Towards a complete NNLO Prediction for $\bar{B} \to X_s \gamma$: m_c-dependent matrix elements

Radja Boughezal

ITP University of Zürich, Switzerland

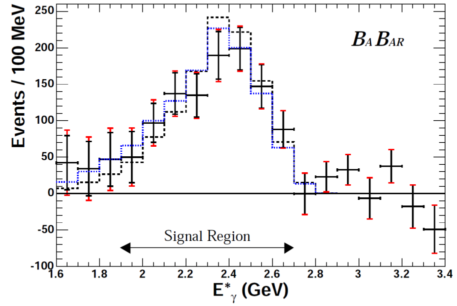
Durham September 2008

Collaborators:
M. Czakon and T. Schutzmeier

Motivations

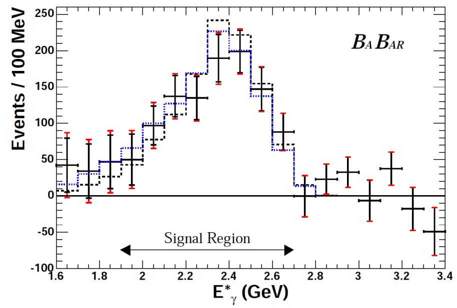
- $\bar{B} \to X_s\gamma$ most precise short-distance information currently available for $\Delta B = 1$ FCNC

$$\mathcal{B}(\bar{B} \to X_s\gamma)^{\rm exp}_{E_\gamma > 1.6\,{\rm GeV}} = (3.55 \pm 0.26) \times 10^{-4}$$ [HFAG2006]



Motivations

- $\bar{B} \to X_s \gamma$ most precise short-distance information currently available for $\Delta B = 1$ FCNC



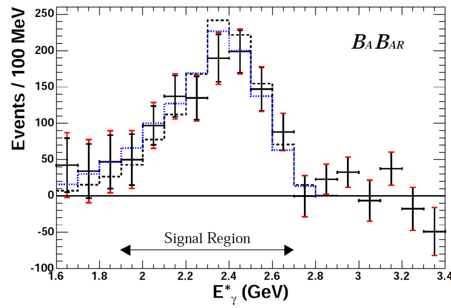
$$\boxed{\mathcal{B}(\bar{B} \to X_s \gamma)^{\text{exp}}_{E_\gamma > 1.6\,\text{GeV}} = (3.55 \pm 0.26) \times 10^{-4}}$$ [HFAG2006]

  - less sensitive to non-perturbative effects
    dominant ones: $\mathcal{O}(\Lambda^2/m_b^2)$, $\mathcal{O}(\Lambda^2/m_c^2)$, $\mathcal{O}(\alpha_s \Lambda/m_b)$

$$\Longrightarrow \quad \Gamma(\bar{B} \to X_s \gamma) \approx \Gamma(b \to X_s^{parton} \gamma)$$
$$= \Gamma(b \to s\gamma) + \Gamma(b \to s\gamma g) + \ldots$$

Motivations

- $\bar{B} \to X_s\gamma$ most precise short-distance information currently available for $\Delta B = 1$ FCNC

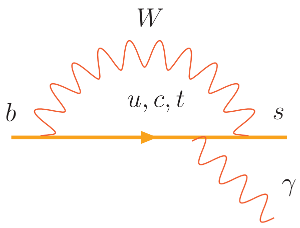


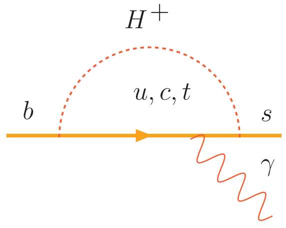
$$\mathcal{B}(\bar{B} \to X_s\gamma)^{\text{exp}}_{E_\gamma > 1.6\,\text{GeV}} = (3.55 \pm 0.26) \times 10^{-4}$$ [HFAG2006]

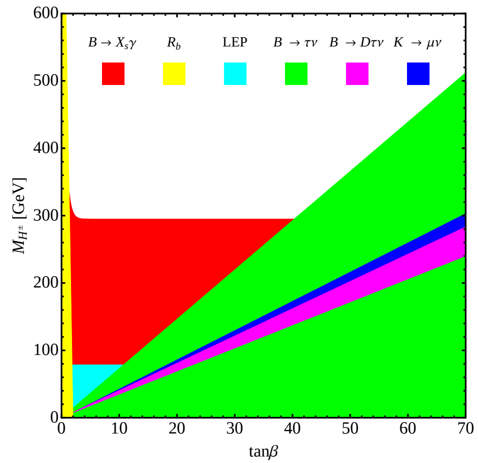
  - less sensitive to non-perturbative effects
    dominant ones: $\mathcal{O}(\Lambda^2/m_b^2)$, $\mathcal{O}(\Lambda^2/m_c^2)$, $\mathcal{O}(\alpha_s \Lambda/m_b)$

$$\Longrightarrow \quad \Gamma(\bar{B} \to X_s\gamma) \approx \Gamma(b \to X_s^{parton}\gamma)$$
$$= \Gamma(b \to s\gamma) + \Gamma(b \to s\gamma g) + \ldots$$

  - loop induced in SM and highly sensitive to new physics which is not suppressed by factors of α as compared to SM contributions







Motivations

- Theoretical error vs. experimental one:

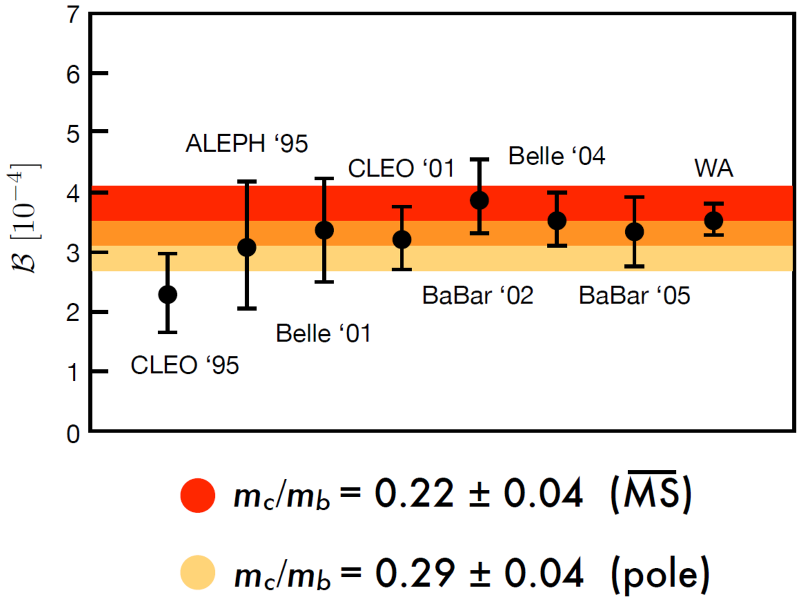
  - $\mathcal{B}(\bar{B} \to X_s\gamma)^{\rm th,NLO}_{E_\gamma > 1.6\,{\rm GeV}} = (3.57 \pm 0.30) \times 10^{-4}$

    [Misiak et al 2001,Buras et al 2002]

  - $\mathcal{B}(\bar{B} \to X_s\gamma)^{\rm exp}_{E_\gamma > 1.6\,{\rm GeV}} = (3.55 \pm 0.26) \times 10^{-4}$

    [HFAG 2006]

    Super-B factory: 5% uncertainty possible
    (more statistics, lower E_γ)



Motivations

- Theoretical error vs. experimental one:

