



University
of Glasgow

THE PRICE OF A LARGE ELECTRON YUKAWA MODIFICATION

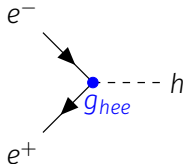
...and what we would learn from a Higgs pole run at FCC-ee

Ben Smith

(based on arXiv:2511.02642 w/ L. Allwicher, M. McCullough, S. Renner, D. Rocha)

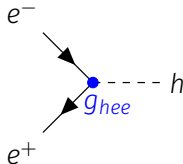
17th December 2025, **YTF 2025**

THE ELECTRON YUKAWA COUPLING



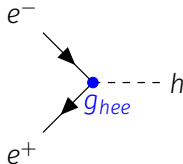
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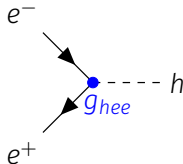
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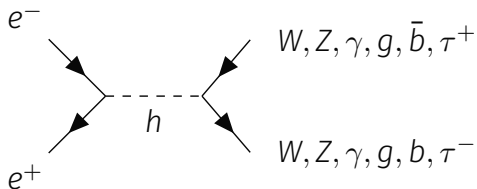
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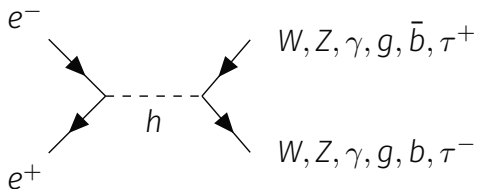


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- Parametrise deviations in terms of $\kappa_e = \frac{g_{hee}}{g_{hee}^{SM}}$
- Constraints:
 - $\kappa_e^{\text{LHC}} < 240$ (Tumasyan et al. 2023)
 - $\kappa_e^{\text{HL-LHC}} < 120$ (Cepeda et al. 2019)

- Projected $|\kappa_e| < 1.6$ @ 95% C.L from a dedicated run at the Higgs pole ($\sqrt{s} = m_h$) (d'Enterria, Poldaru, and Wojcik [2022](#)) .

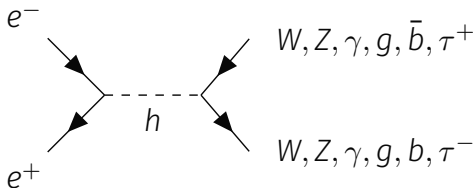


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 - Need high precision on Higgs mass (few MeV).
 - Need monochromatised beams.
 - Large backgrounds.

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➡ What do we learn (EFT/Models)?
Other observables?

EFT PERSPECTIVE

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{k,D>4} c_k^{(D)} \mathcal{O}_k^{(D)}$$

LEPTON YUKAWAS IN THE SMEFT

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- New contributions to lepton flavour conserving and violating Higgs couplings.

$$[g_{h\ell\ell}]_{ij} = \frac{1}{v}[M_\ell]_{ij} - \frac{v^2}{\sqrt{2}}[c_{eH}^*]_{ji} \Rightarrow g_{H\ell\ell} \not\propto M_\ell!$$

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- Assume real and ‘**electrophilic**’ flavour structure:
 $\rightarrow$ only $\text{Re}([c_{eH}]_{11}) \neq 0$

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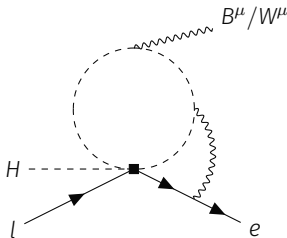
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- Leading connection $\mathcal{O}_{eH} \rightarrow \mathcal{O}_{eW}/\mathcal{O}_{eB}$ is at the two-loop level.



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Working the SMEFT alone, Δa_e is insufficient to constrain κ_e to $\mathcal{O}(1)$.

