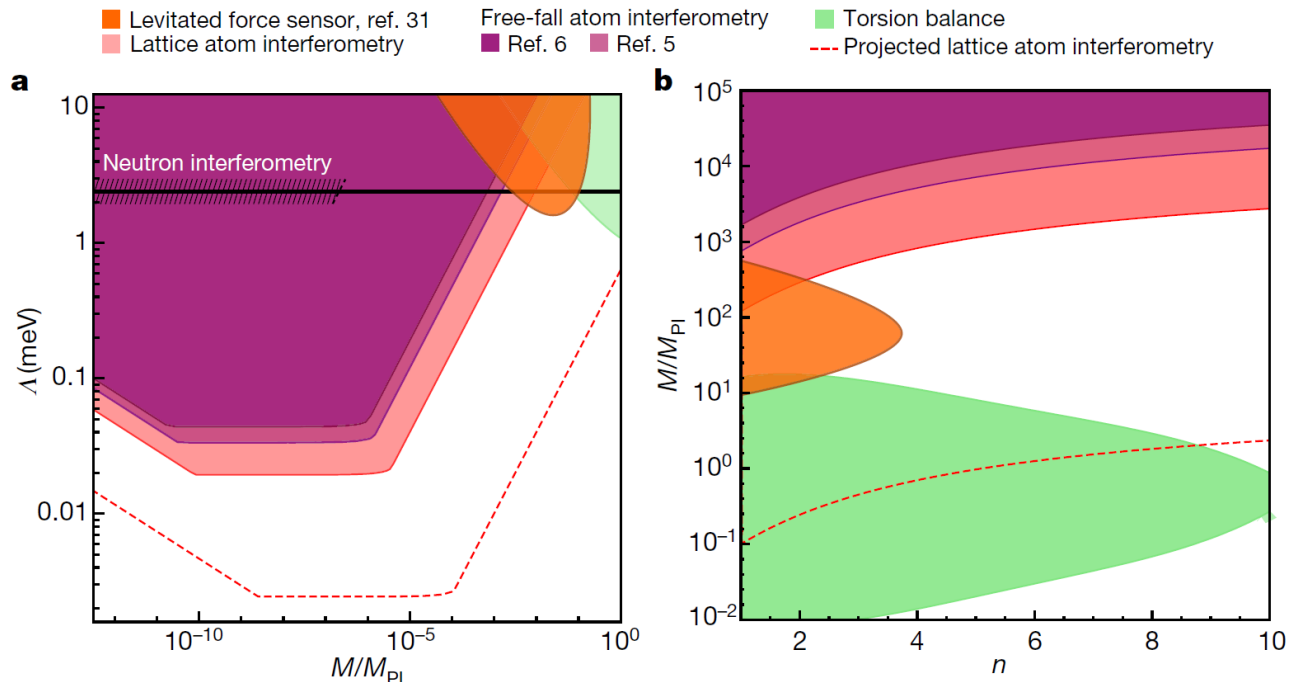
High density

Field is suppressed inside compact objects

Chameleon: Combined Constraints

Bare potential: $V(\phi) = \Lambda^{n+4}/\phi^n$

Bare matter coupling: M



Anomalous acceleration $< 6.2 \text{ nm s}^{-2}$

Panda, Tao, Ceja, Khoury, Tino, Muller. Nature 631, 515–520 (2024)

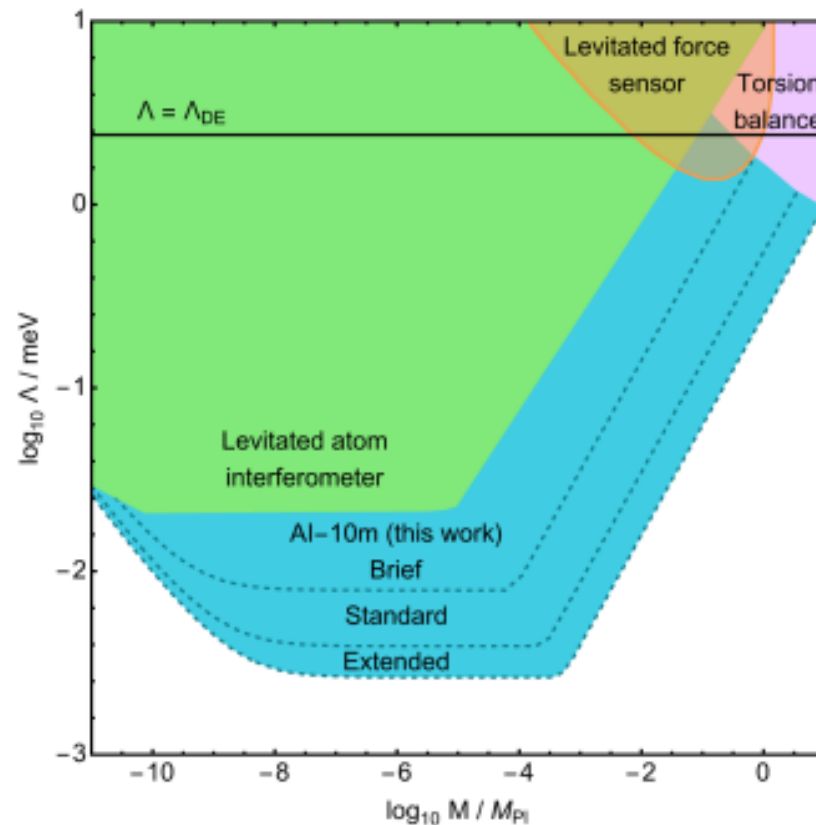
See also: Sabulsky, Dutta, Hinds, Elder, CB, Copeland. Phys.Rev.Lett. 123 (2019) 6, 061102