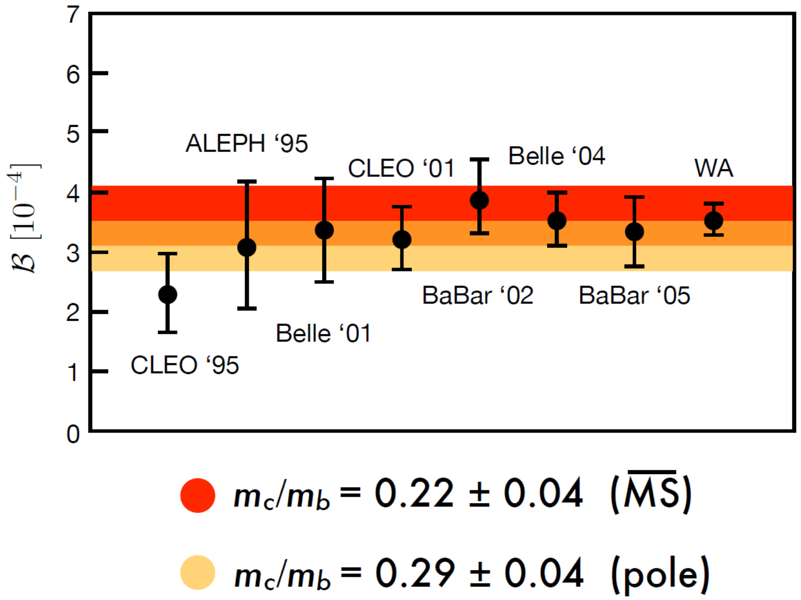
  - $\mathcal{B}(\bar{B} \to X_s\gamma)^{\mathrm{th,NLO}}_{E_\gamma > 1.6\,\mathrm{GeV}} = (3.57 \pm 0.30) \times 10^{-4}$

    [Misiak et al 2001,Buras et al 2002]

  - $\mathcal{B}(\bar{B} \to X_s\gamma)^{\mathrm{exp}}_{E_\gamma > 1.6\,\mathrm{GeV}} = (3.55 \pm 0.26) \times 10^{-4}$

    [HFAG 2006]

  Super-B factory: 5% uncertainty possible
  (more statistics, lower E_γ)



$\Longrightarrow$ strong constraints on new physics require better theoretical precision

Motivations

$$\mathcal{B}(\bar{B} \to X_s \gamma)^{\rm exp}_{E_\gamma > 1.6\,{\rm GeV}} = (3.55 \pm 0.26) \times 10^{-4}$$

[HFAG 2006]

- Contributions to the theory prediction

$$\mathcal{B}(\bar{B} \to X_s\gamma)_{E_\gamma > 1.6\,{\rm GeV}} = \mathcal{B}(\bar{B} \to X_c e \bar{\nu})_{\rm exp} \left[\frac{\Gamma(b \to s\gamma)}{\Gamma(b \to c e \bar{\nu})}\right]_{\rm LO\ EW} f\left(\frac{\alpha_s(M_W)}{\alpha_s(m_b)}\right) \times$$

$$\times \left\{ 1 + \underbrace{\underset{\substack{\textbf{NLO} \\ \sim 25\%}}{\mathcal{O}(\alpha_s)} + \underset{\substack{\textbf{NNLO} \\ \sim 7\%}}{\mathcal{O}(\alpha_s^2)} + \underset{\sim 4\%}{\mathcal{O}(\alpha_{\rm em})}}_{\textbf{perturbative corrections}} + \underbrace{\underset{\sim 1\%}{\mathcal{O}\left(\frac{\Lambda^2}{m_b^2}\right)} + \underset{\sim 3\%}{\mathcal{O}\left(\frac{\Lambda^2}{m_c^2}\right)} + \underset{<\sim 5\%}{\mathcal{O}\left(\frac{\Lambda}{m_b}\alpha_s\right)}}_{\textbf{non-perturbative corrections}} \right\}$$

Motivations

$$\mathcal{B}(\bar{B} \to X_s\gamma)^{\rm exp}_{E_\gamma > 1.6\,{\rm GeV}} = (3.55 \pm 0.26) \times 10^{-4}$$

[HFAG 2006]

- Contributions to the theory prediction

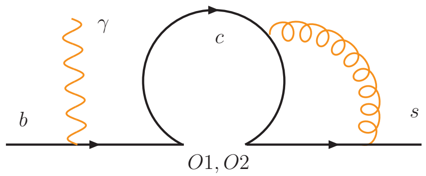
$$\mathcal{B}(\bar{B} \to X_s\gamma)_{E_\gamma>1.6\,{\rm GeV}} = \mathcal{B}(\bar{B} \to X_c e\bar{\nu})_{\rm exp} \left[\frac{\Gamma(b\to s\gamma)}{\Gamma(b\to ce\bar{\nu})}\right]_{\rm LO\ EW} f\left(\frac{\alpha_s(M_W)}{\alpha_s(m_b)}\right) \times$$

$$\times \Big\{ \underbrace{1 + \underset{\substack{\textbf{NLO}\\ \sim 25\%}}{\mathcal{O}(\alpha_s)} + \underset{\substack{\textbf{NNLO}\\ \sim 7\%}}{\mathcal{O}(\alpha_s^2)} + \underset{\sim 4\%}{\mathcal{O}(\alpha_{\rm em})}}_{\textbf{perturbative corrections}} + \underbrace{\underset{\sim 1\%}{\mathcal{O}\left(\frac{\Lambda^2}{m_b^2}\right)} + \underset{\sim 3\%}{\mathcal{O}\left(\frac{\Lambda^2}{m_c^2}\right)} + \underset{<\sim 5\%}{\mathcal{O}\left(\frac{\Lambda}{m_b}\alpha_s\right)}}_{\textbf{non-perturbative corrections}} \Big\}$$

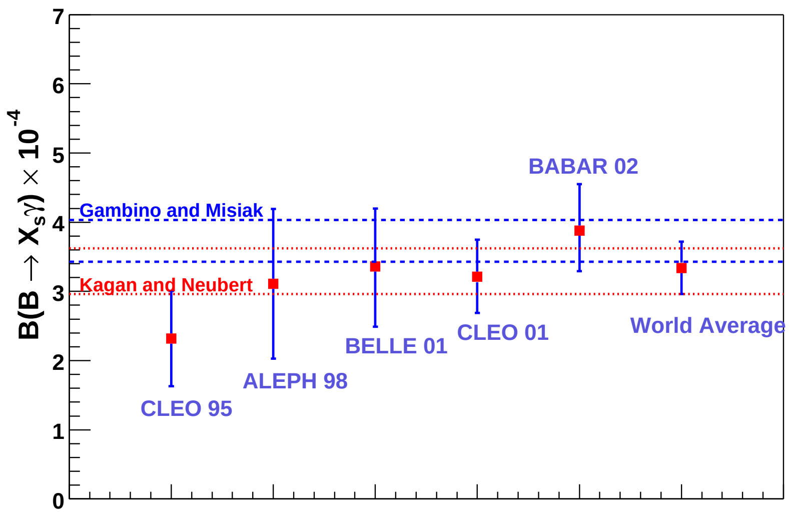
expected NNLO corrections to $\mathcal{B}$ ($\sim 7\%$) are of the same size as the experimental error

Motivations

- Charm quark mass definition ambiguity
  - dependence of $\mathcal{B}(\bar{B} \to X_s\gamma)^{theo}$ on m_c enters through the $\langle s\gamma|O_{1,2}|b\rangle$ which start contributing at $\mathcal{O}(\alpha_s)$
  - $m_c^{pole}/m_b^{pole} = 0.29 \pm 0.02$
    $\mathcal{B}(\bar{B} \to X_s\gamma)^{theo} = (3.32 \pm 0.30) \times 10^{-4}$
  - $\overline{m}_c(m_b/2)/m_b^{pole} = 0.22 \pm 0.04$
    $\mathcal{B}(\bar{B} \to X_s\gamma)^{theo} = (3.70 \pm 0.30) \times 10^{-4}$

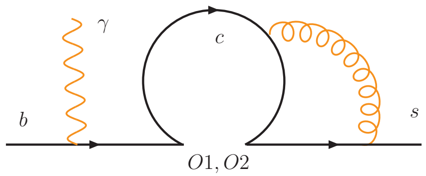


[Aubertetal 02]