UV MODEL PERSPECTIVE

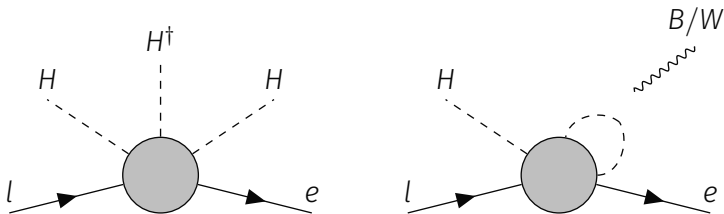
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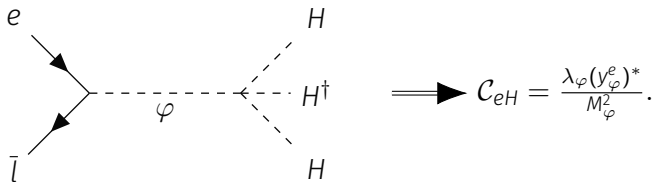
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Grey blob = diagram of arbitrary loop order involving heavy state exchange.

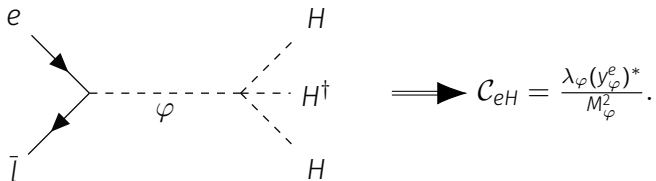
THE φ (2HDM) EXCEPTION

- If extension is a scalar doublet $\varphi \sim (1, 2, \frac{1}{2})$, can generate $\mathcal{O}_{eH}$ at tree level.

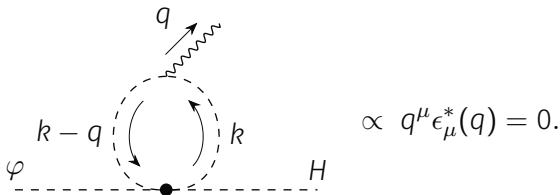


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- However, we do *not* generate $\mathcal{O}_{eB}/\mathcal{O}_{eW}$ at one loop.



$$\mathcal{C}_{e\gamma} = (\cos(\theta_W)\mathcal{C}_{eB} - \sin(\theta_W)\mathcal{C}_{eW}) = \frac{e}{16\pi^2} \left(\frac{g^2}{16\pi^2} \right)^{N_{\text{loops}}-1} \mathcal{C}_{eH}$$

N_{loops}	$ \kappa_e (\Delta a_e^{\text{future}})$
0	< 1.01
1	< 3
2	< 300

cf. $|\kappa_e^{\text{FCC-ee}}| < 1.6$

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→ Two loop suppression between κ_e and Δa_e required for large enhancements.

