

Ultralight dark matter haloes in galaxies

Kateryna Korshynska

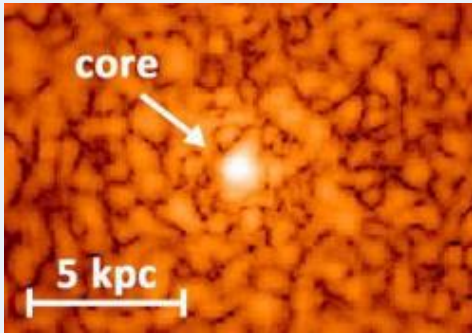
*Fundamentale Physik für Metrologie (FPM) PTB Braunschweig,
TU Braunschweig*

Durham 2026

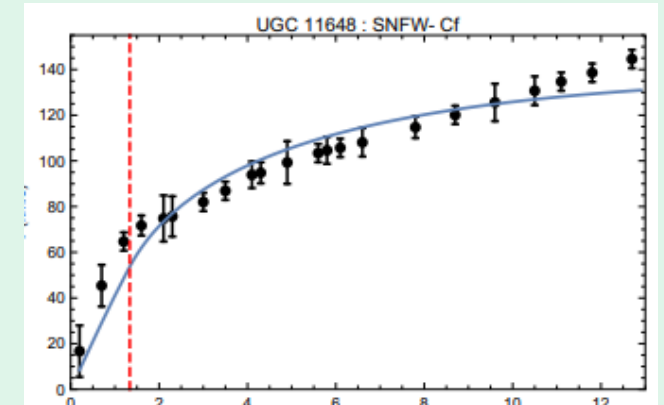
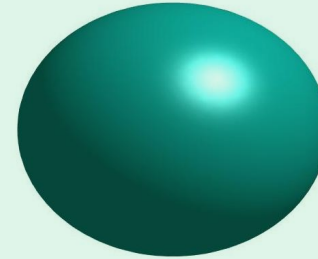
"High Precision Experiments for Light Dark Matter"

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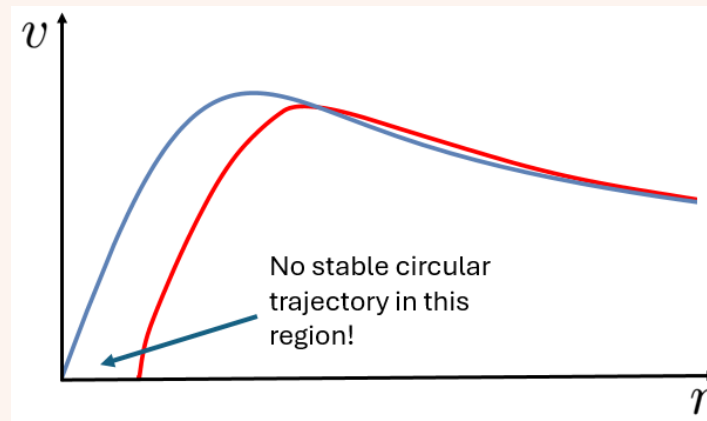
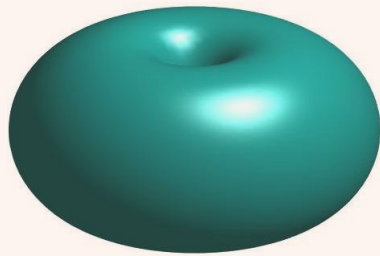
1. Ultralight dark matter (ULDM): introduction



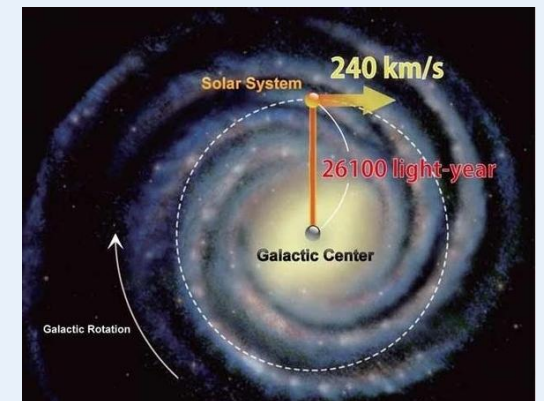
2. ULDM soliton and its astrophysical implications



3. ULDM vortex and its astrophysical implications

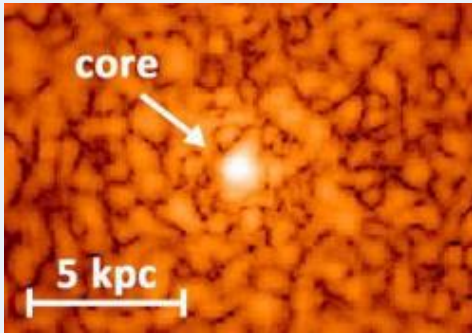


4. Conclusions

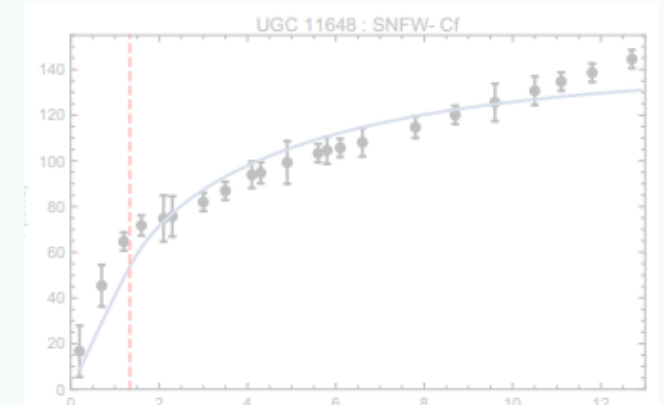
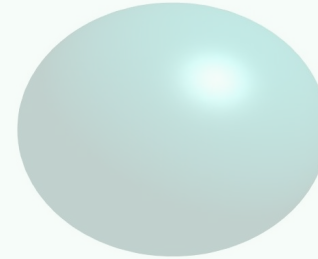


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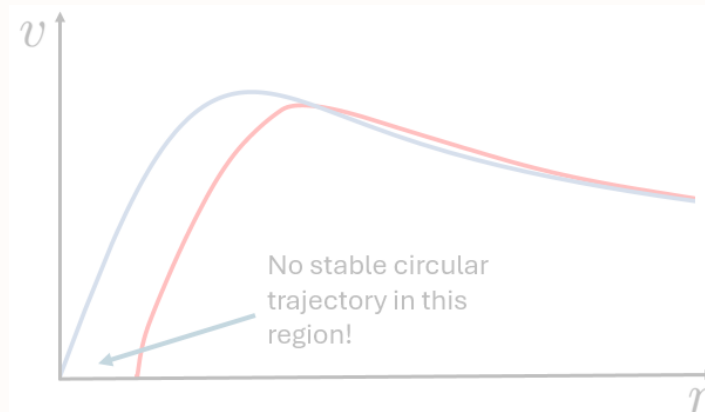
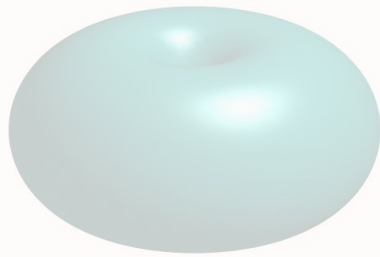
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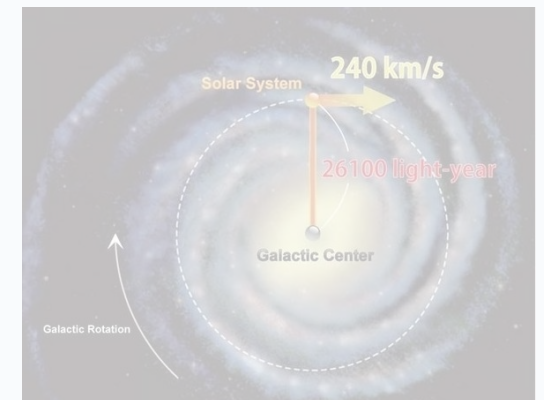
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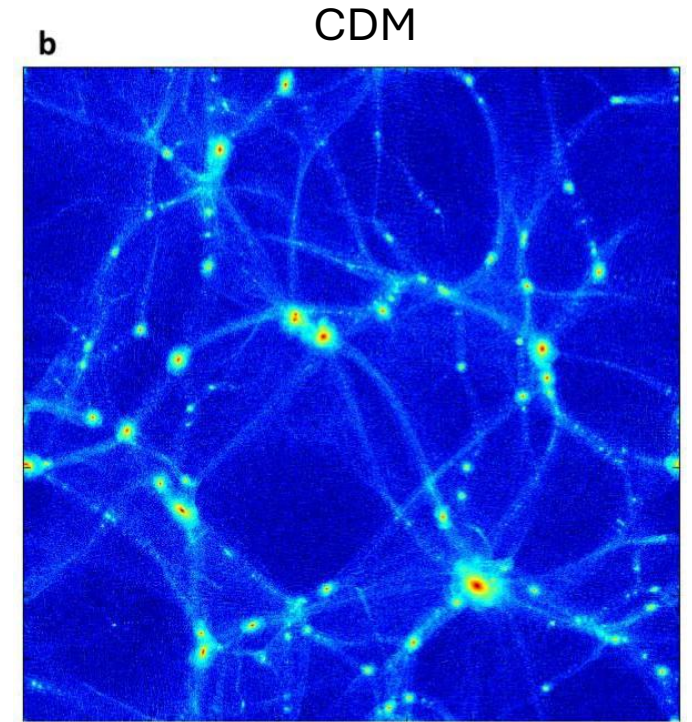


4. Conclusions



Cold dark matter (CDM)

On large scales dark matter is well-described by cold dark matter (CDM) model as a non-relativistic collisionless perfect fluid.

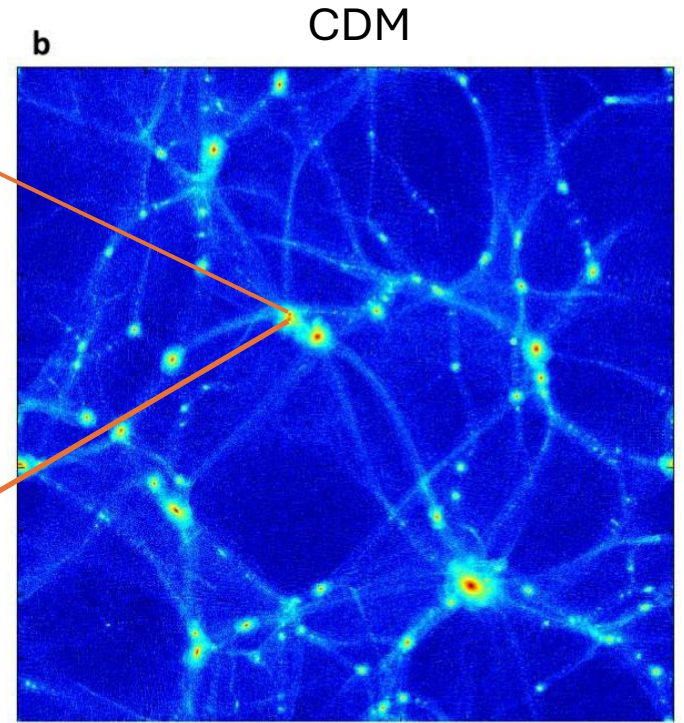
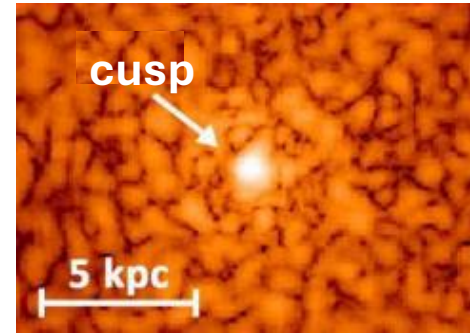
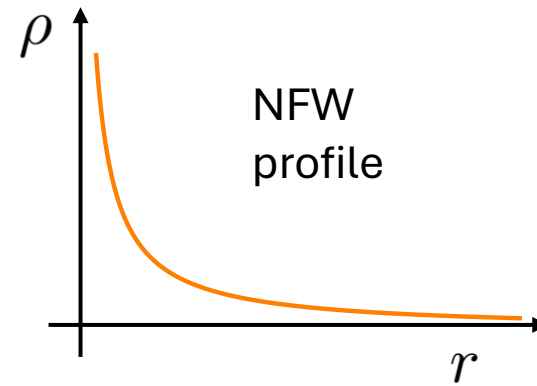


Hsi-Yu Schive, Tzihong Chiueh, Tom Broadhurst, NaturePhysics (2014)

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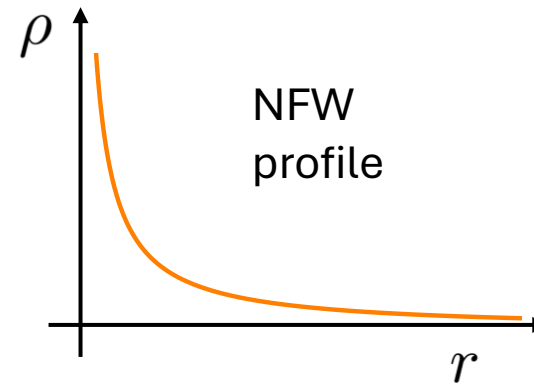
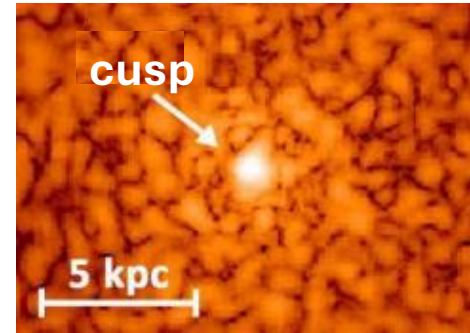
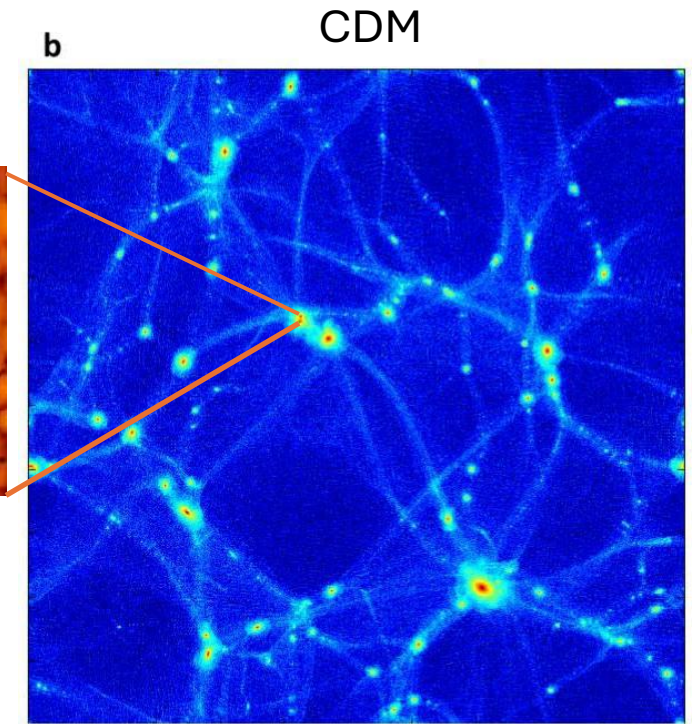
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Include the impact of baryonic matter in galactic simulations

Introduce modified gravity, based on the scaling relations in galaxies

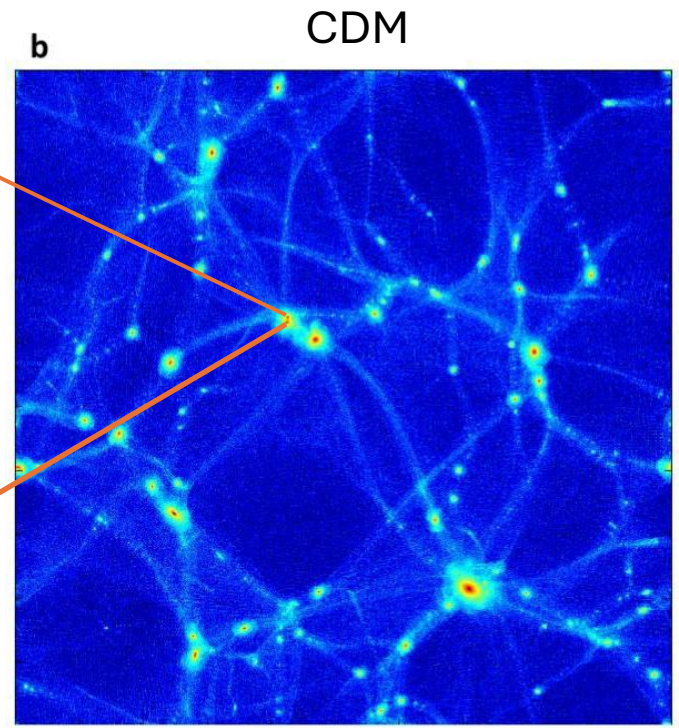
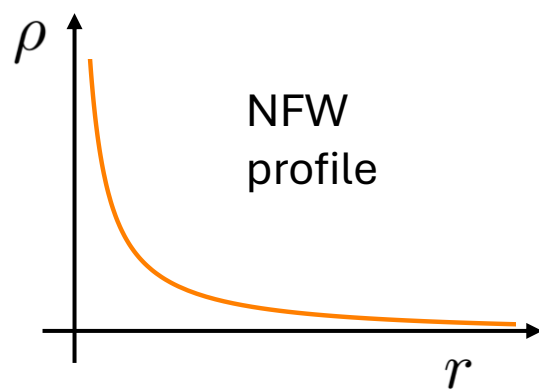
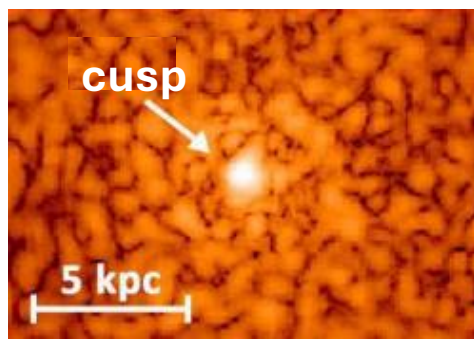


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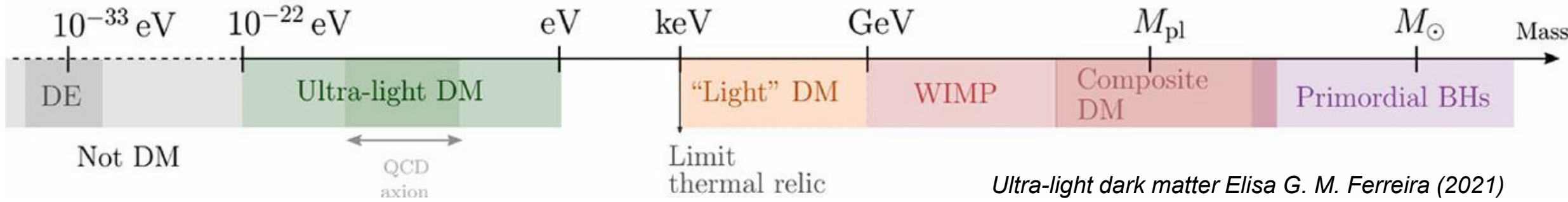
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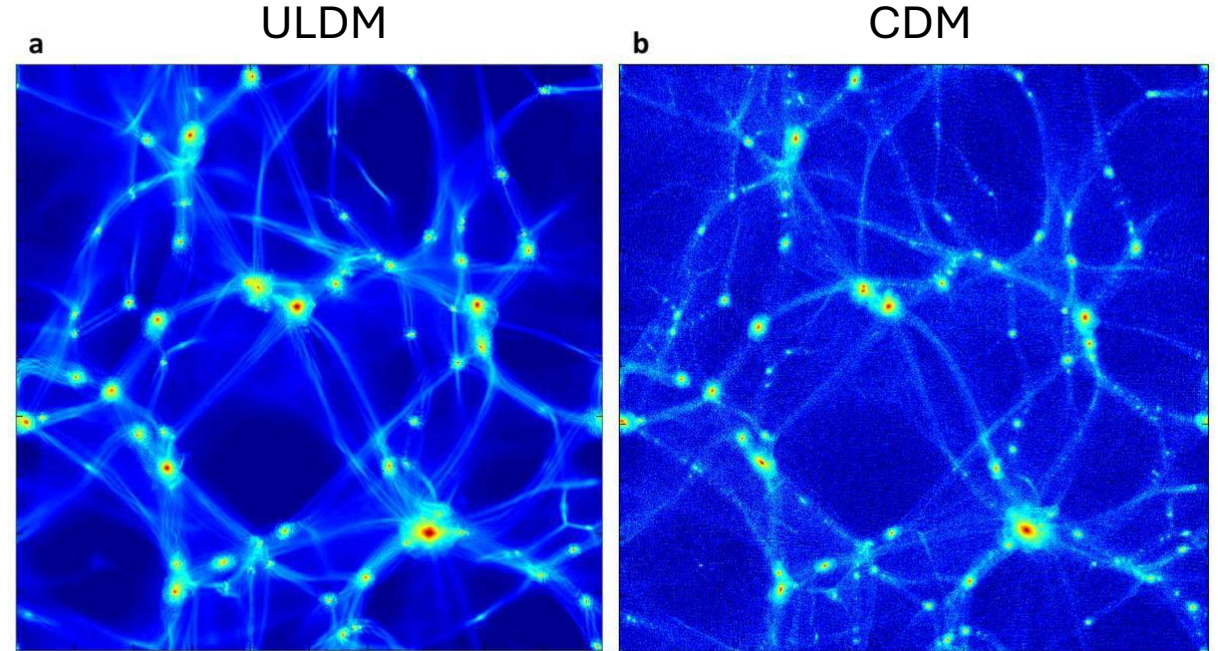
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Ultra-light dark matter Elisa G. M. Ferreira (2021)

