

Issues with jets

Mrinal Dasgupta

University of Manchester

Durham, February 20, 2008

u-logo

ur-logo



- **Jets and IRC safety**
- Non-perturbative effects at hadron colliders.
- **Analytical** studies of hadronisation contribution to jet energy.
- Monte Carlo studies.
- Possible tests and future measurements.

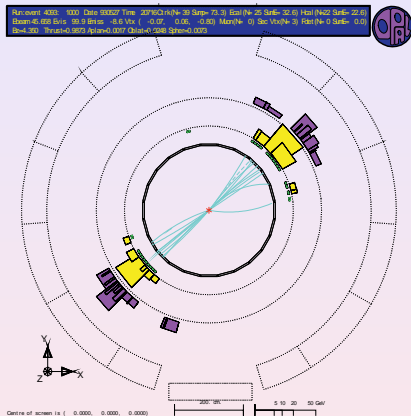
- Jets and IRC safety
- Non-perturbative effects at hadron colliders.
- Analytical studies of hadronisation contribution to jet energy.
- Monte Carlo studies.
- Possible tests and future measurements.

- Jets and IRC safety
- Non-perturbative effects at hadron colliders.
- **Analytical** studies of hadronisation contribution to jet energy.
- Monte Carlo studies.
- Possible tests and future measurements.

- Jets and IRC safety
- Non-perturbative effects at hadron colliders.
- **Analytical** studies of hadronisation contribution to jet energy.
- Monte Carlo studies.
- Possible tests and future measurements.

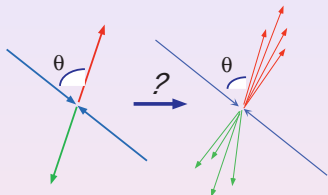
- Jets and IRC safety
- Non-perturbative effects at hadron colliders.
- **Analytical** studies of hadronisation contribution to jet energy.
- Monte Carlo studies.
- Possible tests and future measurements.

Prehistory - Birth of jet algorithms



Experiments observe sprays of hadrons. What does theory say ?

Theory fails unless...

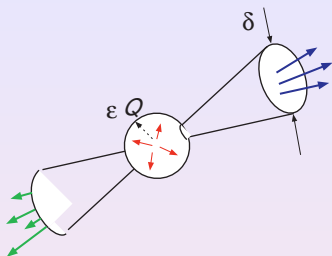


Beyond parton model one obtains divergent results for three and more particle production. Need to introduce energy resolution (like in QED) but also **angular resolution**.

ur-logo

ur-logo

Jet definition cures problem



A jet algorithm is an **agreement**! Given this one can predict jet observables such as rates and distributions :

$$f_2 = 1 - 8C_F \frac{\alpha_S}{2\pi} \left(\ln \frac{1}{\delta} \left[\ln \left(\frac{1}{2\epsilon} - 1 \right) - \frac{3}{4} + 3\epsilon \right] + \frac{\pi^2}{12} - \frac{7}{16} - \epsilon + \frac{3}{2}\epsilon^2 + \mathcal{O}(\delta^2 \ln \epsilon) \right)$$
$$f_3 = 1 - f_2$$

Things are much more sophisticated. Jet algorithms fall into two main categories

- **Sequential recombination** algorithms (k_t , Cambridge/Aachen). Cluster particles repeatedly according to some distance measure with some stopping criterion. IRC safe.
- Cone type algorithms (seeded, seedless). Look for energy flow into limited angular regions given by stable cones. Stable cone is one whose axis is centre of momentum of particles in it. Seeded cones are IRC unsafe. Salam and Soyez 2007

tu-logo

ur-logo

Things are much more sophisticated. Jet algorithms fall into two main categories

- **Sequential recombination** algorithms (k_t , Cambridge/Aachen). Cluster particles repeatedly according to some distance measure with some stopping criterion. IRC safe.
- Cone type algorithms (seeded, seedless). Look for energy flow into limited angular regions given by stable cones. Stable cone is one whose axis is centre of momentum of particles in it. Seeded cones are IRC unsafe. Salam and Soyez 2007

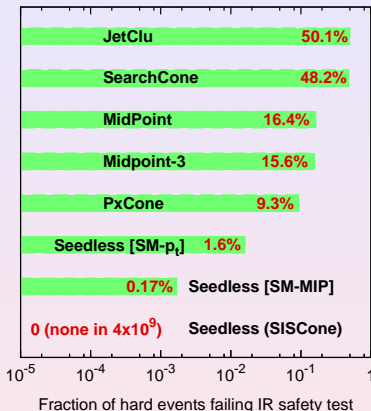
Things are much more sophisticated. Jet algorithms fall into two main categories

- **Sequential recombination** algorithms (k_t , Cambridge/Aachen). Cluster particles repeatedly according to some distance measure with some stopping criterion. IRC safe.
- Cone type algorithms (seeded, seedless). Look for energy flow into limited angular regions given by stable cones. Stable cone is one whose axis is centre of momentum of particles in it. Seeded cones are IRC unsafe. Salam and Soyez 2007

Things are much more sophisticated. Jet algorithms fall into two main categories

- **Sequential recombination** algorithms (k_t , Cambridge/Aachen). Cluster particles repeatedly according to some distance measure with some stopping criterion. IRC safe.
- Cone type algorithms (seeded, seedless). Look for energy flow into limited angular regions given by stable cones. Stable cone is one whose axis is centre of momentum of particles in it. Seeded cones are IRC unsafe. Salam and Soyez 2007

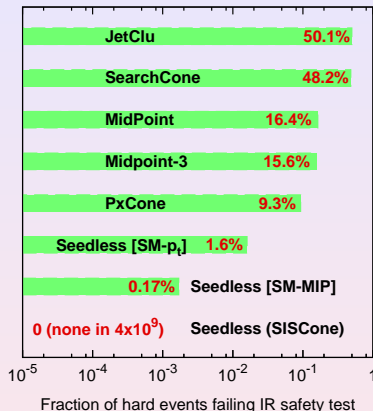
Unsafe algorithms commonly used



Discussion ongoing since 1990s ! Various fixes proposed e.g midpoint. Adopted for Tevatron run 2. However **still IRC unsafe** ! New **fast practical** seedless algorithm on the market **SISCONE**.

