

# Spectral Networks

& their role in 2d conformal field theory

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# Background

# Definition

Let  $C$  be a genus  $g$  Riemann surface with  $n \geq 1$  punctures, and  $\iota : \Sigma \hookrightarrow T^*C$  be a branched double cover,

$$\pi : \Sigma \xrightarrow{2:1} C.$$

For a curve  $\gamma \subset C$ , let  $\gamma_i$  denote the lift of  $\gamma$  to the  $i$ th sheet of  $\Sigma$ .

## Definition

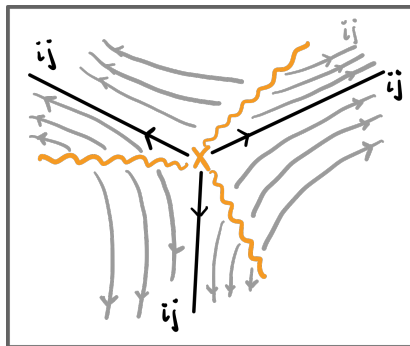
For a choice of phase  $\theta$ , an  $ij$ -trajectory is a curve  $w \subset C$  such that

$$e^{-i\theta} \int_{w_{ij}} \iota^* \lambda \in \mathbb{R}_{\geq 0}$$

where  $w_{ij} := w_i - w_j$  and  $\lambda \in \Omega^1(T^*C)$  is the canonical 1-form.

# Definition

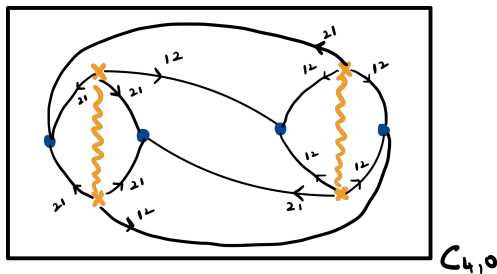
This defines a foliation of  $C$  (known as the horizontal foliation or WKB foliation). At the branch points of  $\Sigma$ , the foliation has a three-pronged structure.



These are known as the **critical  $ij$ -trajectories**. We orient all the critical  $ij$ -trajectories away from the branch point.

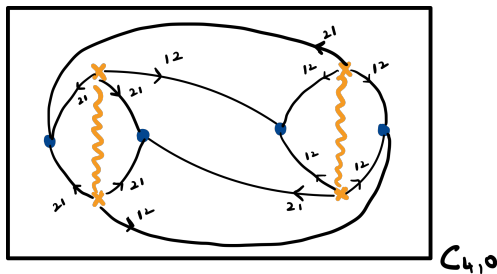
# Definition

The collection of all the oriented critical  $ij$  trajectories, the branch points, and the  $ij$  labels is known as a (rank 2) **WKB spectral network**  $\mathcal{W}_\theta(\Sigma)$ .



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## Working definition

A (rank 2) spectral network  $\mathcal{W}$  is a directed graph on  $C$  with a trivalent vertex at each branch point of  $\Sigma$ , and all the edges terminate on either punctures or branch points.

# The Physics

# The 4d story

The surface  $C$  is the moduli space of (canonical) UV surface defects in a 4d class S theory, or equivalently, the **UV curve** of the class S theory.

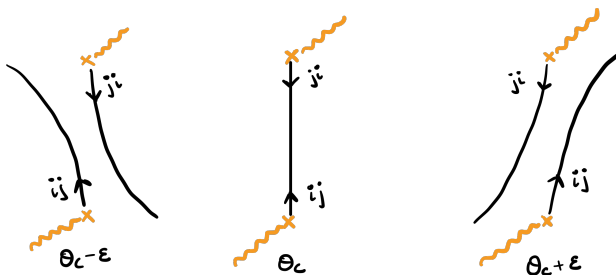
The surface  $\Sigma$  is the moduli space of IR surface defects in the 4d theory, or equivalently, the **Seiberg-Witten curve** of the class S theory.

A spectral network  $\mathcal{W}_\theta$  at phase  $\theta$  is the locus of all the surface defects  $z \in C$  that admit a BPS soliton at phase  $\theta$ . The analytical condition from before is the BPS equation. [GMN13], [BHM25]



# The 4d story

At certain phases, the topology of the network might degenerate.  
Such degenerate topologies encode the 4d BPS states.



This follows from the string/M-theory construction of class S theories.

# The 2d story

The AGT correspondence states that  $4d \mathcal{N} = 2$   $SU(2)$  gauge theories are dual to 2d Liouville conformal field theory on  $C$ . [AGT10]

**4d instanton partition functions**



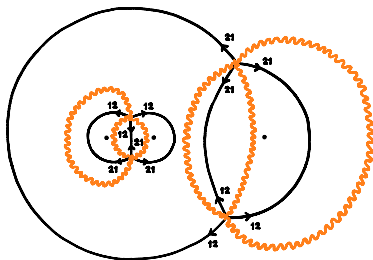
**2d Liouville conformal blocks**

These conformal blocks admit an integral representation for a suitable choice of contours on  $C$ . [DF84]

**Punchline:** spectral networks can supply the required contours for these integral representations. [HN24], [HM\*]

# The 2d story

In particular, the integral representations  $\mathcal{Z}_{\mathcal{W}_{\text{FN}}}$  defined using Fenchel-Nielsen networks reproduce the *standard* Liouville conformal blocks. [HM\*]



For a Fenchel-Nielsen network, the complement  $C \setminus \mathcal{W}_{\text{FN}}$  is a union of punctured discs. Note that none of the critical  $ij$ -trajectories of a Fenchel-Nielsen network terminate on a puncture.

# Outlook

# Related ideas

- We may obtain the algebra of loop operators in the theory using  $q$ -nonabelianisation. [NY20]
- We may extend this to higher-rank spectral networks and higher-rank conformal field theories.
- We would like to extend this to 3d Chern-Simons theory.

Thank you!



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