Brax, Davis, Elder. Phys. Rev. D 107 (2023) 084025

Chameleon: Combined Constraints

Bare potential: $V(\phi) = \Lambda^{n+4}/\phi^n$

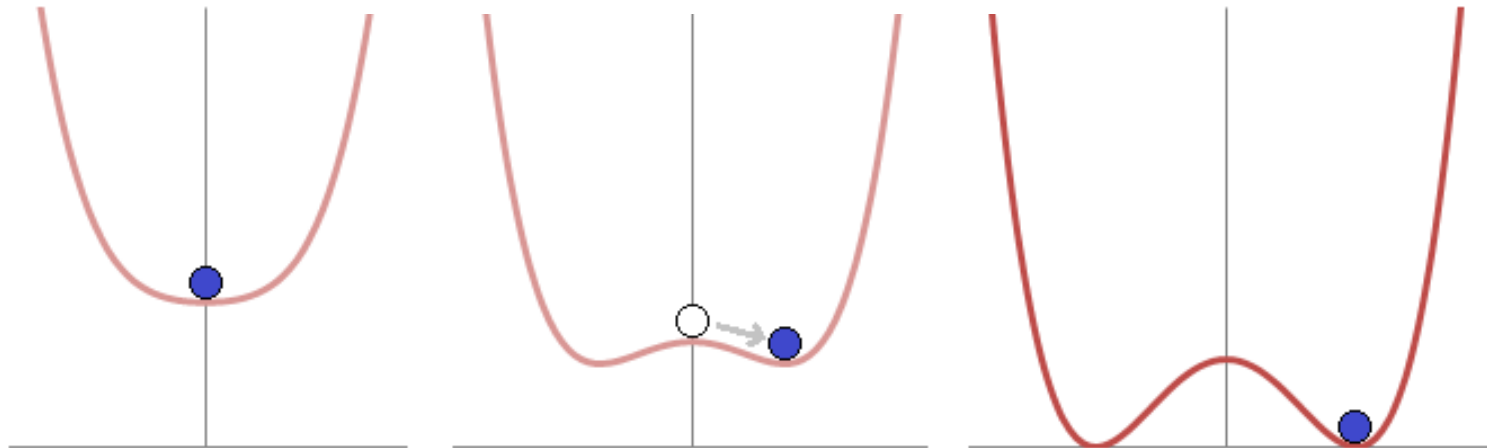
Bare matter coupling: M



Symmetron Screening

Effective potential

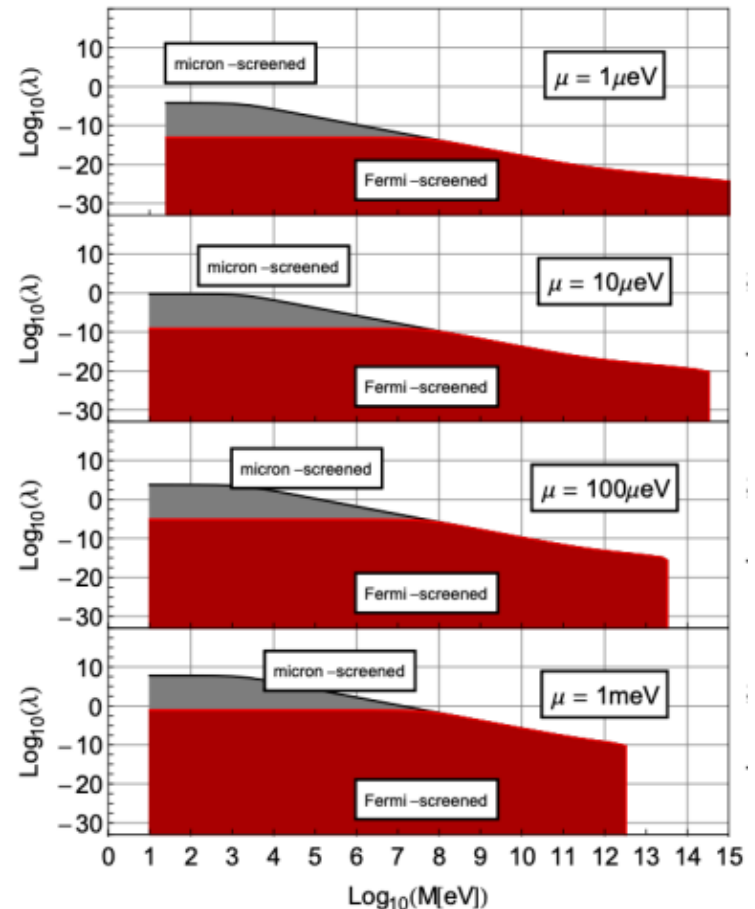
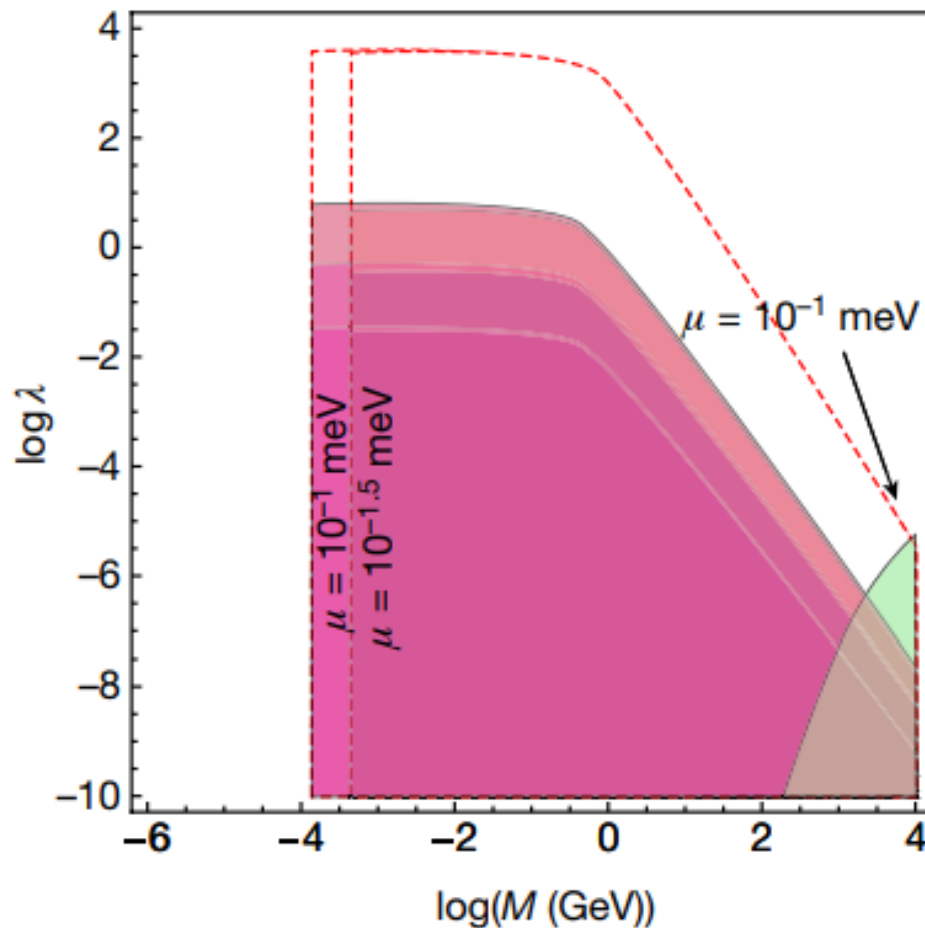
$$V_{\text{eff}}(\phi) = \frac{1}{2} \left(\frac{\rho}{M^2} - \mu^2 \right) \phi^2 + \frac{1}{4} \lambda \phi^4$$



Fifth force vanishes when symmetry is restored

$$F = \phi \nabla \phi / M^2$$

Symmetron: Laboratory Constraints



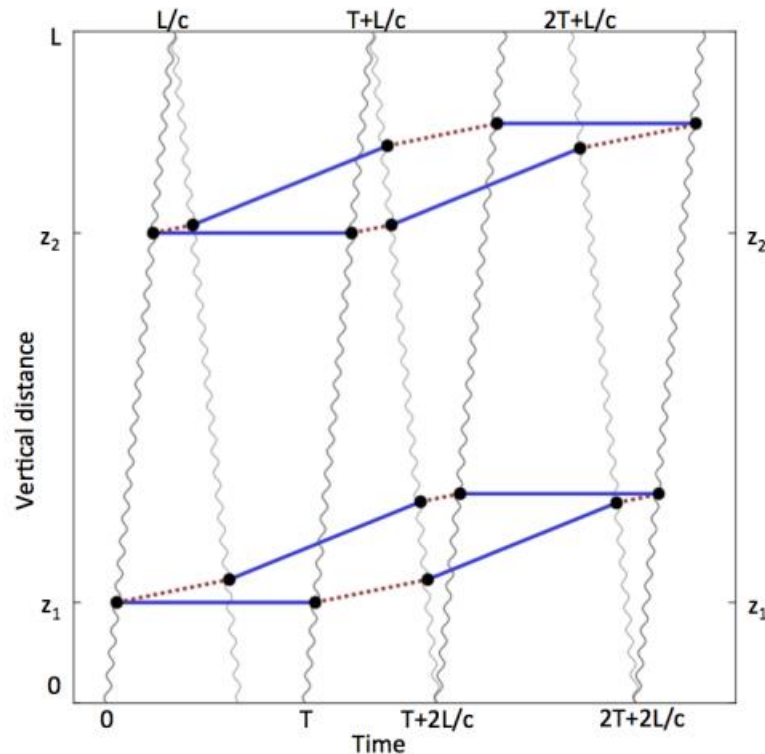
Panda, Tao, Ceja, Khoury, Tino, Muller. Nature 631, 515–520 (2024), Jenke, Bosina, Micko, Pitschmann, Sedmik, Abele, Eur. Phys. J. ST 230, 1131 (2021), 2012.07472.

Constraints do not contain corrections from Millington, Udemba. arXiv:2508.16726₃₁

Detecting Oscillations

Transition frequency may also oscillate with time

Differential interferometer, one laser interacts with (at least) two spatially separated clouds of atoms