Ultralight dark matter

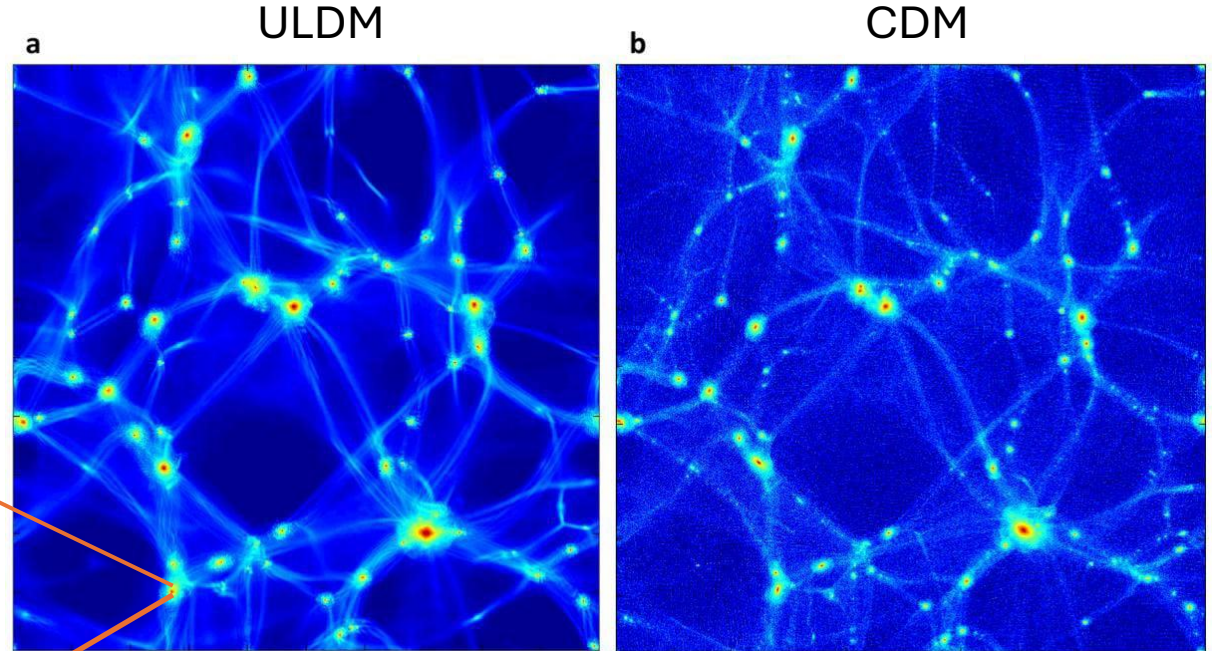
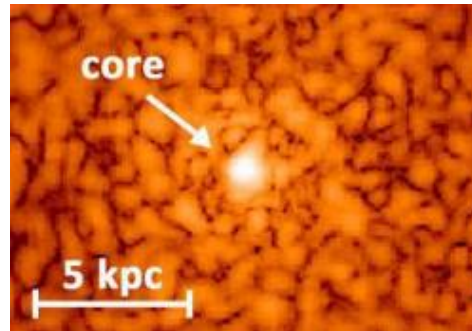
We assume that the DM consists of **ultralight bosons** from Standard Model extensions – axions or axionlike particles



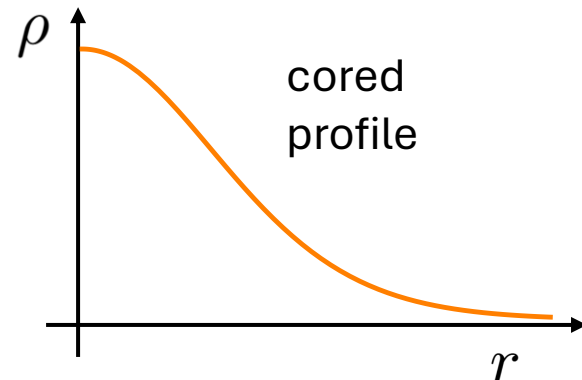
Cosmic Structure as the Quantum Interference of a Coherent Dark Wave, Hsi-Yu Schive et al. (2014)

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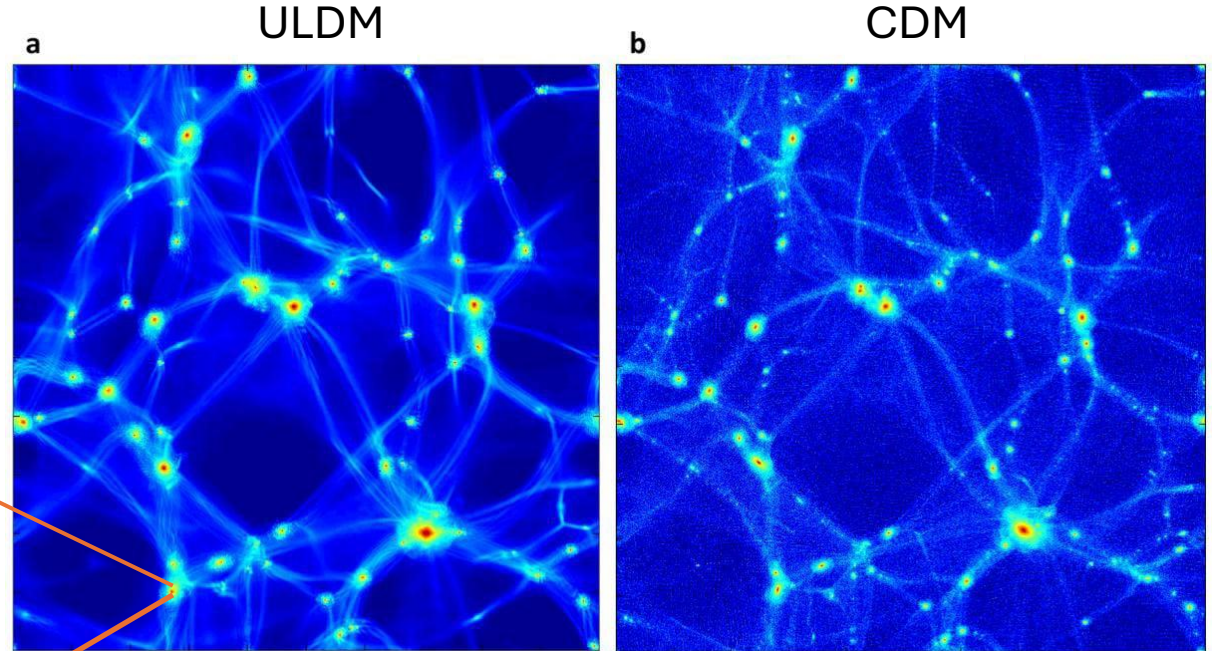
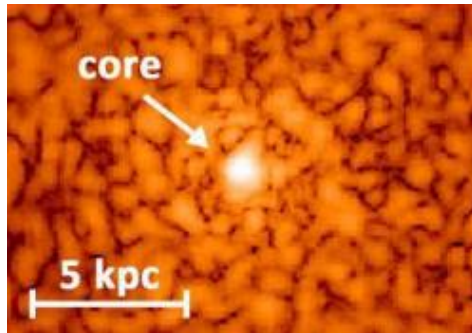
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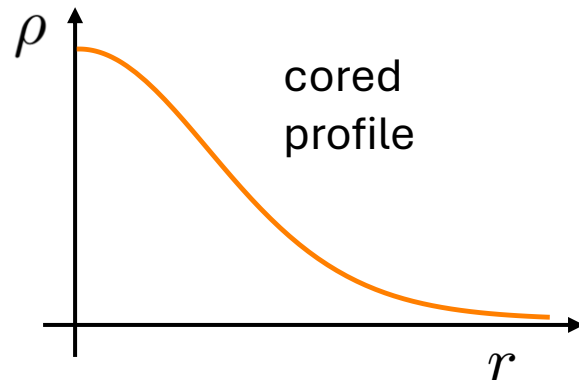
On the large scales ($>$ galaxy) ULDM reproduces CDM, but it behaves differently on galactic scales!

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On the large scales ($>$ galaxy) ULDM reproduces CDM, but it behaves differently on galactic scales!

Why?

GPP equations

Ultralight dark matter is described by the system of Gross-Pitaevskii-Poisson equations

$$\begin{cases} i\hbar\partial_t\psi(\mathbf{r}, t) = \left(-\frac{\hbar^2}{2m}\Delta + m\Phi(\mathbf{r}, t) + gN|\psi(\mathbf{r}, t)|^2\right) \psi(\mathbf{r}, t) \\ \Delta\Phi(\mathbf{r}, t) = 4\pi GmN|\psi(\mathbf{r}, t)|^2 \end{cases}$$

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Names in literature: BECDM, ψ DM, ULDM, wave DM, fuzzy DM (FDM)

Extremely small
boson mass: $m \sim 10^{-22} \text{ eV}/c^2 \rightarrow \lambda \sim 1 \text{ kpc}$

Huge occupation numbers: $\frac{10^8 M_\odot}{10^{-22} \text{ eV}/c^2} \sim 10^{96}$

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Let's transform GPP to equivalent hydrodynamic equations

$$\psi = |\psi|e^{\frac{i}{\hbar}S(\mathbf{r}, t)} \quad \rho = mN|\psi|^2 \text{ - density} \quad \mathbf{u} = \frac{1}{m}\nabla S \text{ - velocity}$$

$$\begin{aligned} \partial_t\rho + \nabla(\rho\mathbf{u}) &= 0 \\ \partial_t\mathbf{u} + (\mathbf{u} \cdot \nabla)\mathbf{u} &= -\frac{1}{\rho}\nabla p - \nabla\Phi - \frac{1}{m}\nabla Q \\ \Delta\Phi &= 4\pi G\rho \end{aligned}$$

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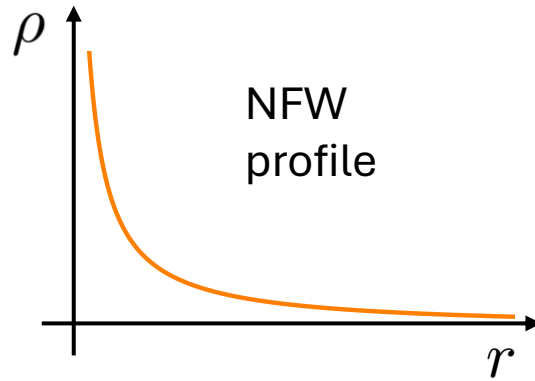
CDM is pressureless, but ULDM is not, since it has

- quantum pressure $Q = -\frac{\hbar^2}{2m} \frac{\Delta\sqrt{\rho}}{\sqrt{\rho}}$
- pressure due to self-interaction $p = \frac{g}{2m^2} \rho^2$

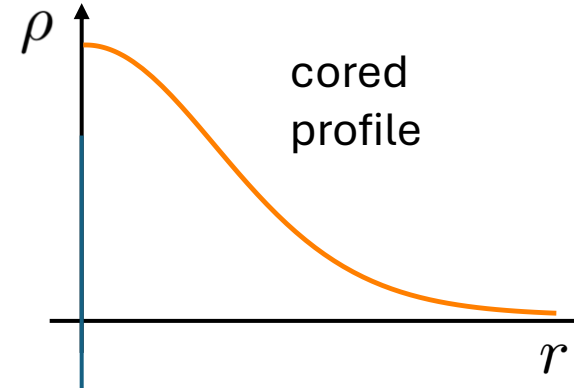
Absent in the case of CDM!

Ultralight boson and halo structure

CDM: gravity attracts matter to the center



ULDM: quantum pressure and self-interaction (if present) counteract gravity



$$\partial_t \rho + \nabla(\rho \mathbf{u}) = 0$$

$$\partial_t \mathbf{u} + (\mathbf{u} \cdot \nabla) \mathbf{u} = -\frac{1}{\rho} \nabla p - \nabla \Phi - \frac{1}{m} \nabla Q$$

$$\Delta \Phi = 4\pi G \rho,$$

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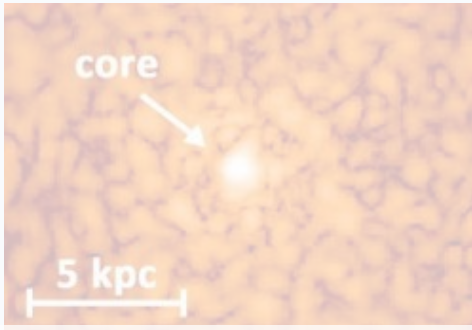
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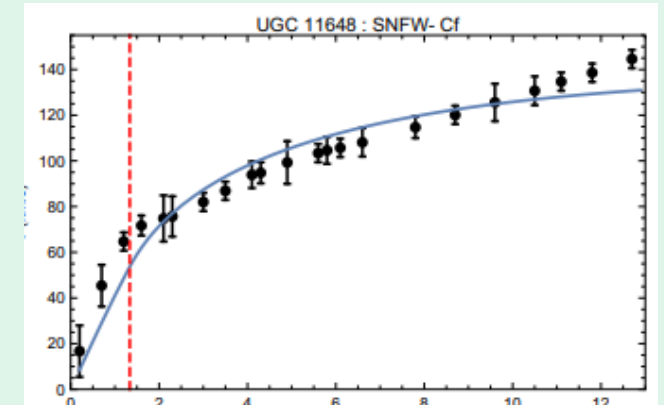
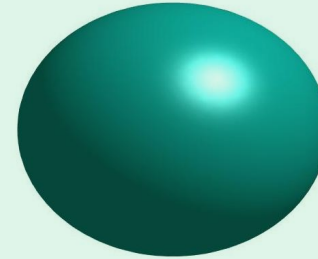
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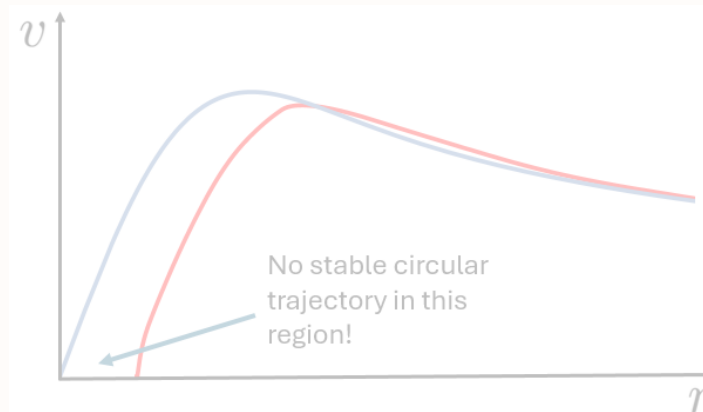
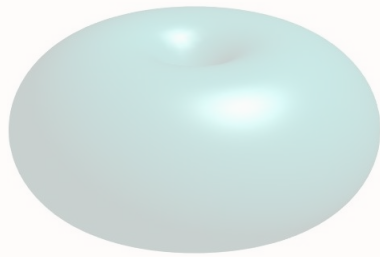
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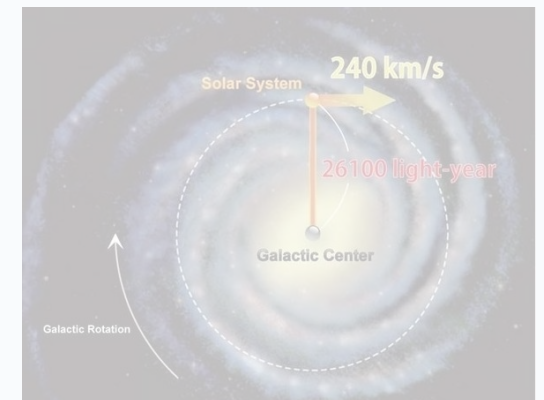
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4. Conclusions



Simple model to describe a ULDM core

One can fit the numerical solution of the GPP equations

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„Understanding the Core-Halo Relation of Quantum Wave Dark Matter from 3D Simulations“ H.-Y. Schive et al. 2014

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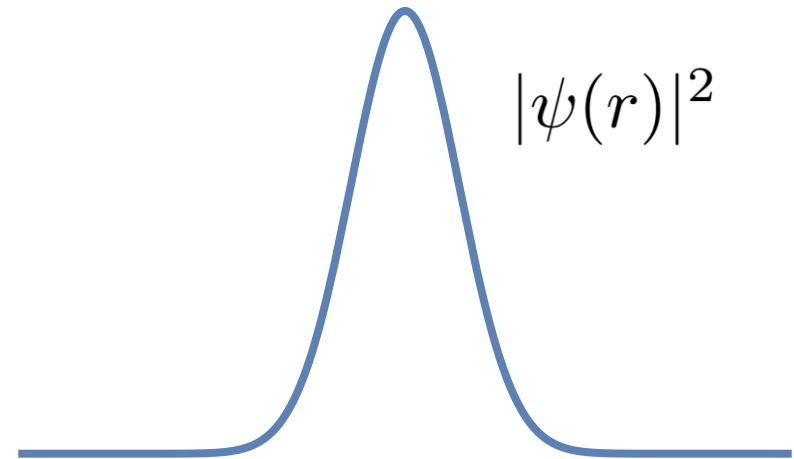
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A complementary approach: variational ansatz

$$\psi(\mathbf{r}, t) = \frac{\sqrt{M}}{\pi^{3/4} R^{3/2}} e^{-r^2/(2R^2) + i\beta r^2}$$

variational parameters

$$S_{\text{GPP}} = \frac{1}{m} \int dt \int d\mathbf{r} \psi^* \left[i\hbar \frac{\partial}{\partial t} + \frac{\hbar^2}{2m} \nabla^2 - \frac{m}{2} \Phi \right] \psi$$



Chavanis P. Phys Rev D (2011)

Simple model to describe a ULDM core

$$\psi(\mathbf{r}, t) = \frac{\sqrt{M}}{\pi^{3/4} R^{3/2}} e^{-r^2/(2R^2) + i\beta r^2}$$

From the GPP action we find the Lagrange equations of motion for the variational parameters:

$$\left\{ \begin{array}{l} \beta = \frac{m}{2\hbar} \frac{\dot{R}}{R} \\ \ddot{R} - \frac{\hbar^2}{m^2} \frac{1}{R^3} + \frac{\sqrt{2}}{3\sqrt{\pi}} \frac{GM}{R^2} = 0 \end{array} \right.$$

quantum pressure gravity

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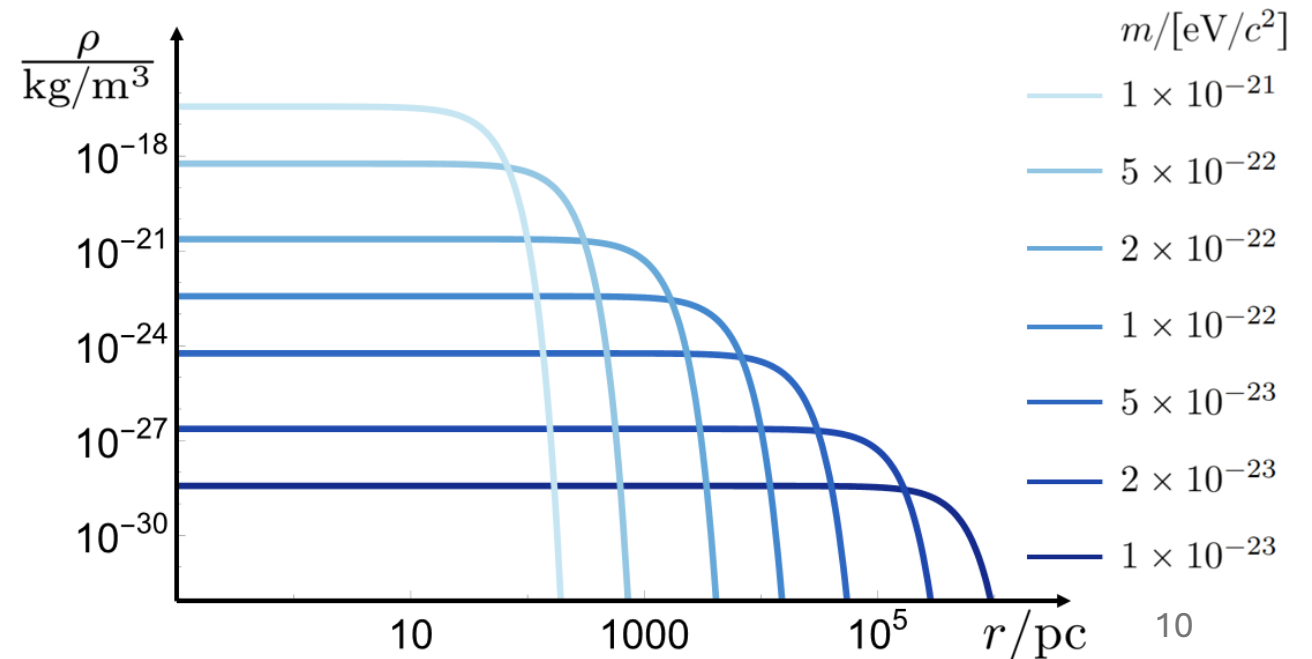
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*CDM limit: nothing would counteract gravity and we would get $R \rightarrow 0$