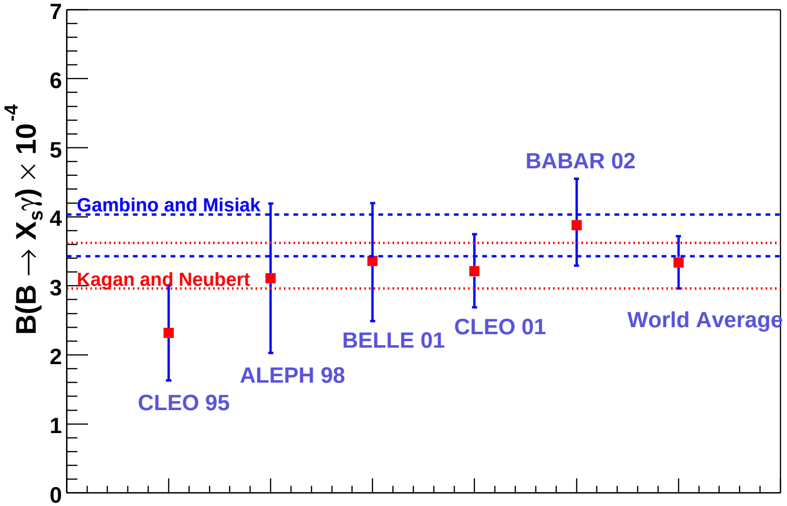
Motivations

- Charm quark mass definition ambiguity
  - dependence of $\mathcal{B}(\bar{B} \to X_s \gamma)^{theo}$ on m_c enters through the $\langle s\gamma | O_{1,2} | b \rangle$ which start contributing at $\mathcal{O}(\alpha_s)$



[Aubertetal 02]

  - $m_c^{pole}/m_b^{pole} = 0.29 \pm 0.02$
    $\mathcal{B}(\bar{B} \to X_s \gamma)^{theo} = (3.32 \pm 0.30) \times 10^{-4}$
  - $\overline{m}_c(m_b/2)/m_b^{pole} = 0.22 \pm 0.04$
    $\mathcal{B}(\bar{B} \to X_s \gamma)^{theo} = (3.70 \pm 0.30) \times 10^{-4}$



- difference between using $\overline{m}_c(\mu)$ and m_c^{pole} is a NNLO effect in the branching ratio
  $\Longrightarrow$ resolving the ambiguity requires going to the NNLO level

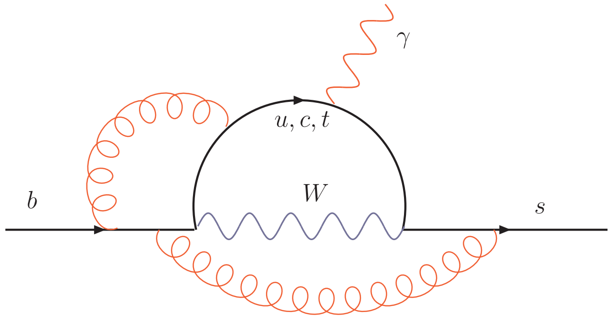
Theoretical framework

- diagrams involve scales with large hierarchy

  $M_W\,, M_t \gg m_b \gg m_s \implies$ large $\log\left(\frac{M_W^2}{m_b^2}\right)$

  $\longrightarrow$ resummation of $\alpha_s \log\left(\frac{M_W^2}{m_b^2}\right)$ is necessary

  using RG techniques



- start by introducing an effective theory without the heavy fields

$$\mathcal{L}_{\rm eff} = \mathcal{L}_{\rm QCD\times QED}(u,d,s,c,b) + \frac{4G_F}{\sqrt{2}} V_{ts}^* V_{tb} \sum_i C_i(\mu)\, O_i(\mu)$$

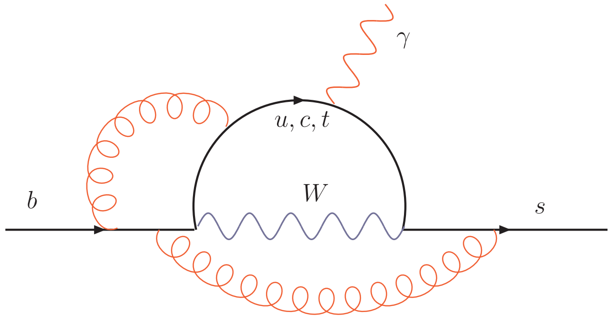
Theoretical framework

- diagrams involve scales with large hierarchy

  $M_W, M_t \gg m_b \gg m_s \Longrightarrow$ large $\log\left(\frac{M_W^2}{m_b^2}\right)$

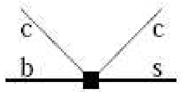
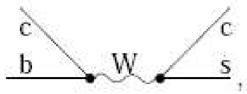
  $\longrightarrow$ resummation of $\alpha_s \log\left(\frac{M_W^2}{m_b^2}\right)$ is necessary

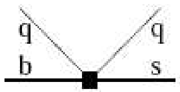
  using RG techniques

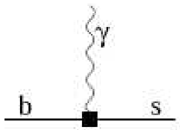


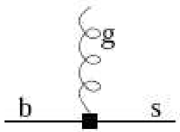
- start by introducing an effective theory without the heavy fields

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{QCD}\times\text{QED}}(u,d,s,c,b) + \frac{4G_F}{\sqrt{2}} V_{ts}^* V_{tb} \sum_i C_i(\mu)\, O_i(\mu)$$

$O_{1,2} =$  $= (\bar{s}\Gamma_i c)(\bar{c}\Gamma'_i b)$, **from** , $|C_i(m_b)| \sim 1$

$O_{3,4,5,6} =$  $= (\bar{s}\Gamma_i b)\sum_q(\bar{q}\Gamma'_i q)$, $|C_i(m_b)| < 0.07$

$O_7 =$  $= \frac{e m_b}{16\pi^2}\, \bar{s}_L \sigma^{\mu\nu} b_R F_{\mu\nu}$, $C_7(m_b) \simeq -0.3$

$O_8 =$  $= \frac{g m_b}{16\pi^2}\, \bar{s}_L \sigma^{\mu\nu} T^a b_R G^a_{\mu\nu}$, $C_8(m_b) \simeq -0.15$

Theoretical framework

Calculation done in three steps:

- Matching find the Wilson coefficients $C_i(\mu)$ by comparing the full and the effective theory at the mass scale $\mu \approx M_W$
  $\Rightarrow$ no large logarithms and only vacuum diagrams

- Mixing compute the anomalous dimensions of the operators and solve the renormalization group equations to go down with the Wilson coefficients to $\mu \approx m_b$

$$\frac{d}{d\mu} C_j(\mu) = C_i(\mu)\gamma_{ij}(\mu)$$

- Matrix elements calculate the matrix elements of all the operators at $\mu \approx m_b \Rightarrow$ no large logarithms as no heavy masses are present

Current state of the art for NNLO corrections

1. Matching
   - 2-loop matching for (O_1, ...,O_6) [Bobeth,Misiak,Urban 00]
   - 3-loop matching for O_7 and O_8 [Misiak,Steinhauser 04]
2. Mixing
   - 3-loop: $(O_1, \ldots, O_6)$ and $(O_7,\ O_8)$ sectors [Gorbahn,Haisch 05] [Gorbahn,Haisch,Misiak 05]
   - 4-loop $(O_1, \ldots, O_6) \longrightarrow (O_7,\ O_8)$ [Czakon,Haisch,Misiak 06]

