Quantum Corrections to Symmetron Fifth Forces

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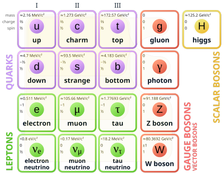

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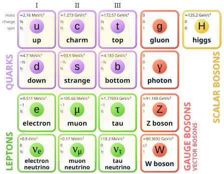

- Inflation
- Dark energy
- Dark matter

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Observational effects: Contribution to vacuum energy (dark energy or inflation), fifth forces (dark matter).

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Example: The chamelon model. The mass m of the scalar is proportional to ambient density of matter. The range of the fifth force goes like m^{-1}.

The Symmetron Model

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$$A(\phi) = 1 + \frac{\phi^2}{2M^2}$$

$$V(\phi) = -\frac{1}{2}\mu^2\phi^2 + \frac{1}{4}\lambda\phi^4$$

$$V_{\text{eff}}(\phi) = \frac{1}{2}\left(\frac{\rho}{M^2} - \mu^2\right)\phi^2 + \frac{1}{4}\lambda\phi^4$$

$\rho \rightarrow$ background matter density
$M \rightarrow$ matter coupling
$\mu \rightarrow$ mass term
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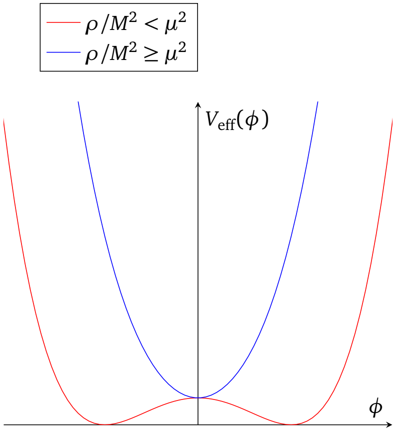

Figure 1: The symmetron effective potential.

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Low density $\Longrightarrow$ spontaneous symmetry breaking, nonzero vacuum expectation value v, unscreened force with coupling strength v/M.

High density $\Longrightarrow$ symmetry is restored, $\phi \to 0$, fifth force is screened.

Constraints

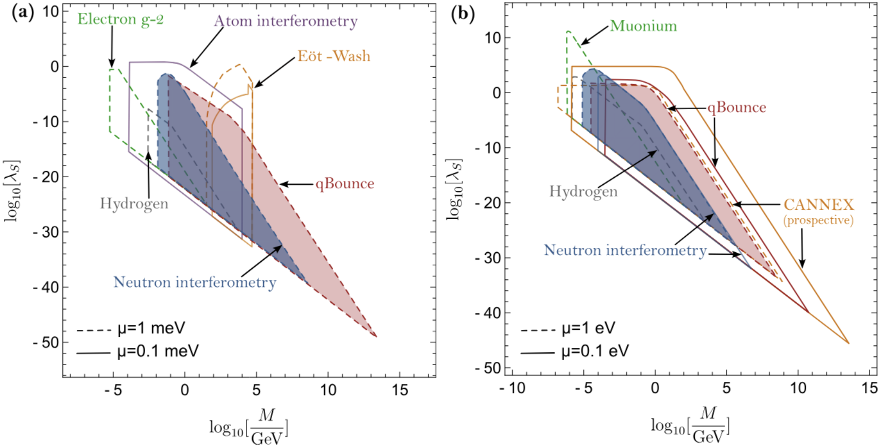

Figure 2: Constraints from H. Fischer, C. Käding and M. Pitschmann, 2024

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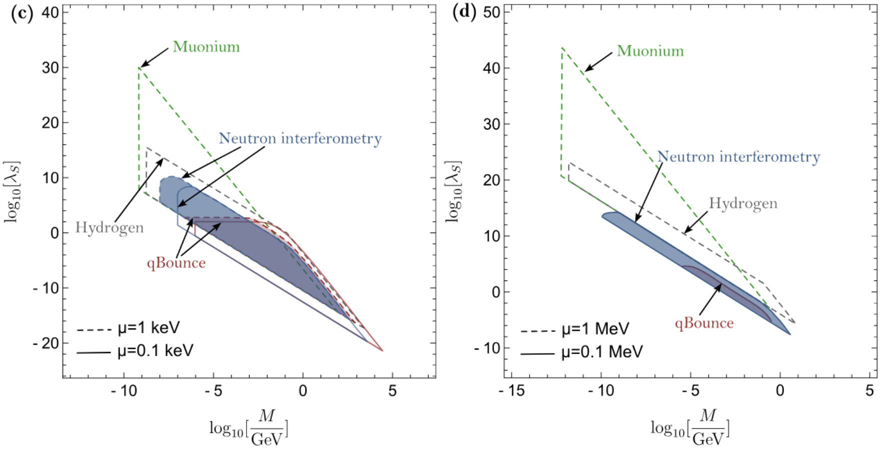