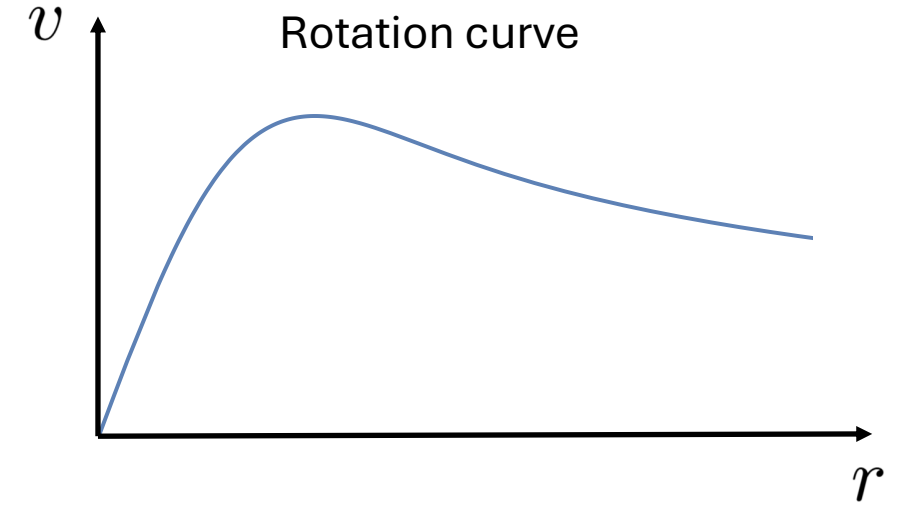
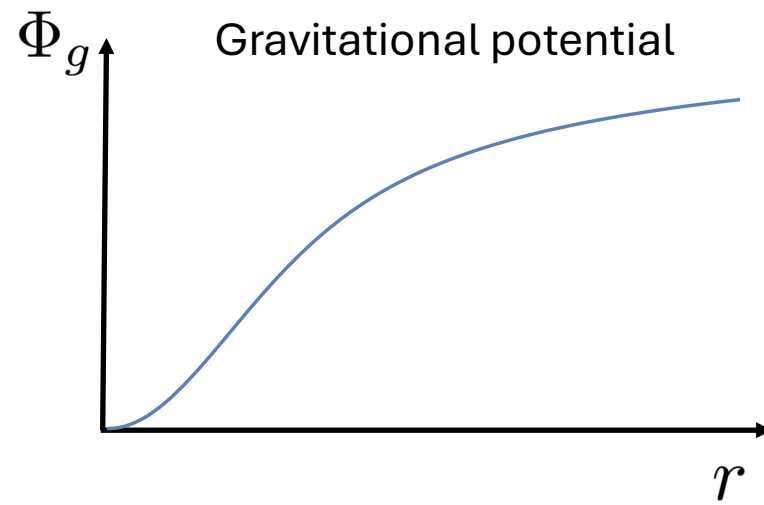
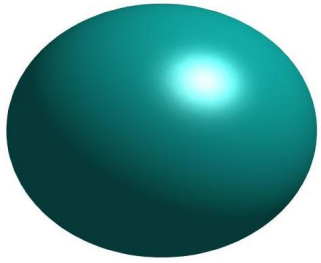
Equilibrium stationary solution: $\dot{R} = 0, \beta = 0$

$$R = \frac{3\sqrt{\pi}}{\sqrt{2}} \frac{\hbar^2}{GMm^2}$$



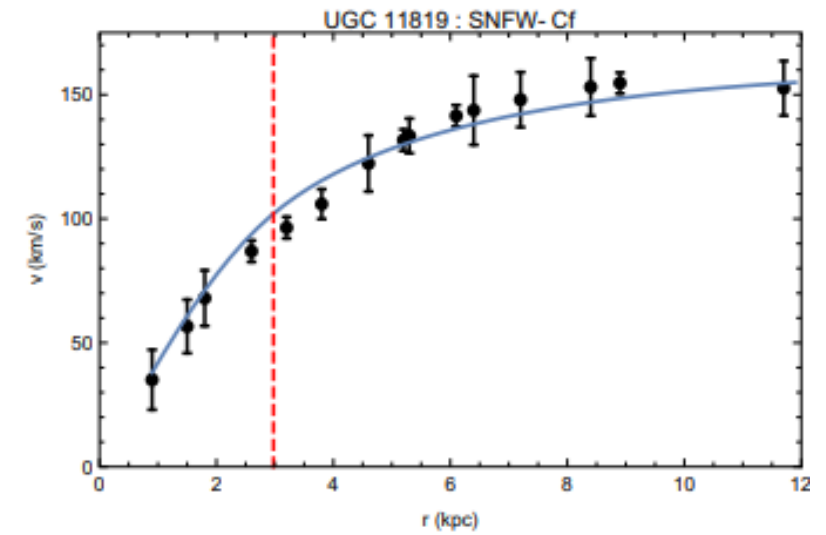
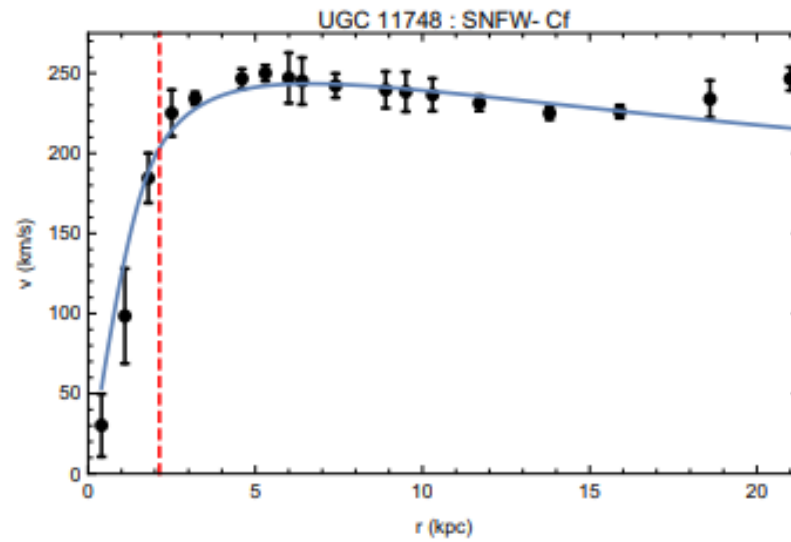
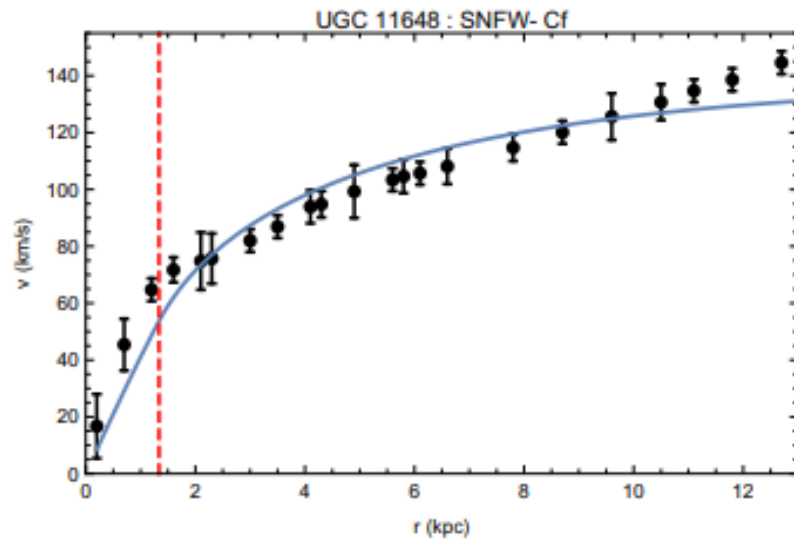
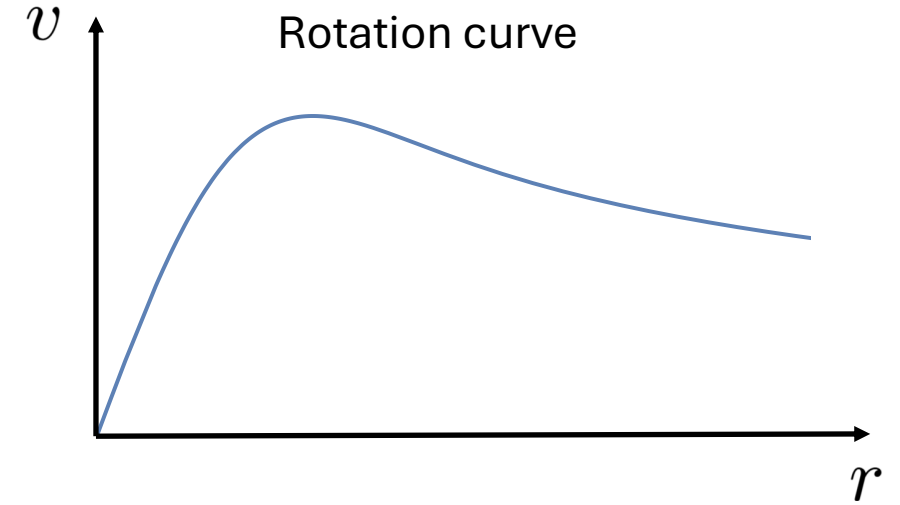
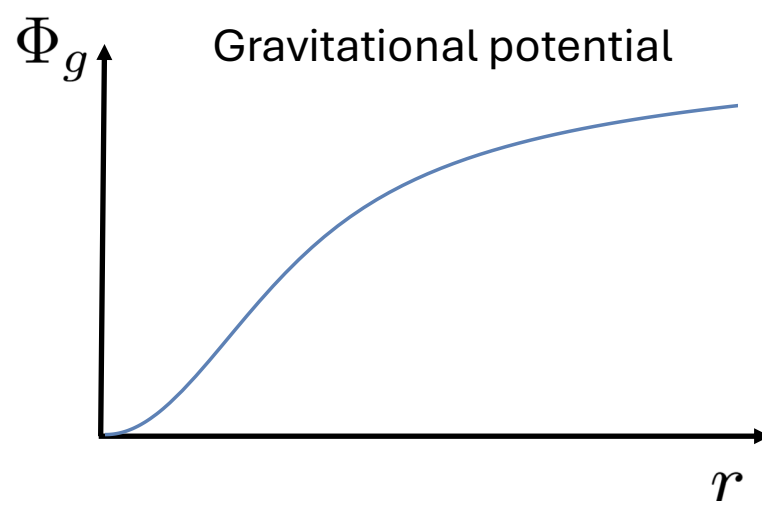
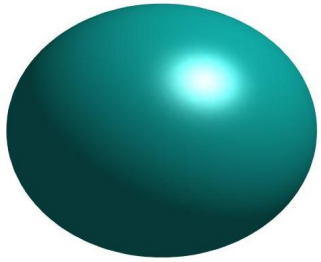
Rotation curves

ULDM soliton



Rotation curves

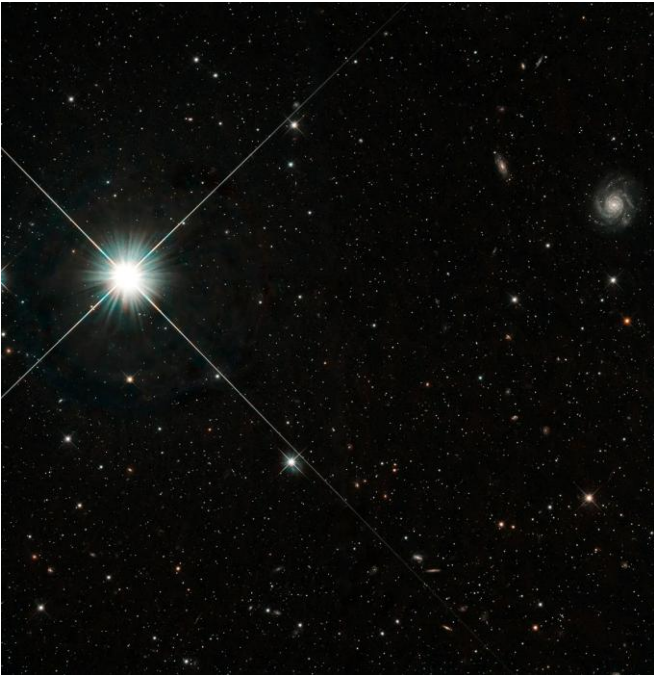
ULDM soliton



Fit of observed rotation curves for LSB galaxies.

Velocity dispersion in dwarf galaxies

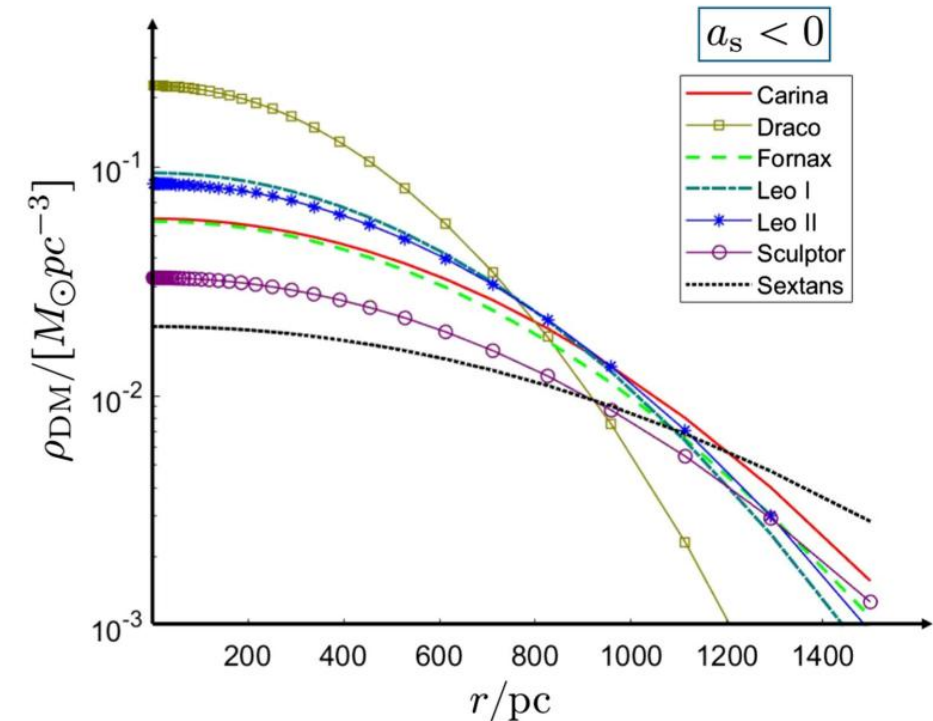
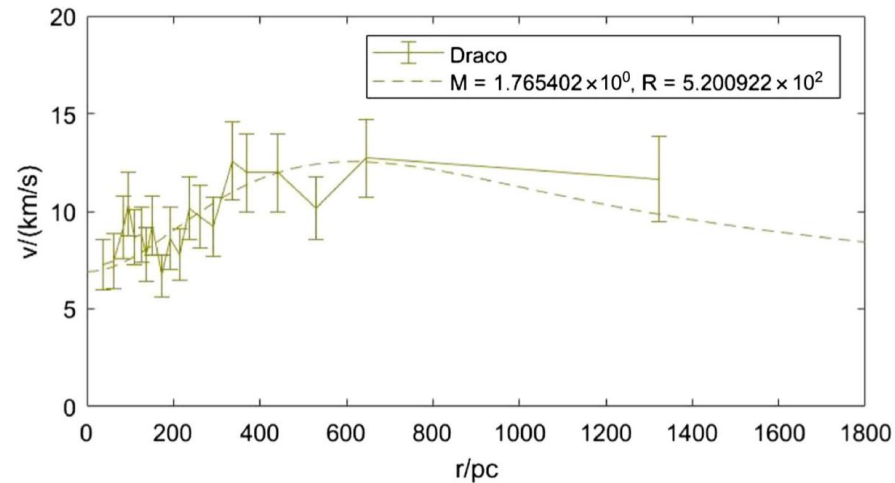
Idea: fit velocity dispersion data for seven Milky Way dwarf spheroidals with ULDM halo model



Hubble Space Telescope view showing part of the Draco Dwarf Galaxy. Credits: NASA, ESA, Eduardo Vitral, Roeland van der Marel, and Sangmo Tony Sohn (STScI), DSS; Image processing: Joseph DePasquale (STScI)

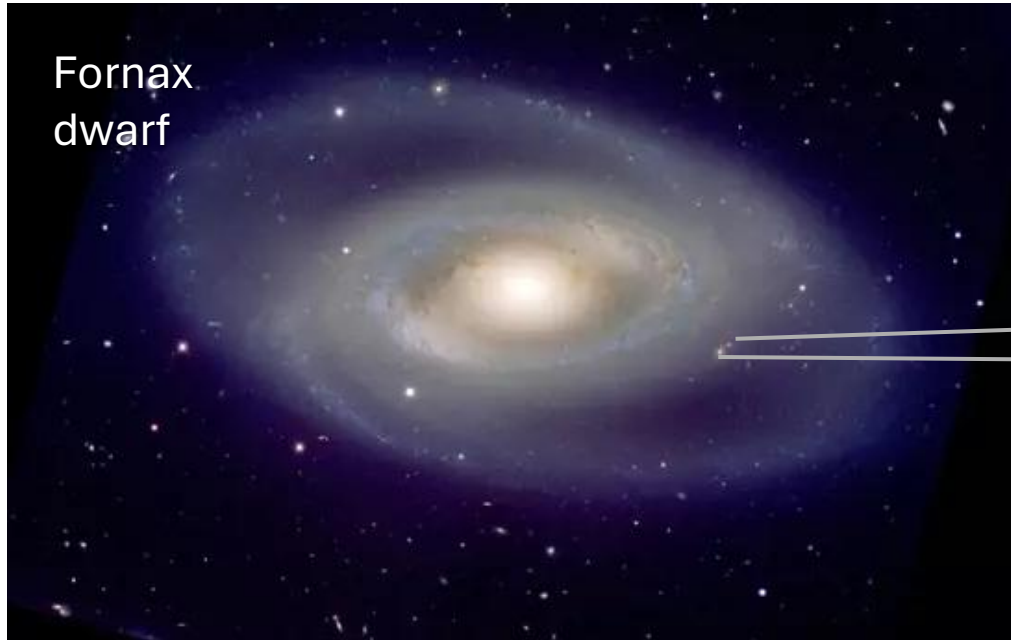
Velocity dispersion $\overline{v_r^2} = \frac{1}{\rho_b(r)} \int_r^\infty dr' \rho_b(r') \frac{d\Phi}{dr'}$

Gravitational potential, dominated by ULDM halo

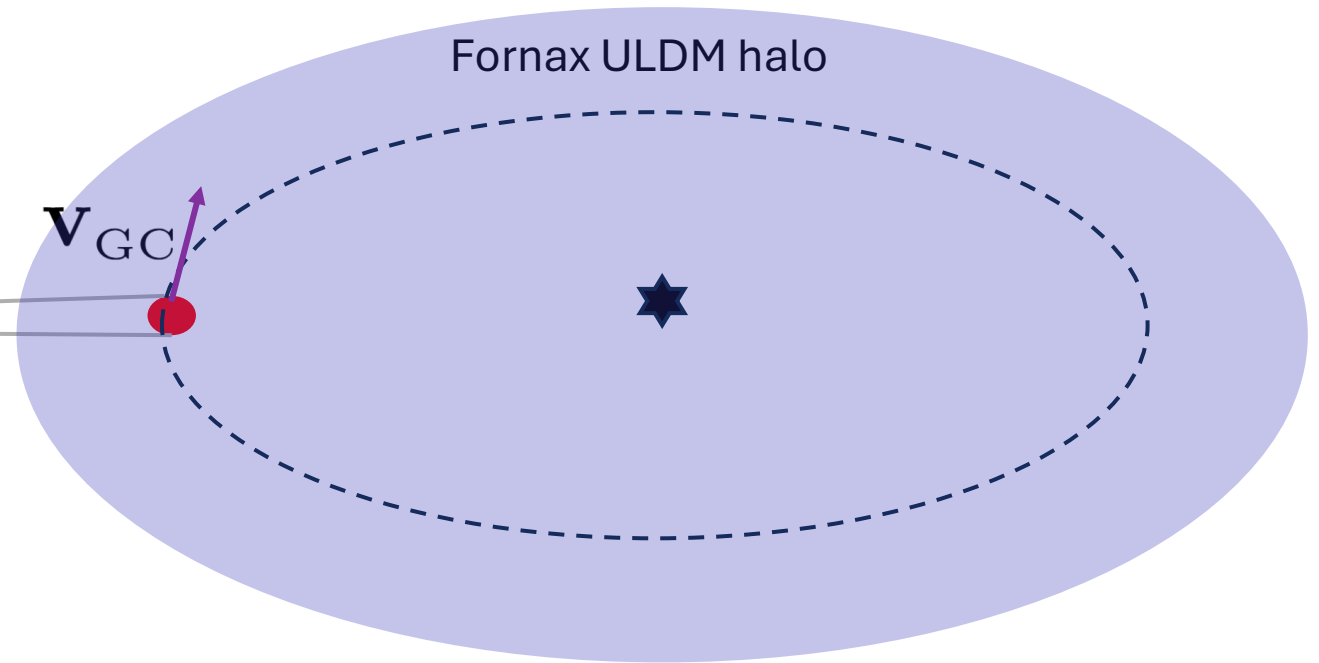


K Korshynska, EV Gorbar, YM Bidasyuk, AI Yakimenko, Y Revaz Phys Rev D (2026)