Current state of the art for NNLO corrections

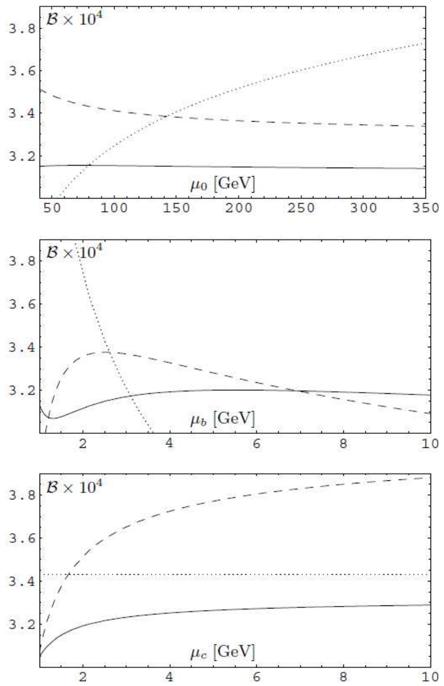
1. Matching
   - 2-loop matching for ($O_1, \ldots, O_6$) [Bobeth,Misiak,Urban 00]
   - 3-loop matching for O_7 and O_8 [Misiak,Steinhauser 04]
2. Mixing
   - 3-loop: $(O_1, \ldots, O_6)$ and $(O_7,\, O_8)$ sectors [Gorbahn,Haisch 05] [Gorbahn,Haisch,Misiak 05]
   - 4-loop $(O_1, \ldots, O_6) \longrightarrow (O_7,\, O_8)$ [Czakon,Haisch,Misiak 06]
3. Matrix elements
   - $O_1,\, O_2,\, O_7,\, O_8$ large β_0 [Bieri,Greub,Steinhauser 03]
   - O_7 [Blokland,Czarnecki,Misiak,Slusarczyk,Tkachov 05] [Asatrian,Hovhannisyan,Poghosyan,Ewerth,Greub,Hurth 06]
   - O_7, photon spectrum [Melnikov,Mitov 05] [Asatrian,Ewerth,Ferroglia,Gambino,Greub 06]
   - O_1, O_2 leading term for $m_c \gg m_b$ [Misiak,Steinhauser 06]

The NNLO estimated Branching Ratio

$$\mathcal{B}(\bar{B} \to X_s \gamma)^{\rm theo}_{E_\gamma > 1.6\,{\rm GeV}} = (3.15 \pm 0.23) \times 10^{-4}$$

[Misiak et al 06] [Misiak,Steinhauser 06]

- Decomposition of Uncertainty
  - non-perturbative 5% $\mathcal{O}(\alpha_s \Lambda / m_b)$
  - parametric 3% $\alpha_s(M_Z)$, $\mathcal{B}^{exp}_{SL}$, $m_c \ldots$
  - $\boldsymbol{m_c}$ interpolation 3% ($O_{1,2}$ matrix elements)
  - higher order 3% (μ_b, μ_c, μ_0 dependence)

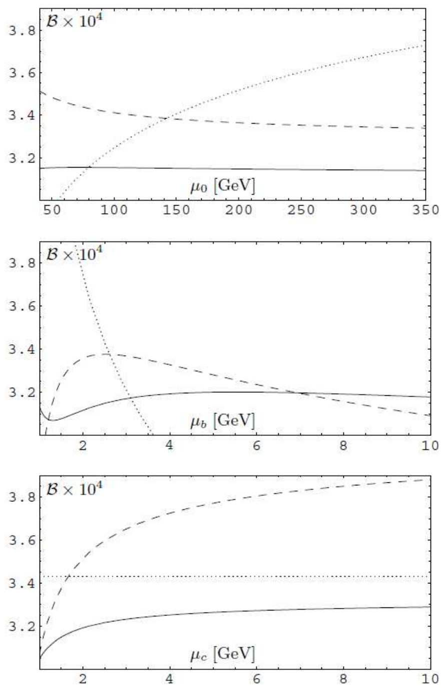


The NNLO estimated Branching Ratio

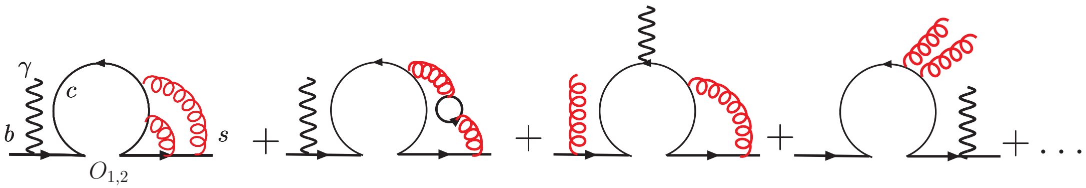
$$\mathcal{B}(\bar{B} \to X_s\gamma)^{\rm theo}_{E_\gamma > 1.6\,{\rm GeV}} = (3.15 \pm 0.23) \times 10^{-4}$$

[Misiak et al 06] [Misiak,Steinhauser 06]

- Decomposition of Uncertainty
  - non-perturbative 5% $\mathcal{O}(\alpha_s \Lambda/m_b)$
  - parametric 3% $\alpha_s(M_Z)$, $\mathcal{B}^{exp}_{SL}$, $m_c \ldots$
  - m_c interpolation 3% ($O_{1,2}$ matrix elements)
  - higher order 3% (μ_b, μ_c, μ_0 dependence)



- source of the interpolation uncertainty is the missing $\boldsymbol{\mathcal{O}}\,(\boldsymbol{\alpha_s^2})$ correction to $\langle s\gamma|O_{1,2}|b\rangle$

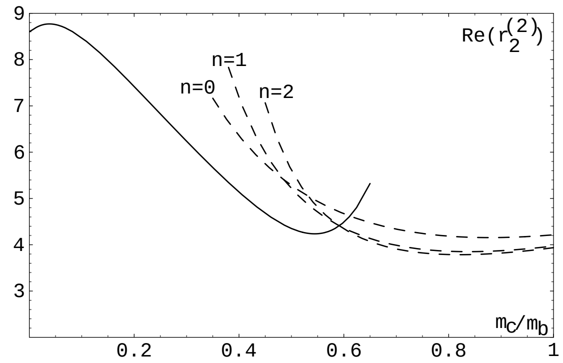


More about the interpolation uncertainty

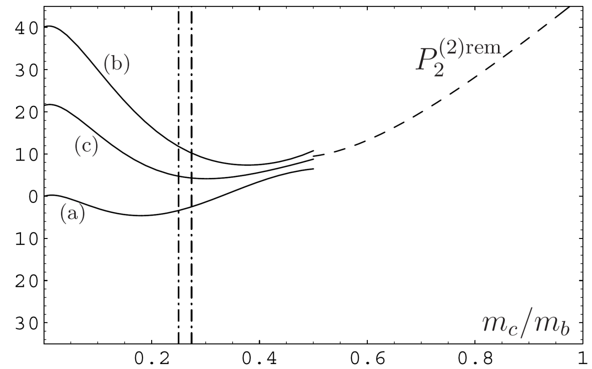
- $\mathcal{O}\left(\alpha_s^2\right)$ perturbative contribution to $\mathcal{B}(\bar{B} \to X_s\gamma)$: $\qquad P_2^{(2)} = \sum_{i,j=1}^{8} C_i^{(0)} C_j^{(0)} \left(n_f\, A_{ij} + B_{ij}\right)$

- using large β_0 approx. $\qquad P_2^{(2)} = \sum_{i,j=1}^{8} C_i^{(0)} C_j^{(0)} \left(\frac{-3}{2}\beta_0\, A_{ij} + B'_{ij}\right) = P_2^{(2),\beta_0} + P_2^{(2),rem}$
  - $P_2^{(2),\beta_0}$ known for $\langle s\gamma|O_{1,2,7,8}|b\rangle$

- expansions in limits $m_c/m_b \to 0$ and $m_c \gg m_b$ match nicely for Re$\langle s\gamma|O_2|b\rangle^{\beta_0}$
- good approximation already for $n = 0$
- no large $c\bar{c}$ threshold effects at $m_c = m_b/2$