Salam and Soyez 2007

Unsafe algorithms commonly used



Discussion ongoing since 1990s ! Various fixes proposed e.g midpoint. Adopted for Tevatron run 2. However **still IRC unsafe** ! New **fast practical** seedless algorithm on the market **SISCONE**.
[Salam and Soyez 2007](#).

Non-perturbative contributions

At hadron colliders two **distinct** NP effects contribute:

- **hadronisation** (ubiquitous, familiar from LEP , HERA.)
- **underlying event** (UE, specific to hadron colliders, messy, large at the LHC)

Significant progress in understanding of hadronisation at LEP, HERA. Can we use it for jets at hadron colliders ?

tu-logo

ur-logo

Non-perturbative contributions

At hadron colliders two **distinct** NP effects contribute:

- **hadronisation** (ubiquitous, familiar from LEP , HERA.)
- **underlying event** (UE, specific to hadron colliders, messy, large at the LHC)

Significant progress in understanding of hadronisation at LEP, HERA. Can we use it for jets at hadron colliders ?

tu-logo

ur-logo

Non-perturbative contributions

At hadron colliders two **distinct** NP effects contribute:

- **hadronisation** (ubiquitous, familiar from LEP , HERA.)
- **underlying event** (UE, specific to hadron colliders, messy, large at the LHC)

Significant progress in understanding of hadronisation at LEP, HERA. Can we use it for jets at hadron colliders ?

tu-logo

ur-logo

Non-perturbative contributions

At hadron colliders two **distinct** NP effects contribute:

- **hadronisation** (ubiquitous, familiar from LEP , HERA.)
- **underlying event** (UE, specific to hadron colliders, messy, large at the LHC)

Significant progress in understanding of hadronisation at LEP, HERA. Can we use it for jets at hadron colliders ?

tu-logo

ur-logo

Hadronisation from LEP and HERA

Learnt a fair amount from studies of event shapes. Related to flow of energy-momentum in final state –linear in momenta.

The summary of LEP and HERA experience : Event shapes receive large hadronisation corrections $\propto 1/Q$. Associated to hadronisation. **BUT** possible to **analytically** estimate size of these effects. Experimentally popular is DW approach . Extends perturbation theory into infrared with **infrared finite universal coupling**.

tu-logo

ur-logo

Hadronisation from LEP and HERA

Learnt a fair amount from studies of event shapes. Related to flow of energy-momentum in final state –linear in momenta. The summary of LEP and HERA experience : Event shapes receive large hadronisation corrections $\propto 1/Q$. Associated to hadronisation. BUT possible to *analytically* estimate size of these effects. Experimentally popular is DW approach . Extends perturbation theory into infrared with *infrared finite universal coupling*.

tu-logo

ur-logo

Hadronisation from LEP and HERA

Learnt a fair amount from studies of event shapes. Related to flow of energy-momentum in final state –linear in momenta. The summary of LEP and HERA experience : Event shapes receive large hadronisation corrections $\propto 1/Q$. Associated to hadronisation. **BUT** possible to **analytically** estimate size of these effects. Experimentally popular is DW approach . Extends perturbation theory into infrared with **infrared finite universal coupling**.

tu-logo

ur-logo

Hadronisation from LEP and HERA

Learnt a fair amount from studies of event shapes. Related to flow of energy-momentum in final state –linear in momenta. The summary of LEP and HERA experience : Event shapes receive large hadronisation corrections $\propto 1/Q$. Associated to hadronisation. **BUT** possible to **analytically** estimate size of these effects. Experimentally popular is DW approach . Extends perturbation theory into infrared with **infrared finite universal coupling**.

tu-logo

ur-logo

Hadronisation from LEP and HERA

Power corrections linked to soft gluon emission with $k_t \sim \Lambda_{\text{QCD}}$.
Soft gluons modify event shapes as

$$\delta V^{\text{NP}} = \frac{k_t}{Q} f_V(\eta)$$

Change to event shape mean values is

$$\int_0^{\mu_I} \frac{dk_t}{k_t} \delta\alpha_s(k_t) \frac{k_t}{Q} d\eta f_V(\eta) = A(\mu_I) \frac{C_V}{Q}$$

Here one has the **universal** quantity

$$A(\mu_I) = \int_0^{\mu_I} \frac{dk_t}{k_t} k_t \delta\alpha_s(k_t)$$

Fit different event shapes at LEP and HERA...try to check universality.

Hadronisation from LEP and HERA

Power corrections linked to soft gluon emission with $k_t \sim \Lambda_{\text{QCD}}$.
Soft gluons modify event shapes as

$$\delta V^{\text{NP}} = \frac{k_t}{Q} f_V(\eta)$$

Change to event shape mean values is

$$\int_0^{\mu_I} \frac{dk_t}{k_t} \delta\alpha_s(k_t) \frac{k_t}{Q} d\eta f_V(\eta) = A(\mu_I) \frac{C_V}{Q}$$

Here one has the **universal** quantity

$$A(\mu_I) = \int_0^{\mu_I} \frac{dk_t}{k_t} k_t \delta\alpha_s(k_t)$$

Fit different event shapes at LEP and HERA...try to check universality.