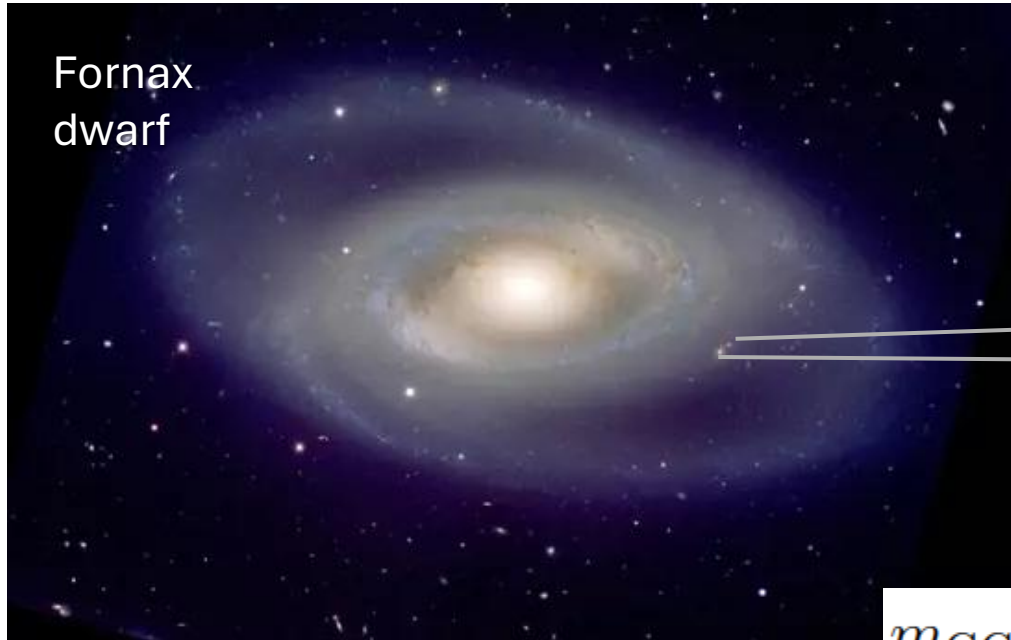
Trajectories of globular clusters: dynamical friction



<https://www.constellation-guide.com/constellation-list/fornax-constellation/>

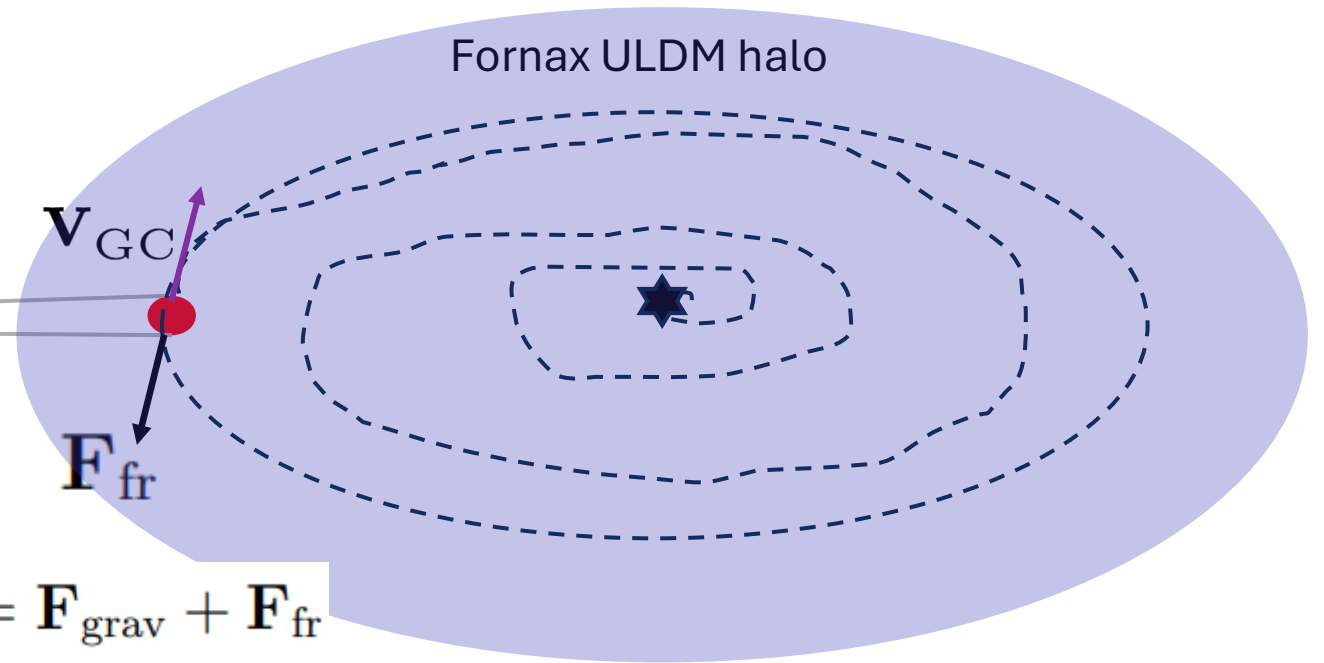


Trajectories of globular clusters: dynamical friction



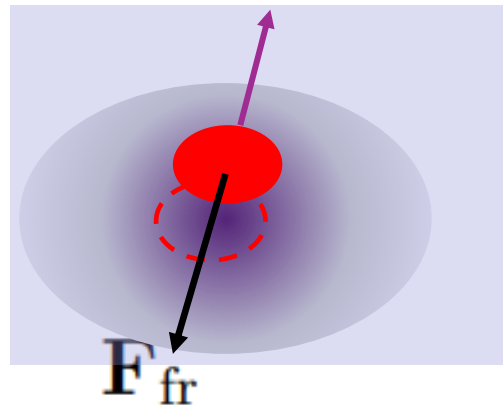
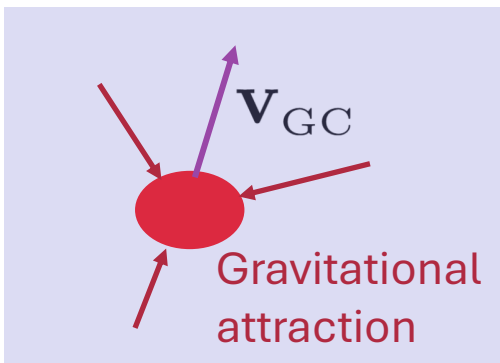
Fornax
dwarf

<https://www.constellation-guide.com/constellation-list/fornax-constellation/>

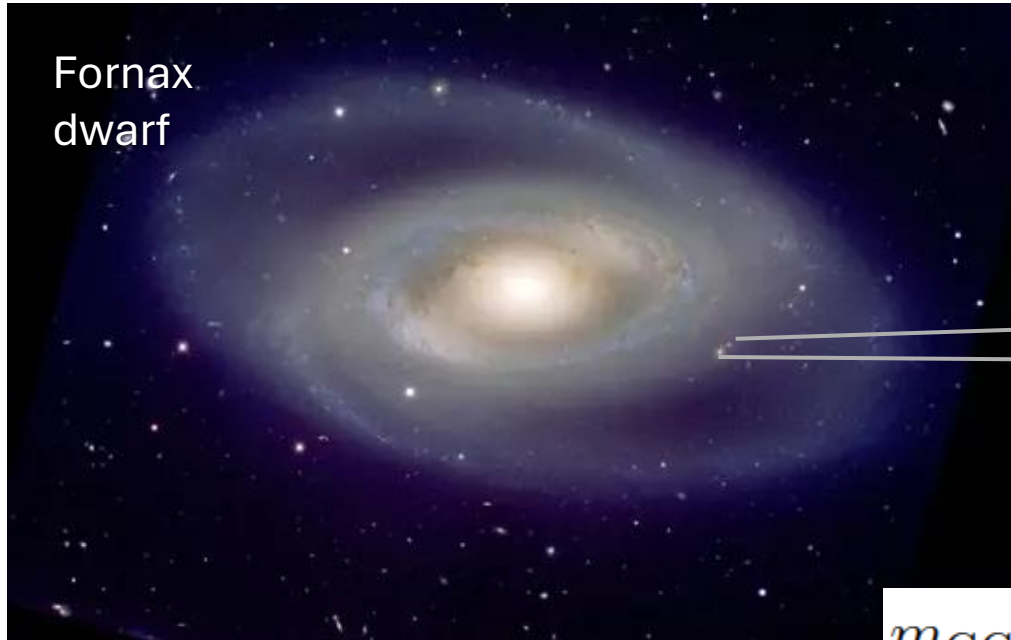


$$m_{GC}\ddot{\mathbf{r}} = \mathbf{F}_{grav} + \mathbf{F}_{fr}$$

For the supersonic velocity of a globular cluster

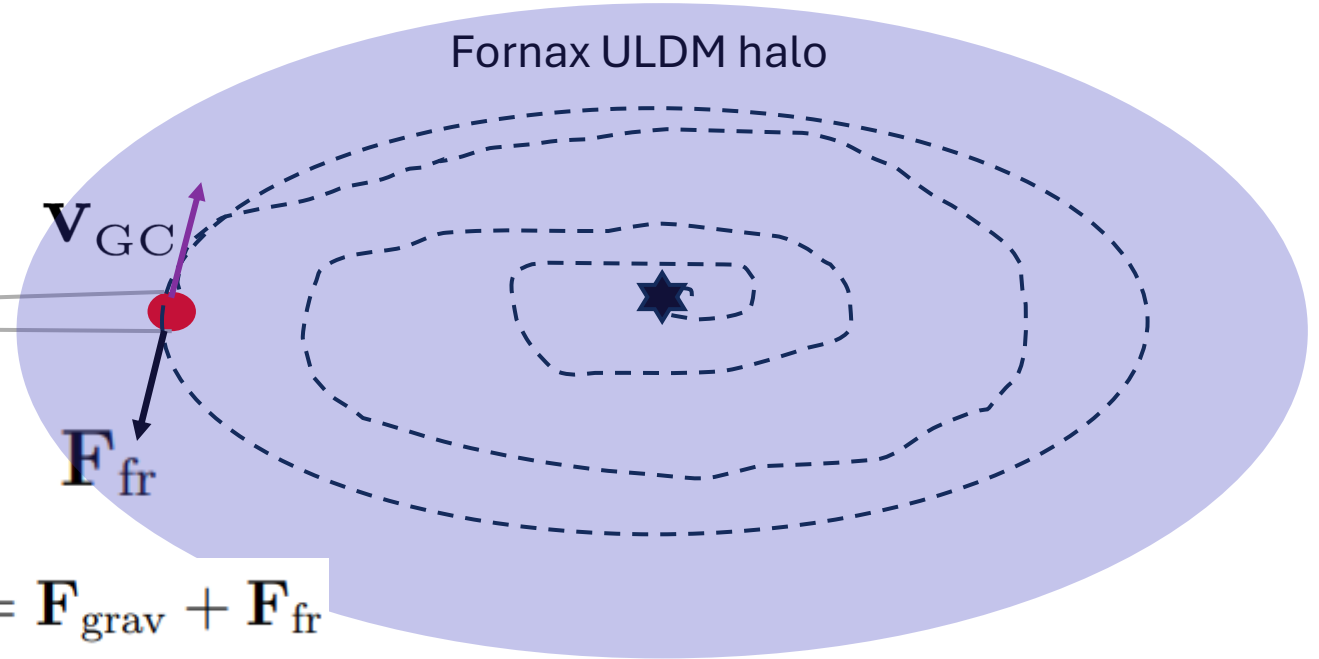


Trajectories of globular clusters: dynamical friction

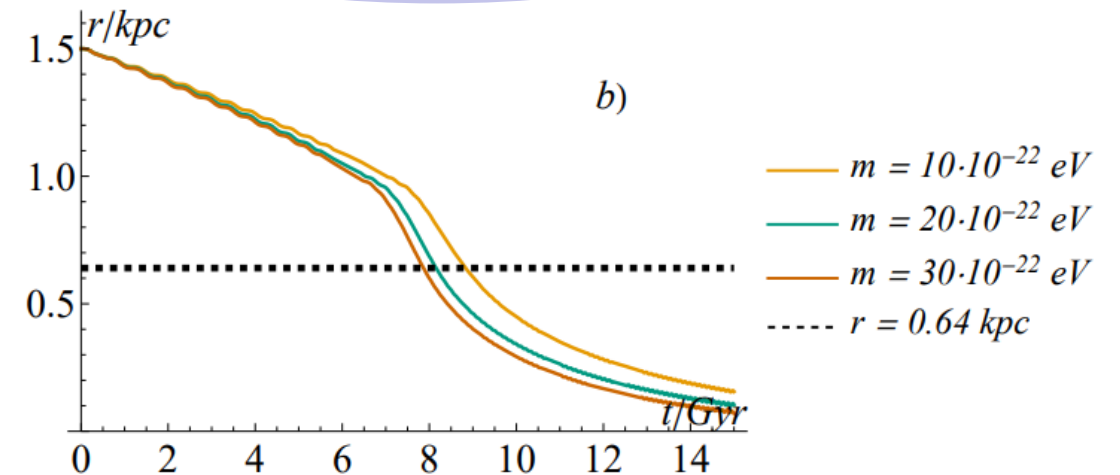
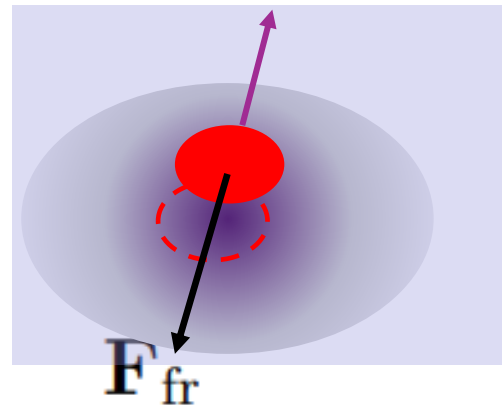
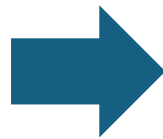
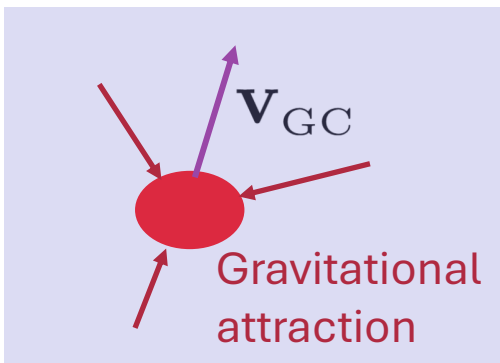


Fornax dwarf

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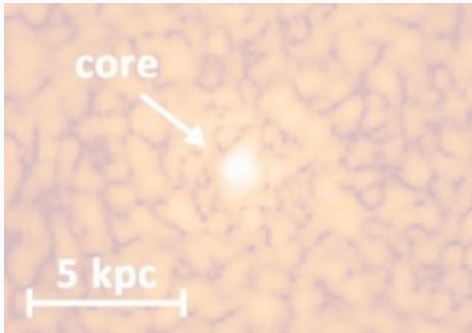
For the supersonic velocity of a globular cluster



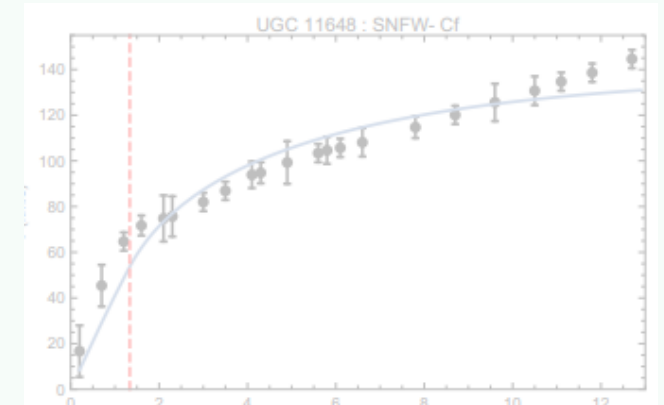
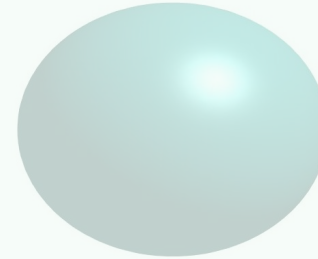
V.M. Gorkavenko, A.I. Yakimenko, A.O. Zaporozhchenko, E.V. Gorbar, *Physica Scripta* (2025)
E Gorbar, O Barabash, V Gorkavenko, K Korshynska, Andriy Momot, A Zaporozhchenko
Classical and Quantum Gravity (2026)

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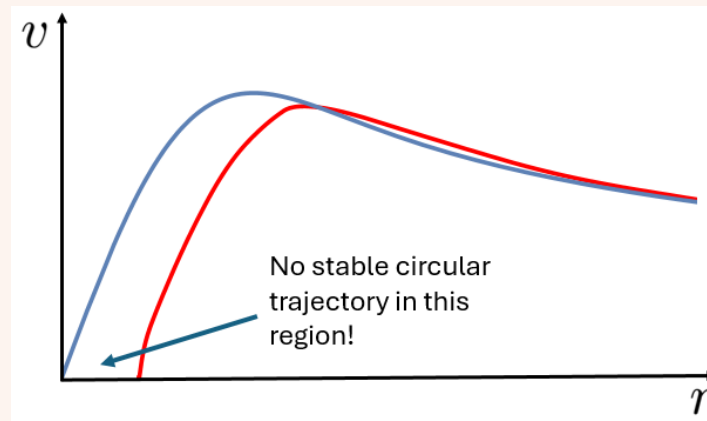
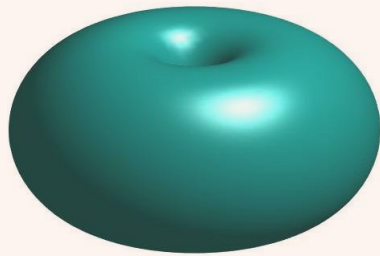
1. Ultralight dark matter (ULDM): introduction



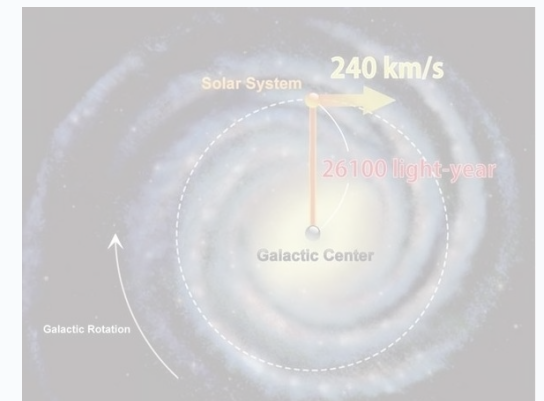
2. ULDM soliton and its astrophysical implications



3. ULDM vortex and its astrophysical implications



4. Conclusions



ULDM vortices

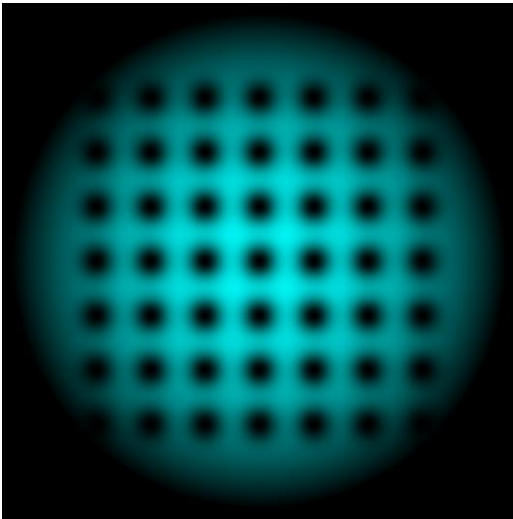
Could be formed due to

ULDM co-rotates with baryonic matter
in spiral galaxies

$$\omega \gg \omega_{cr}$$

Baryon
rotational
velocity

Critical velocity
for vortex
formation



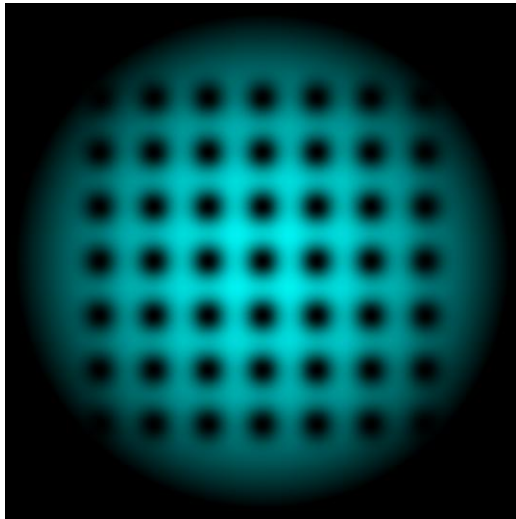
*M. P. Silverman and R. L. Mallett, General
Relativity and Gravitation (2002)*

