

pySecDec TUTORIAL

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$$\sigma = \sigma_{\text{LO}} + \sigma_{\text{NLO}} = \int_n d\sigma^{\text{B}} + \int_n d\sigma^{\text{V}} + \int_{n+1} d\sigma^{\text{R}}$$

 **Finite**  *Infinite*  *Infinite*

$d\sigma^{\text{B}}$ = born cross section

$d\sigma^{\text{V}}$ = **virtual** correction

$d\sigma^{\text{R}}$ = real correction

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σ $\swarrow$ **Finite**

$\int_n d\sigma^{\text{V}}$ $\swarrow$ *Infinite*

$\int_{n+1} d\sigma^{\text{R}}$ $\swarrow$ *Infinite*

$d\sigma^{\text{V}} = \textbf{virtual}$ correction

$d\sigma^{\text{B}} = \text{born cross section}$

$d\sigma^{\text{R}} = \text{real correction}$

Virtual correction involves **loop integrals**

$$\sim \int \frac{d^4 k}{(2\pi)^4} \frac{1}{D_1 D_2 D_3 D_4}$$

$$\begin{aligned} D_1 &= k^2 - m^2 \\ D_2 &= (k + p_1)^2 - m^2 \\ D_3 &= (k + p_1 + p_2)^2 - m^2 \\ D_4 &= (k + p_1 + p_2 - p_3)^2 - m^2 \end{aligned}$$

- These integrals can be *divergent*; require **regularisation** procedure..

Divergences & Regularisation

- **Dimensional regularisation**

- Shift dimension such that integrals are well defined: $D = 4 - 2\epsilon$
- Restore four-dimensional theory by: $\epsilon \rightarrow 0$

Integration measure: $\int \frac{d^D k}{(2\pi)^D}$

- Infinities “captured” as **poles in ϵ**

- They must cancel with other parts of the calculation (UV/IR subtraction)

Divergences & Regularisation

- **Dimensional regularisation**

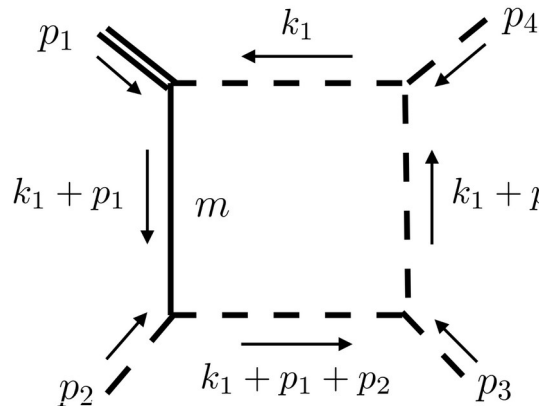
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A typical result can look like:



Box with massless propagators

Series expansion in ϵ

$$\sim \frac{A(s, t)}{\epsilon^2} + \frac{B(s, t)}{\epsilon} + C(s, t) + \mathcal{O}(\epsilon)$$

- **Infinities extracted as 1/eps**
- A,B,C are *finite* $\rightarrow$ numerical integration

How to find this series expansion?

Feynman Parametrisation

General loop integral:

$$G = \int \prod_{l=1}^L [d^D k_l] \frac{1}{\prod_{j=1}^N P_j^{\nu_j}(\{k\}, \{p\}, m_j^2)}$$

 **L loops**
 **N propagators**

$$[d^D k_l] = \frac{\mu^{4-D}}{i\pi^{\frac{D}{2}}} d^D k_l$$

Feynman Parametrisation

General loop integral:

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$$[d^D k_l] = \frac{\mu^{4-D}}{i\pi^{\frac{D}{2}}} d^D k_l$$

It is often useful to work in *parameter space*. Feynman's trick: $\frac{1}{AB} = \int_0^1 \frac{du}{[uA + (1-u)B]^2}$

Or more generally:

$$\frac{1}{\prod_{j=1}^N P_j^{\nu_j}} = \frac{\Gamma(N_\nu)}{\prod_{j=1}^N \Gamma(\nu_j)} \int_0^\infty \prod_{j=1}^N dx_j x_j^{\nu_j-1} \delta(1 - \sum_{i=1}^N x_i) \frac{1}{\left[\sum_{j=1}^N x_j P_j\right]^{N_\nu}}$$

 **Product**

 **Feynman Parameters**

 **Sum**

Feynman Parametrisation II

Feynman parameterizing our loop integral:

$$G = \int_{-\infty}^{\infty} \prod_{l=1}^L [d^D k_l] \frac{1}{\prod_{j=1}^N P_j^{\nu_j}} = \frac{\Gamma(N_\nu)}{\prod_{j=1}^N \Gamma(\nu_j)} \int_0^\infty \prod_{j=1}^N dx_j x_j^{\nu_j-1} \delta(1 - \sum_{i=1}^N x_i)$$
$$\times \int_{-\infty}^{\infty} \prod_{l=1}^L [d^D k_l] \left[\sum_{i,j=1}^L k_i^T \overset{\text{From quadratic (in k) terms of propagators}}{\uparrow} M_{ij} k_j - 2 \sum_{j=1}^L k_j^T \overset{\text{Linear (in k) terms}}{\uparrow} \cdot Q_j + J + i\delta \right]^{-N_\nu}$$

