

## QUESTION 1: HALO

The halo is the region populated by dark matter with a galaxy at its centre. One can show that the rotation curves of galaxies can be reproduced if the dark matter mass distribution scales like

$$\rho(r) = \frac{v_c^2}{4\pi G} \frac{1}{r^2} \quad \Rightarrow \quad M(r) = \frac{v_c^2 r}{G_N}$$

where  $v_c$  is a constant circular velocity and  $G_N$  is Newtons constant. Here the goal is to show that an ideal, isotropic gas in hydrostatic equilibrium with gravity with a Maxwell velocity distribution

$$f(\vec{v}) = N \exp\left(-\frac{1}{2} \frac{|\vec{v}|^2}{\sigma^2}\right),$$

predicts such a mass distribution.

- Show that the mean square velocity squared  $\langle v^2 \rangle$  is proportional to the temperature  $T$  of an ideal gas with average kinetic energy  $\langle K \rangle = 3/2 k_B T$ .
- Calculate the mean speed for the Maxwell Boltzmann distribution

$$\langle v^2 \rangle = \int_0^\infty v^2 f(v) dv$$

and solve for the velocity distribution  $\sigma$  using the result from a) (Determine the normalisation constant  $N$  first, then calculate  $\langle v^2 \rangle$ ). Using the ideal gas law you should find the relation between the pressure  $P(r)$  and the density  $\rho(r)$ ,

$$P(r) = \sigma^2 \rho(r).$$

Hint:  $\int_{-\infty}^\infty e^{-x^2/b} dx = \sqrt{b\pi}$

- Now we want to know the density profile for such a dark matter halo. To derive that, consider a thin volume of the dark matter halo with surface  $A$  perpendicular to  $r$  and infinitesimal thickness  $dr$ . If you balance the force exerted by the hydrostatic pressure on a that volume element with the gravitational force, you can derive a differential equation for the matter-energy density  $\rho(r)$ . Show that the ansatz  $\rho(r) = a/r^2$  solves this differential equation and therefore a gravitating ideal gas produces the density profile observed in dark matter halos. *Hint: You can use*

$$\begin{aligned} F_{\text{hydro}} &= F_{\text{grav}} \\ (P(r+dr) - P(r))A &= -\rho(r)A dr \frac{G_N M(r)}{r^2} \\ \Rightarrow \frac{dP(r)}{dr} &= -\rho(r) \frac{G_N M(r)}{r^2}. \end{aligned}$$

and the relation for spherically symmetric mass distributions

$$\frac{dM(r)}{dr} = 4\pi r^2 \rho(r).$$

## QUESTION 2: CMB

The cosmic microwave background (CMB) radiation is the relic of the recombination era when photons for the first time were free to travel through the universe without constantly being scattered off free electrons and protons. It emerged when the Universe was  $t_{\text{CMB}}$  380000 years old. Today the CMB is an almost perfect black body radiation with a temperature of  $T_{\text{CMB}} = 2.725 \text{ K}$ . During most of the time between  $t_{\text{CMB}}$  and  $t_{\text{now}}$ , the energy-matter density of the Universe was matter-dominated. This means that the relation between temperature and time is given to good approximation by

$$T(t_{\text{past}}) = T(t_{\text{now}}) \left( \frac{t_{\text{past}}}{t_{\text{now}}} \right)^{-2/3}. \quad (1)$$

- What is the wavelength and frequency of the CMB today? *hint: The CMB is black-body radiation. Use Wiens displacement law  $\lambda_{\text{peak}} = b/T$  with  $b = 2.9 \times 10^{-3} \text{ mK}$ .*
- Show that the temperature of the CMB when the photons decoupled was 3000 K. What was its wavelength then? Would it have been visible to the human eye and can you guess its color?
- How old was the universe when the night sky turned black? *Hint: This should happen at the latest when the black-body spectrum has an average temperature of  $T \approx 700 \text{ K}$  (the Draper point).*
- Go to <https://chrisnorth.github.io/planckapps/Simulator/> and solve the CMB challenge: Find values that agree 100% with our Universe. Understand what the parameter mean and try to get the fit right (don't click answer yet!) Why does the flatness of the universe depend on the first three parameters, but not on the Hubble constant, the spectral index and the re-ionization redshift?