ULDM vortices

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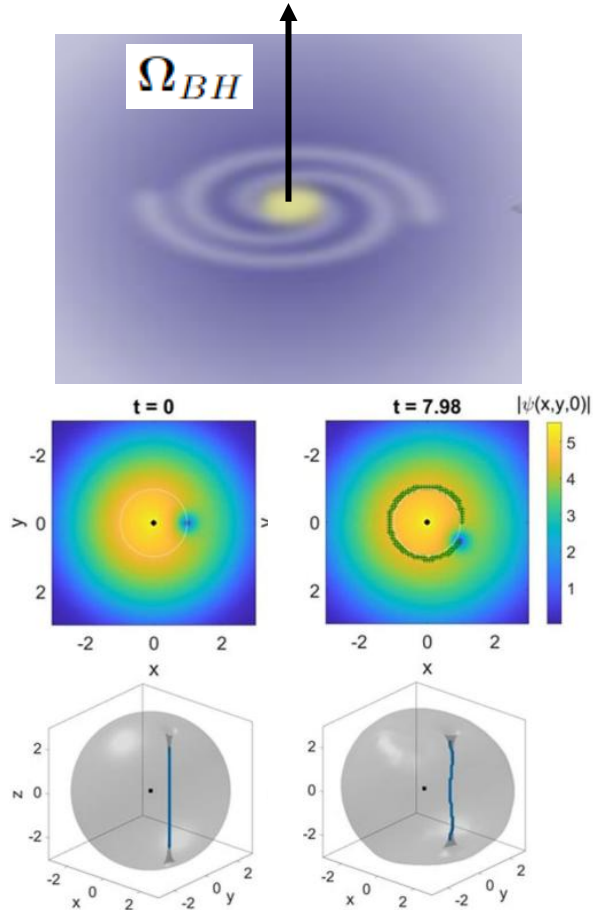
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M. P. Silverman and R. L. Mallett, General Relativity and Gravitation (2002)

Rotating black hole



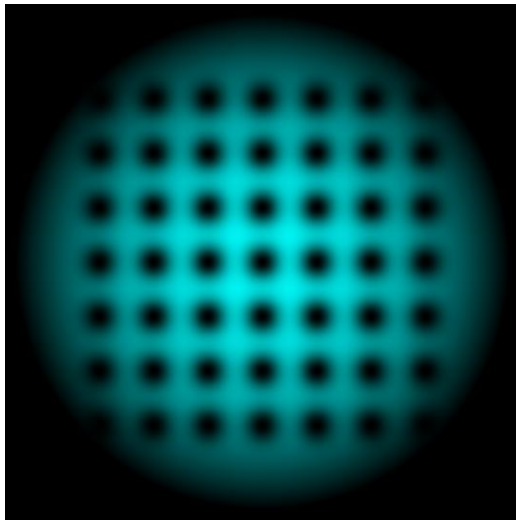
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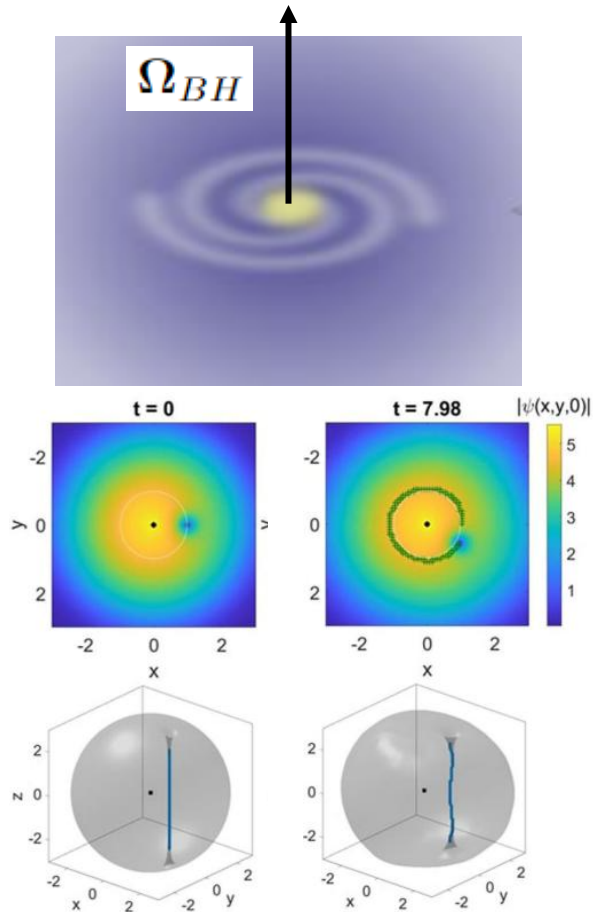
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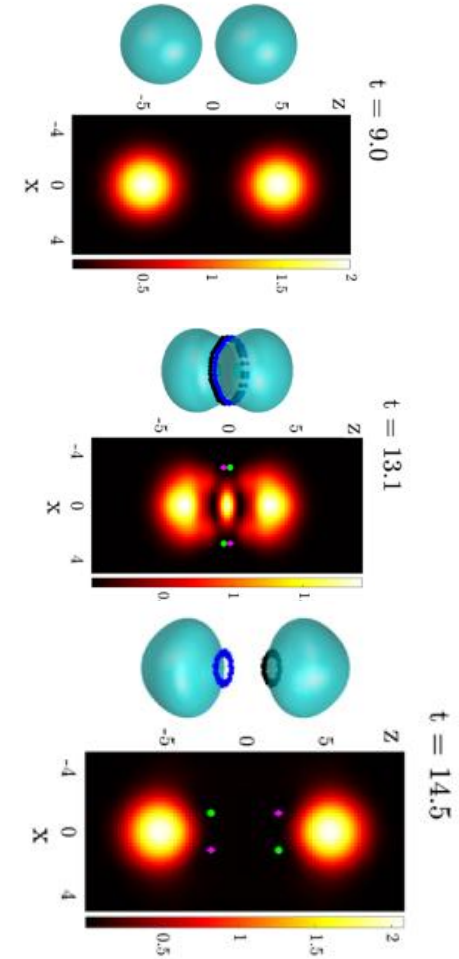
M. P. Silverman and R. L. Mallett, General Relativity and Gravitation (2002)

Rotating black hole



K Korshynska, OO Prykhodko, EV Gorbar, Junji Jia, AI Yakimenko PhysRev D (2025)

Collisions of halos



YO Nikolaieva, YM Bidasyuk, K Korshynska, EV Gorbar, Junji Jia, AI Yakimenko Phys Rev D (2023)

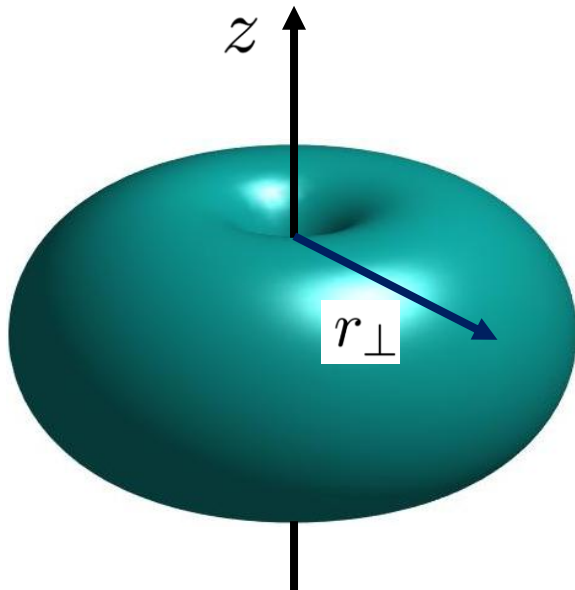
ULDM vortices

GPP equations

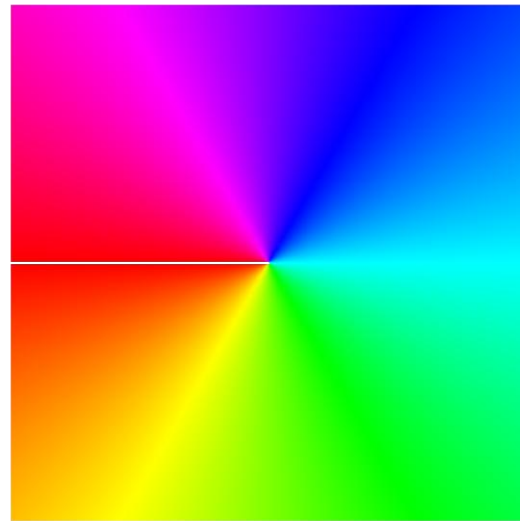
$$\left\{ \begin{aligned} i\hbar\partial_t\psi(\mathbf{r}, t) &= \left(-\frac{\hbar^2}{2m}\Delta + m\Phi(\mathbf{r}, t) + gN|\psi(\mathbf{r}, t)|^2 \right) \psi(\mathbf{r}, t) \\ \Delta\Phi(\mathbf{r}, t) &= 4\pi GmN|\psi(\mathbf{r}, t)|^2 \end{aligned} \right.$$

... have another type of stationary solution – vortices!

$$\psi(r_{\perp}, \theta, z) = A \left(\frac{r_{\perp}}{R} \right)^s e^{-\frac{r_{\perp}^2}{2R^2} - \frac{z^2}{2(R\eta)^2} + is\theta}$$



density



phase (s=1)

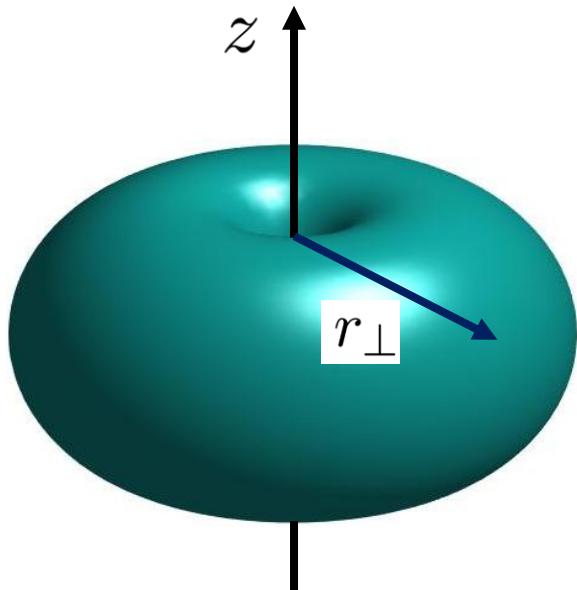
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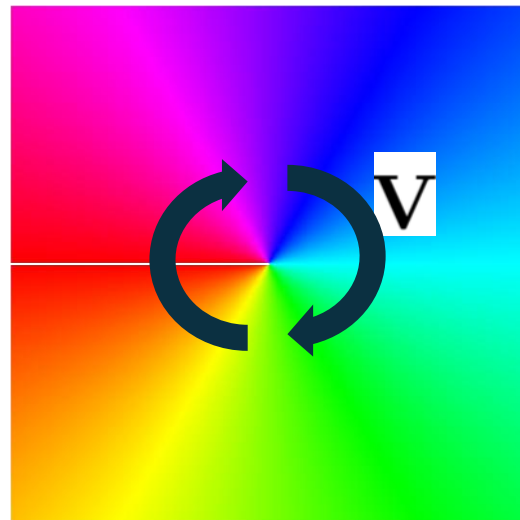
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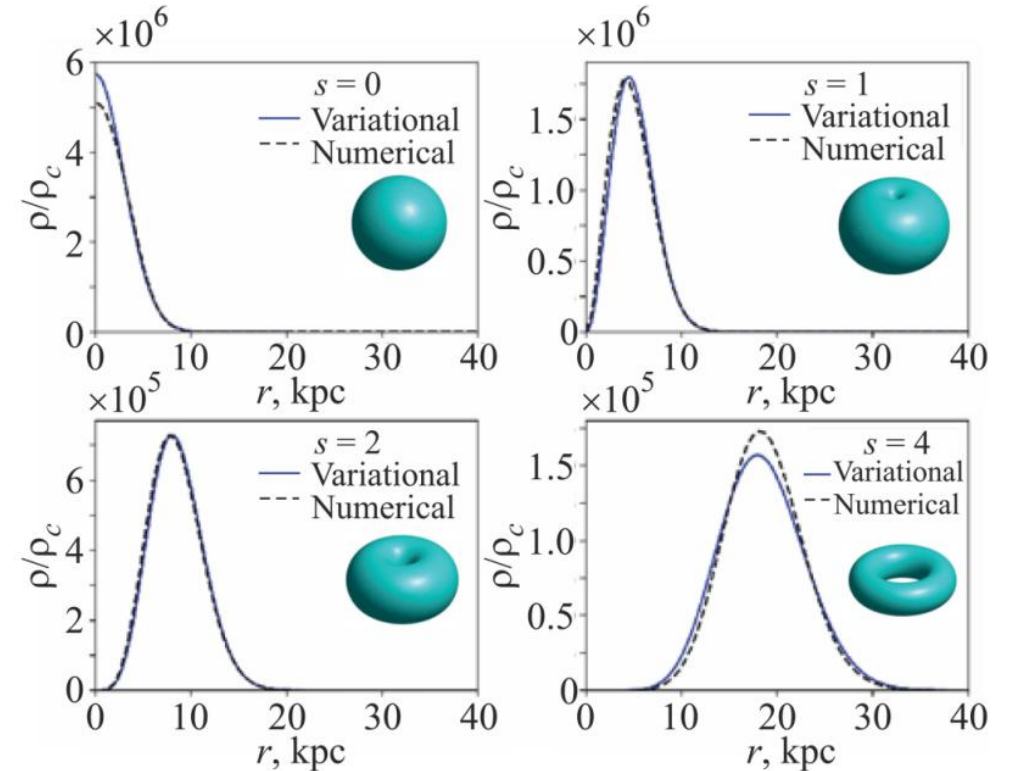
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Y. O. Nikolaieva, A. O. Olashyn, Y. I. Kuriatnikov, S. I. Vilchynskii, A. I. Yakimenko et al. *Low Temp Physics* (2021)

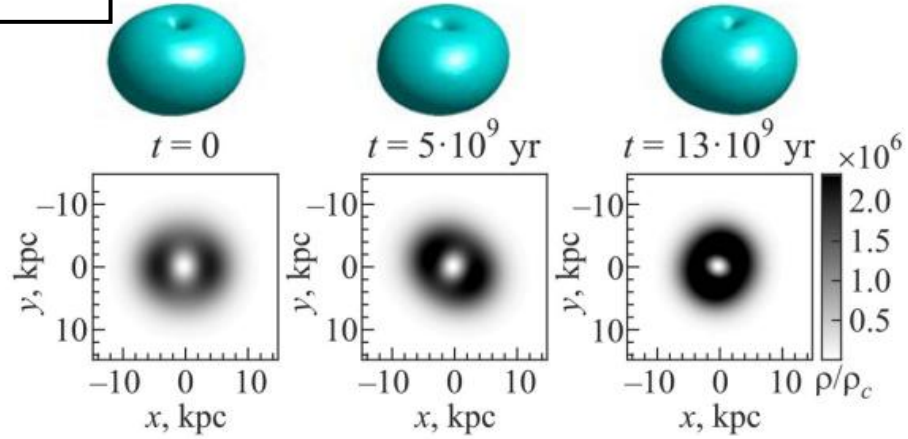
Vortices have a superfluid flow

$$\mathbf{V} = \frac{\mathbf{j}}{|\psi|^2} = -\frac{i\hbar}{2m} \frac{\psi^* \nabla \psi - \psi \nabla \psi^*}{|\psi|^2} = \frac{s\hbar}{m} \frac{\mathbf{e}_\phi}{r_\perp}$$

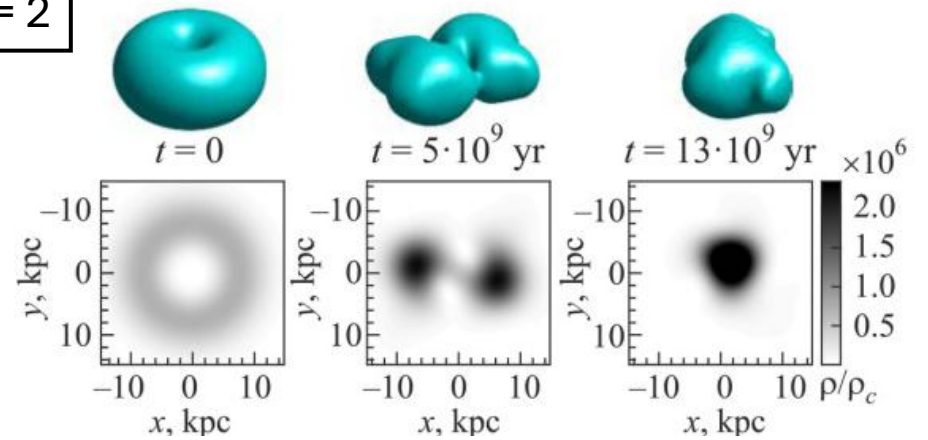
ULDM vortices

Are stable with respect to perturbations:

$S = 1$



$S = 2$

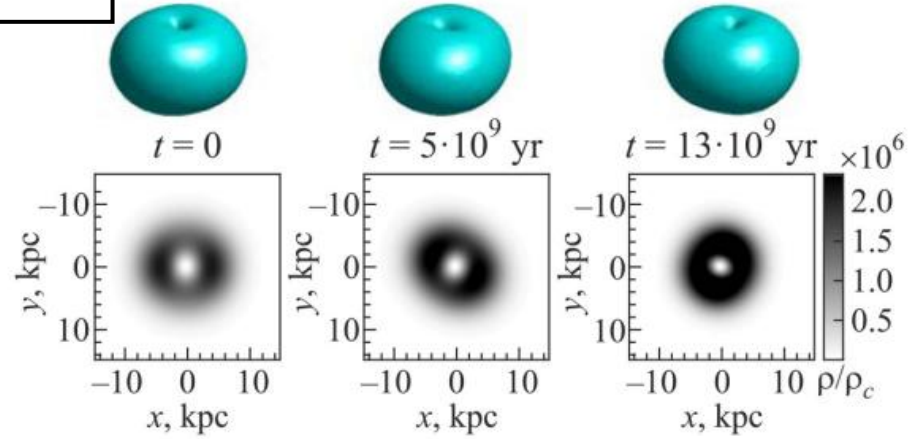


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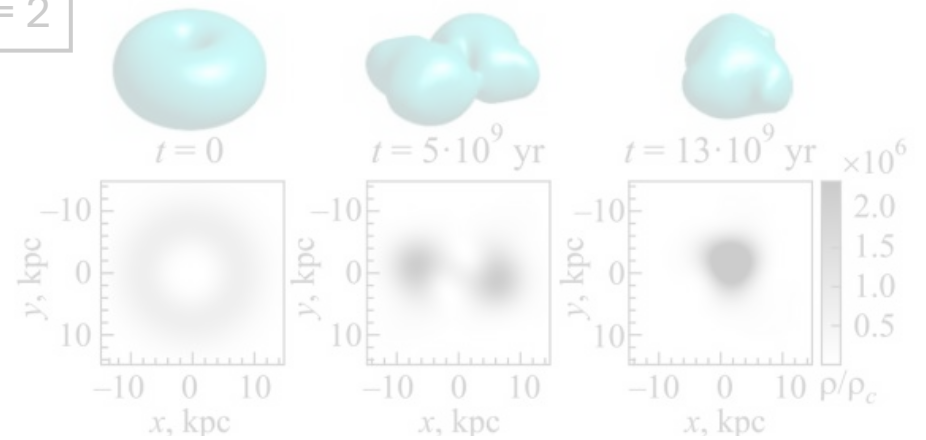
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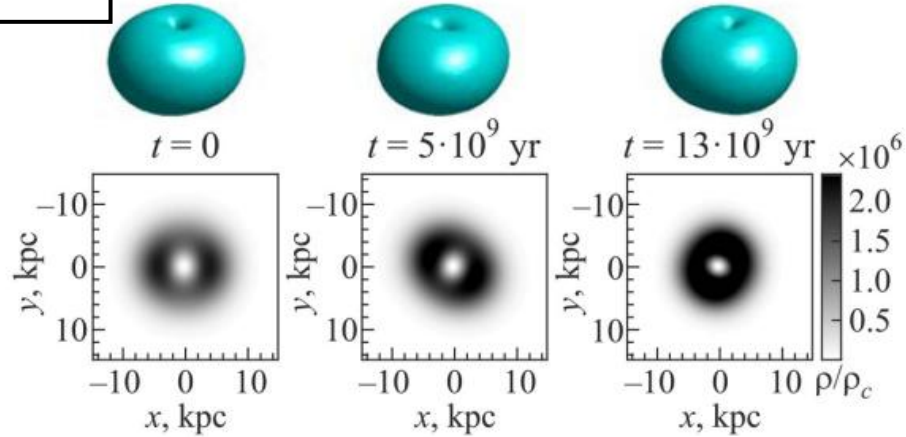