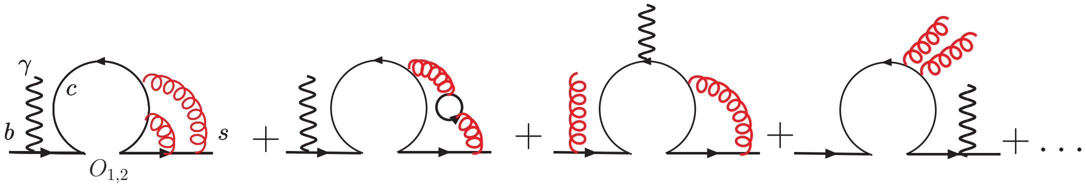
- calculate the leading term of large m_c expansion for $P_2^{(2),rem}$ and interpolate to physical m_c [Misiak & Steinhauser 06]
- making assumptions for $P_2^{(2),rem}$ at $m_c = 0$ is the source of the interpolation uncertainty



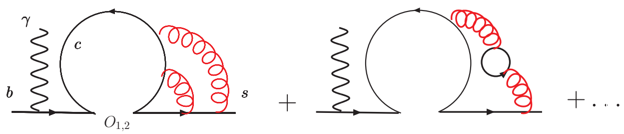
Reducing the overall uncertainty of $\mathcal{B}(\bar{B} \to X_s\gamma)^{\rm theo,NNLO}_{E_\gamma > 1.6\,{\rm GeV}}$

- removing the interpolation uncertainty

$\Longrightarrow$ need a complete calculation of $\langle s\gamma|O_{1,2}|b\rangle$ at $m_c \neq 0$



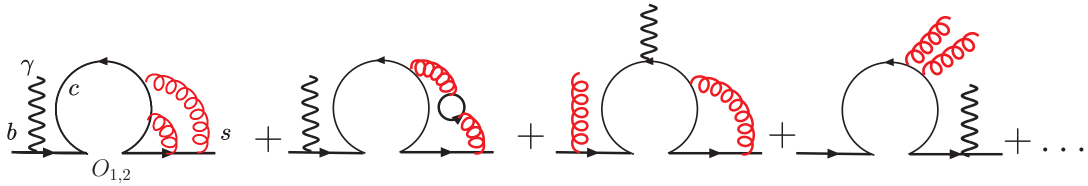
$\longrightarrow$ working on the virtual part [R. B, Czakon, Schutzmeier]



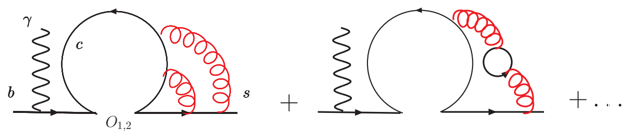
Reducing the overall uncertainty of $\mathcal{B}(\bar{B} \to X_s\gamma)^{\text{theo,NNLO}}_{E_\gamma > 1.6\,\text{GeV}}$

- removing the interpolation uncertainty

$\Longrightarrow$ need a complete calculation of $\langle s\gamma|O_{1,2}|b\rangle$ at $m_c \neq 0$

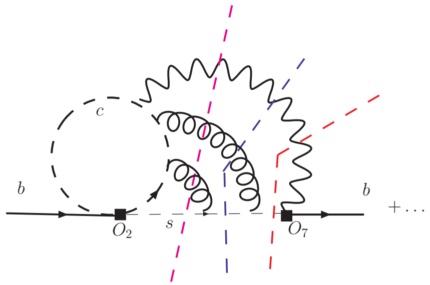


$\longrightarrow$ working on the virtual part [R. B, Czakon, Schutzmeier]



- reducing the interpolation uncertainty

$\Longrightarrow$ need a complete calculation of $\langle s\gamma|O_{1,2}|b\rangle$ at $m_c = 0$



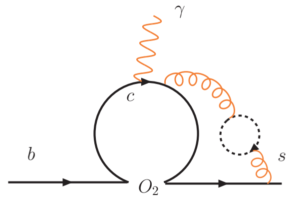
in progress [R. B, Czakon, Schutzmeier]

Removing the interpolation uncertainty: virtual part

- approx. 400 3-loop on-shell vertex diagrams with two scales m_b & m_c
- around 500 masters are involved in the bare amplitude
- symbolic reduction down to masters is not yet complete for the full 3-loop vertex
- $\mathcal{O}\left(\alpha_s^2 n_f\right)$ correction to $\langle s\gamma|O_{1,2}|b\rangle$: [R. B, Czakon, Schutzmeier 07]

Removing the interpolation uncertainty: virtual part

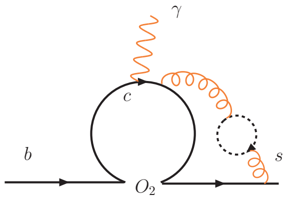
- approx. 400 3-loop on-shell vertex diagrams with two scales m_b & m_c
- around 500 masters are involved in the bare amplitude
- symbolic reduction down to masters is not yet complete for the full 3-loop vertex
- $\mathcal{O}\left(\alpha_s^2 n_f\right)$ correction to $\langle s\gamma|O_{1,2}|b\rangle$: [R. B, Czakon, Schutzmeier 07]

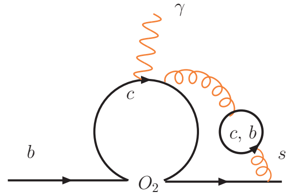


- masters were calculated with Mellin Barnes
  - first way: a numerical integration of the MB representations is performed for specific values of z using the MB package [MB : Czakon 05] ,
  - second way:
    - perform an expansion in $z = m_c^2/m_b^2$ by closing contours
    - coefficients of the expansion are given by at most a 1-dimensional MB integral expressed as a sum over residues
    - sum these infinite series using **XSummer** [Moch & Uwer 05]

Removing the interpolation uncertainty: virtual part

- approx. 400 3-loop on-shell vertex diagrams with two scales m_b & m_c
- around 500 masters are involved in the bare amplitude
- symbolic reduction down to masters is not yet complete for the full 3-loop vertex
- $\mathcal{O}\left(\alpha_s^2 n_f\right)$ correction to $\langle s\gamma|O_{1,2}|b\rangle$: [R. B, Czakon, Schutzmeier 07]





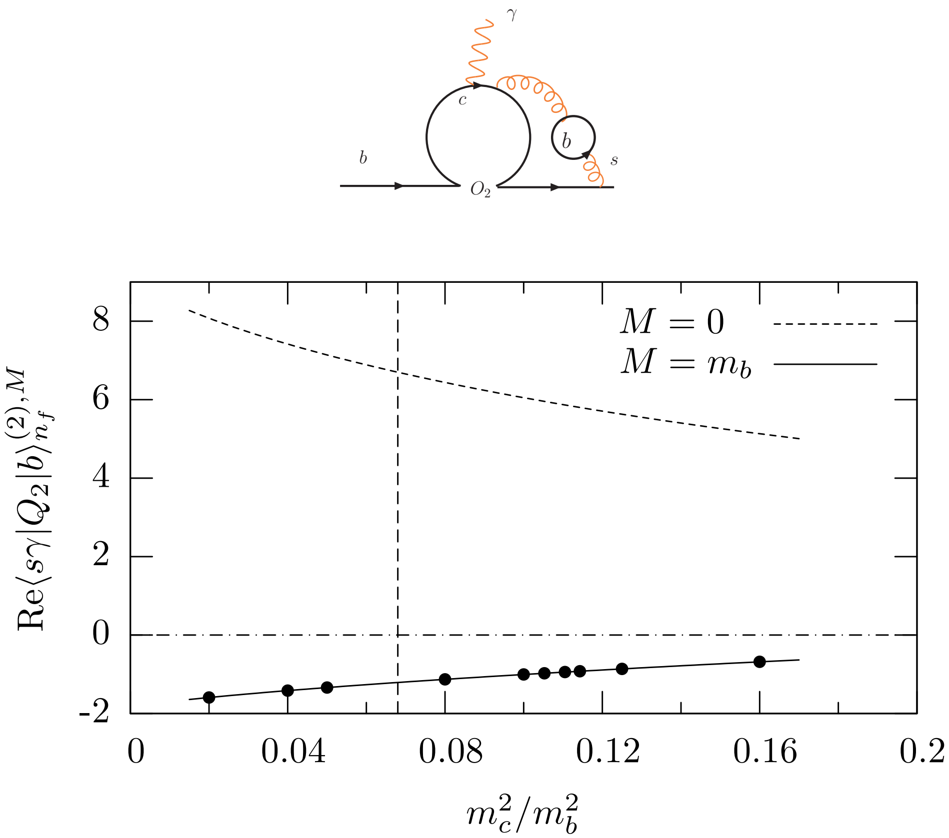
- masters were calculated with Mellin Barnes
  - first way: a numerical integration of the MB representations is performed for specific values of z using the MB package [MB : Czakon 05] ,
  - second way:
    - perform an expansion in $z = m_c^2/m_b^2$ by closing contours
    - coefficients of the expansion are given by at most a 1-dimensional MB integral expressed as a sum over residues
    - sum these infinite series using **XSummer** [Moch & Uwer 05]