Hadronisation from LEP and HERA

Power corrections linked to soft gluon emission with $k_t \sim \Lambda_{\text{QCD}}$.
Soft gluons modify event shapes as

$$\delta V^{\text{NP}} = \frac{k_t}{Q} f_V(\eta)$$

Change to event shape mean values is

$$\int_0^{\mu_I} \frac{dk_t}{k_t} \delta\alpha_s(k_t) \frac{k_t}{Q} d\eta f_V(\eta) = A(\mu_I) \frac{C_V}{Q}$$

Here one has the **universal** quantity

$$A(\mu_I) = \int_0^{\mu_I} \frac{dk_t}{k_t} k_t \delta\alpha_s(k_t)$$

Fit different event shapes at LEP and HERA...try to check universality.

Hadronisation from LEP and HERA

Power corrections linked to soft gluon emission with $k_t \sim \Lambda_{\text{QCD}}$.
Soft gluons modify event shapes as

$$\delta V^{\text{NP}} = \frac{k_t}{Q} f_V(\eta)$$

Change to event shape mean values is

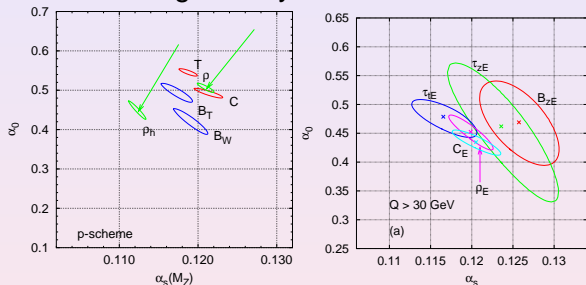
$$\int_0^{\mu_I} \frac{dk_t}{k_t} \delta\alpha_s(k_t) \frac{k_t}{Q} d\eta f_V(\eta) = A(\mu_I) \frac{C_V}{Q}$$

Here one has the **universal** quantity

$$A(\mu_I) = \int_0^{\mu_I} \frac{dk_t}{k_t} k_t \delta\alpha_s(k_t)$$

Fit different event shapes at LEP and HERA...try to check universality.

Observed to generally work well at LEP and HERA



Can we take over to hadron collider jets ?

Why bother when we have MC ?

- MC does not reflect understanding of physics of hadronisation. Analytical studies do.
- MC studies do not provide any detailed parametric understanding of NP effects or insight on how to separate hadronisation from UE.
- Lack of parametric understanding leads to myths: e.g. cone jets said to suffer from large hadronisation while k_t jets from UE. But $R_{\text{cone}} = 0.4$ whereas $R_{k_t} = 1$.
- MC hadronisation taken from difference between hadron level and parton shower, then added to NLO sometimes without cross-checks.

Analytical insight **sorely needed** !

ur-logo

ur-logo

Why bother when we have MC ?

- MC does not reflect understanding of physics of hadronisation. Analytical studies do.
- MC studies do not provide any detailed parametric understanding of NP effects or insight on how to separate hadronisation from UE.
- Lack of parametric understanding leads to myths: e.g. cone jets said to suffer from large hadronisation while k_t jets from UE. But $R_{\text{cone}} = 0.4$ whereas $R_{k_t} = 1$.
- MC hadronisation taken from difference between hadron level and parton shower, then added to NLO sometimes without cross-checks.

Analytical insight sorely needed !

ur-logo

ur-logo

Why bother when we have MC ?

- MC does not reflect understanding of physics of hadronisation. Analytical studies do.
- MC studies do not provide any detailed parametric understanding of NP effects or insight on how to separate hadronisation from UE.
- Lack of parametric understanding leads to myths: e.g. cone jets said to suffer from large hadronisation while k_t jets from UE. But $R_{\text{cone}} = 0.4$ whereas $R_{k_t} = 1$.
- MC hadronisation taken from difference between hadron level and parton shower, then added to NLO sometimes without cross-checks.

Analytical insight sorely needed !

ur-logo

ur-logo

Why bother when we have MC ?

- MC does not reflect understanding of physics of hadronisation. Analytical studies do.
- MC studies do not provide any detailed parametric understanding of NP effects or insight on how to separate hadronisation from UE.
- Lack of parametric understanding leads to myths: e.g. cone jets said to suffer from large hadronisation while k_t jets from UE. But $R_{\text{cone}} = 0.4$ whereas $R_{k_t} = 1$.
- MC hadronisation taken from difference between hadron level and parton shower, then added to NLO sometimes without cross-checks.

Analytical insight **sorely needed** !

ur-logo

ur-logo

Why bother when we have MC ?

- MC does not reflect understanding of physics of hadronisation. Analytical studies do.
- MC studies do not provide any detailed parametric understanding of NP effects or insight on how to separate hadronisation from UE.
- Lack of parametric understanding leads to myths: e.g. cone jets said to suffer from large hadronisation while k_t jets from UE. But $R_{\text{cone}} = 0.4$ whereas $R_{k_t} = 1$.
- MC hadronisation taken from difference between hadron level and parton shower, then added to NLO sometimes without cross-checks.

Analytical insight **sorely needed** !

ur-logo

ur-logo

Why bother when we have MC ?

- MC does not reflect understanding of physics of hadronisation. Analytical studies do.
- MC studies do not provide any detailed parametric understanding of NP effects or insight on how to separate hadronisation from UE.
- Lack of parametric understanding leads to myths: e.g. cone jets said to suffer from large hadronisation while k_t jets from UE. But $R_{\text{cone}} = 0.4$ whereas $R_{k_t} = 1$.
- MC hadronisation taken from difference between hadron level and parton shower, then added to NLO sometimes without cross-checks.

Analytical insight **sorely needed** !

ur-logo

ur-logo

Jet transverse energy scale

Important ingredient of experimental jet studies.

Let's focus on this. Use LEP and HERA methods + universality to try to estimate hadronisation. Consider effects of soft emission and relate to IR universal coupling.

tu-logo

ur-logo



Jet transverse energy scale

Important ingredient of experimental jet studies.