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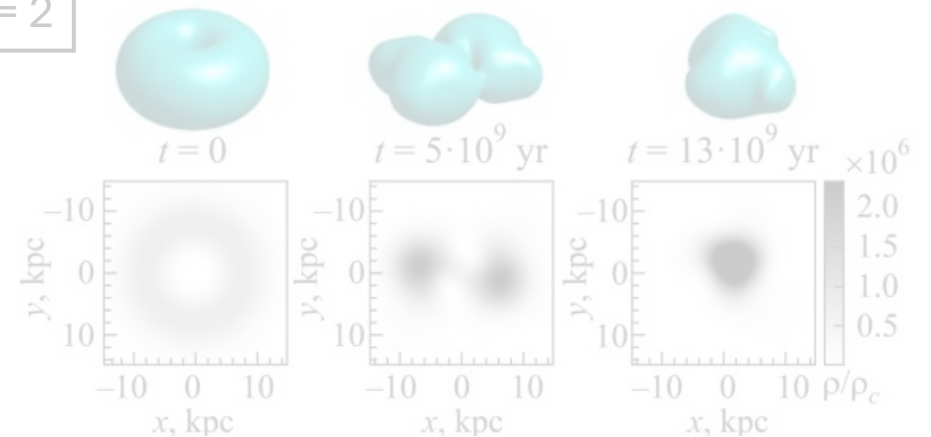
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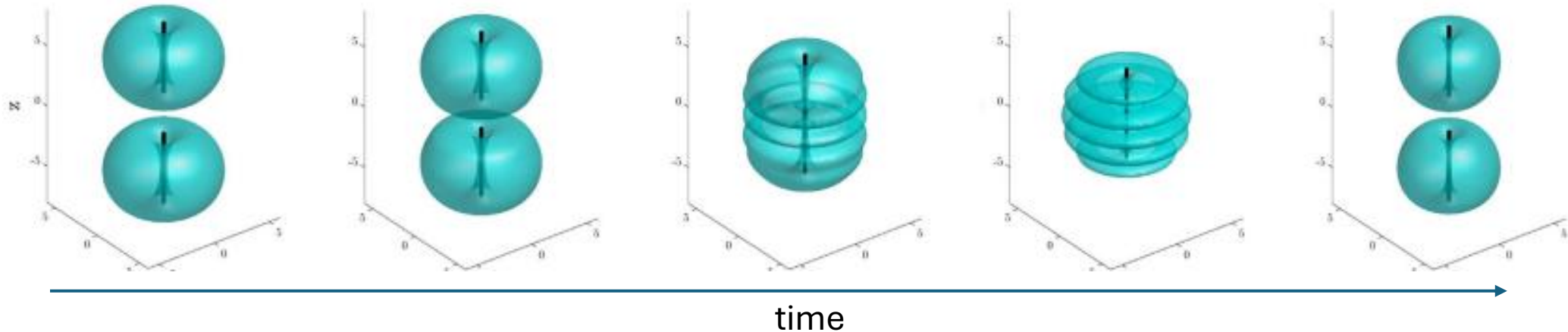


$S = 2$



Y. O. Nikolaieva, A. O. Olashyn, Y. I. Kuriatnikov, S. I. Vilchynskii, A. I. Yakimenko Low Temp Physics (2021)

Are stable with respect to collisions:



YO Nikolaieva, YM Bidasyuk, K Korshynska, EV Gorbar, Junji Jia, AI Yakimenko Phys Rev D (2023)

ULDM vortices: gravielectric potential

The presence of a vortex changes the spacetime metric:

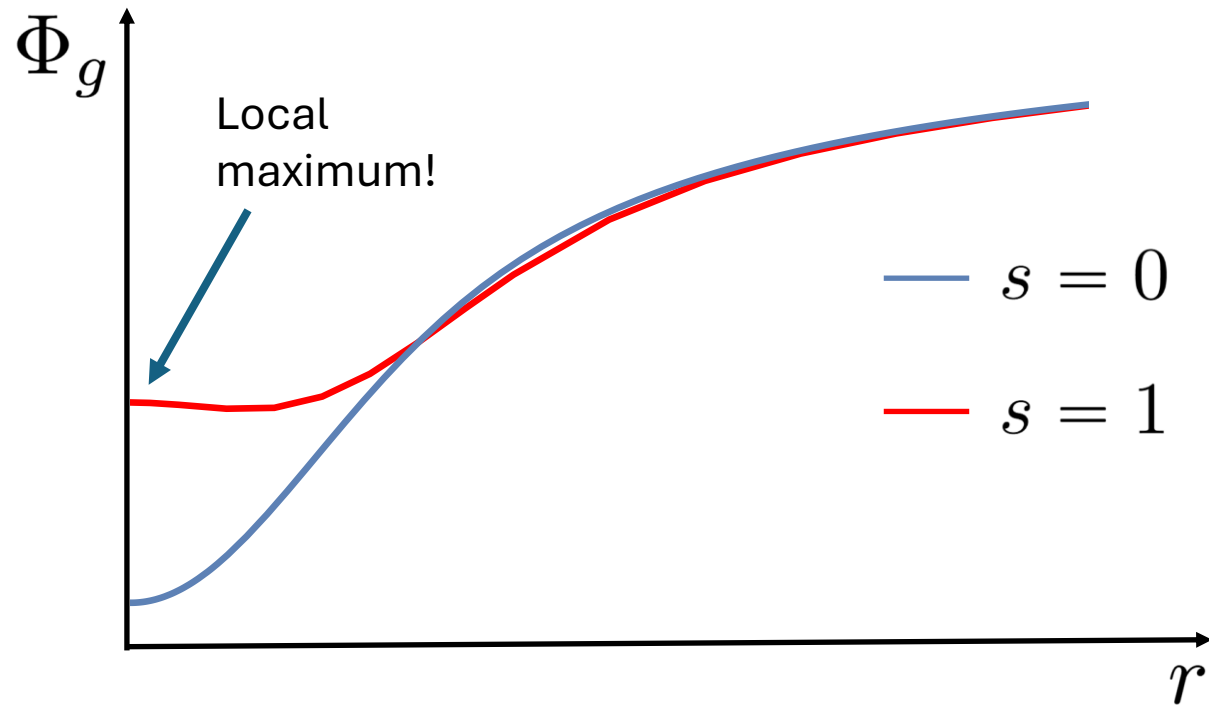
$$dS^2 = g_{\mu\nu} dx^\mu dx^\nu = \left(1 - \frac{2\Phi_g}{c^2}\right) (dx^0)^2 + \frac{4}{c^2} (\mathbf{A}_g \mathbf{d}\mathbf{x}) dx^0 + \left(-1 - \frac{2\Phi_g}{c^2}\right) \delta_{ij} dx^i dx^j$$

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Gravielectric (Newtonian) potential and rotation curves

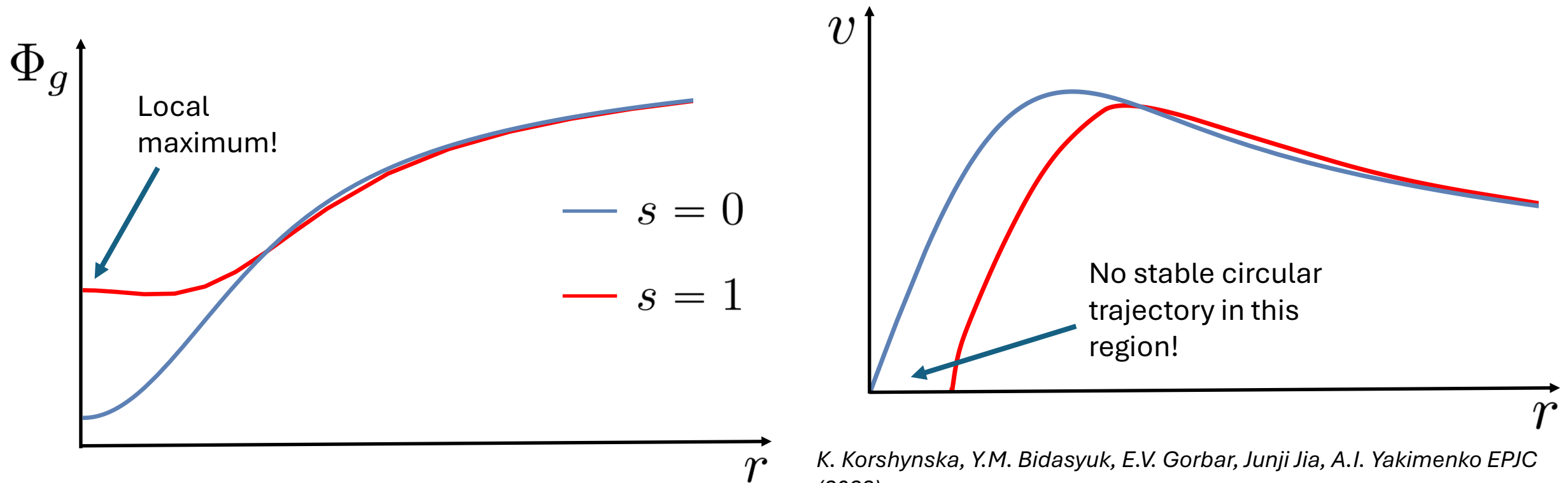


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Gravielectric (Newtonian) potential and rotation curves



K. Korshynska, Y.M. Bidasyuk, E.V. Gorbar, Junji Jia, A.I. Yakimenko EPJC (2023)

ULDM vortices: gravimagnetic effect

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The geodesic trajectory of a test particle is determined by a gravielectromagnetic force

$$\begin{aligned} \mathbf{F}_g &= m \left(\nabla \Phi - \frac{2}{c} \mathbf{v} \times [\nabla \times \mathbf{A}_g] \right) \\ &= m(\mathbf{a}_E + \mathbf{a}_B) \end{aligned}$$

ULDM vortices: gravimagnetic effect

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Gravimagnetic force: Lense-Thirring effect

The geodesic trajectory of a test particle is determined by a gravielectromagnetic force

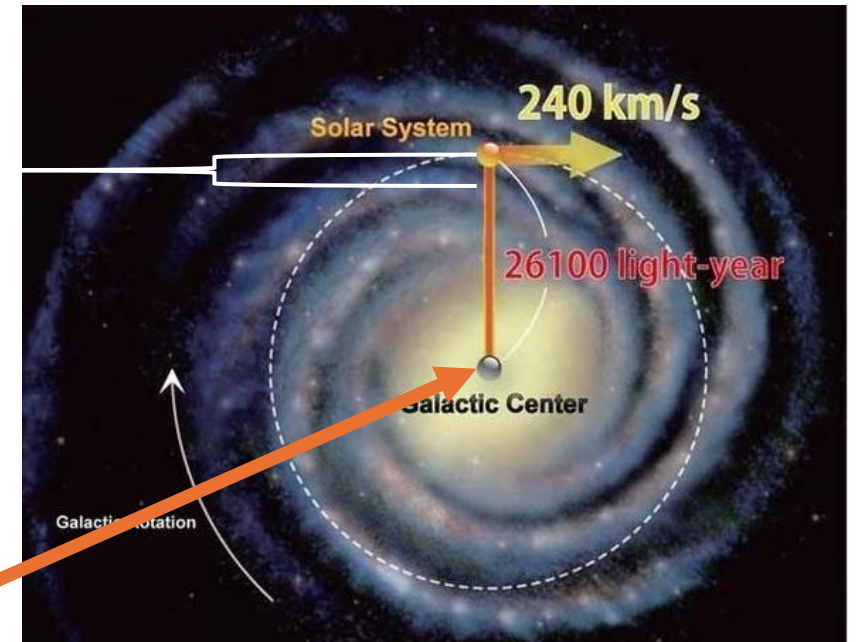
$$\mathbf{F}_g = m \left(\nabla \Phi - \frac{2}{c} \mathbf{v} \times [\nabla \times \mathbf{A}_g] \right) = m(\mathbf{a}_E + \mathbf{a}_B)$$

$$\frac{|\mathbf{a}_B|}{|\mathbf{a}_E|} \sim 10^{-6}$$

For a distant star (8 kpc from the MW center):

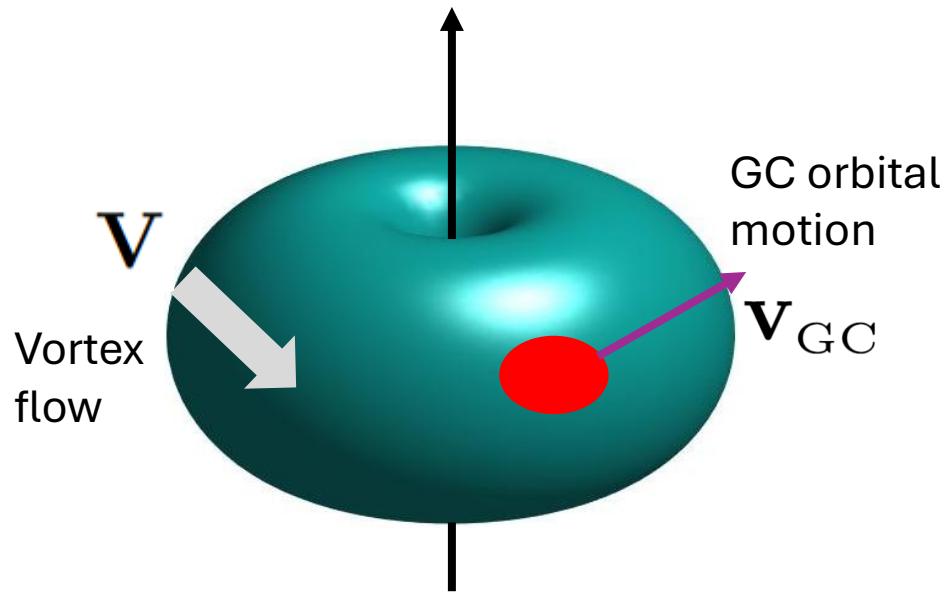
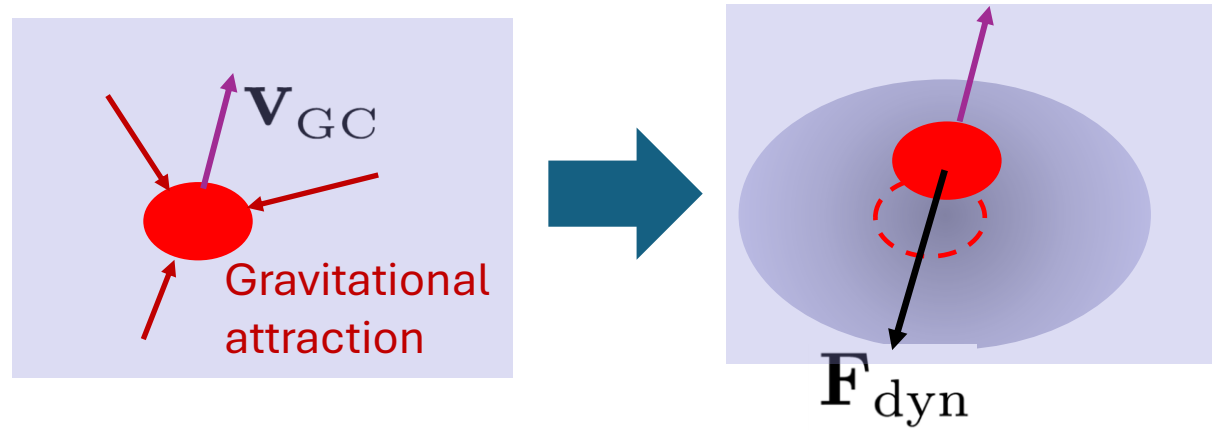
gravimagnetic force would displace the Sun trajectory by only 80 solar radii!

Here the vortex model could be tested better!

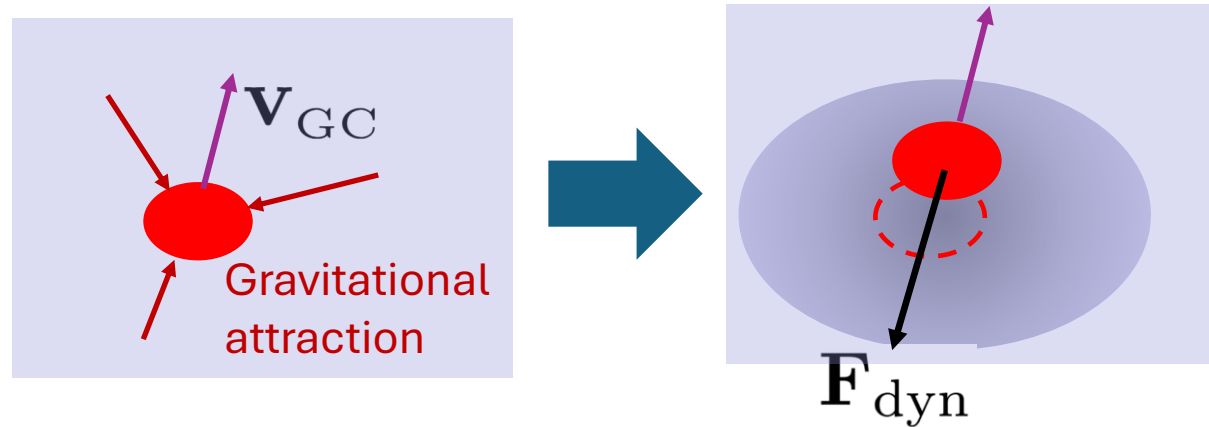


<https://cosmosatyourdoorstep.com/2023/01/18/how-massive-is-the-milky-way/>

ULDM vortices: dynamical friction

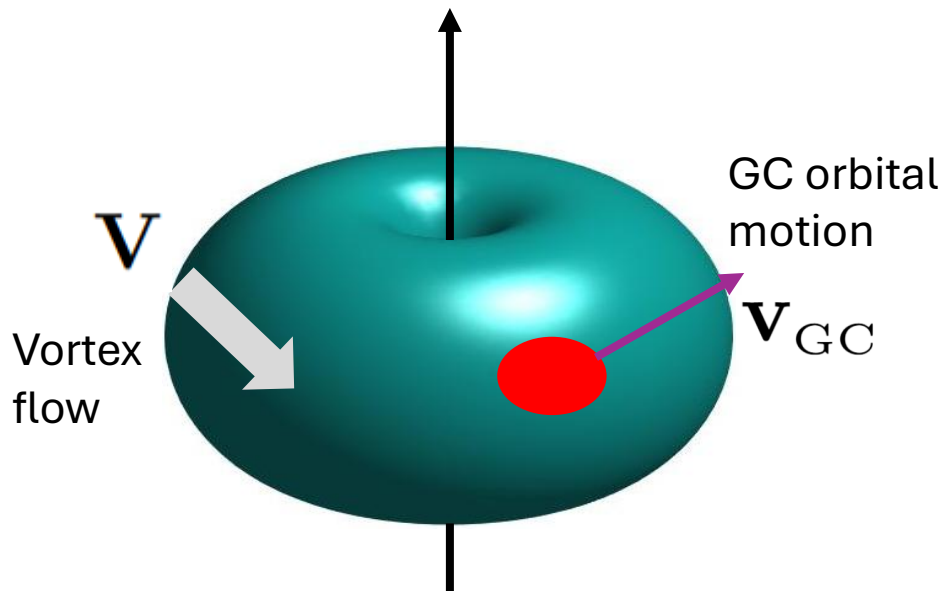


ULDM vortices: dynamical friction



Now relative velocity between the Globular Cluster and the ULDM medium is:

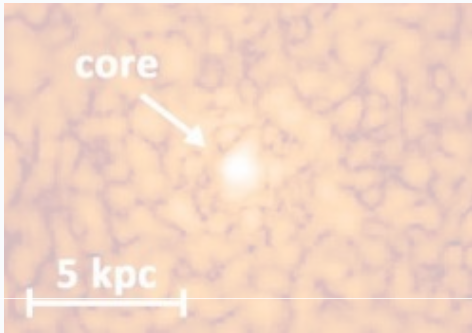
$$\mathbf{V}_{GC} - \mathbf{V}$$



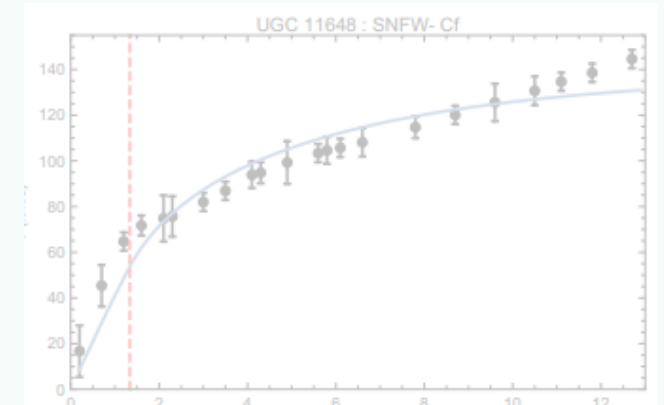
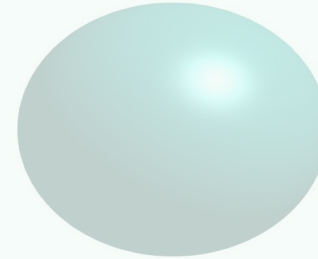
co-directional rotation leads to a pronounced suppression of dynamical friction