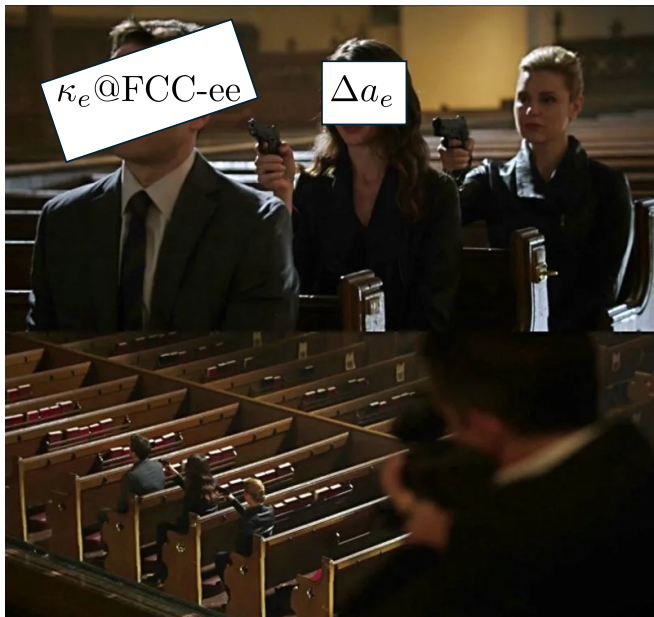
VISUAL SUMMARY



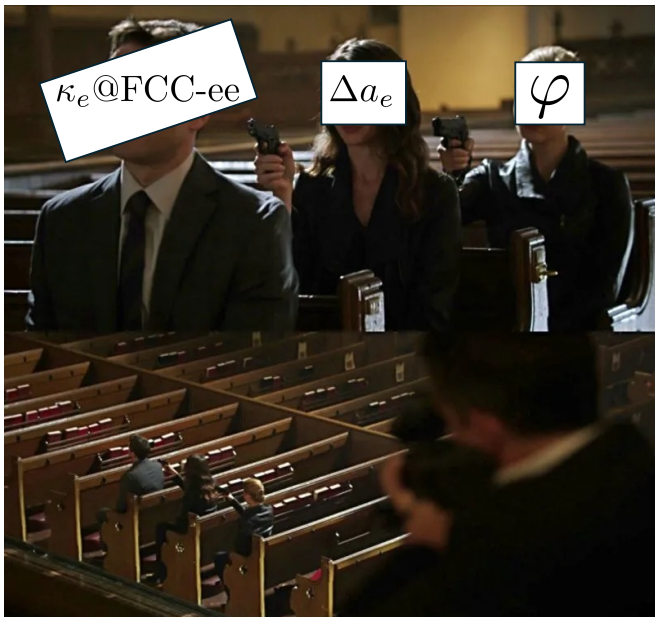
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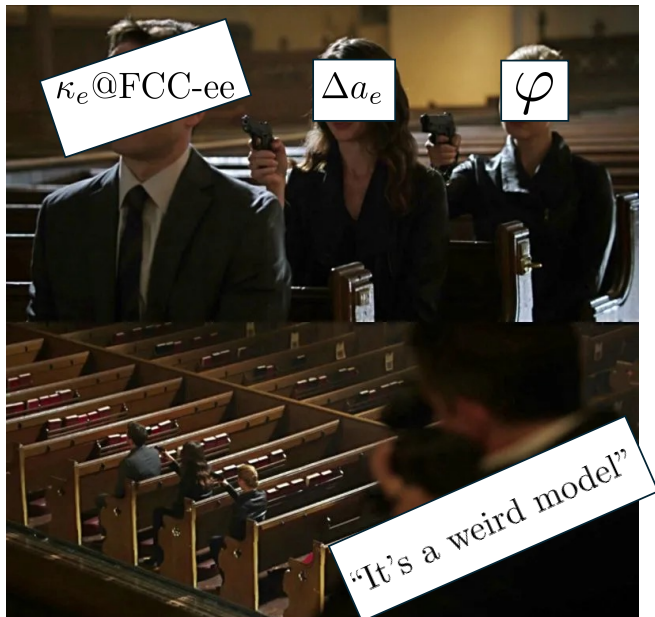
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Thank you!

BACKUP SLIDES

UV MODEL PERSPECTIVE

State	Spin	SM charges	$\mathcal{C}_{eH}$	$\mathcal{C}_{eB}/\mathcal{C}_{eW}$
$\mathcal{S}$	0	(1, 1, 0)	tree	1 loop
φ (with Higgs coupling)	0	(1, 2, $\frac{1}{2}$)	tree	2 loop
Ξ	0	(1, 3, 0)	tree	1 loop
Ξ_1	0	(1, 3, 1)	tree	1 loop
E	$\frac{1}{2}$	(1, 1, -1)	tree	1 loop
Δ_1	$\frac{1}{2}$	(1, 2, $-\frac{1}{2}$)	tree	1 loop
Δ_3	$\frac{1}{2}$	(1, 2, $-\frac{3}{2}$)	tree	1 loop
Σ	$\frac{1}{2}$	(1, 3, 0)	tree	1 loop
Σ_1	$\frac{1}{2}$	(1, 3, -1)	tree	1 loop
φ (with top coupling)	0	(1, 2, $\frac{1}{2}$)	1 loop	2 loop
ω_1	0	(3, 1, $-\frac{1}{3}$)	1 loop	1 loop
Π_7	0	(3, 2, $\frac{7}{6}$)	1 loop	1 loop
$\mathcal{U}_2$	1	(3, 1, $\frac{2}{3}$)	1 loop	1 loop
$\mathcal{Q}_5$	1	(3, 2, $-\frac{5}{6}$)	1 loop	1 loop