Let's focus on this. Use LEP and HERA methods + universality to try to estimate hadronisation. Consider effects of soft emission and relate to IR universal coupling.

tu-logo

ur-logo

Jet transverse energy scale

Important ingredient of experimental jet studies.

Let's focus on this. Use LEP and HERA methods + universality to try to estimate hadronisation. Consider effects of soft emission and relate to IR universal coupling.

tu-logo

ur-logo

Jet transverse energy scale

Important ingredient of experimental jet studies.
Let's focus on this. Use LEP and HERA methods + universality to try to estimate hadronisation. Consider effects of soft emission and relate to IR universal coupling.

tu-logo

ur-logo

One gluon calculation

Work with dijets near threshold. At Born level

$$p_1 = \frac{\sqrt{s}}{2}(1, 0, 0, 1)$$

$$p_2 = \frac{\sqrt{s}}{2}(1, 0, 0, -1)$$

$$p_j = p_t(1, 1, 0, 0)$$

$$p_r = p_t(1, -1, 0, 0)$$

$$p_t = \frac{\sqrt{s}}{2}$$

Then we need one gluon correction. Two possibilities:

- Gluon recombined into jet.
- Gluon outside jet.

At one gluon level all algorithms work the same.

tu-logic

ur-logic

One gluon calculation

Work with dijets near threshold. At Born level

$$p_1 = \frac{\sqrt{s}}{2}(1, 0, 0, 1)$$

$$p_2 = \frac{\sqrt{s}}{2}(1, 0, 0, -1)$$

$$p_j = p_t(1, 1, 0, 0)$$

$$p_r = p_t(1, -1, 0, 0)$$

$$p_t = \frac{\sqrt{s}}{2}$$

Then we need one gluon correction. Two possibilities:

- Gluon recombined into jet.
- Gluon outside jet.

At one gluon level all algorithms work the same.

tu-logic

ur-logic

One gluon calculation

Work with dijets near threshold. At Born level

$$p_1 = \frac{\sqrt{s}}{2}(1, 0, 0, 1)$$

$$p_2 = \frac{\sqrt{s}}{2}(1, 0, 0, -1)$$

$$p_j = p_t(1, 1, 0, 0)$$

$$p_r = p_t(1, -1, 0, 0)$$

$$p_t = \frac{\sqrt{s}}{2}$$

Then we need one gluon correction. Two possibilities:

- Gluon recombined into jet.
- Gluon outside jet.

At one gluon level all algorithms work the same.

tu-logic

ur-logic

One gluon calculation

Work with dijets near threshold. At Born level

$$p_1 = \frac{\sqrt{s}}{2}(1, 0, 0, 1)$$

$$p_2 = \frac{\sqrt{s}}{2}(1, 0, 0, -1)$$

$$p_j = p_t(1, 1, 0, 0)$$

$$p_r = p_t(1, -1, 0, 0)$$

$$p_t = \frac{\sqrt{s}}{2}$$

Then we need one gluon correction. Two possibilities:

- Gluon recombined into jet.
- Gluon outside jet.

At one gluon level all algorithms work the same.

tu-logo

ur-logo

Soft gluon contribution

Gluon + initiating parton = massive jet

$$p_1 = \frac{\sqrt{s}}{2}(1, 0, 0, 1)$$

$$p_2 = \frac{\sqrt{s}}{2}(1, 0, 0, -1)$$

$$p_j = (\sqrt{p_t^2 + M_j^2}, p_t, 0, 0)$$

$$p_r = (p_t, -p_t, 0, 0)$$

This leads to

$$p_t = \frac{\sqrt{s}}{2} \left(1 - \frac{M_j^2}{s} \right) \rightarrow \delta p_t^+ = \frac{-M_j^2}{2\sqrt{s}} = \frac{-p_j \cdot k}{\sqrt{s}}$$

Similarly when gluon is not recombined, recoil system is massive :

$$\delta p_t^- = \frac{-M_r^2}{2\sqrt{s}} = \frac{-p_r \cdot k}{\sqrt{s}}$$

Soft gluon contribution

Gluon + initiating parton = massive jet

$$p_1 = \frac{\sqrt{s}}{2}(1, 0, 0, 1)$$

$$p_2 = \frac{\sqrt{s}}{2}(1, 0, 0, -1)$$

$$p_j = (\sqrt{p_t^2 + M_j^2}, p_t, 0, 0)$$

$$p_r = (p_t, -p_t, 0, 0)$$

This leads to

$$p_t = \frac{\sqrt{s}}{2} \left(1 - \frac{M_j^2}{s} \right) \rightarrow \delta p_t^+ = \frac{-M_j^2}{2\sqrt{s}} = \frac{-p_j \cdot k}{\sqrt{s}}$$

Similarly when gluon is not recombined, recoil system is massive :

$$\delta p_t^- = \frac{-M_r^2}{2\sqrt{s}} = \frac{-p_r \cdot k}{\sqrt{s}}$$

Integrating the one gluon contribution

2 \rightarrow 2 process means several dipoles are involved in the antenna pattern for emission of soft gluon. Several copies of $e^+ e^-$. One has

$$\langle \delta p_t \rangle = \sum_{ij} C_{ij} I_{ij}$$

$$I_{ij} = I_{ij}^+ + I_{ij}^-$$

$$I^\pm(R) \equiv \int_{\pm} d\eta \frac{d\phi}{2\pi} d\kappa_T^{(ij)} \delta\alpha_s \left(\kappa_T^{(ij)} \right) k_T \left| \frac{\partial k_T}{\partial \kappa_T^{(ij)}} \right| \frac{p_i \cdot p_j}{p_i \cdot k p_j \cdot k} \delta p_t^\pm,$$

with

$$\left(\kappa_T^{(ij)} \right)^2 = \frac{2 p_i \cdot k p_j \cdot k}{p_i \cdot p_j},$$