Detecting Oscillations

Over the interferometer

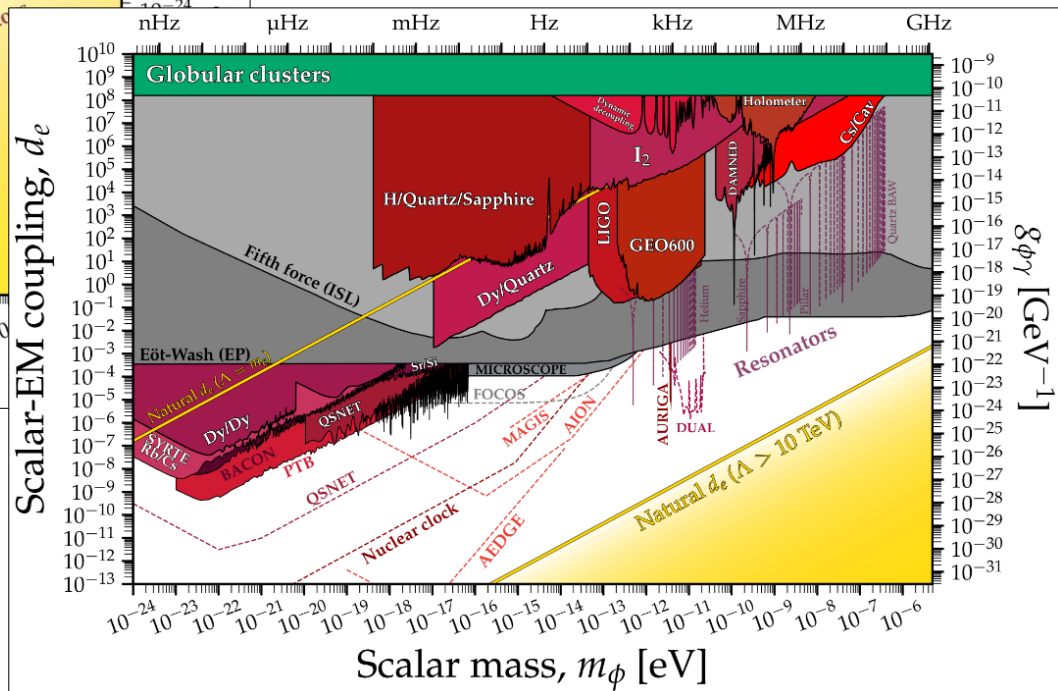
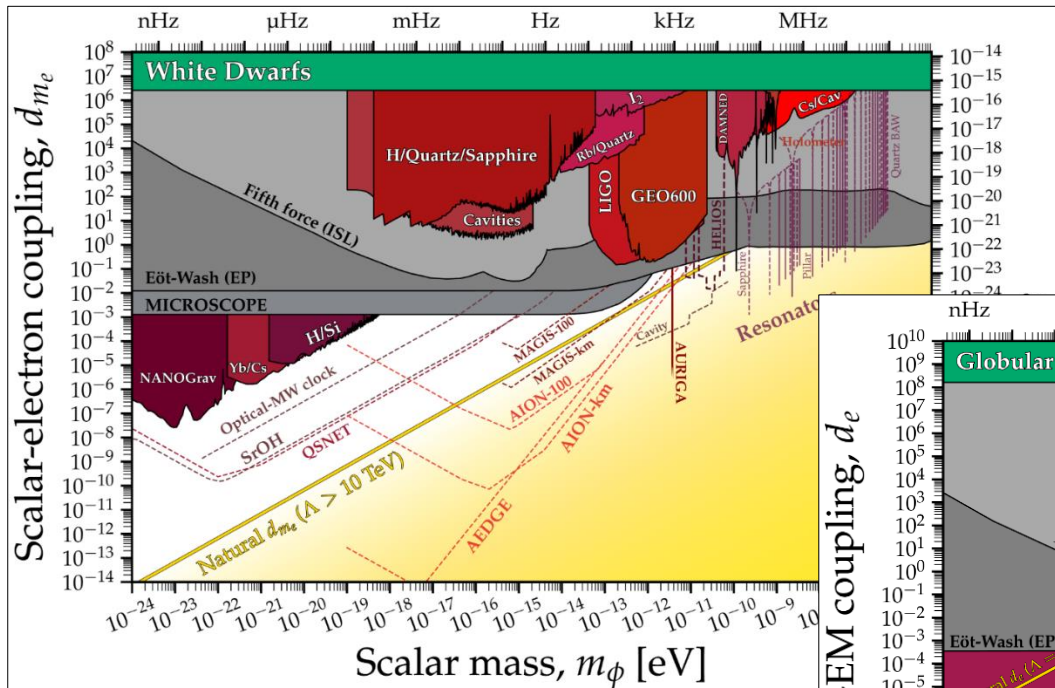
$$\overline{\Phi}_s = 8 \frac{\Delta\omega_A}{m_\phi} \times \left| \sin \left[\frac{m_\phi n L}{2} \right] \sin \left[\frac{m_\phi (T - (n-1)L)}{2} \right] \sin \left[\frac{m_\phi T}{2} \right] \right|$$

Terrestrial Very Long Baseline AI, e.g. AION, Magis,
100m and 1km baseline

Space based AEDGE, $\sim 10^5$ km baseline



ULDM Sensitivity



<https://cajohare.github.io/AxionLimits/>

Also sensitive to other dark matter models e.g. spin 2.

Blas, Carlton, McCabe. Phys. Rev. D. 111 (2025) 11 11 5020

Also sensitive to dark matter gradients. Geraci, Derevianko. Phys. Rev. Lett. 117, 261301 (2016)

AI GW Detector

GW distortion of space time

$$ds^2 = c^2 dt^2 - (1 + h \sin(\omega(t - z) + \phi_0)) dz^2 - (1 - h \sin(\omega(t - z) + \phi_0)) dy^2 - dx^2$$

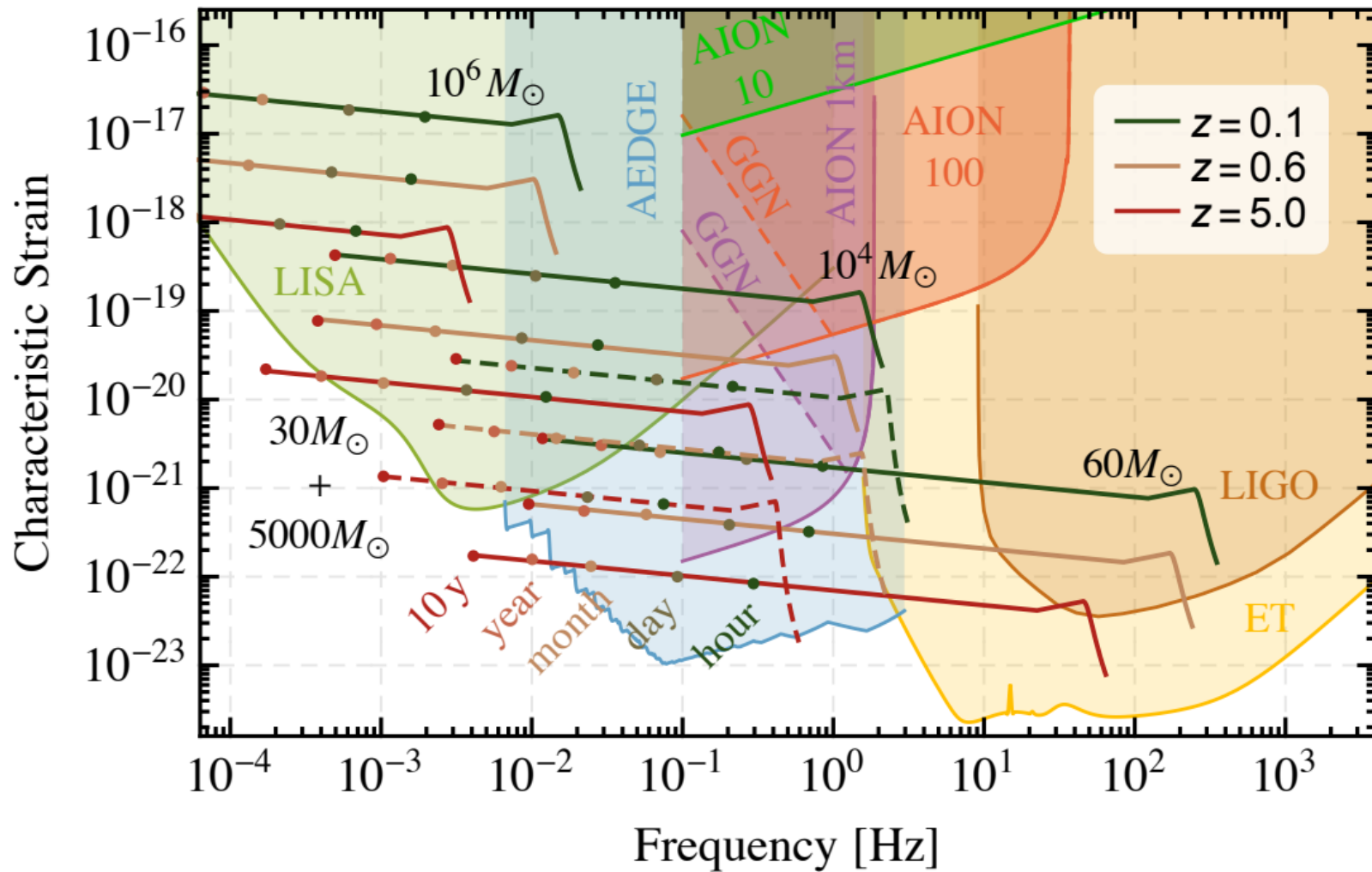
will change the space-time coordinates of the
interaction points