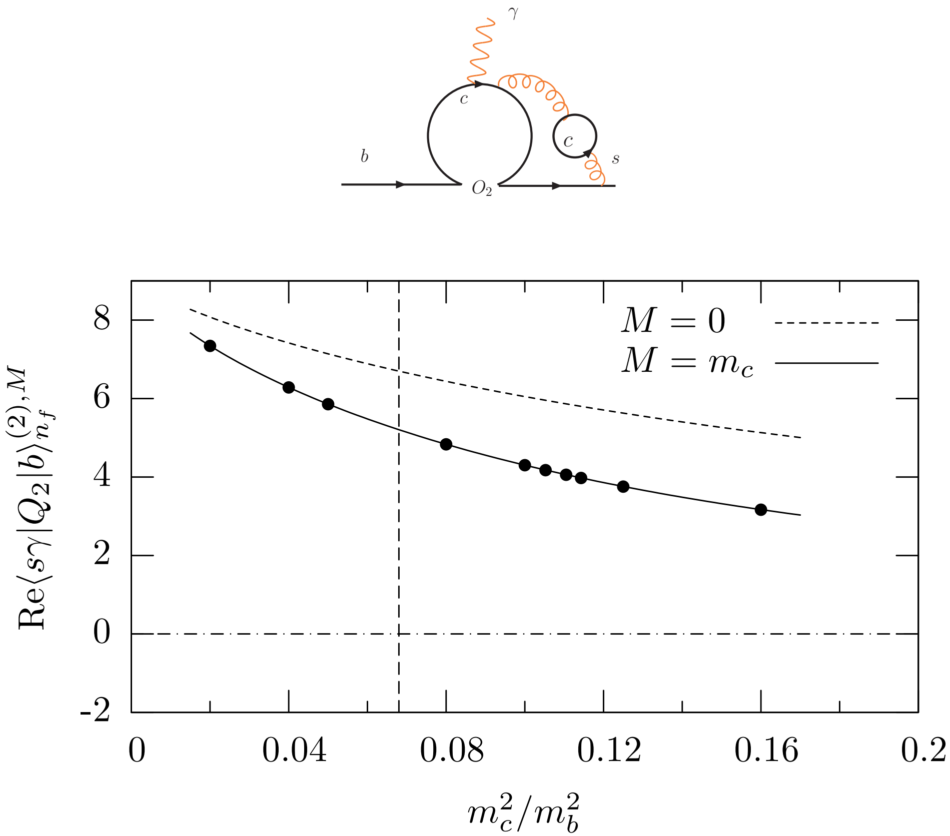
- MB alone was not enough to calculate all the masters due to poor convergence
- use differential equations solved numerically
  - boundaries were obtained using diagrammatic large mass expansion for $m_c \gg m_b$

    $\longrightarrow$ more about this method later

$\langle s\gamma|O_2|b\rangle_{\mathcal{O}(\alpha_s^2 n_f)}$

- Results for the massive fermionic contributions: [R. B, Czakon, Schutzmeier 07]



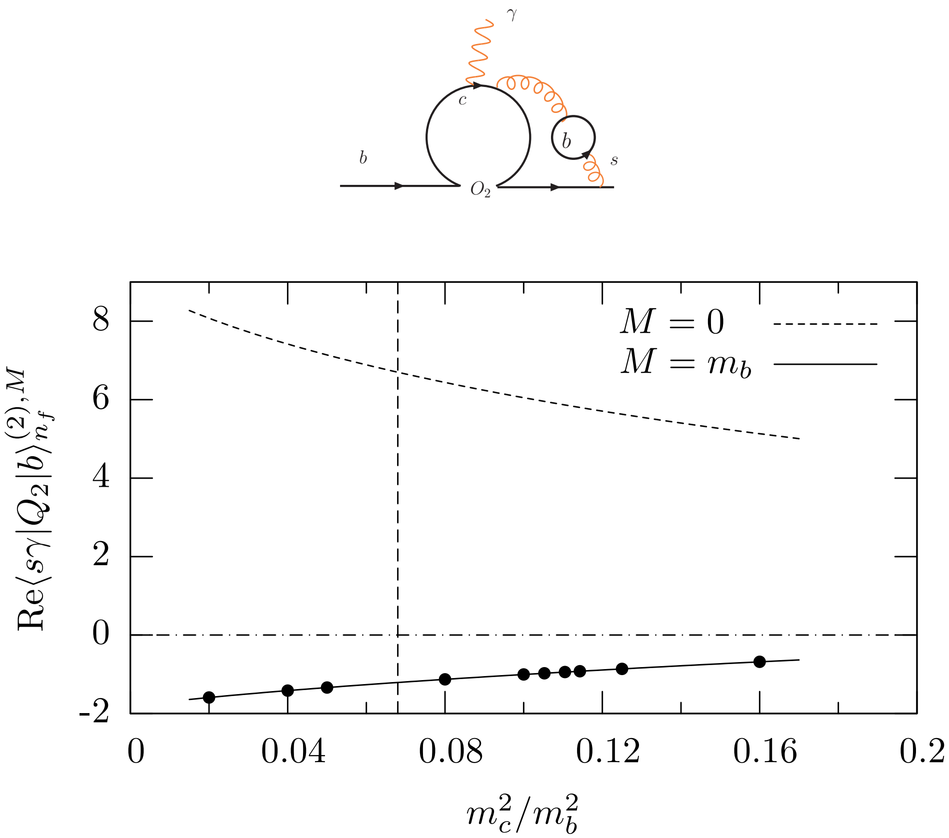


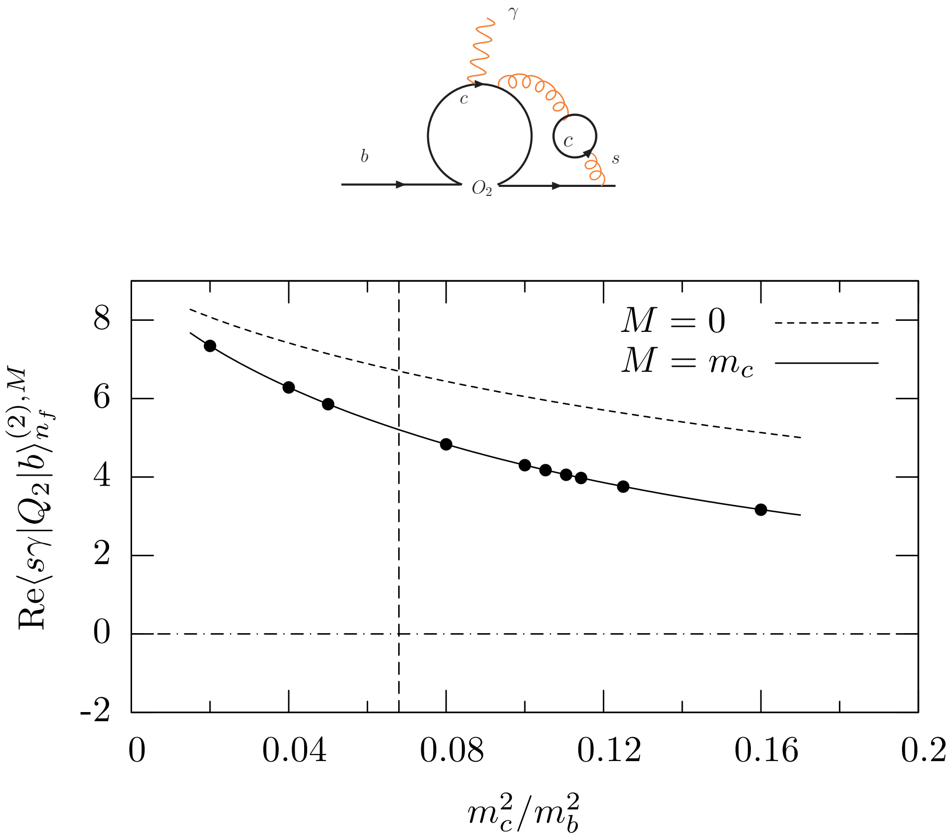
- massless approximation overestimates the massive b result and has the opposite sign !