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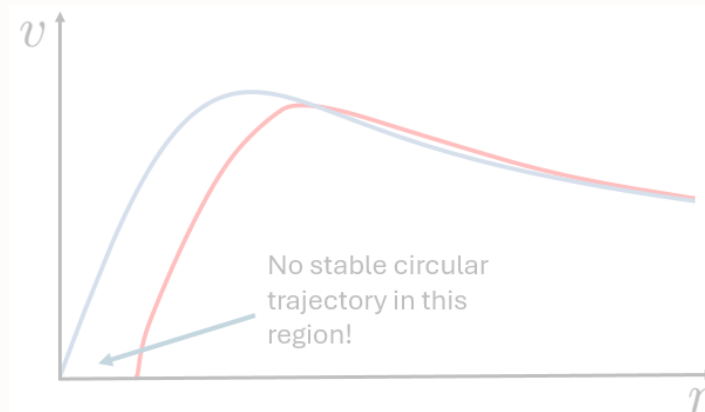
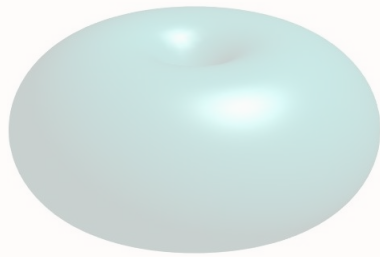
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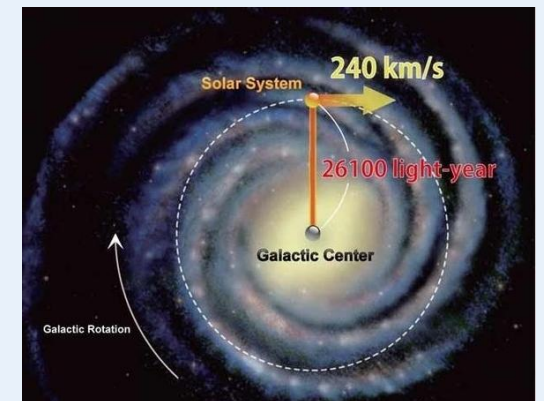
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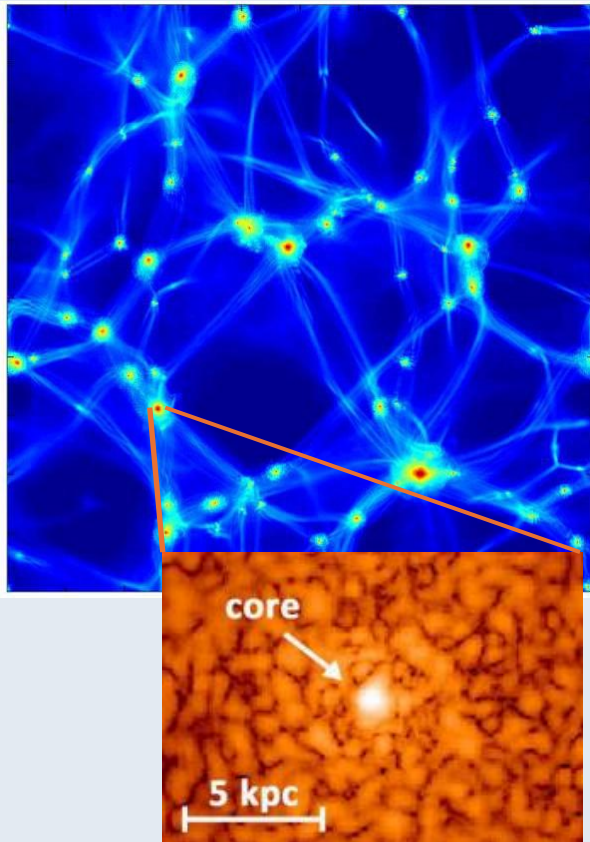
4. Conclusions



Conclusions

Ultralight dark matter model:

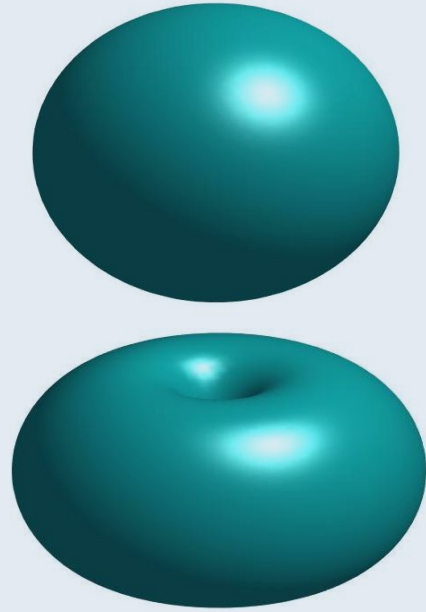
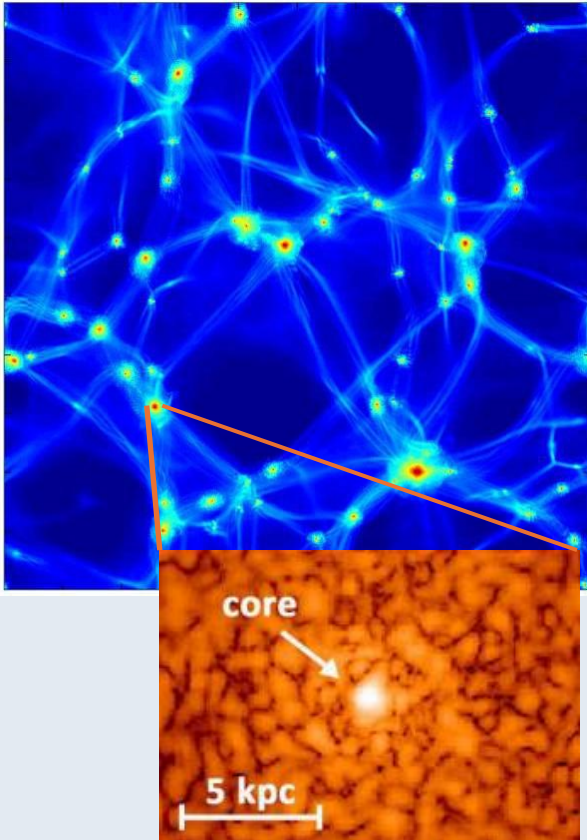
- axionlike bosons with small masses;
- cores in galactic centers



Conclusions

Ultralight dark matter model:

- axionlike bosons with small masses;
- cores in galactic centers



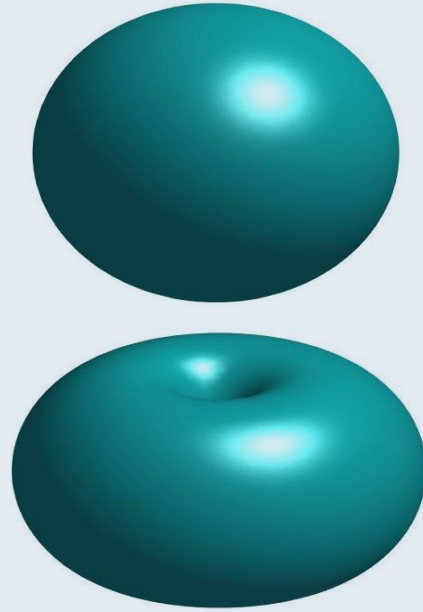
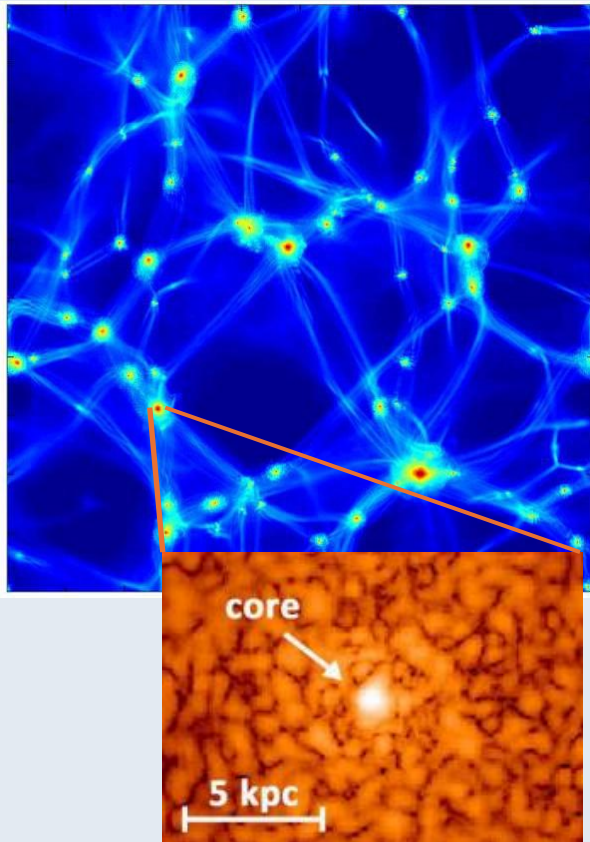
Balance between
gravitational attraction
and quantum pressure
-> cores in ULDM

They can also form
stable vortices

Conclusions

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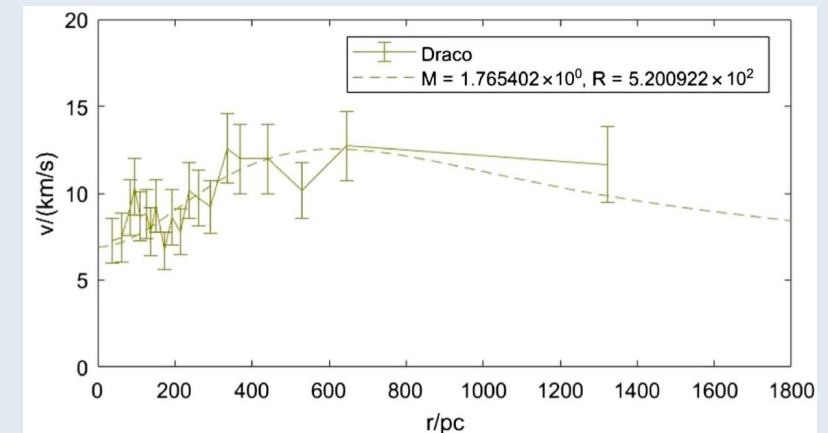
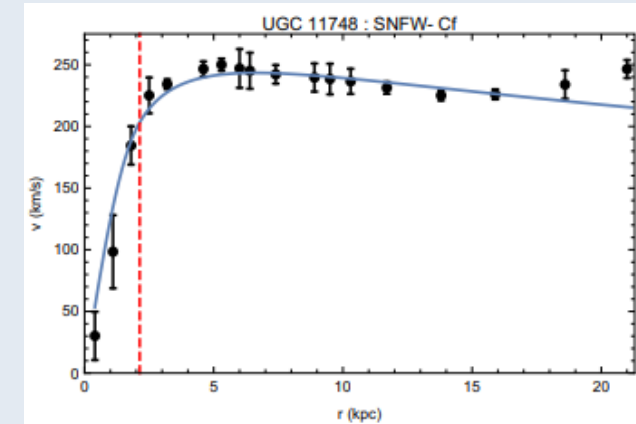
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Balance between gravitational attraction and quantum pressure
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They can also form stable vortices

ULDM model test: comparison with astrophysical observations
(rotation curves, velocity dispersion and globular clusters' trajectories)



I would like to acknowledge



Olexander
Yakymenko



Eduard
Gorbar



Yurii
Bidasyuk



Volodymyr
Gorkavenko

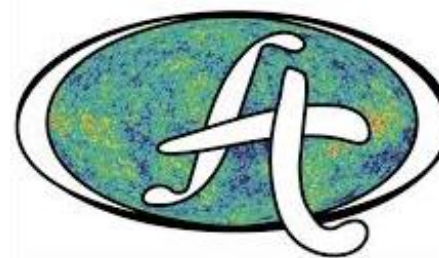


Oleh
Barabash



Yves
Revaz

...and many others for their collaboration.



QuantumFrontiers
Cluster of Excellence



I would like to acknowledge



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Yakymenko



Eduard
Gorbar



Yurii
Bidasyuk



Volodymyr
Gorkavenko



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Barabash



Yves
Revaz

...and many others for their collaboration.



EPFL

**And especially FPM
institute at PTB=)**

Literature

- „**Dynamical galactic effects induced by solitonic vortex structure in bosonic dark matter**“

K. Korshynska, Y.M. Bidasyuk, E.V. Gorbar, Junji Jia, A.I. Yakimenko EPJC (2023)

- „**Stable vortex structures in colliding self-gravitating Bose-Einstein condensates**“

YO Nikolaieva, YM Bidasyuk, K Korshynska, EV Gorbar, Junji Jia, AI Yakimenko Phys Rev D (2023)

- „**Vortex lines in ultralight bosonic dark matter around rotating supermassive black holes**“

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- „**Velocity dispersion profiles of dwarf spheroidal galaxies with self-interacting ultralight dark matter**“

K Korshynska, EV Gorbar, YM Bidasyuk, AI Yakimenko, Y Revaz Phys Rev D (2026)

- „**Damping of dynamical friction force in self-interacting ultralight dark matter and Fornax timing problem**“

E Gorbar, O Barabash, V Gorkavenko, K Korshynska, Andriy Momot, A Zaporozhchenko Classical and Quantum Gravity (2026)

- „**The Vortex State of Ultralight Dark Matter and the Fornax Timing Problem**“

V Gorkavenko, O Barabash, T Gorkavenko, K Korshynska, O Teslyk, A Zaporozhchenko, E Gorbar, Universe (2026)

- „**Cosmological structure formation in ultralight dark matter**“

K. Korshynska, A. I. Yakimenko, Y. Revaz and E. V. Gorbar [in preparation]

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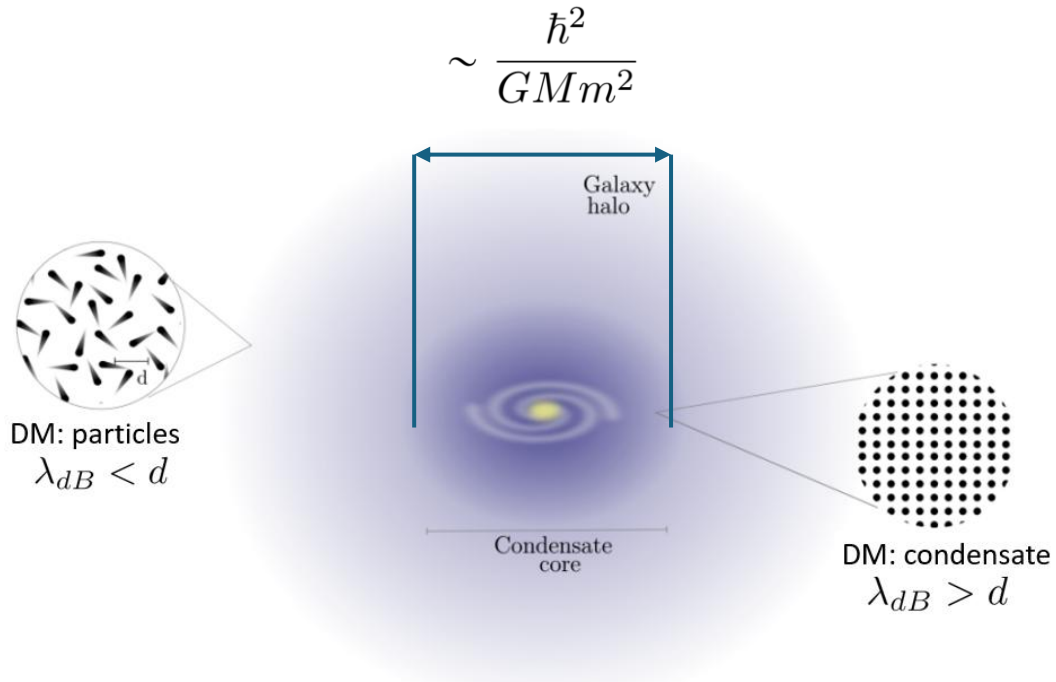
- **„Cosmological structure formation in ultralight dark matter“**

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Thank you for your attention!

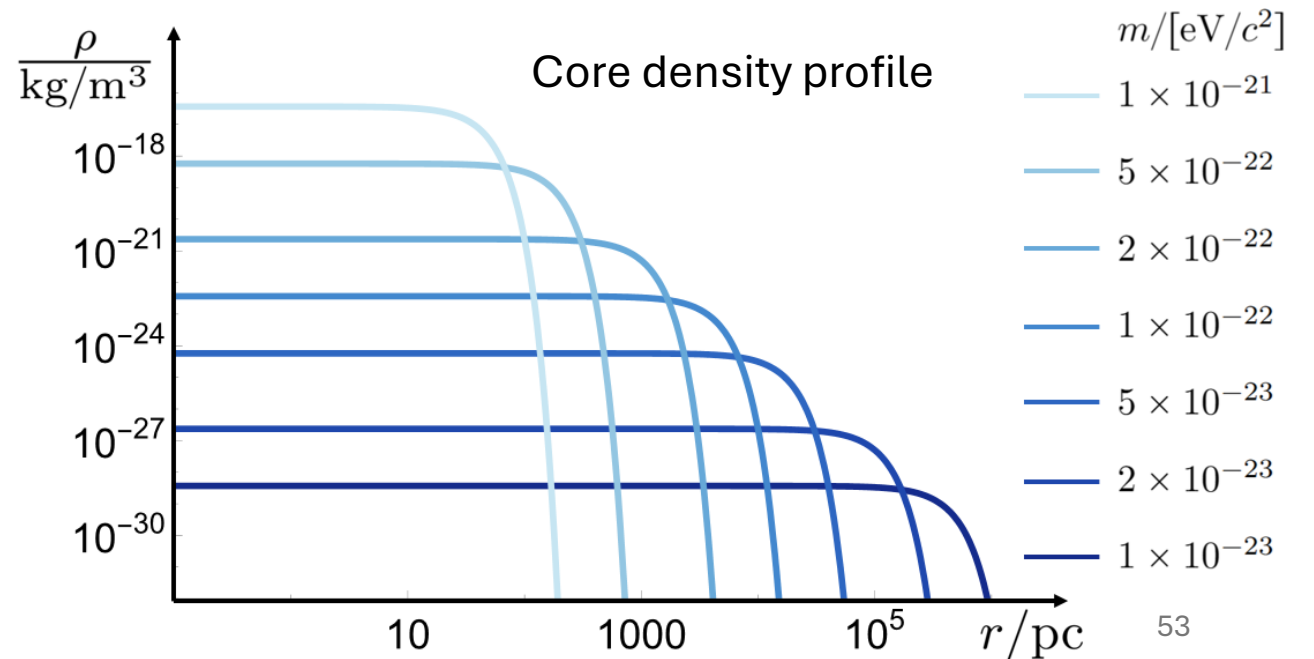
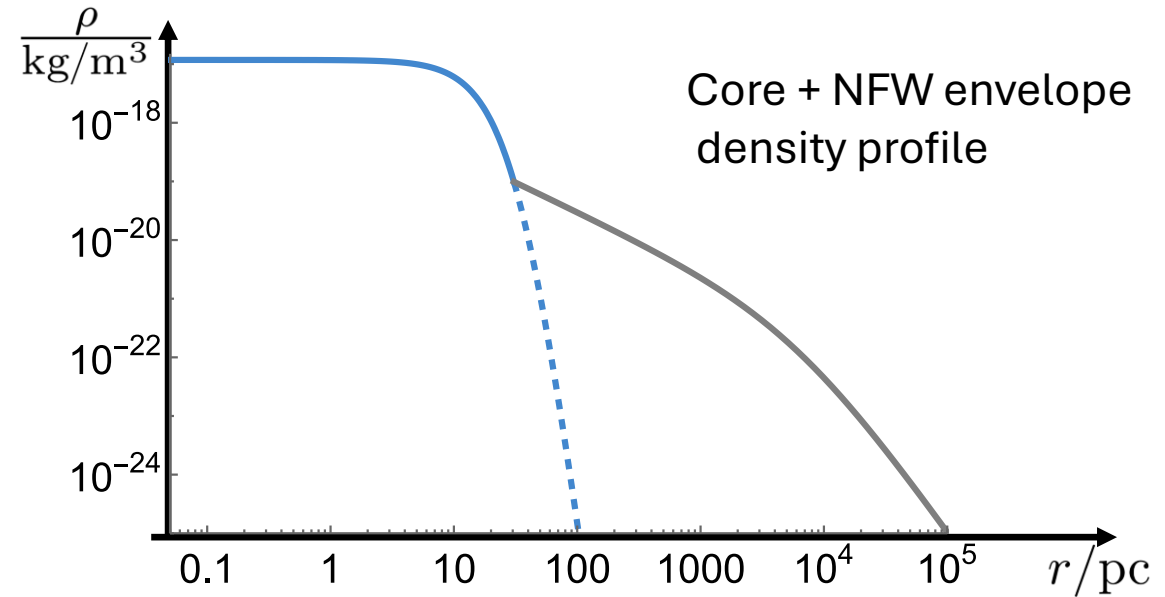
Backup slides

Ultralight boson and halo structure



Ultra-light dark matter Elisa G. M. Ferreira (2021)