Spatial separation of the two interferometers much
larger than the separation of the wavepackets in each
interferometer

Phase shift difference between interferometers

$$\Delta\phi_{\text{diff}} = 4 \frac{h \omega_A}{\omega} \sin^2 \left(\frac{\omega T}{2} \right) \sin \left(\frac{\omega (z_2 - z_1)}{2c} \right) \sin \left(\omega T + \frac{\omega (z_1 + z_2)}{2c} + \phi_0 \right)$$

Sensitivity

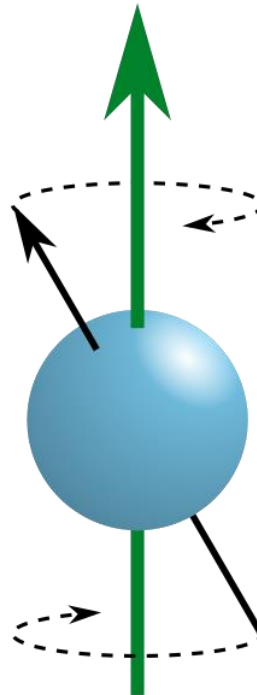


Spin Precession

In the presence of a magnetic field the spin dipole axis precesses

- Can be detected with NMR

Can search for things that look like magnetic fields



Axions – Pseudo-scalar ULDM

$$\mathcal{L}_{\text{axion}} = \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{DS}} + \frac{a}{4f_\gamma} F_{\mu\nu} \tilde{F}_{\mu\nu} + \frac{a}{4f_G} \text{Tr} G_{\mu\nu} \tilde{G}_{\mu\nu} + \frac{\partial_\mu a}{f_l} \sum_\alpha \bar{l}_\alpha \gamma_\mu \gamma_5 l_\alpha + \frac{\partial_\mu a}{f_q} \sum_\beta \bar{q}_\beta \gamma_\mu \gamma_5 q_\beta$$

Coupling to photons

Coupling to Gluons

Coupling to leptons

Coupling to quarks

Axion solution to the strong CP problem*

Axion-like particles – pseudo scalars with similar couplings

* But see also: Ai, Cruz, Garbrecht, Tamarit. Phys.Lett.B 822 (2021) 136616

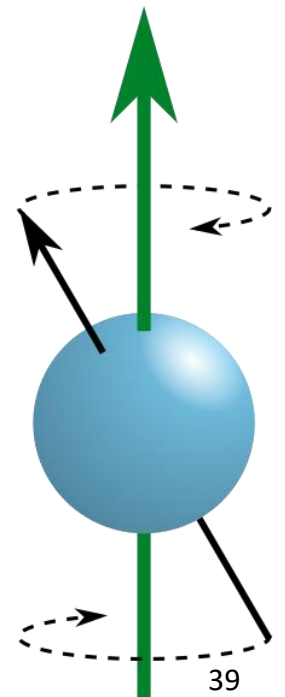
Spin Precession

Axions and ALPS couple to CP odd quantities

$$\mathcal{L}_{\text{ax}} = g_a \bar{\psi} \psi \partial_\mu a \bar{\psi} \gamma^\mu \gamma^5 \psi$$

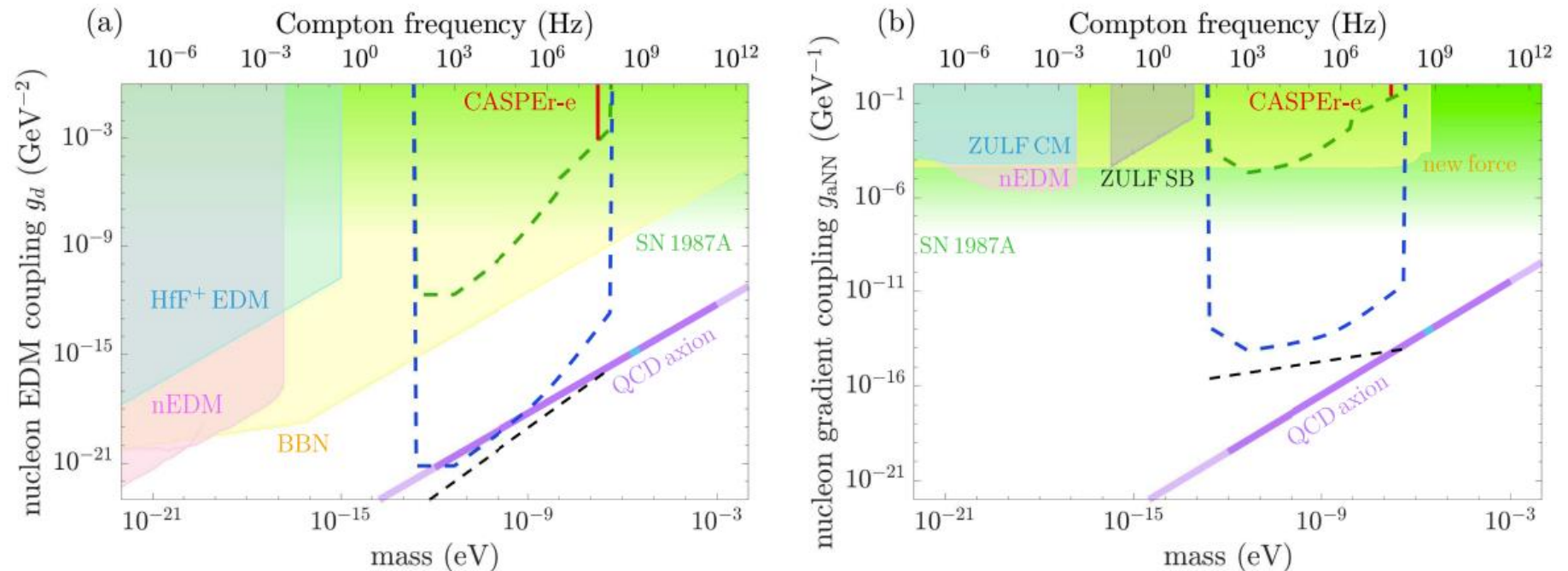
The axion gradient acts like a (pseudo-) magnetic field