### QUESTION 3: DIRECT DETECTION

Direct searches for Dark Matter are shielded detectors placed underground, which try to measure interactions of Dark Matter with targets made from different elements. A recent summary plot of direct detection bounds is shown in Fig. 1. Experimental bounds are shown as solid lines, projected bounds are shown as dashed lines. Possible signal candidates are shown as colored regions, while the thick, dashed orange line demarks the neutrino floor.

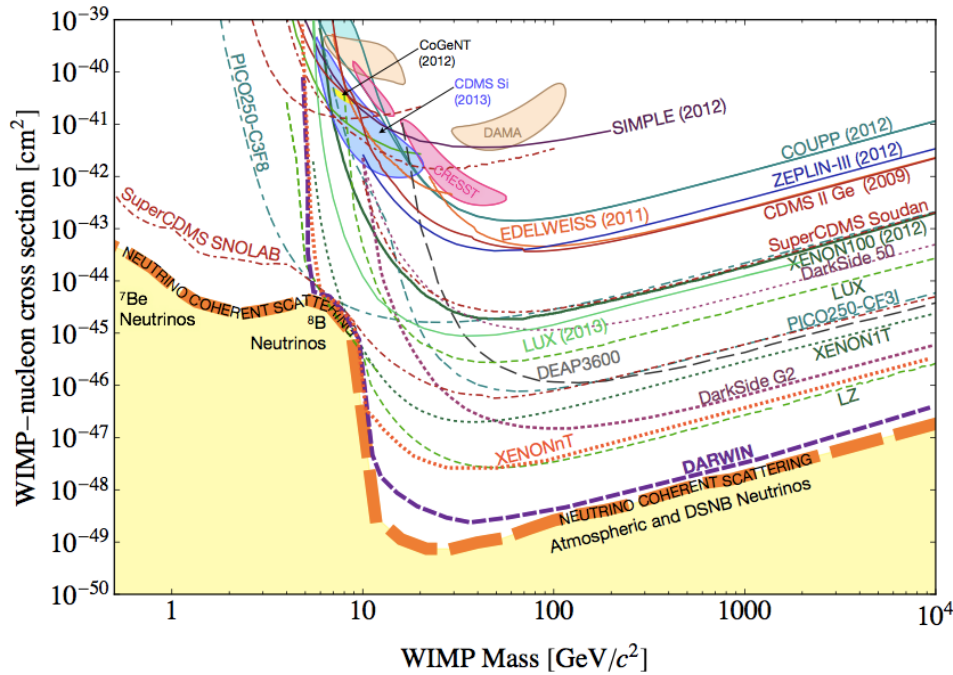


Fig 1: Direct Detection bounds from various experiments. (adapted from Snowmass, [1310.8327](#) )

- The local dark matter density is  $\rho = 0.3 \text{ GeV/cm}^3$ . How many dark matter particles would you expect in a 2L bottle if the dark matter had a mass of 5 GeV or 500 GeV?
- Explain the general shape of the exclusion limits (the sharp drop at a  $\sim 10 \text{ GeV}$  and the linear slope for higher masses)! Can you explain the differences in the excluded regions from different experiments? What would you guess is the scaling for the sensitivity with mass and runtime of these experiments?

You are now an experimentalist working in a Direct Detection collaboration considering two detectors. The first is based on a Germanium target and the second has a Xenon target:

|                     |            | Germanium  | Xenon      |
|---------------------|------------|------------|------------|
| Energy Threshold    | $E_t$      | 1 keV      | 5 keV      |
| Energy Interval     | $\Delta E$ | (1-40) keV | (5-40) keV |
| Target Mass         | $M$        | 1 kg       | 35 kg      |
| Target Element Mass | $m_A$      | 65 GeV     | 122 GeV    |
| Mass Number         | $A$        | 73         | 131        |

The goal is to provide the best possible limit on a theory predicting a

- 1) light dark matter candidate  $M_\chi = 5 \text{ GeV}$  ,
- 2) heavy dark matter candidate  $M_\chi = 500 \text{ GeV}$  .

In both cases the commissioned runtime is  $T = 100 \text{ days}$ .

- c) Direct detection experiments are limited by the nuclear recoil energy threshold of the target material  $E_t$ . In terms of the velocity  $v$  of the Dark Matter particle and the scattering angle  $\theta$ , the recoil energy is given by

$$E_R = v^2 \frac{\mu_N^2}{m_A} (1 - \cos \theta) , \quad (2)$$

in which the reduced mass is given by  $\mu_N = \frac{m_A M_\chi}{m_A + M_\chi}$ . Compute the minimal velocity  $v_{\min}$  for the Germanium and Xenon detector and the two Dark Matter masses given above (note, that  $v$  is given as a fraction of  $c = 3 \times 10^5 \text{ km/s}$  using the units in the table.)