Blue = non-renormalisable interaction required to match to $\mathcal{C}_{eH}, \mathcal{C}_{eB}, \mathcal{C}_{eW} \Rightarrow$ See (Erdelyi, Gröber, and Selimovic 2025) .

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$$\kappa_e = \frac{g_{hee}}{\frac{m_e}{v}} = \frac{\frac{1}{2}(y_e - \frac{3v^2}{\sqrt{2}}\mathcal{C}_{eH})}{\frac{1}{\sqrt{2}}(y_e - \frac{v^2}{2}\mathcal{C}_{eH})} = 1 - \frac{\frac{v^2\mathcal{C}_{eH}}{y_e}}{1 - \frac{v^2\mathcal{C}_{eH}}{2y_e}} = 1 - \frac{\zeta}{1 - \zeta/2}$$

Large κ_e requires tuning ζ ($= \frac{v^2\mathcal{C}_{eH}}{2y_e}$) close to 1.

$$\Delta = \left| \frac{\partial \ln \kappa}{\partial \ln \zeta} \right| = \left| \frac{(1 - \kappa)(3 - \kappa)}{2\kappa} \right| ,$$

which behaves as

$$\lim_{\kappa \gg 1} \Delta = \frac{\kappa}{2} .$$

$\Rightarrow \kappa_e = 10$ requires $\sim 20\%$ tuning.

BIBLIOGRAPHY I



Cepeda, M. et al. (2019). “Report from Working Group 2: Higgs Physics at the HL-LHC and HE-LHC”. In: *CERN Yellow Rep. Monogr.* 7. Ed. by Andrea Dainese et al., pp. 221–584. DOI: [10.23731/CYRM-2019-007.221](#). arXiv: [1902.00134 \[hep-ph\]](#).



d’Enterria, David, Andres Poldaru, and George Wojcik (2022). “Measuring the electron Yukawa coupling via resonant s-channel Higgs production at FCC-ee”. In: *Eur. Phys. J. Plus* 137.2, p. 201. DOI: [10.1140/epjp/s13360-021-02204-2](#). arXiv: [2107.02686 \[hep-ex\]](#).



Di Luzio, Luca et al. (2025). “Model-Independent Tests of the Hadronic Vacuum Polarization Contribution to the Muon $g-2$ ”. In: *Phys. Rev. Lett.* 134.1, p. 011902. DOI: [10.1103/PhysRevLett.134.011902](#). arXiv: [2408.01123 \[hep-ph\]](#).



Erdelyi, Barbara Anna, Ramona Gröber, and Nudzeim Selimovic (2025). “Probing new physics with the electron Yukawa coupling”. In: *JHEP* 05, p. 135. DOI: [10.1007/JHEP05\(2025\)135](#). arXiv: [2501.07628 \[hep-ph\]](#).



Panico, Giuliano, Alex Pomarol, and Marc Riembau (2019). “EFT approach to the electron Electric Dipole Moment at the two-loop level”. In: *JHEP* 04, p. 090. DOI: [10.1007/JHEP04\(2019\)090](#). arXiv: [1810.09413 \[hep-ph\]](#).



Tumasyan, Armen et al. (2023). “Search for the Higgs boson decay to a pair of electrons in proton-proton collisions at $\sqrt{s}=13\text{TeV}$ ”. In: *Phys. Lett. B* 846, p. 137783. DOI: [10.1016/j.physletb.2023.137783](#). arXiv: [2208.00265 \[hep-ex\]](#).