Spins will precess around the direction of the axion gradient



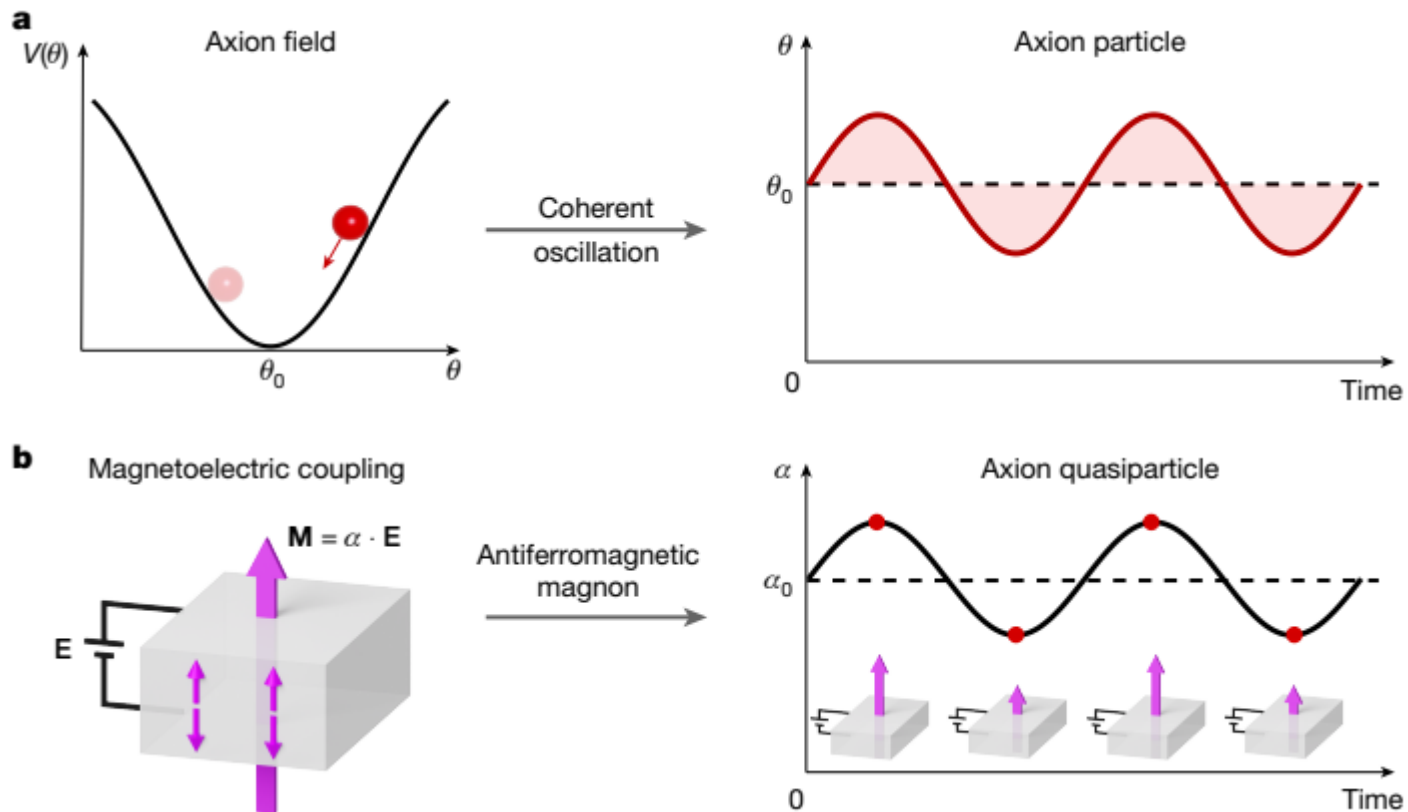
CASPER Constraints

Precision measurement of ^{207}Pb solid state NMR in polarised crystal (4.6 mm) – extension to 80cm sample



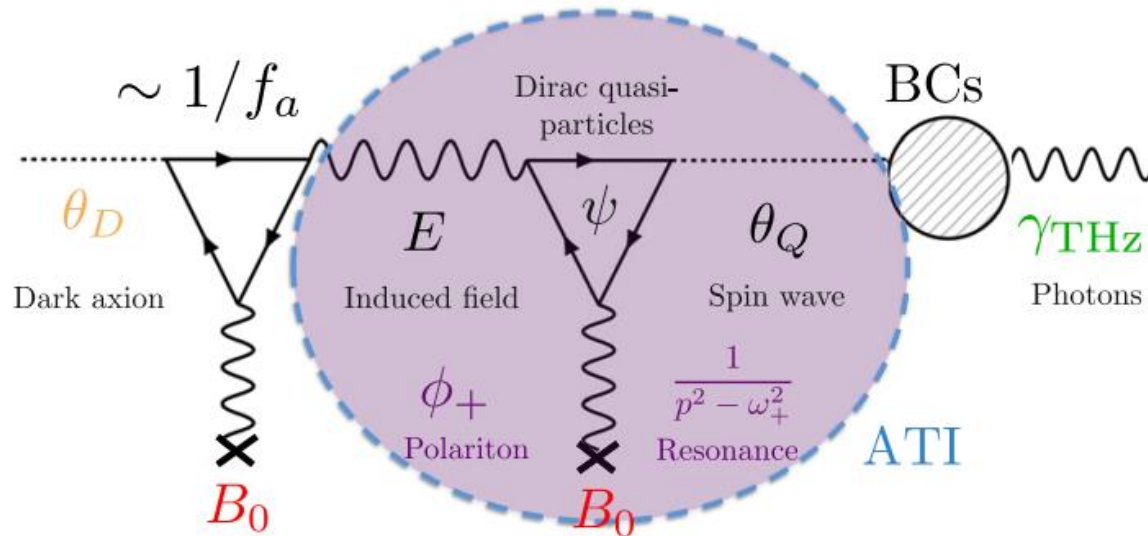
Quasi-particles

Axion quasi-particle observed in MnBi_2Te_4 (2D material)
- magnons in an anti-ferromagnetic material



Quasi-particles

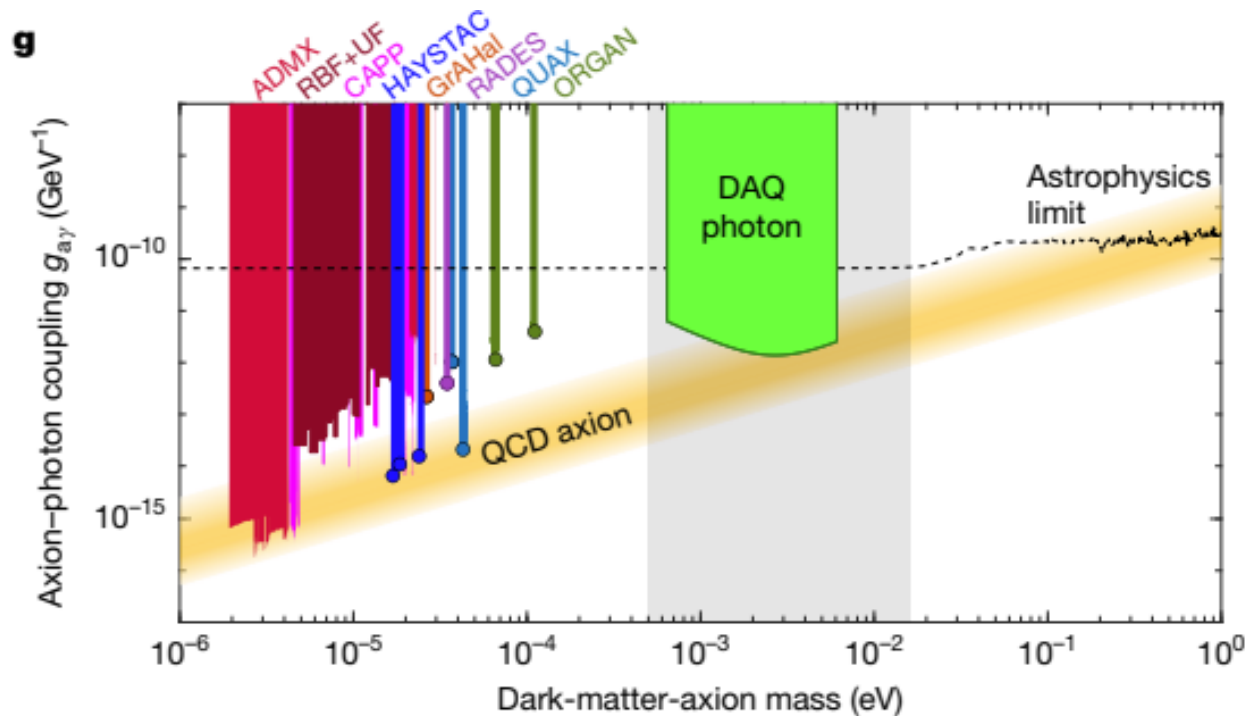
Axion quasi-particle observed in MnBi_2Te_4 (2D material)
Couple to a dark matter axion field