- d) The Dark Matter velocity distribution follows a Maxwell-Boltzmann distribution

$$f(\mathbf{v}) = N e^{-\mathbf{v}^2/v_0^2} , \quad (3)$$

with  $v_0 = 220 \text{ km/s}$  the circular velocity of the Dark Matter halo and  $N = 1/(\sqrt{\pi}v_0)^3$ . Integrate over the solid angle to obtain the velocity distribution  $f(|\mathbf{v}|)$ . Plot this function and discuss what it implies for the velocities you have derived in c)? Does it even make sense to consider very fast Dark Matter particles or should the velocity distribution be cut off at a certain speed  $v_{\max}$ ?

- e) The expected rate for WIMP interactions can be expressed as

$$R \approx \frac{A^2}{2\mu_P^2 M_\chi} \sigma_0 \rho_\chi \int_{v_{\min}}^{v_{\max}} \frac{f(v)}{v} dv \cdot \Delta E , \quad (4)$$

in which  $\mu_P = \frac{m_N M_\chi}{m_N + M_\chi}$  is the reduced mass of the Dark Matter and a Nucleon (either proton or neutron)  $m_N \approx 1 \text{ GeV}$ . The local dark matter density  $\rho_\chi = 0.3 \text{ GeV/cm}^3$  and the velocity distribution given above are astrophysical inputs. The mass of the Dark Matter candidate and the cross section  $\sigma_0 = 1 \cdot 10^{-38} \text{ cm}^2$  are quantities provided by your particle physics colleague. Compute the expected number of events  $N = R \cdot T \cdot M$  for the two detectors and both the heavy and light Dark Matter candidate. (Make sure you are using consistent units!)

TABLE IV. 90% C.L. intervals for the Poisson signal mean  $\mu$ , for total events observed  $n_0$ , for known mean background  $b$  ranging from 0 to 5.