Figure 3: Constraints from H. Fischer, C. Käding and M. Pitschmann, 2024

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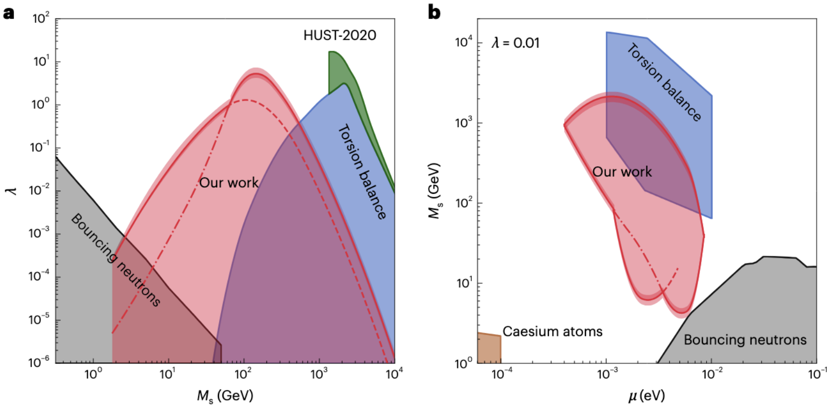

Figure 4: Constraints from P. Yin, X. Xu, K. Tian, et al., 2025

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Have we been missing significant quantum corrections to symmetron field profiles?

Leading-Order Quantum Correction

(B. Garbrecht and P. Millington, 2015)

Let $\phi_{\text{qu}}(x) = \phi_{\text{cl}}(x) + \delta\phi(x)$. Then

$$\delta\phi(x) = \int \mathrm{d}^4y\, G(x,y)\Pi^R(y)\phi_{\text{cl}}(y)\,, \tag{2}$$

where $G(x,y)$ is the propagator, defined by

$$\left.\frac{\delta^2 S[\phi]}{\delta\phi(x)\delta\phi(y)}\right|_{\phi=\phi_{\text{cl}}} G(x,y) = -\delta(x-y) \tag{3}$$

and $\Pi^R(x) = 3i\lambda G(x,x) + \delta m^2 + \delta\lambda\phi_{\text{cl}}(x)^2$ is the renormalised tadpole contribution.

Classical Solution

Static configuration,

$$\nabla^2 \phi = \frac{\mathrm{d}V_{\mathrm{eff}}}{\mathrm{d}\phi} = \left(\frac{\rho(\mathbf{x})}{M^2} - \mu^2 \right) \phi + \lambda \phi^3 \,. \tag{4}$$

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Spherical source $\rho(\mathbf{x}) \to \rho(r) = \rho_0\Theta(R-r)$,

$$\frac{\mathrm{d}^2\phi}{\mathrm{d}r^2} + \frac{2}{r}\frac{\mathrm{d}\phi}{\mathrm{d}r} = \left(\frac{\rho(r)}{M^2} - \mu^2\right)\phi + \lambda\phi^3\,. \tag{5}$$

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Thin-wall approximation: Assume $R \gg \mu^{-1}$. Then $2r^{-1}\mathrm{d}\phi/\mathrm{d}r \to 0$. Let $s = R - r$,

$$\frac{\mathrm{d}^2\phi}{\mathrm{d}s^2} = \left(\frac{\rho(s)}{M^2} - \mu^2\right)\phi + \lambda\phi^3\,. \tag{6}$$

Classical Solution

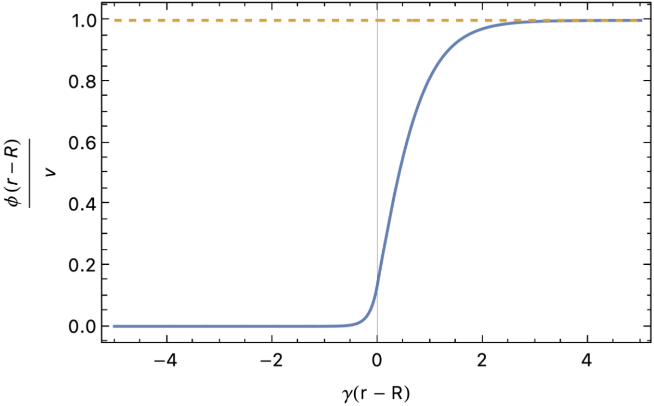

Figure 5: The classical symmetron field profile in the thin-wall approximation.