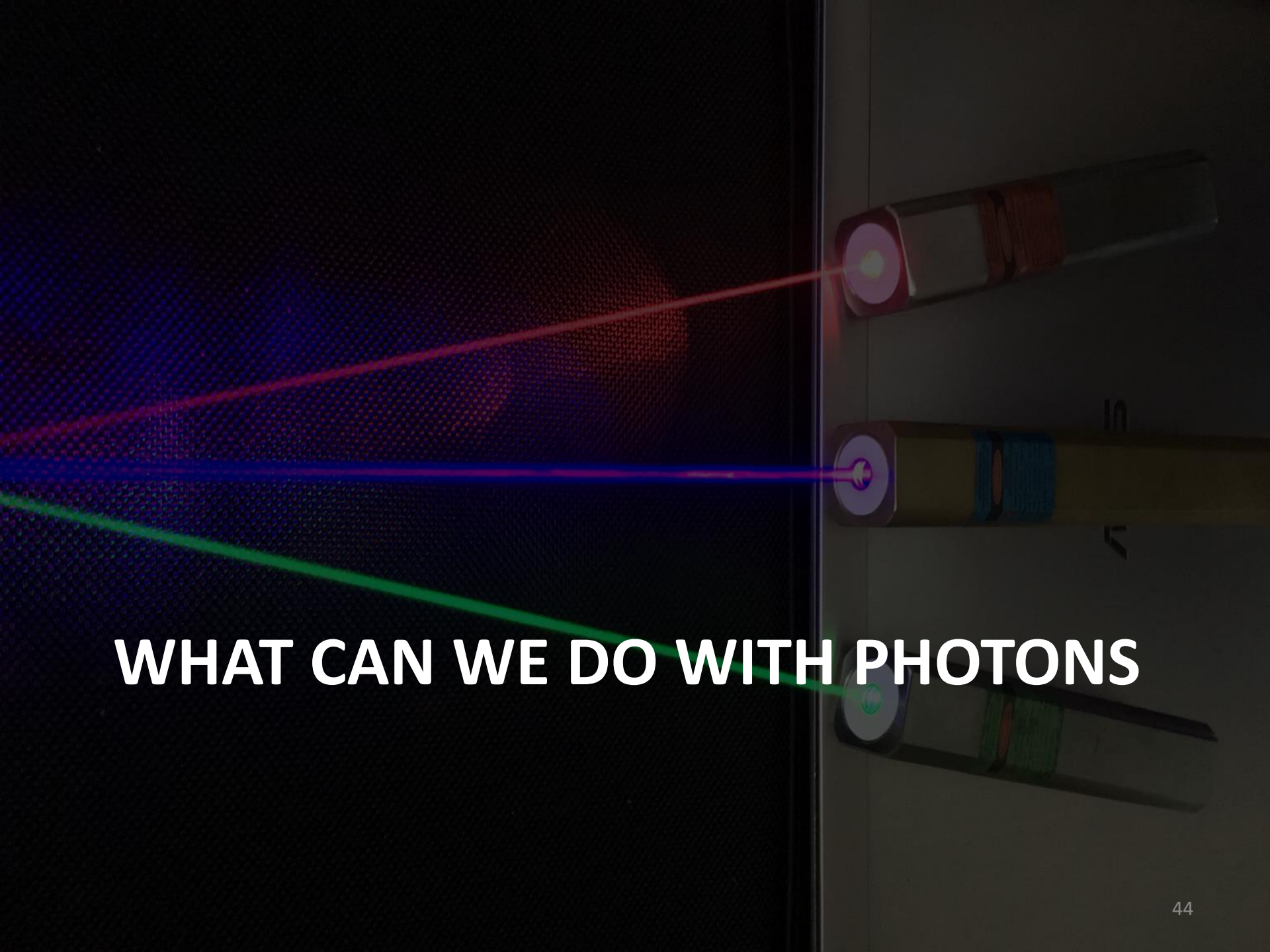


Quasi-particles

Collective excitations in MnBi_2Te_4 (2D material)

Couple to a dark matter axion field



An optical experiment setup is shown. Three fiber optic cables are positioned vertically on the right side of the frame. The top cable is red and emits a red laser beam. The middle cable is yellow and emits a yellow laser beam. The bottom cable is green and emits a green laser beam. The beams are directed towards a dark, textured surface on the left, which appears to be a diffraction grating. The red beam is deflected upwards, the yellow beam is deflected downwards, and the green beam is deflected downwards. The text "WHAT CAN WE DO WITH PHOTONS" is overlaid in white capital letters across the center of the image.

WHAT CAN WE DO WITH PHOTONS

Axions – Pseudo-scalar ULDM

$$\mathcal{L}_{\text{axion}} = \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{DS}} + \frac{a}{4f_\gamma} F_{\mu\nu} \tilde{F}_{\mu\nu} + \frac{a}{4f_G} \text{Tr} G_{\mu\nu} \tilde{G}_{\mu\nu} + \frac{\partial_\mu a}{f_l} \sum_\alpha \bar{l}_\alpha \gamma_\mu \gamma_5 l_\alpha + \frac{\partial_\mu a}{f_q} \sum_\beta \bar{q}_\beta \gamma_\mu \gamma_5 q_\beta$$

Coupling to photons

Coupling to Gluons

Coupling to leptons

Coupling to quarks

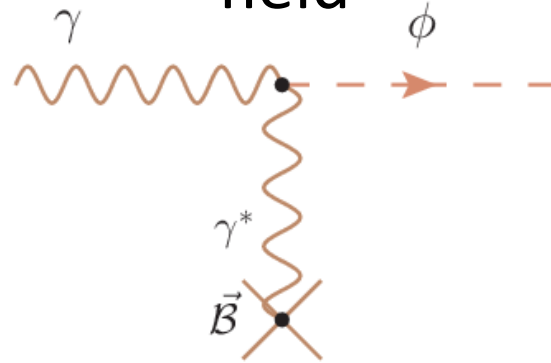
Axion solution to the strong CP problem*

Axion-like particles – pseudo scalars with similar couplings

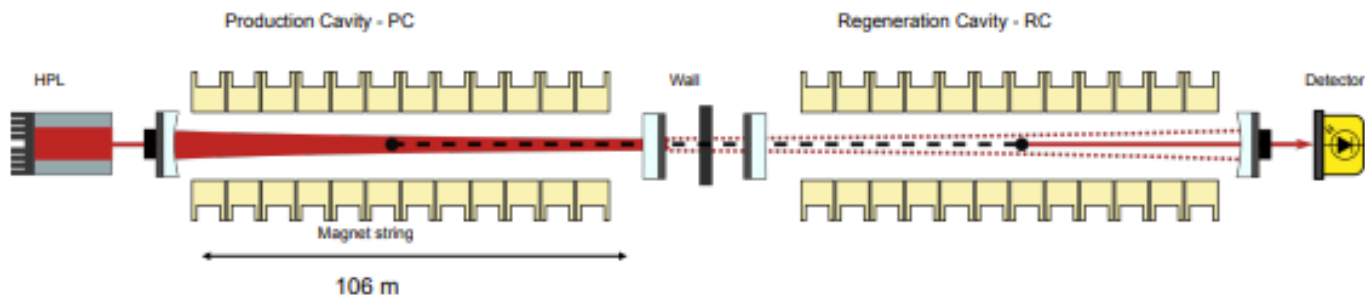
* But see also: Ai, Cruz, Garbrecht, Tamarit. Phys.Lett.B 822 (2021) 136616