Integrating the one gluon contribution

2 \rightarrow 2 process means several dipoles are involved in the antenna pattern for emission of soft gluon. Several copies of $e^+ e^-$. One has

$$\langle \delta p_t \rangle = \sum_{ij} C_{ij} I_{ij}$$

$$I_{ij} = I_{ij}^+ + I_{ij}^-$$

$$I^\pm(R) \equiv \int_{\pm} d\eta \frac{d\phi}{2\pi} d\kappa_T^{(ij)} \delta\alpha_s \left(\kappa_T^{(ij)} \right) k_T \left| \frac{\partial k_T}{\partial \kappa_T^{(ij)}} \right| \frac{p_i \cdot p_j}{p_i \cdot k p_j \cdot k} \delta p_t^\pm,$$

with

$$\left(\kappa_T^{(ij)} \right)^2 = \frac{2 p_i \cdot k p_j \cdot k}{p_i \cdot p_j},$$

Integrating the one gluon contribution

2 \rightarrow 2 process means several dipoles are involved in the antenna pattern for emission of soft gluon. Several copies of $e^+ e^-$. One has

$$\langle \delta p_t \rangle = \sum_{ij} C_{ij} I_{ij}$$

$$I_{ij} = I_{ij}^+ + I_{ij}^-$$

$$I^\pm(R) \equiv \int_{\pm} d\eta \frac{d\phi}{2\pi} d\kappa_T^{(ij)} \delta\alpha_s(\kappa_T^{(ij)}) k_T \left| \frac{\partial k_T}{\partial \kappa_T^{(ij)}} \right| \frac{p_i \cdot p_j}{p_i \cdot k p_j \cdot k} \delta p_t^\pm,$$

with

$$\left(\kappa_T^{(ij)} \right)^2 = \frac{2 p_i \cdot k p_j \cdot k}{p_i \cdot p_j},$$

Different dipoles have different qualitative behaviour :

- In-In dipole

$$\begin{aligned} I_{12}(R) &= 2 \mathcal{A}(\mu_I) \int_{\eta^2 + \phi^2 < R^2} d\eta \frac{d\phi}{2\pi} \cos \phi \\ &= 2 \mathcal{A}(\mu_I) R J_1(R) \\ &= 2 \mathcal{A}(\mu_I) \left(\frac{1}{2} R^2 - \frac{1}{16} R^4 + \frac{1}{384} R^6 + \mathcal{O}(R^8) \right). \end{aligned}$$

- In-Recoil dipole

$$I_{1R}(R) = -\mathcal{A}(\mu_I) \left(\frac{1}{8} R^2 - \frac{9}{512} R^4 + \frac{73}{24576} R^6 + \mathcal{O}(R^8) \right).$$

Different dipoles have different qualitative behaviour :

- In-In dipole

$$\begin{aligned}I_{12}(R) &= 2 \mathcal{A}(\mu_I) \int_{\eta^2 + \phi^2 < R^2} d\eta \frac{d\phi}{2\pi} \cos \phi \\ &= 2 \mathcal{A}(\mu_I) R J_1(R) \\ &= 2 \mathcal{A}(\mu_I) \left(\frac{1}{2} R^2 - \frac{1}{16} R^4 + \frac{1}{384} R^6 + \mathcal{O}(R^8) \right) .\end{aligned}$$

- In-Recoil dipole

$$I_{1R}(R) = -\mathcal{A}(\mu_I) \left(\frac{1}{8} R^2 - \frac{9}{512} R^4 + \frac{73}{24576} R^6 + \mathcal{O}(R^8) \right) .$$

Singular R dependence

In dipoles involving the triggered jet - collinear enhancement of integrand reflected in a linear power correction at jet boundary.

- Incoming-Jet dipole

$$I_{1J}(R) = \mathcal{A}(\mu_I) \left(-\frac{2}{R} + \frac{5}{8}R - \frac{23}{1536}R^3 - \frac{95}{73728}R^5 + \mathcal{O}(R^7) \right),$$

- Recoil-Jet dipole

$$I_{RJ} = -\mathcal{A}(\mu_I) \left(\frac{2}{R} + \frac{1}{4}R - \frac{1}{192}R^3 + \frac{5}{2304}R^5 + \mathcal{O}(R^7) \right).$$

Singular R dependence

In dipoles involving the triggered jet - collinear enhancement of integrand reflected in a linear power correction at jet boundary.

- Incoming-Jet dipole

$$I_{1J}(R) = \mathcal{A}(\mu_I) \left(-\frac{2}{R} + \frac{5}{8}R - \frac{23}{1536}R^3 - \frac{95}{73728}R^5 + \mathcal{O}(R^7) \right),$$

- Recoil-Jet dipole

$$I_{RJ} = -\mathcal{A}(\mu_I) \left(\frac{2}{R} + \frac{1}{4}R - \frac{1}{192}R^3 + \frac{5}{2304}R^5 + \mathcal{O}(R^7) \right).$$

Parametric separation from underlying event

Our results are

$$\langle \delta p_t \rangle^h = -C_i \frac{2}{R} A(\mu_I) + \mathcal{O}(R)$$

Value for $2C_F A(\mu_I) \approx 0.5$ GeV from e^+e^- event shapes.

Testable prediction (more cleanly at HERA).

$$\langle \delta p_t \rangle^{\text{UE}} = \frac{\Lambda_{\text{UE}}}{2} R J_1(R) = \frac{\Lambda_{\text{UE}}}{2} \left(R^2 - \frac{R^4}{8} + \mathcal{O}(R^6) \right)$$

No handle on Λ_{UE} except MC studies.

ur-logo

ur-logo

Parametric separation from underlying event

Our results are

$$\langle \delta p_t \rangle^h = -C_i \frac{2}{R} A(\mu_I) + \mathcal{O}(R)$$

Value for $2C_F A(\mu_I) \approx 0.5$ GeV from e^+e^- event shapes.