Green's Function

$G(x,x')$ determined by weighted sum over l and integral over E of $G_l(r,r';E)$, satisfying

$$\left(\frac{\mathrm{d}^2}{\mathrm{d}r^2} - \frac{l(l+1)}{r^2} + E^2 - \left.\frac{\mathrm{d}^2 V_{\mathrm{eff}}}{\mathrm{d}\phi^2}\right|_{\phi=\phi_{\mathrm{cl}}}\right) G_l(r,r';E) = \frac{\delta(r-r')}{r^2} \tag{7}$$

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Planar limit

$$\frac{l(l+1)}{R^2} \approx p^2 \Longrightarrow \sum_l (2l+1) \to \int \mathrm{d}p\, p \tag{9}$$

Potential

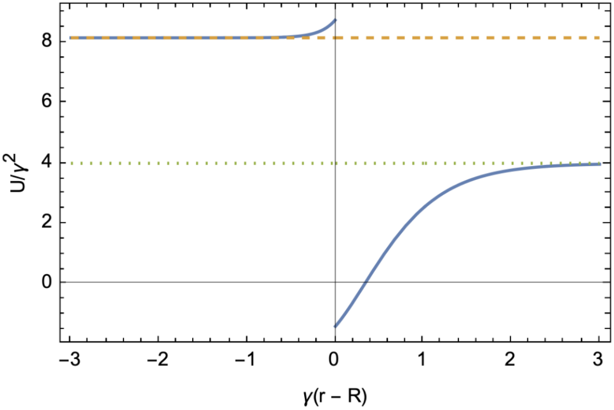

Figure 6: The potential for the Green's function and eigenfunction problem.

Decomposition

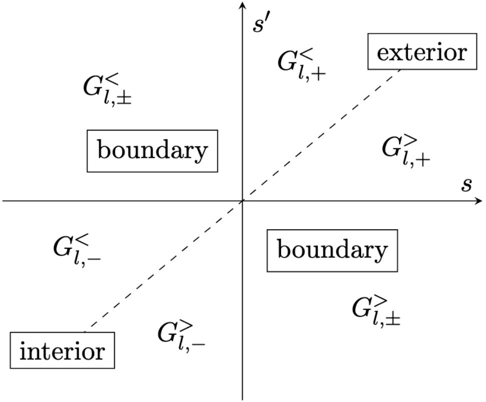

Figure 7: A graphical depiction of the decomposition of the Green's function in the s, s'-plane.

Coincident Green's Function

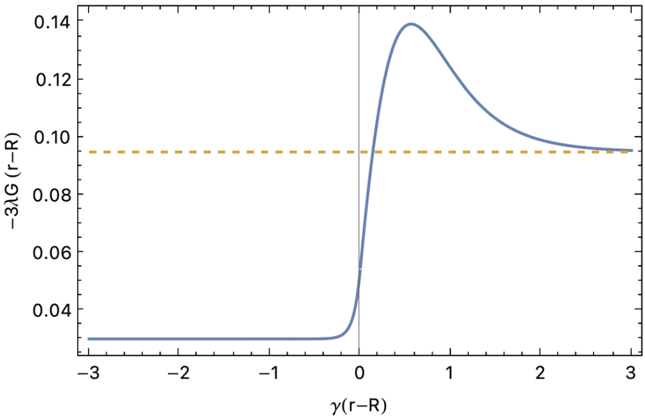

Figure 8: The coincident Green's function in position space.

Tangent: What if we don't take the thin wall approximation?

The 1D equation of motion also describes very long rods or very large planes.

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The 1D equation of motion also describes very long rods or very large planes. In this way, the system of equations

$$\frac{\mathrm{d}^2\phi}{\mathrm{d}x^2} = \left(\frac{\rho_0 \operatorname{rect}\left(\frac{x}{2R}\right)}{M^2} - \mu^2\right)\phi + \lambda\phi^3\,;\ \phi'(0) = 0, \phi(\pm\infty) = v \tag{10}$$

captures cylindrical-planar symmetry, as opposed to the spherical-planar symmetry of the thin-wall approximation.