Light Shining Through Walls

Axion-photon mixing in the presence of a magnetic field

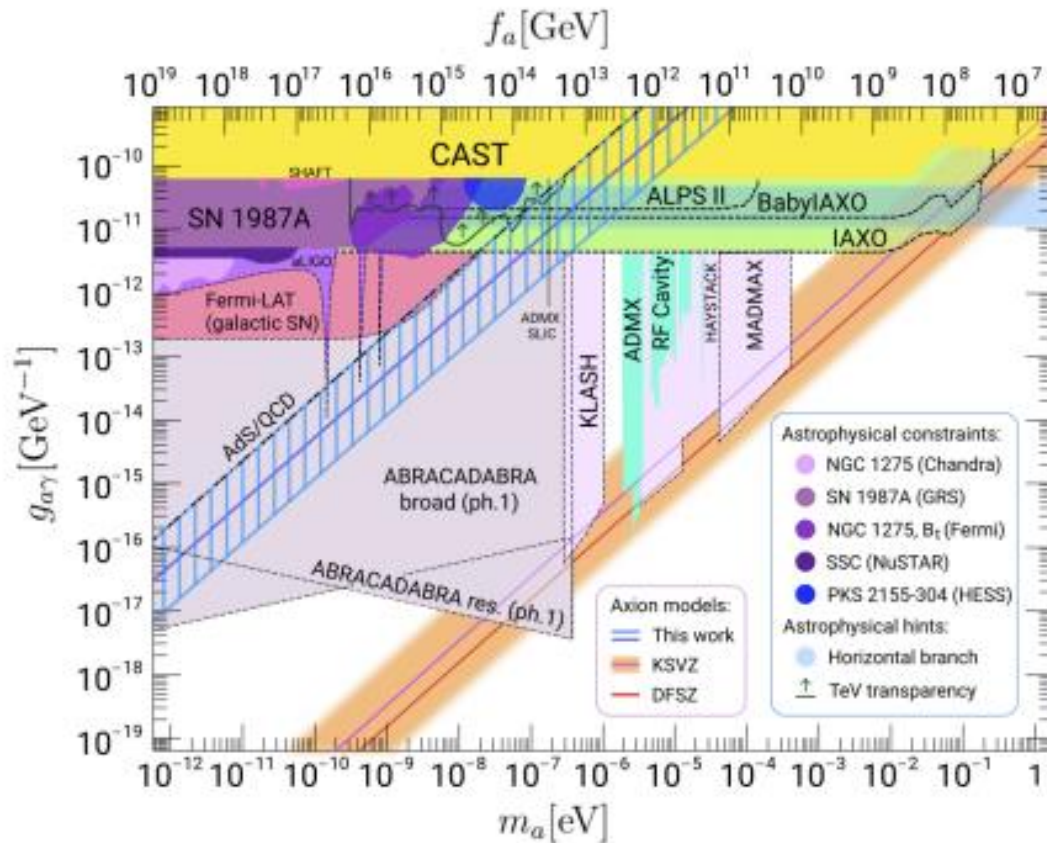


ALPS II experiment at DESY, using recycled HERA magnets



Light Shining Through Walls

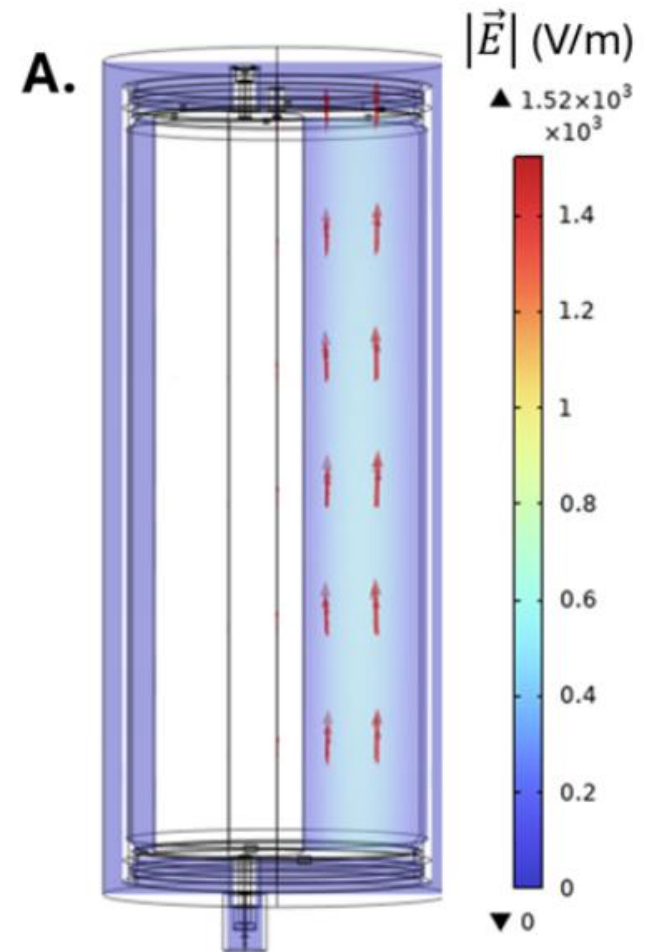
Axion-photon mixing in the presence of a magnetic field



Haloscope

Axions in the dark matter halo,
converted into photons in a
magnetic field

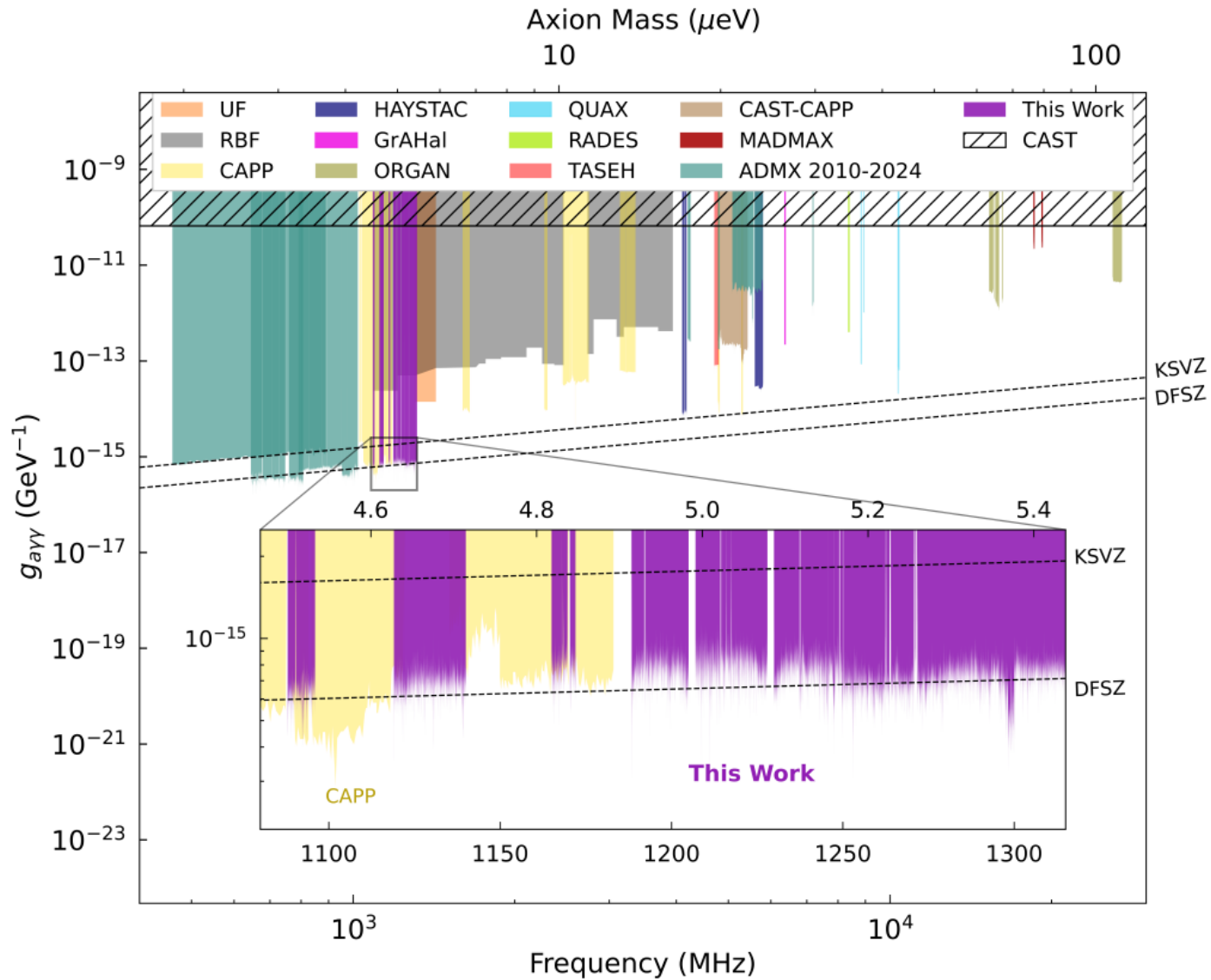
- resonant enhancement at
cavity frequency
- Signal proportional to cavity
volume, magnetic field and
cavity Q-factor
- Noise proportional to black
body radiation, amplifier
noise



Sikivie. PRL 51 16 1415 (1983)

Carosi et al. Phys. Rev. Lett. 135, 191001 (2025)

Haloscope



Gravitational Wave Detection

Coupling of the form $\sim h \mathbf{B} \cdot \mathbf{B}_0$

Inverse Gertsenshtein effect induces effective currents that couple to the EM modes of a cavity



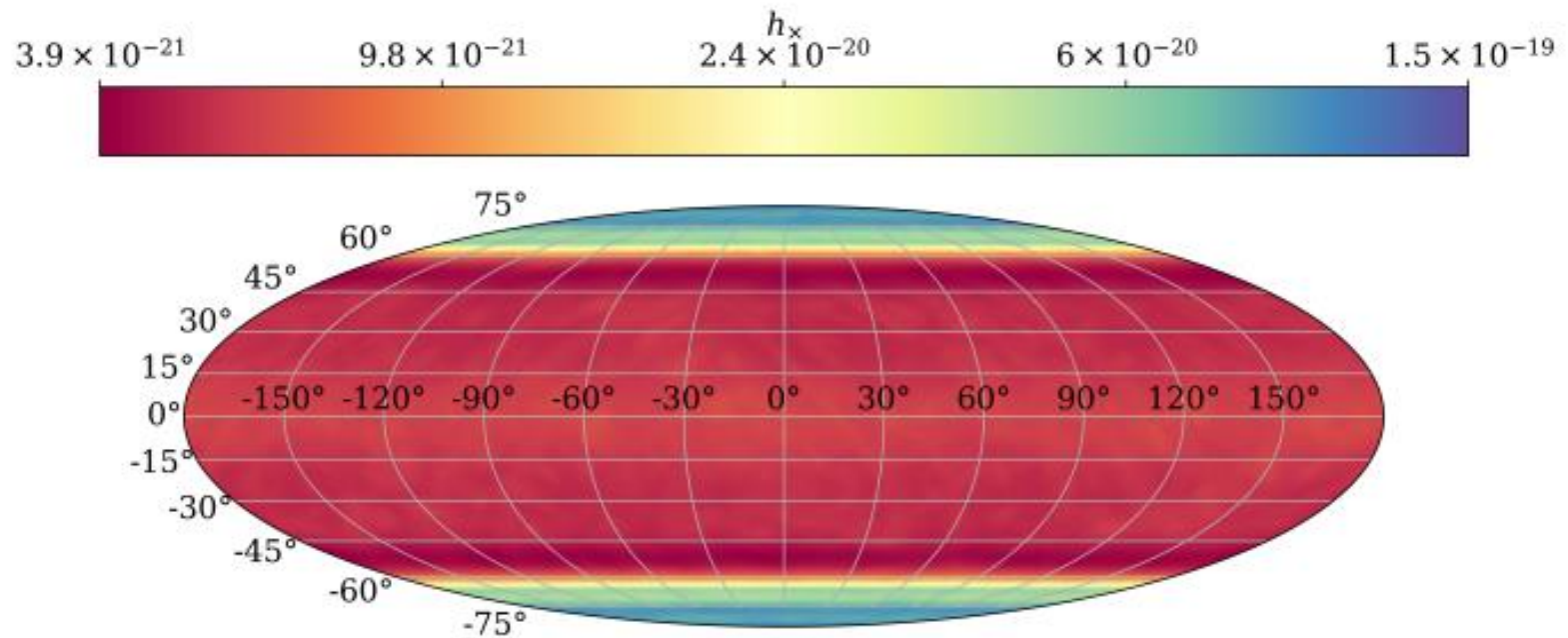
Enables narrow band high-frequency gravitational wave detection