While for small galaxies (dwarf spheroidals or ultra-faint dwarfs) the core can be used to model the total halo, the bigger halos (spiral, like Milky Way) require core+envelope model

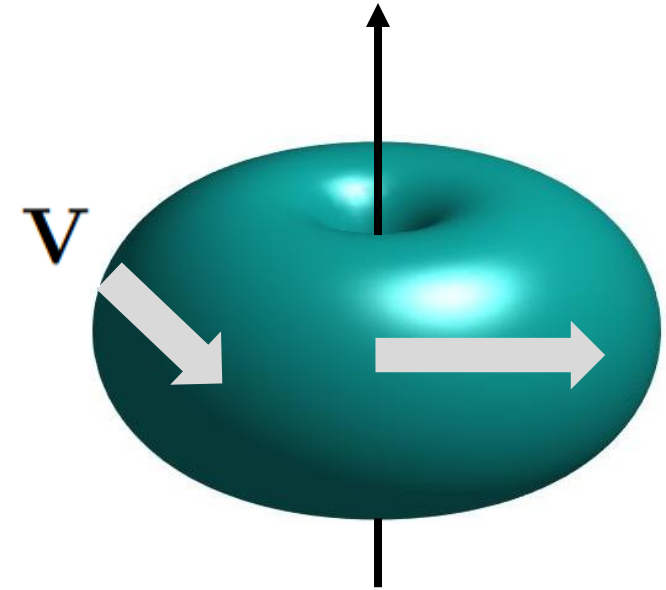
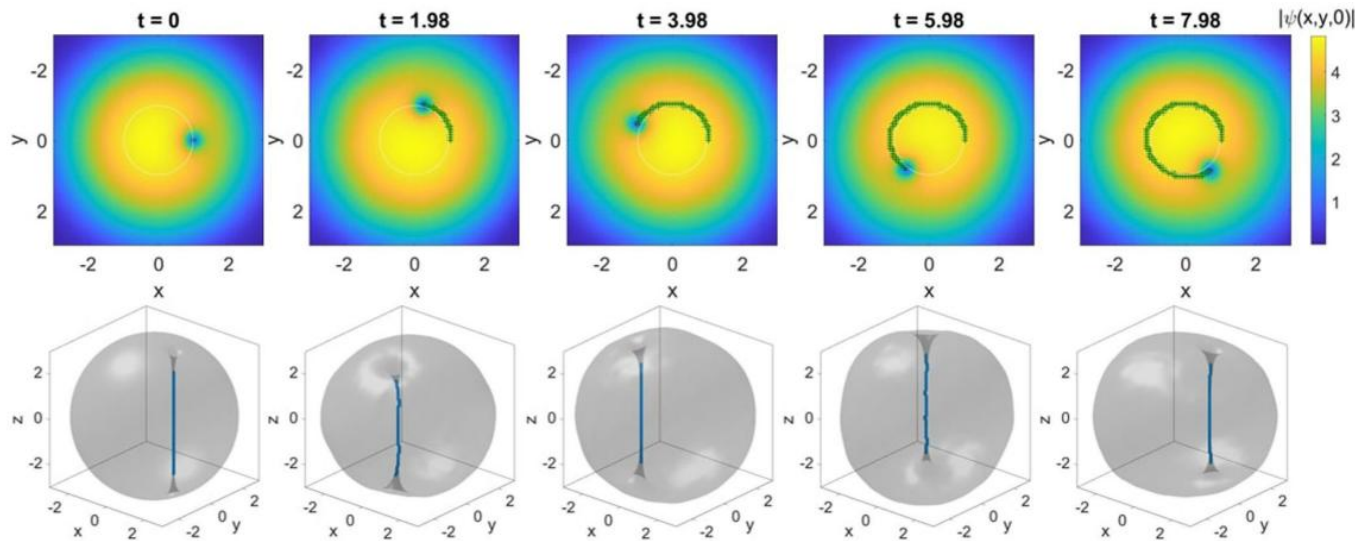


ULDM vortices

1) Have density hole and the superfluid flow

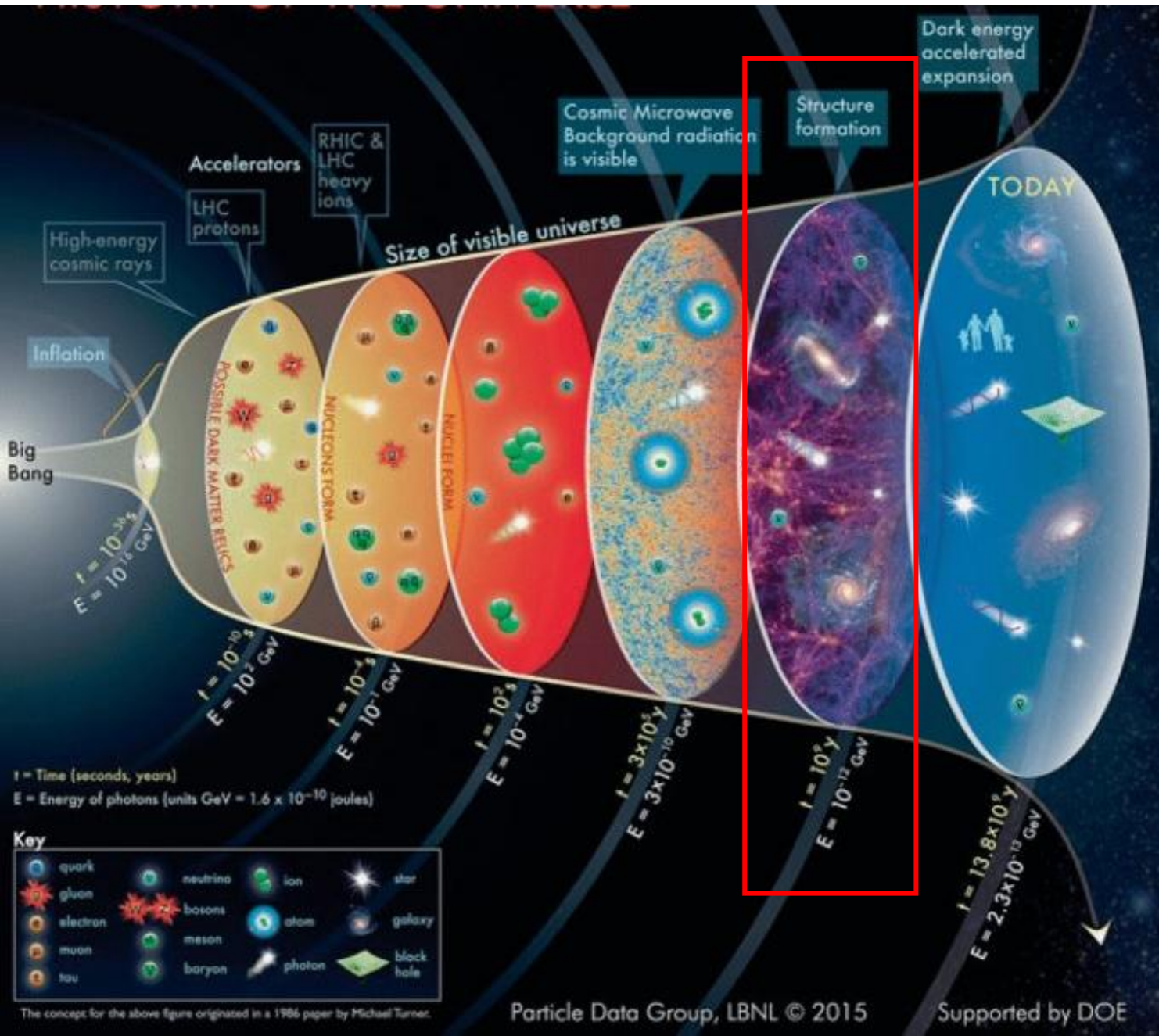
$$\mathbf{V} = \frac{\mathbf{j}}{|\psi|^2} = -\frac{i\hbar}{2m|\psi|^2}(\psi^*\nabla\psi - \psi\nabla\psi^*) = \frac{s\hbar}{m} \frac{1}{r} \mathbf{e}_\phi$$

2) Interesting dynamics



3) Have interesting implications on astrophysical observations...

Cosmological perspective: structure formation



“The emergence of stars, galaxies, groups and clusters of galaxies, and even larger inhomogeneities from the nearly homogeneous early universe is referred to as structure formation”.

Galactic Dynamics J. Binney and S. Tremaine (2011)

ULDM structure formation

Investigated in

- „Fuzzy dark matter simulations“ Schive 2026
- „Cosmic structure as the quantum interference of a coherent dark wave “ Schive et al. 2014
- „Jeans-type instability of a complex self-interacting scalar field in general relativity“ A. Suarez et al., Phys Rev D 2018
- “Structure formation in large-volume cosmological simulations of fuzzy dark matter: impact of the non-linear dynamics” S. May et al. 2021
- “Coherent and incoherent structures in fuzzy dark matter haloes” I.-K. Liu et al. 2023....
- and many more.

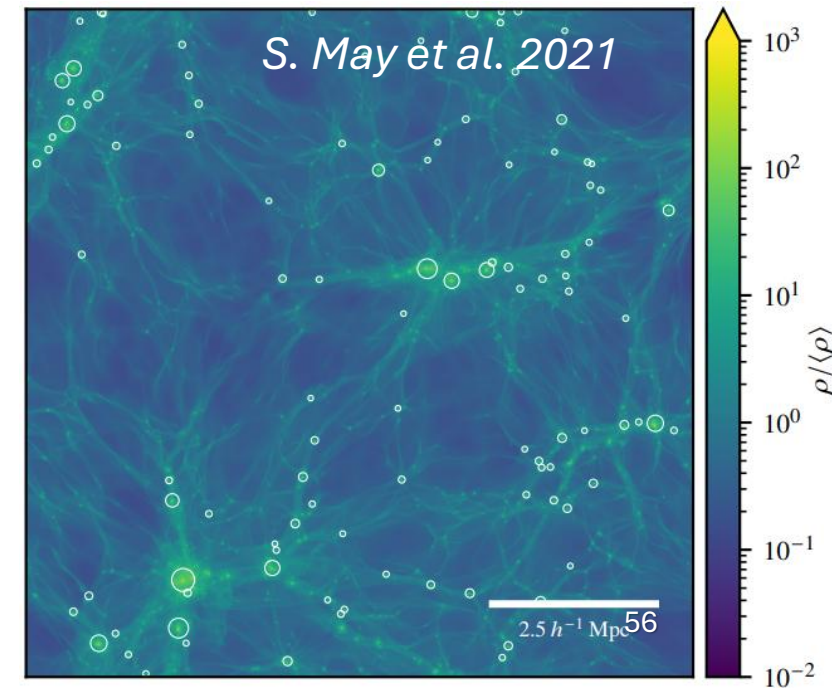
Usual approach:

GPP equations + Universe expansion

$$\left[\begin{aligned} i\hbar \left(\frac{\partial}{\partial t} + \frac{3}{2} H \right) \psi &= -\frac{\hbar^2}{2m a^2} \nabla^2 \psi + m \Phi \psi \\ \frac{1}{a^2} \Delta \Phi &= 4\pi G \rho \end{aligned} \right.$$

... and then a large-scale numerical simulation

Can we try to understand this process with a simple model?



ULDM structure formation

Perturbed FLRW spacetime:

$$ds^2 = \left(1 + \frac{2\Phi}{c^2}\right) c^2 dt^2 + a^2(t) \left(-1 + \frac{2\Phi}{c^2}\right) \delta_{ij} dx^i dx^j$$

$$\square\psi + \left(\frac{mc}{\hbar}\right)^2 \psi = 0 \quad \xrightarrow{\text{nonrelativistic limit}} \quad i\hbar \left(\frac{\partial}{\partial t} + \frac{3}{2}H\right) \psi = -\frac{\hbar^2}{2ma^2} \nabla^2 \psi + m\Phi\psi$$

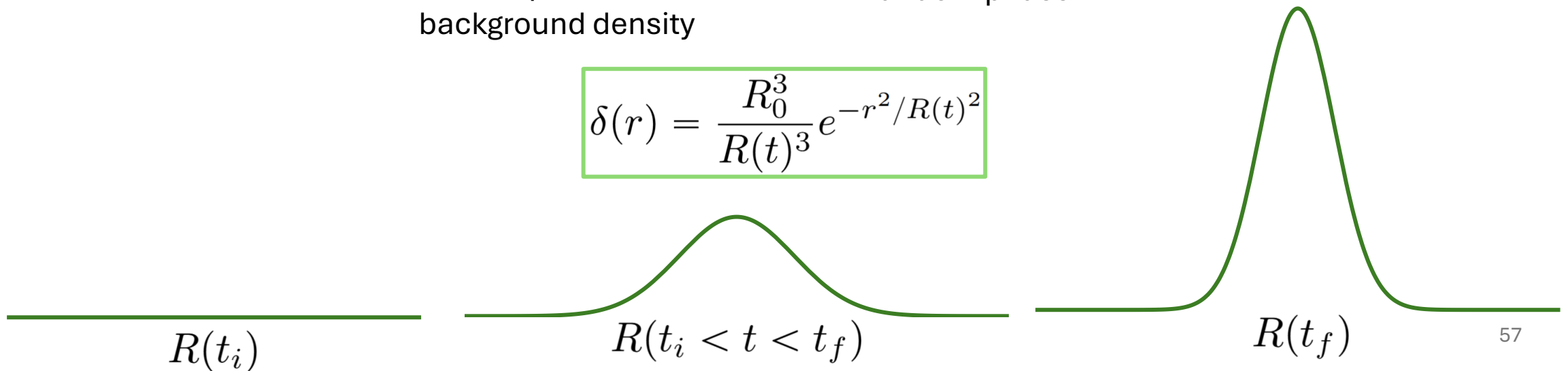
Variational approach: $\psi(\mathbf{r}, t) = \sqrt{\rho_0(t)} \left(e^{i\alpha(\mathbf{r})} + e^{i\theta(\mathbf{r}, t)} \sqrt{\delta(\mathbf{r}, t)} \right)$

time-dependent uniform background density $\rightarrow \rho_0(t)$

random phase $\rightarrow \alpha(\mathbf{r})$

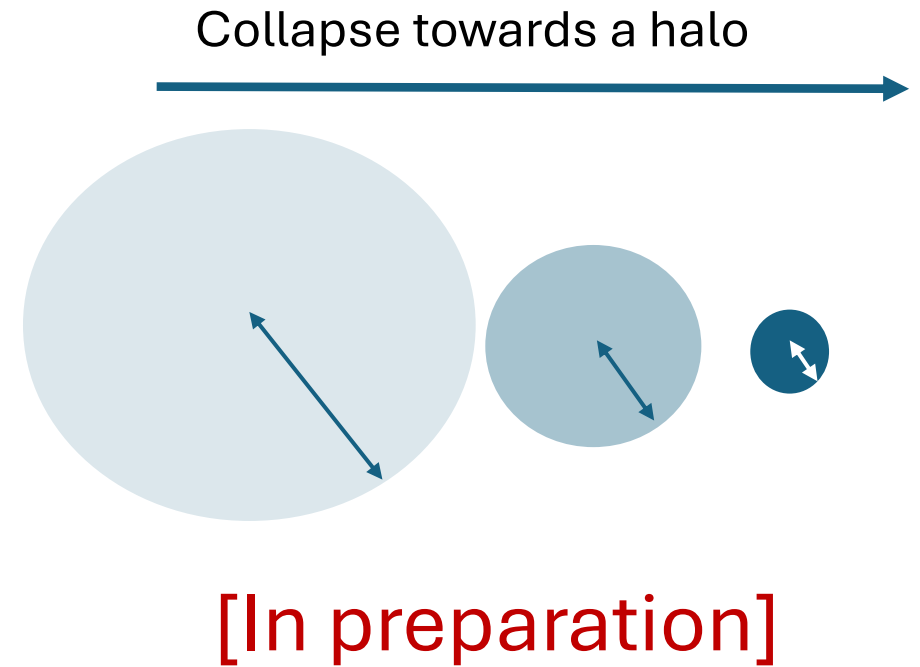
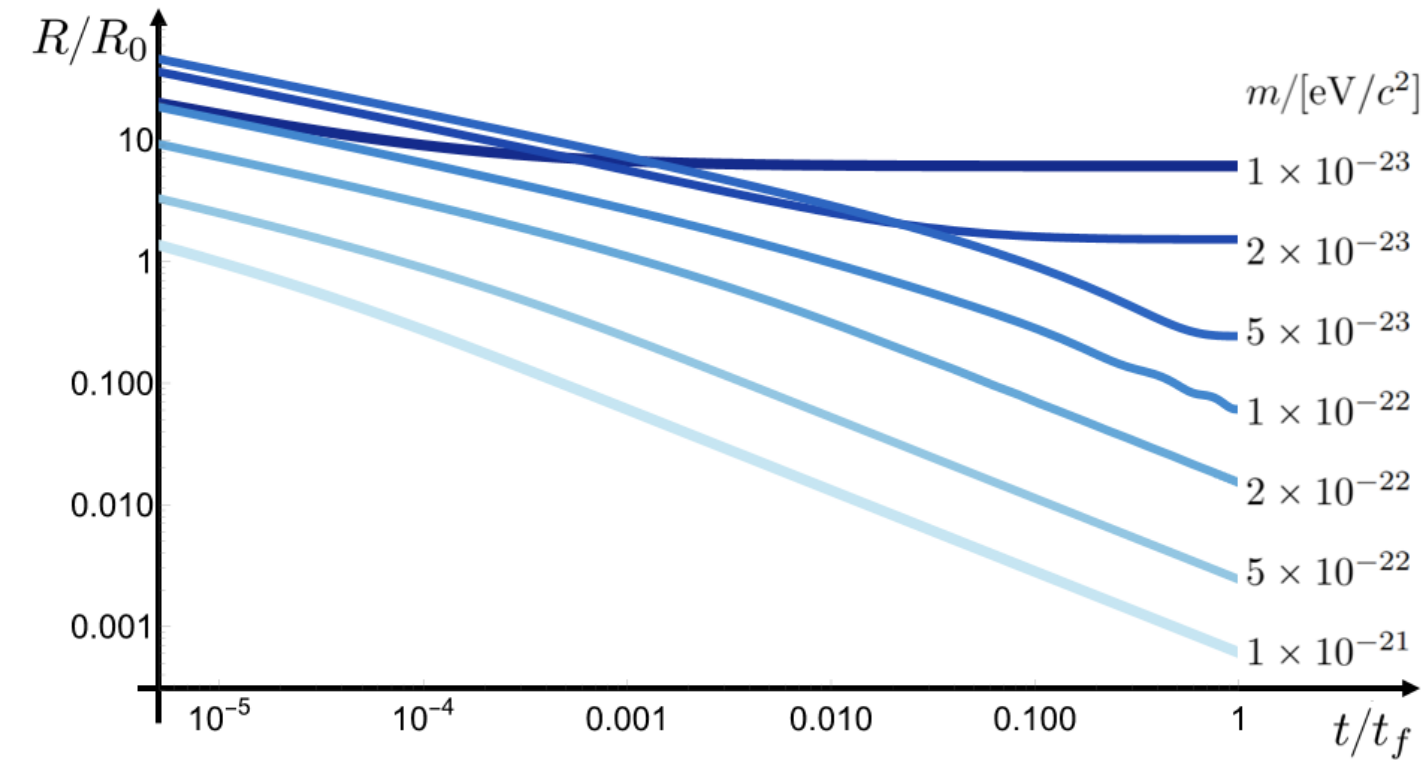
Inhomogeneity (halo) wavefunction $\rightarrow \delta(\mathbf{r}, t)$

$$\delta(r) = \frac{R_0^3}{R(t)^3} e^{-r^2/R(t)^2}$$



Evolution of the halo size in expanding Universe

$$\ddot{R} + \frac{4}{3t}\dot{R} - \frac{\hbar^2 t_f^{8/3}}{m^2 t^{8/3}} \frac{1}{R^3} + \frac{2c^2 R_0^3 t_f^{4/3} R}{3t^2} \int_{t_i}^t \frac{\tau^{-1/3} d\tau}{(R^2(t) + R^2(\tau) + 4C(\tau, t))^{5/2}} = 0$$



ULDM gravitational potential

Einstein's equations in the
perturbed FLRW spacetime:

$$G_{\mu\nu} = 8\pi G T_{\mu\nu}$$

$$\frac{1}{a^2} \Delta \Phi - \frac{3\dot{a}}{a} \frac{\dot{\Phi}}{c^2} - 8\pi G \rho_0 \frac{\Phi}{c^2} = 4\pi G \rho_0 \delta$$

$$\begin{aligned} \mathcal{L}(t) = & \frac{1}{G m t_f^2} \frac{\sqrt{\pi}}{4} \left[-\hbar R_0^3 R^2 \dot{\beta} - \frac{\hbar^2 t_f^{4/3}}{m t^{4/3}} \frac{R_0^3 (1 + 4R^4 \beta^2)}{2R^2} \right] \int_u^\eta dv \frac{a}{a'} \\ \Phi(\eta, \mathbf{x}) = & + \frac{m c^2 R_0^6 t_f^{2/3}}{6 t^{2/3}} \int_{t_i}^t \frac{\dot{a}(\tau) d\tau}{(R^2(t) + R^2(\tau) + 4C(\tau, t))^{3/2}} \end{aligned}$$