1D is Harder Than 3D (plus approximations)

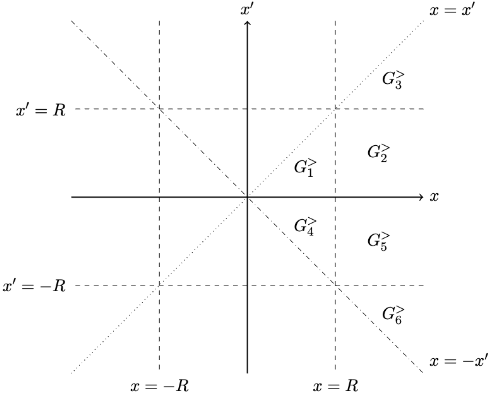

Figure 9: A graphical depiction of the decomposition of the Green's function in the x, x'-plane.

Renormalisation

Renormalisation conditions

$$\left.\frac{\mathrm{d}^2 V_{\text{1-loop}}}{\mathrm{d}\phi^2}\right|_{\phi=v} = 2\mu^2 \text{ and } \left.\frac{\mathrm{d}^4 V_{\text{1-loop}}}{\mathrm{d}\phi^4}\right|_{\phi=v} = 6\lambda \tag{11}$$

determine counterterms δm^2 and $\delta\lambda$.

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$$\frac{1}{2\pi^2}\int_0^\Lambda \mathrm{d}p\, p^2 \Delta m^2(p) = \delta m^2 \text{ and } \frac{1}{2\pi^2}\int_0^\Lambda \mathrm{d}p\, p^2 \Delta\lambda(p) = \delta\lambda\,. \tag{12}$$

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Pseudocounterterms facilitate numerical computation of the tadpole contribution,

$$\Pi^R(x) = \frac{1}{2\pi^2}\int_0^\Lambda \mathrm{d}p\, p^2 \left(-3\lambda G(s,s) + \Delta m^2 + \Delta\lambda\phi^2\right)\,. \tag{13}$$

Quantum-Corrected Field

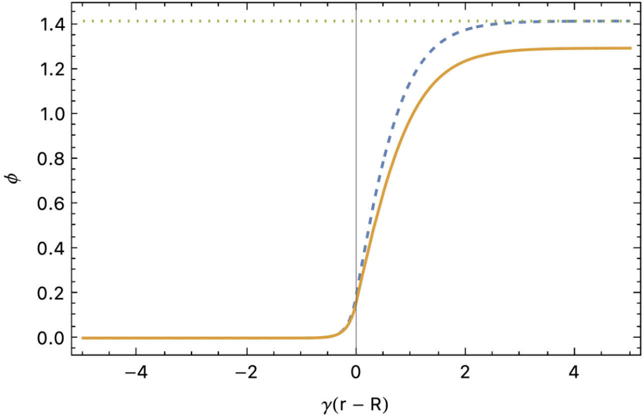

Figure 10: The classical (blue, dashed) and quantum (orange, solid) symmetron field profiles.

Some Analytical Understanding

Shift in the vev Δv,

$$\left.\frac{\partial V_{\text{1-loop}}}{\partial \phi}\right|_{v+\Delta v} = 0 \Rightarrow \Delta v = -\frac{27\lambda}{16\pi^2} v + O(\lambda^2)\,, \tag{14}$$

changes the asymptotic value.

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Shift in the mass,

$$m_{\text{1-loop}} = \left.\frac{\partial^2 V_{\text{1-loop}}}{\partial \phi^2}\right|_{v+\Delta v} = 2\mu^2 \left(1 - \frac{81\lambda}{16\pi^2}\right) + O(\lambda^2)\,, \tag{15}$$

changes the slope.

Slope Shift

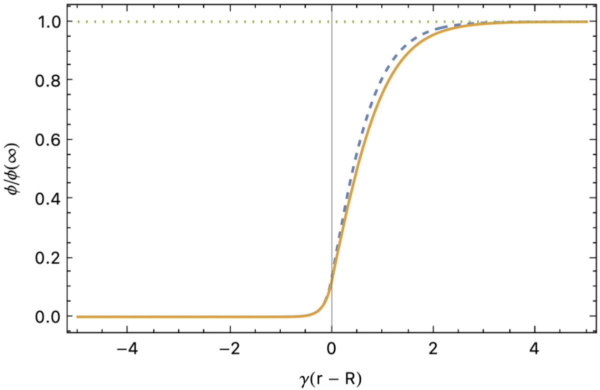

Figure 11: The classical (blue, dashed) and quantum (orange, solid) symmetron field profiles, normalised to their asymptotic values.