Testable prediction (more cleanly at HERA).

$$\langle \delta p_t \rangle^{\text{UE}} = \frac{\Lambda_{\text{UE}}}{2} R J_1(R) = \frac{\Lambda_{\text{UE}}}{2} \left(R^2 - \frac{R^4}{8} + \mathcal{O}(R^6) \right)$$

No handle on Λ_{UE} except MC studies.

ur-logo

ur-logo

Parametric separation from underlying event

Our results are

$$\langle \delta p_t \rangle^h = -C_i \frac{2}{R} A(\mu_I) + \mathcal{O}(R)$$

Value for $2C_F A(\mu_I) \approx 0.5$ GeV from e^+e^- event shapes.

Testable prediction (more cleanly at HERA).

$$\langle \delta p_t \rangle^{\text{UE}} = \frac{\Lambda_{\text{UE}}}{2} R J_1(R) = \frac{\Lambda_{\text{UE}}}{2} \left(R^2 - \frac{R^4}{8} + \mathcal{O}(R^6) \right)$$

No handle on Λ_{UE} except MC studies.

ur-logo

ur-logo

Parametric separation from underlying event

Our results are

$$\langle \delta p_t \rangle^h = -C_i \frac{2}{R} A(\mu_I) + \mathcal{O}(R)$$

Value for $2C_F A(\mu_I) \approx 0.5$ GeV from e^+e^- event shapes.

Testable prediction (more cleanly at HERA).

$$\langle \delta p_t \rangle^{\text{UE}} = \frac{\Lambda_{\text{UE}}}{2} R J_1(R) = \frac{\Lambda_{\text{UE}}}{2} \left(R^2 - \frac{R^4}{8} + \mathcal{O}(R^6) \right)$$

No handle on Λ_{UE} except MC studies.

ur-logo

ur-logo

Parametric separation from underlying event

Our results are

$$\langle \delta p_t \rangle^h = -C_i \frac{2}{R} A(\mu_I) + \mathcal{O}(R)$$

Value for $2C_F A(\mu_I) \approx 0.5$ GeV from e^+e^- event shapes.

Testable prediction (more cleanly at HERA).

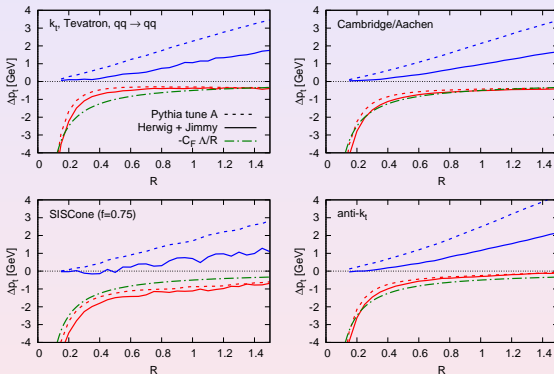
$$\langle \delta p_t \rangle^{\text{UE}} = \frac{\Lambda_{\text{UE}}}{2} R J_1(R) = \frac{\Lambda_{\text{UE}}}{2} \left(R^2 - \frac{R^4}{8} + \mathcal{O}(R^6) \right)$$

No handle on Λ_{UE} except MC studies.

ur-logo

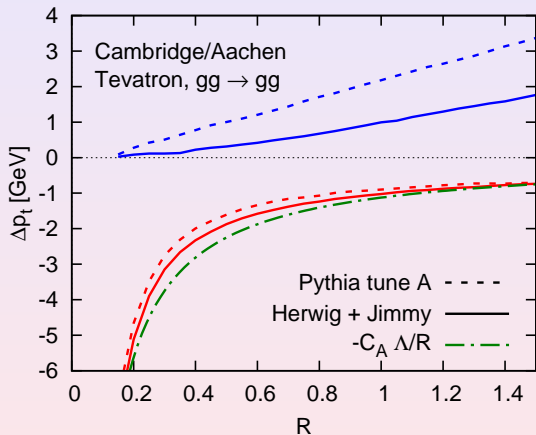
ur-logo

Comparisons to Monte Carlo

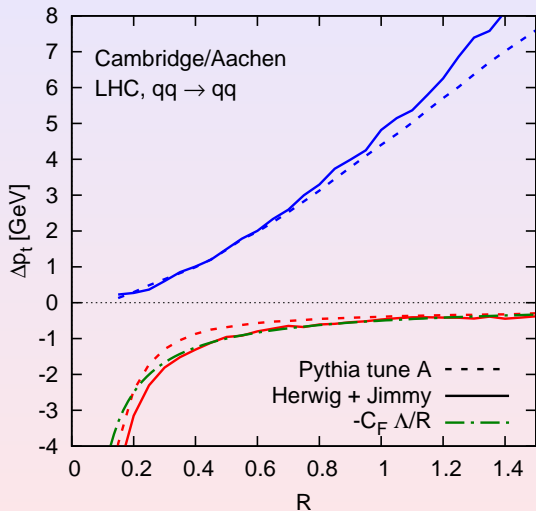


Similar behaviour for all algorithms. Differences in UE between MC's.

Comparisons to Monte Carlo models (contd.)



Comparisons to Monte Carlo models (contd.)



LHC underlying event is enormous effect.