| $n_0 \backslash b$ | 0.0         | 0.5         | 1.0         | 1.5         | 2.0         | 2.5         | 3.0         | 3.5         | 4.0        | 5.0        |
|--------------------|-------------|-------------|-------------|-------------|-------------|-------------|-------------|-------------|------------|------------|
| 0                  | 0.00, 2.44  | 0.00, 1.94  | 0.00, 1.61  | 0.00, 1.33  | 0.00, 1.26  | 0.00, 1.18  | 0.00, 1.08  | 0.00, 1.06  | 0.00, 1.01 | 0.00, 0.98 |
| 1                  | 0.11, 4.36  | 0.00, 3.86  | 0.00, 3.36  | 0.00, 2.91  | 0.00, 2.53  | 0.00, 2.19  | 0.00, 1.88  | 0.00, 1.59  | 0.00, 1.39 | 0.00, 1.22 |
| 2                  | 0.53, 5.91  | 0.03, 5.41  | 0.00, 4.91  | 0.00, 4.41  | 0.00, 3.91  | 0.00, 3.45  | 0.00, 3.04  | 0.00, 2.67  | 0.00, 2.33 | 0.00, 1.73 |
| 3                  | 1.10, 7.42  | 0.60, 6.92  | 0.10, 6.42  | 0.00, 5.92  | 0.00, 5.42  | 0.00, 4.92  | 0.00, 4.42  | 0.00, 3.95  | 0.00, 3.53 | 0.00, 2.78 |
| 4                  | 1.47, 8.60  | 1.17, 8.10  | 0.74, 7.60  | 0.24, 7.10  | 0.00, 6.60  | 0.00, 6.10  | 0.00, 5.60  | 0.00, 5.10  | 0.00, 4.60 | 0.00, 3.60 |
| 5                  | 1.84, 9.99  | 1.53, 9.49  | 1.25, 8.99  | 0.93, 8.49  | 0.43, 7.99  | 0.00, 7.49  | 0.00, 6.99  | 0.00, 6.49  | 0.00, 5.99 | 0.00, 4.99 |
| 6                  | 2.21,11.47  | 1.90,10.97  | 1.61,10.47  | 1.33, 9.97  | 1.08, 9.47  | 0.65, 8.97  | 0.15, 8.47  | 0.00, 7.97  | 0.00, 7.47 | 0.00, 6.47 |
| 7                  | 3.56,12.53  | 3.06,12.03  | 2.56,11.53  | 2.09,11.03  | 1.59,10.53  | 1.18,10.03  | 0.89, 9.53  | 0.39, 9.03  | 0.00, 8.53 | 0.00, 7.53 |
| 8                  | 3.96,13.99  | 3.46,13.49  | 2.96,12.99  | 2.51,12.49  | 2.14,11.99  | 1.81,11.49  | 1.51,10.99  | 1.06,10.49  | 0.66, 9.99 | 0.00, 8.99 |
| 9                  | 4.36,15.30  | 3.86,14.80  | 3.36,14.30  | 2.91,13.80  | 2.53,13.30  | 2.19,12.80  | 1.88,12.30  | 1.59,11.80  | 1.33,11.30 | 0.43,10.30 |
| 10                 | 5.50,16.50  | 5.00,16.00  | 4.50,15.50  | 4.00,15.00  | 3.50,14.50  | 3.04,14.00  | 2.63,13.50  | 2.27,13.00  | 1.94,12.50 | 1.19,11.50 |
| 11                 | 5.91,17.81  | 5.41,17.31  | 4.91,16.81  | 4.41,16.31  | 3.91,15.81  | 3.45,15.31  | 3.04,14.81  | 2.67,14.31  | 2.33,13.81 | 1.73,12.81 |
| 12                 | 7.01,19.00  | 6.51,18.50  | 6.01,18.00  | 5.51,17.50  | 5.01,17.00  | 4.51,16.50  | 4.01,16.00  | 3.54,15.50  | 3.12,15.00 | 2.38,14.00 |
| 13                 | 7.42,20.05  | 6.92,19.55  | 6.42,19.05  | 5.92,18.55  | 5.42,18.05  | 4.92,17.55  | 4.42,17.05  | 3.95,16.55  | 3.53,16.05 | 2.78,15.05 |
| 14                 | 8.50,21.50  | 8.00,21.00  | 7.50,20.50  | 7.00,20.00  | 6.50,19.50  | 6.00,19.00  | 5.50,18.50  | 5.00,18.00  | 4.50,17.50 | 3.59,16.50 |
| 15                 | 9.48,22.52  | 8.98,22.02  | 8.48,21.52  | 7.98,21.02  | 7.48,20.52  | 6.98,20.02  | 6.48,19.52  | 5.98,19.02  | 5.48,18.52 | 4.48,17.52 |
| 16                 | 9.99,23.99  | 9.49,23.49  | 8.99,22.99  | 8.49,22.49  | 7.99,21.99  | 7.49,21.49  | 6.99,20.99  | 6.49,20.49  | 5.99,19.99 | 4.99,18.99 |
| 17                 | 11.04,25.02 | 10.54,24.52 | 10.04,24.02 | 9.54,23.52  | 9.04,23.02  | 8.54,22.52  | 8.04,22.02  | 7.54,21.52  | 7.04,21.02 | 6.04,20.02 |
| 18                 | 11.47,26.16 | 10.97,25.66 | 10.47,25.16 | 9.97,24.66  | 9.47,24.16  | 8.97,23.66  | 8.47,23.16  | 7.97,22.66  | 7.47,22.16 | 6.47,21.16 |
| 19                 | 12.51,27.51 | 12.01,27.01 | 11.51,26.51 | 11.01,26.01 | 10.51,25.51 | 10.01,25.01 | 9.51,24.51  | 9.01,24.01  | 8.51,23.51 | 7.51,22.51 |
| 20                 | 13.55,28.52 | 13.05,28.02 | 12.55,27.52 | 12.05,27.02 | 11.55,26.52 | 11.05,26.02 | 10.55,25.52 | 10.05,25.02 | 9.55,24.52 | 8.55,23.52 |

Tab.1: Lower and upper bound on the expected events for a 90% confidence level and a given number of signal  $n_0$  and background  $b$  events. (Feldman, Cousins, [9711021](#))

- f) Now assume that you measured 3 signal events during the runtime, but also 2 background events. Table 1 gives the lower and upper bound on the expected events for a 90% confidence level and a given number of signal  $n_0$  and background  $b$  events. Use these numbers and the formulas in d) to derive a bound on the cross section  $\sigma_0$  for the two experiments and scenarios described above.