Quantum-Corrected Force

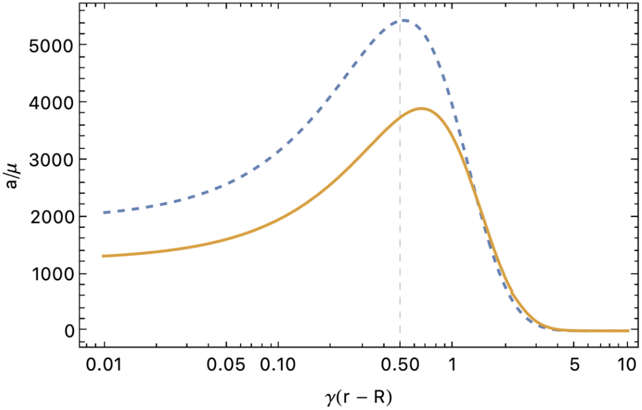

Figure 12: The classical (blue, dashed) and quantum (orange, solid) symmetron fifth force profiles. $\mu = 1\text{GeV}$, $M = 10\text{MeV}$, $\lambda = 0.5$, $\rho_0 = 2.45 \times 10^{-3}\text{GeV}^4$ (hydrogen spectroscopy)

Quantum-Corrected Force

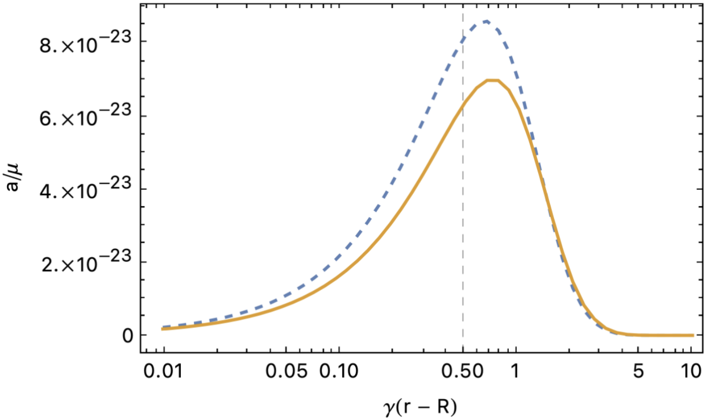

Figure 13: The classical (blue, dashed) and quantum (orange, solid) symmetron fifth force profiles. $\mu = 10^{-1}$meV, $M = 10^{-2}$GeV, $\lambda = 10^{-0.5}$, $\rho_0 = 8.178 \times 10^{-5}$MeV^4 (atom interferometry)

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In other words, we should expect virtually no μ-dependence and approximately linear λ-dependence in ΔF.

μ-dependence

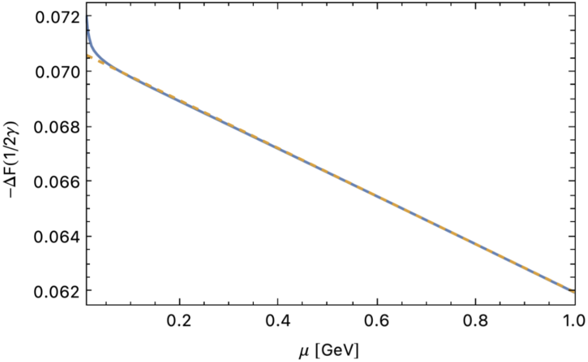

Figure 14: The dependence of the relative shift in the force on the mass μ, at one Compton wavelength from the surface of the source.

λ-dependence

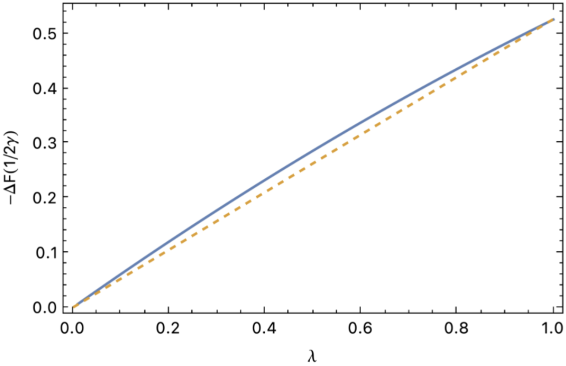

Figure 15: The dependence of the relative shift in the force on the self-coupling λ, at one Compton wavelength from the surface of the source.

Conclusions

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The magnitude of the correction to the force scales almost linearly with the self-coupling.

Current constraints may be weaker than we thought, especially for nonperturbative self-couplings.