Summary of findings

- Different algorithms show a **similar sensitivity** to NP effects (contradicts folklore)
- UE depends on collider energy and MC model as well as R .
- Hadronization on jet “colour factor” and differently on R .
- $\Lambda_{\text{UE}}(1.96\text{TeV}) \approx 2 - 4 \text{ GeV}$ and $\Lambda_{\text{UE}}(14\text{TeV}) \approx 10 \text{ GeV}$.
Large scale at LHC order of magnitude bigger than hadronisation.
- More info in **variable R analytical studies** than fixed R MC analysis.

tu-logo

ur-logo

Summary of findings

- Different algorithms show a **similar sensitivity** to NP effects (contradicts folklore)
- UE depends on collider energy and MC model as well as R .
- Hadronization on jet “colour factor” and differently on R .
- $\Lambda_{\text{UE}}(1.96\text{TeV}) \approx 2 - 4 \text{ GeV}$ and $\Lambda_{\text{UE}}(14\text{TeV}) \approx 10 \text{ GeV}$.
Large scale at LHC order of magnitude bigger than hadronisation.
- More info in **variable R analytical studies** than fixed R MC analysis.

tu-logo

ur-logo

Summary of findings

- Different algorithms show a **similar sensitivity** to NP effects (contradicts folklore)
- UE depends on collider energy and MC model as well as R .
- Hadronization on jet “colour factor” and differently on R .
- $\Lambda_{\text{UE}}(1.96\text{TeV}) \approx 2 - 4 \text{ GeV}$ and $\Lambda_{\text{UE}}(14\text{TeV}) \approx 10 \text{ GeV}$.
Large scale at LHC order of magnitude bigger than hadronisation.
- More info in **variable R analytical studies** than fixed R MC analysis.

tu-logo

ur-logo

Summary of findings

- Different algorithms show a **similar sensitivity** to NP effects (contradicts folklore)
- UE depends on collider energy and MC model as well as R .
- Hadronization on jet “colour factor” and differently on R .
- $\Lambda_{\text{UE}}(1.96\text{TeV}) \approx 2 - 4 \text{ GeV}$ and $\Lambda_{\text{UE}}(14\text{TeV}) \approx 10 \text{ GeV}$.
Large scale at LHC order of magnitude bigger than hadronisation.
- More info in **variable R analytical studies** than fixed R MC analysis.

ur-logo

ur-logo

Summary of findings

- Different algorithms show a **similar sensitivity** to NP effects (contradicts folklore)
- UE depends on collider energy and MC model as well as R .
- Hadronization on jet “colour factor” and differently on R .
- $\Lambda_{\text{UE}}(1.96\text{TeV}) \approx 2 - 4 \text{ GeV}$ and $\Lambda_{\text{UE}}(14\text{TeV}) \approx 10 \text{ GeV}$. Large scale at LHC order of magnitude bigger than hadronisation.
- More info in **variable R analytical studies** than fixed R MC analysis.

Summary of findings

- Different algorithms show a **similar sensitivity** to NP effects (contradicts folklore)
- UE depends on collider energy and MC model as well as R .
- Hadronization on jet “colour factor” and differently on R .
- $\Lambda_{\text{UE}}(1.96\text{TeV}) \approx 2 - 4 \text{ GeV}$ and $\Lambda_{\text{UE}}(14\text{TeV}) \approx 10 \text{ GeV}$.
Large scale at LHC order of magnitude bigger than hadronisation.
- More info in **variable R analytical studies** than fixed R MC analysis.

tu-logo

ur-logo

Experimental considerations

- Knowing R dependence gives rise to idea of **optimal** R for various studies.
- The ideas presented here can be directly tested in measurement.

tu-logo

ur-logo

Experimental considerations

- Knowing R dependence gives rise to idea of **optimal** R for various studies.
- The ideas presented here can be directly tested in measurement.

tu-logo

ur-logo

For studies reconstructing e.g. mass-peaks want to minimise dispersion

$$\langle \delta p_t^2 \rangle = \langle \delta p_t \rangle_h^2 + \langle \delta p_t \rangle_{\text{UE}}^2 + \langle \delta p_t \rangle_{\text{PT}}^2$$

Perturbative R dependence is $\ln R$ at small R (dominant effect).
For pQCD studies just total NP (UE and hadronisation) :

$$R = \sqrt{2} \left(\frac{C_i A(\mu_f)}{\Lambda} \right)^{1/3}$$

tu-logo

ur-logo

For studies reconstructing e.g. mass-peaks want to minimise dispersion

$$\langle \delta p_t^2 \rangle = \langle \delta p_t \rangle_h^2 + \langle \delta p_t \rangle_{\text{UE}}^2 + \langle \delta p_t \rangle_{\text{PT}}^2$$

Perturbative R dependence is $\ln R$ at small R (dominant effect).
For pQCD studies just total NP (UE and hadronisation) :

$$R = \sqrt{2} \left(\frac{C_i A(\mu_f)}{\Lambda} \right)^{1/3}$$

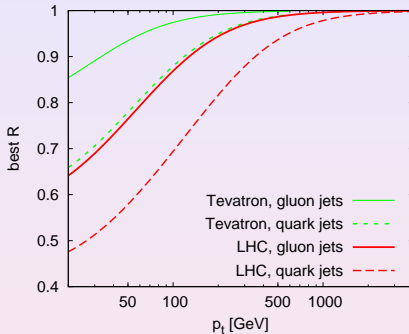
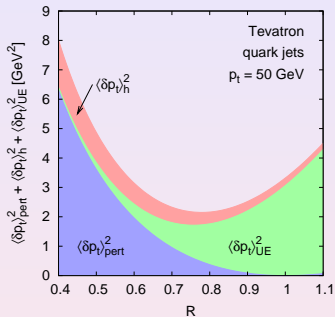
For studies reconstructing e.g. mass-peaks want to minimise dispersion

$$\langle \delta p_t^2 \rangle = \langle \delta p_t \rangle_h^2 + \langle \delta p_t \rangle_{\text{UE}}^2 + \langle \delta p_t \rangle_{\text{PT}}^2$$