*Hint:* Solve

$$N_{\text{limit}} = T M \frac{A^2}{2\mu_P^2 M_\chi} \sigma_0 \rho_\chi \int_{v_{\text{Min}}}^{v_{\text{esc}}} \frac{f(v)}{v} dv \Delta E$$

for the cross section.

| Nucleus           | $Z$ | Odd Nucleon | $J$ | $\langle S_p \rangle$ | $\langle S_n \rangle$ | $C_A^p/C_p$           | $C_A^n/C_n$           |
|-------------------|-----|-------------|-----|-----------------------|-----------------------|-----------------------|-----------------------|
| $^{19}\text{F}$   | 9   | p           | 1/2 | 0.477                 | -0.004                | $9.10 \times 10^{-1}$ | $6.40 \times 10^{-5}$ |
| $^{23}\text{Na}$  | 11  | p           | 3/2 | 0.248                 | 0.020                 | $1.37 \times 10^{-1}$ | $8.89 \times 10^{-4}$ |
| $^{27}\text{Al}$  | 13  | p           | 5/2 | -0.343                | 0.030                 | $2.20 \times 10^{-1}$ | $1.68 \times 10^{-3}$ |
| $^{29}\text{Si}$  | 14  | n           | 1/2 | -0.002                | 0.130                 | $1.60 \times 10^{-5}$ | $6.76 \times 10^{-2}$ |
| $^{35}\text{Cl}$  | 17  | p           | 3/2 | -0.083                | 0.004                 | $1.53 \times 10^{-2}$ | $3.56 \times 10^{-5}$ |
| $^{39}\text{K}$   | 19  | p           | 3/2 | -0.180                | 0.050                 | $7.20 \times 10^{-2}$ | $5.56 \times 10^{-3}$ |
| $^{73}\text{Ge}$  | 32  | n           | 9/2 | 0.030                 | 0.378                 | $1.47 \times 10^{-3}$ | $2.33 \times 10^{-1}$ |
| $^{93}\text{Nb}$  | 41  | p           | 9/2 | 0.460                 | 0.080                 | $3.45 \times 10^{-1}$ | $1.04 \times 10^{-2}$ |
| $^{125}\text{Te}$ | 52  | n           | 1/2 | 0.001                 | 0.287                 | $4.00 \times 10^{-6}$ | $3.29 \times 10^{-1}$ |
| $^{127}\text{I}$  | 53  | p           | 5/2 | 0.309                 | 0.075                 | $1.78 \times 10^{-1}$ | $1.05 \times 10^{-2}$ |
| $^{129}\text{Xe}$ | 54  | n           | 1/2 | 0.028                 | 0.359                 | $3.14 \times 10^{-3}$ | $5.16 \times 10^{-1}$ |
| $^{131}\text{Xe}$ | 54  | n           | 3/2 | -0.009                | -0.227                | $1.80 \times 10^{-4}$ | $1.15 \times 10^{-1}$ |

Tab.2: Input values for spin-dependent cross sections. ( Tovey et al., PLB 488 17(2000), [0005041](#))

- g) You theorist friend comes back and is excited. A new model for dark matter could explain many of the observations. The theorist explains that the model includes a dark matter candidate  $\chi$  with mass  $M_\chi = 15$  GeV and a new particle, a scalar  $S$  that interacts with Standard Model particles and the dark matter and has a mass  $M_S = 300$  GeV. The scalar  $S$  interacts exclusively with neutron spin with some coupling constant  $g$ . Use Table to find the best possible target material for a direct detection experiment to probe such a dark matter model. Now estimate the cross section. Draw a Feynman diagram and use dimensional analysis and your knowledge of Feynman rules to estimate  $\sigma(\chi N \rightarrow \chi N)$ . How long would it take you with a 10 kg target detector to be able to discover the signal of this model.

*Hint:* The general spin-dependent scattering cross section reads

$$\sigma^{\text{SD}} = \frac{32G_F^2\mu^2}{\pi} \frac{J+1}{J} [a_n\langle S_n \rangle + a_p\langle S_p \rangle]^2$$

- h) What is the neutrino floor (the dashed orange line in Fig.1)? Why does it have a similar shape as the exclusion limits? Should it be the same for different target elements? Look into (Billard, Strigari, Figueroa-Feliciano, [1307.5458](#)).