Berlin et al. Phys. Rev. D 105 116011 (2022)

Gertsenshtein. Sov. Phys. JETP 14 84 (1962)

Gravitational Wave Detection

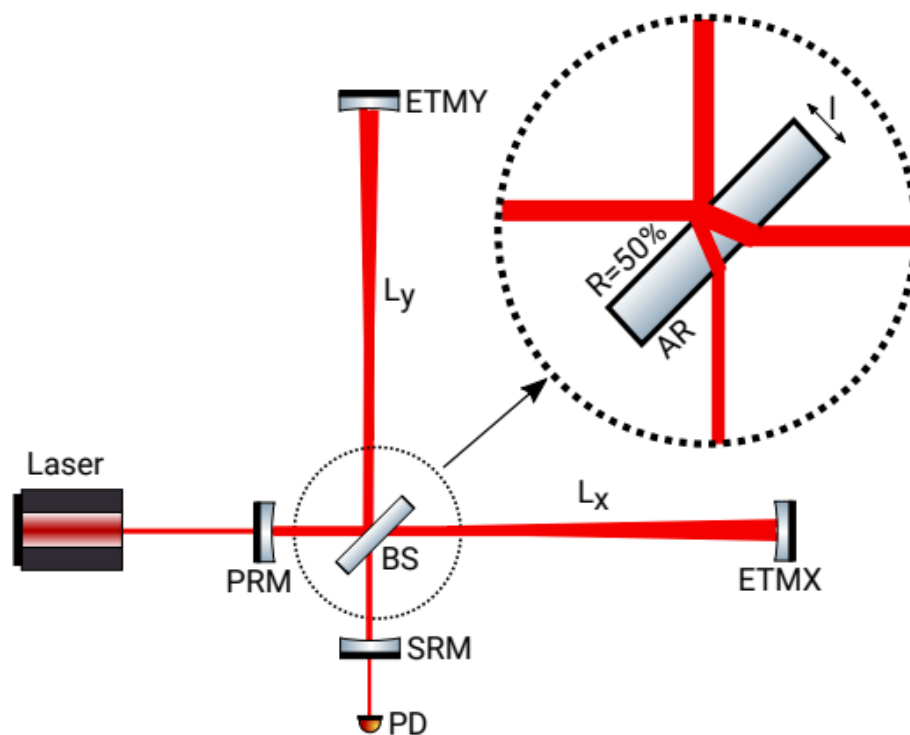
Reanalysis of CAPP-12T (multi-cell) data
for 5.311 GHz GW



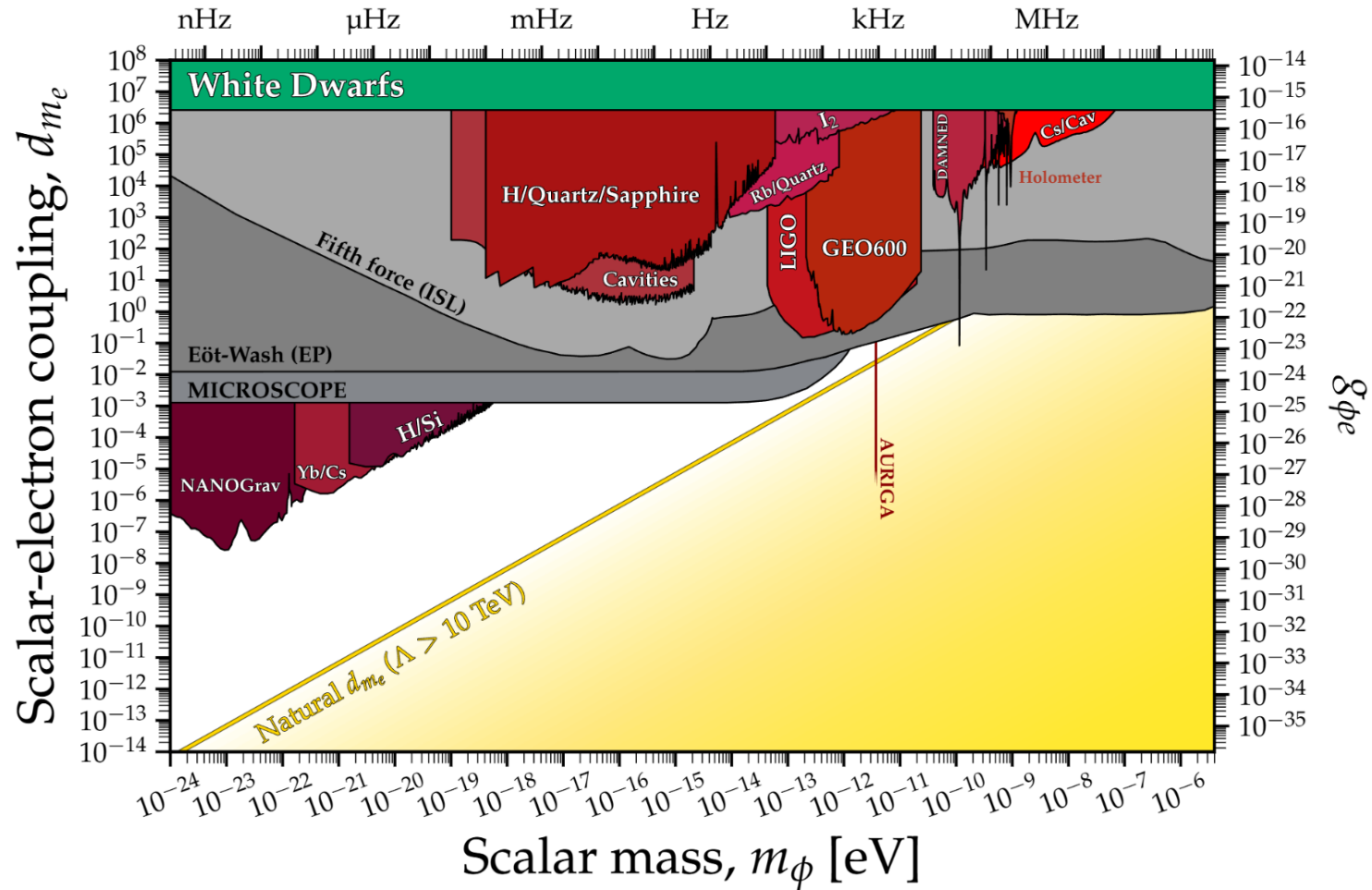
Excludes superradiance from axion clouds around black holes with mass $\sim 10^{-6} M_{\odot}$ within $\sim 10^{-2}$ AU

Michelson Interferometer

Beam splitter and arm mirrors (if freely suspended) produce path length changes if physical constants vary in space / time



Michelson Interferometer



Grote, Stadnik Phys.Rev.Res 1 (2019) 3 033187

<https://cajohare.github.io/AxionLimits/>

Summary

Many opportunities to search for light and feebly coupled new physics

Advances in precision enabled by developing quantum technologies

Careful theoretical modelling needed to match precision of observations

Huge parameter space to explore

- How to best use resources?