- less pronounced differences for the c-quark $\longrightarrow$ moderate negative corrections wrt. massless approximation

$\langle s\gamma|O_2|b\rangle_{\mathcal{O}(\alpha_s^2 n_f)}$

- Results for the massive fermionic contributions: [R. B, Czakon, Schutzmeier 07]





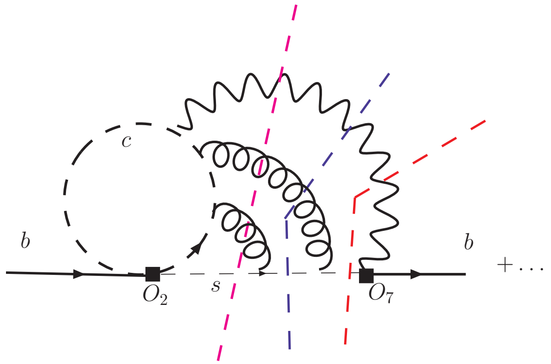
- massless approximation overestimates the massive b result and has the opposite sign !

- less pronounced differences for the c-quark $\longrightarrow$ moderate negative corrections wrt. massless approximation

numerical impact of the mass corrections on $\mathcal{B}(\bar{B} \to X_s\gamma)$ $= +$ 1.1% for $\mu_b = 2.5$ GeV

Reducing the interpolation uncertainty

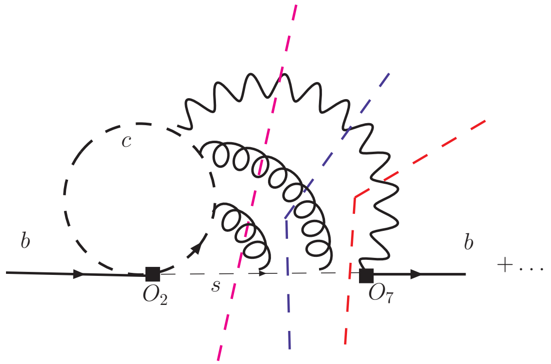
- calculating $\mathcal{O}\left(\alpha_s^2\right)$ correction to $\langle s\gamma|O_{1,2}|b\rangle$ at $m_c = 0$ helps significantly in reducing the interpolation uncertainty

$\Longrightarrow$ 

up to 4-particle cuts: $\gamma s,\ \gamma sg,\ \gamma sgg,\ \gamma sq\bar{q}$

Reducing the interpolation uncertainty

- calculating $\mathcal{O}\left(\alpha_s^2\right)$ correction to $\langle s\gamma|O_{1,2}|b\rangle$ at $m_c = 0$ helps significantly in reducing the interpolation uncertainty

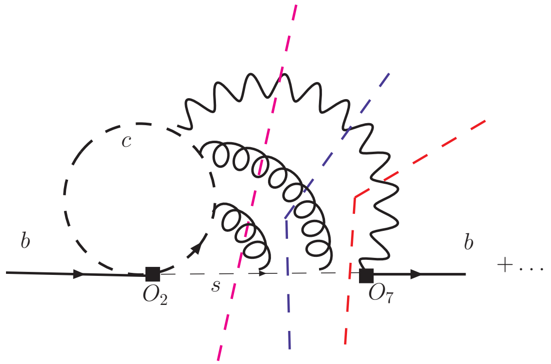
$\Longrightarrow$ 

up to 4-particle cuts: $\gamma s,\, \gamma sg,\, \gamma sgg,\, \gamma sq\bar{q}$

- 506 diagrams expressed through 42093 integrals
- around 300 master integrals have to be calculated BUT HOW ?

Reducing the interpolation uncertainty

- calculating $\mathcal{O}\left(\alpha_s^2\right)$ correction to $\langle s\gamma|O_{1,2}|b\rangle$ at $m_c = 0$ helps significantly in reducing the interpolation uncertainty

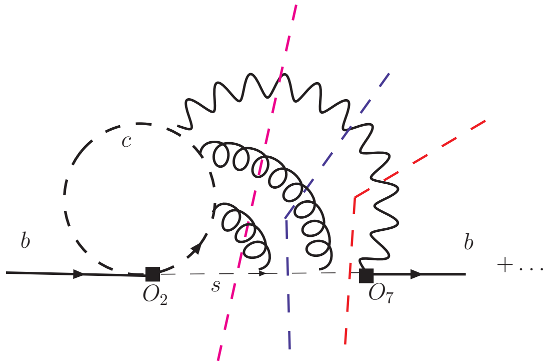
$\Longrightarrow$ 

up to 4-particle cuts: $\gamma s,\ \gamma sg,\ \gamma sgg,\ \gamma sq\bar{q}$

- 506 diagrams expressed through 42093 integrals
- around 300 master integrals have to be calculated   BUT HOW ?
  - sector decomposition: high precision results vs. running time . . .
  - differential equations for $p_b^2 \neq m_b^2$: needs boundaries . . .
  - Mellin Barnes: do we know how to use it for integrals with unitarity cuts ? dimension of the representations for 4-loop cut self energy integrals with up to 4 internal massive lines is an issue

Reducing the interpolation uncertainty

- calculating $\mathcal{O}\left(\alpha_s^2\right)$ correction to $\langle s\gamma|O_{1,2}|b\rangle$ at $m_c = 0$ helps significantly in reducing the interpolation uncertainty

$\Longrightarrow$ 

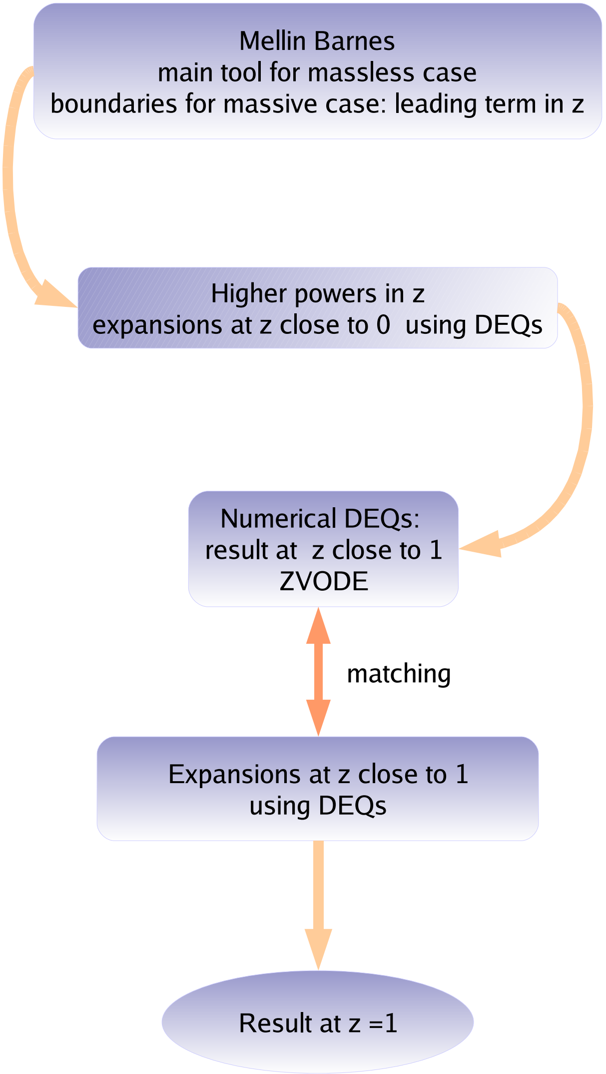
up to 4-particle cuts: $\gamma s,\ \gamma sg,\ \gamma sgg,\ \gamma sq\bar{q}$

- 506 diagrams expressed through 42093 integrals
- around 300 master integrals have to be calculated BUT HOW ?
  - sector decomposition: high precision results vs. running time . . .
  - differential equations for $p_b^2 \neq m_b^2$: needs boundaries . . .
  - Mellin Barnes: do we know how to use it for integrals with unitarity cuts ? dimension of the representations for 4-loop cut self energy integrals with up to 4 internal massive lines is an issue

so what is the way out ?