Feynman Parametrisation II

Feynman parameterizing our loop integral:

$$G = \int_{-\infty}^{\infty} \prod_{l=1}^L [d^D k_l] \frac{1}{\prod_{j=1}^N P_j^{\nu_j}} = \frac{\Gamma(N_\nu)}{\prod_{j=1}^N \Gamma(\nu_j)} \int_0^\infty \prod_{j=1}^N dx_j x_j^{\nu_j-1} \delta(1 - \sum_{i=1}^N x_i) \\ \times \int_{-\infty}^{\infty} \prod_{l=1}^L [d^D k_l] \left[\sum_{i,j=1}^L k_i^T \overset{\substack{\uparrow \\ \text{From quadratic (in k) terms of propagators}}}{M_{ij}}} k_j - 2 \sum_{j=1}^L k_j^T \overset{\substack{\uparrow \\ \text{Linear (in k) terms}}}{Q_j}} + J + i\delta \right]^{-N_\nu}$$

Key Point: In this form we can shift k to eliminate linear terms (obtain spherical symmetry) then do the momentum integrals!

Feynman Parametrisation III

After integration over momenta we obtain:

Master Formula

$$G = (-1)^{N_\nu} \frac{\Gamma(N_\nu - LD/2)}{\prod_{j=1}^N \Gamma(\nu_j)} \int_0^\infty \prod_{j=1}^N dx_j x_j^{\nu_j-1} \delta(1 - \sum_{i=1}^N x_i) \frac{\mathcal{U}^{N_\nu - (L+1)D/2}(\vec{x})}{\mathcal{F}^{N_\nu - LD/2}(\vec{x}, s_{ij})}$$

Graph Polynomials:

1st Symanzik Polynomial: $\mathcal{U}(\vec{x}) = \det(M)$

2nd Symanzik Polynomial: $\mathcal{F}(\vec{x}, s_{ij}) = \det(M) \left[\sum_{i,j=1}^L Q_i M_{ij}^{-1} Q_j - J - i\delta \right]$

We have exchanged L momentum integrals for N parameter integrals

Sector Decomposition

In parameter space we know how to algorithmically obtain series expansions

One technique **Iterated Sector Decomposition** repeat:

Binoth, Heinrich 00

$$\begin{aligned}
 & \int_0^1 dx_1 \int_0^1 dx_2 \frac{1}{(x_1 + x_2)^{2+\epsilon}} \quad \leftarrow \text{Overlapping singularity for } x_1, x_2 \rightarrow 0 \\
 &= \int_0^1 dx_1 \int_0^1 dx_2 \frac{1}{(x_1 + x_2)^{2+\epsilon}} (\theta(x_1 - x_2) + \theta(x_2 - x_1)) \\
 &= \int_0^1 dx_1 \int_0^{x_1} dx_2 \frac{1}{(x_1 + x_2)^{2+\epsilon}} + \int_0^1 dx_2 \int_0^{x_2} dx_1 \frac{1}{(x_1 + x_2)^{2+\epsilon}} \\
 &= \int_0^1 dx_1 \int_0^1 dt_2 \frac{x_1}{(x_1 + x_1 t_2)^{2+\epsilon}} + \int_0^1 dx_2 \int_0^1 dt_1 \frac{x_2}{(x_2 t_1 + x_2)^{2+\epsilon}} \\
 &= \int_0^1 dx_1 \int_0^1 dt_2 \frac{x_1^{-1-\epsilon}}{(1 + t_2)^{2+\epsilon}} + \int_0^1 dx_2 \int_0^1 dt_1 \frac{x_2^{-1-\epsilon}}{(t_1 + 1)^{2+\epsilon}} \quad \leftarrow \text{Singularities factorised}
 \end{aligned}$$

Sector Decomposition II

Consider a **Sector Decomposed** integral (simple case $a = -1$):

$$\begin{aligned} & \int_0^1 dx x^{-1-b\epsilon} f(x) \\ &= \int_0^1 dx x^{-1-b\epsilon} [f(x) - f(0) + f(0)] \\ &= \int_0^1 dx x^{-1-b\epsilon} f(0) + \int_0^1 dx x^{-1-b\epsilon} [f(x) - f(0)] \\ &= \frac{f(0)}{-b\epsilon} + \int_0^1 dx x^{-b\epsilon} \left[\frac{f(x) - f(0)}{x} \right] \end{aligned}$$

 **Poles**  **Finite**

By Definition:

$f(0) \neq 0$

$f(0)$ finite

Key Point: Sector Decomposed integrals can be easily expanded in ϵ and **numerically** integrated!

pySecDec

pySecDec evaluates loop integrals

- Divergences subtracted through **sector decomposition**
- Finite (subtracted) integrands are expanded in the dimensional regulator
- The expansion coefficients (finite integrals) are **integrated numerically**

Python

$$I = \sum_{m=-r}^{2L} \frac{C_m}{\epsilon^m} + \mathcal{O}(\epsilon^{r+1})$$

C++ (with python interface)

Installation

```
$ python3 -m pip install --user --upgrade pySecDec
```



Exercises

```
python3 -m pip install notebook
```

If you want to run the examples locally
(strongly advised)

- Three Jupyter notebooks
 1. Scalar 1-loop box integral (**1_box.ipynb**)
 2. LO amplitude to gg → H (**2_ggh.ipynb**)
 3. LO amplitude to light-by-light scattering (**3_lbl.ipynb**)

Find the notebooks at: <https://yeti-2526.notebooks.danielmaitre.phyip3.dur.ac.uk/>