Perturbative R dependence is $\ln R$ at small R (dominant effect).
For pQCD studies just total NP (UE and hadronisation) :

$$R = \sqrt{2} \left(\frac{C_i A(\mu_f)}{\Lambda} \right)^{1/3}$$

Best R

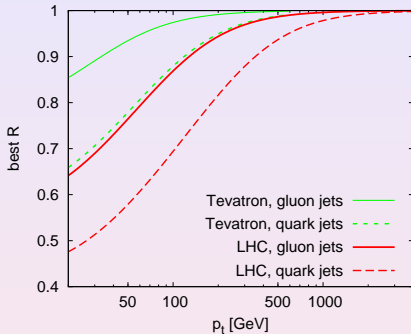
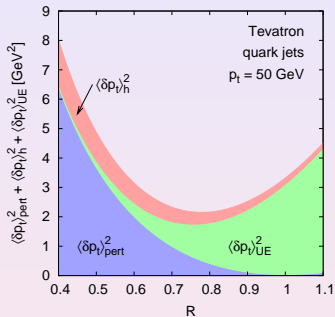


At high p_t one should use a larger R - minimises perturbative effect. Likewise for gluon jets a larger R is suggested. At LHC smaller R values than at Tevatron.

ur-logo

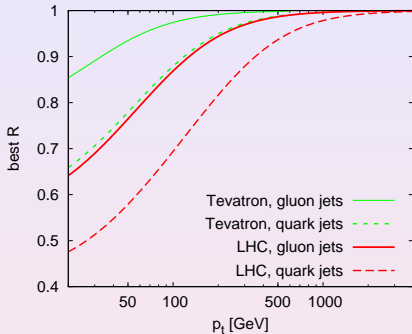
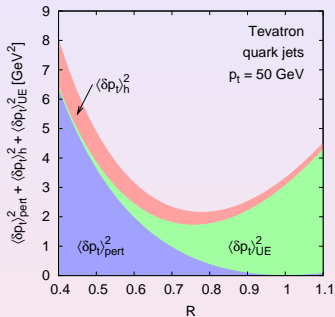
ur-logo

Best R



At high p_t one should use a larger R - minimises perturbative effect. Likewise for gluon jets a larger R is suggested. At LHC smaller R values than at Tevatron.

Best R

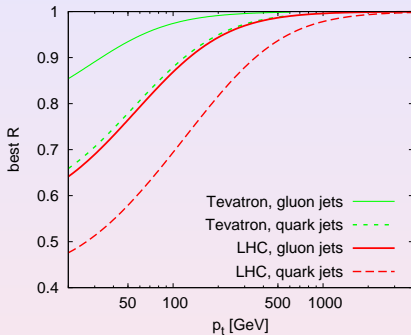
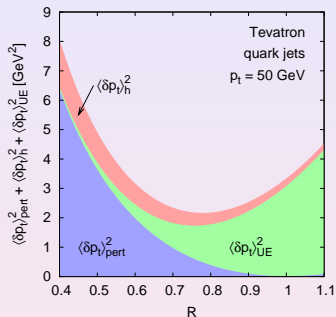


At high p_t one should use a larger R - minimises perturbative effect. Likewise for gluon jets a larger R is suggested. At LHC smaller R values than at Tevatron.

tu-logo

ur-logo

Best R

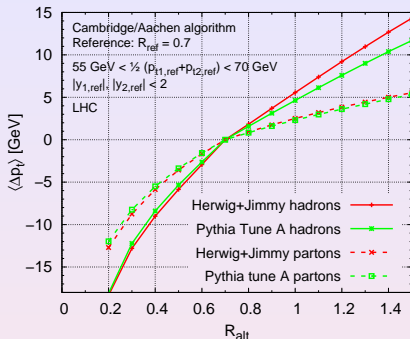


At high p_t one should use a larger R - minimises perturbative effect. Likewise for gluon jets a larger R is suggested. At LHC smaller R values than at Tevatron.

tu-logo

ur-logo

Direct tests



Direct experimental measurements of $\delta p_t(R)$. Can be compared to NLO + NP corrected results. Used to extract Λ_{UE} directly ?

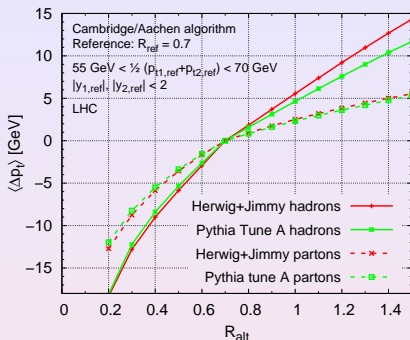
Apply results to single-inclusive jet p_t spectra.

$$\frac{d\sigma}{dp_t} = \frac{d\sigma}{dp_{t,pert}} (p_t - \langle \delta p_t \rangle_{NP})$$

ur-logo

ur-logo

Direct tests



Direct experimental measurements of $\delta p_t(R)$. Can be compared to NLO + NP corrected results. Used to extract Λ_{UE} directly ?

Apply results to single-inclusive jet p_t spectra.

$$\frac{d\sigma}{dp_t} = \frac{d\sigma}{dp_t^{\text{pert}}} (p_t - \langle \delta p_t \rangle_{\text{NP}})$$

- A set of IRC safe jet algorithms now available. Anti k_t recently introduced. **Cacciari Salam and Soyez 2008**
- Features of jet algorithms being analytically and systematically understood. Radius dependence an important aspect of NP effects.
- Can expect development of optimal tools to handle jets at hadron colliders.

- A set of IRC safe jet algorithms now available. Anti k_t recently introduced. Cacciari Salam and Soyez 2008
- Features of jet algorithms being analytically and systematically understood. Radius dependence an important aspect of NP effects.
- Can expect development of optimal tools to handle jets at hadron colliders.

- A set of IRC safe jet algorithms now available. Anti k_t recently introduced. Cacciari Salam and Soyez 2008
- Features of jet algorithms being analytically and systematically understood. Radius dependence an important aspect of NP effects.
- Can expect development of optimal tools to handle jets at hadron colliders.