Combining methods

Merging methods is the way to go, but a long chain of steps:



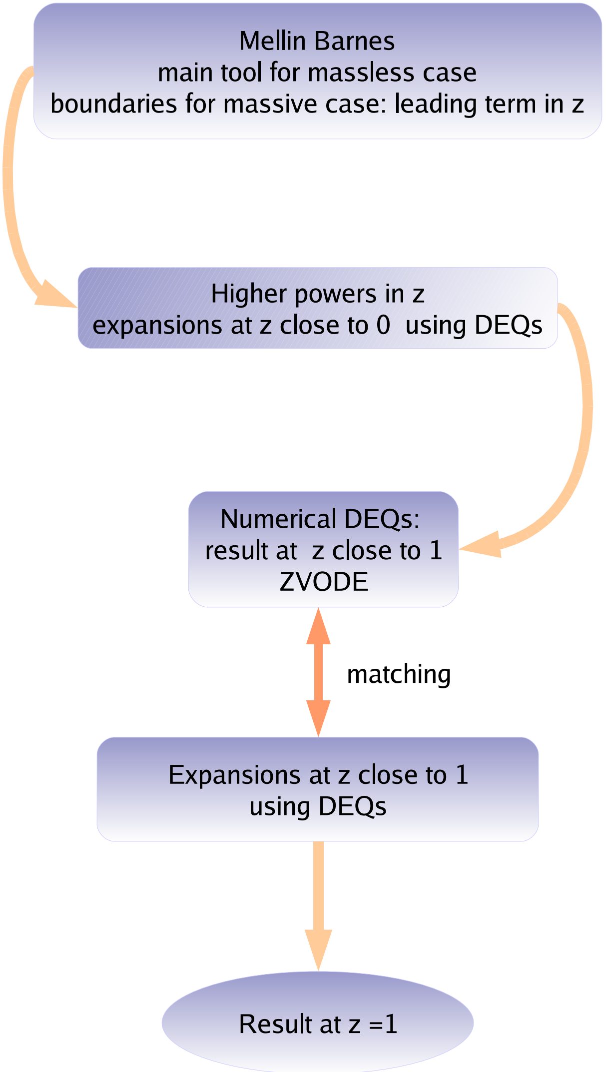
- evaluation of off-shell master integrals $V_i(z,\epsilon)$ with help of numerical differential equations (deqns) [Caffo, Czyz, Remiddi 2002]

$$\frac{d}{dz} V_i(z,\epsilon) = A_{ij}(z,\epsilon) V_j(z,\epsilon), \quad z = p_b^2/m_b^2$$

- Idea:
  - calculate integrals at some "simple" point (e.g. $p_b^2 \ll m_b^2$)
  - Integrate system of deqns starting at this limit up to the on-shell condition $z = 1$

Combining methods

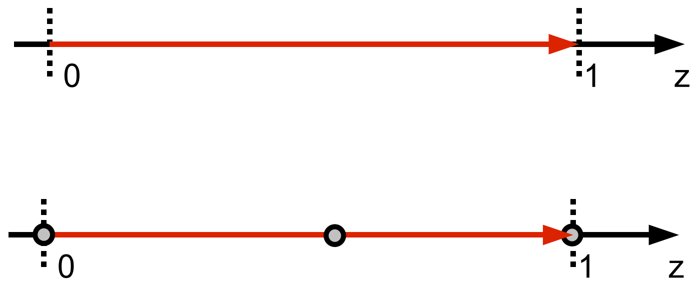
Merging methods is the way to go, but a long chain of steps:



- evaluation of off-shell master integrals $V_i(z, \epsilon)$ with help of numerical differential equations (deqns) [Caffo, Czyz, Remiddi 2002]

$$\frac{d}{dz} V_i(z, \epsilon) = A_{ij}(z, \epsilon) V_j(z, \epsilon), \quad z = p_b^2 / m_b^2$$

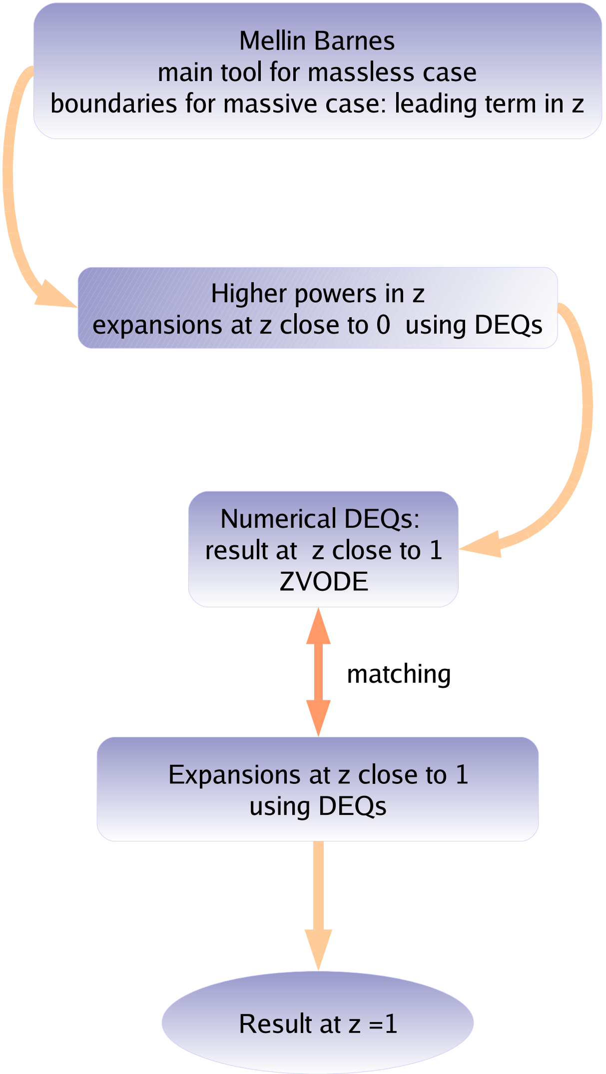
- Idea:
  - calculate integrals at some "simple" point (e.g. $p_b^2 \ll m_b^2$)
  - Integrate system of deqns starting at this limit up to the on-shell condition $z = 1$



$\rightarrow$ but:deqns singular in both endpoints! (and on naive contour $z \in \mathbb{R}$)
$\Rightarrow$ solution:combine expansions with numerical integration in complex plane

Combining methods

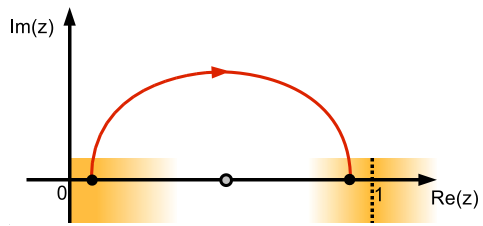
Merging methods is the way to go, but a long chain of steps:



- expand in ϵ and z in the limit $z \to 0$ with ansatz:

$$V_i(z, \epsilon) = \sum_{nmk} c^0_{inmk} \epsilon^n z^m \log^k z$$

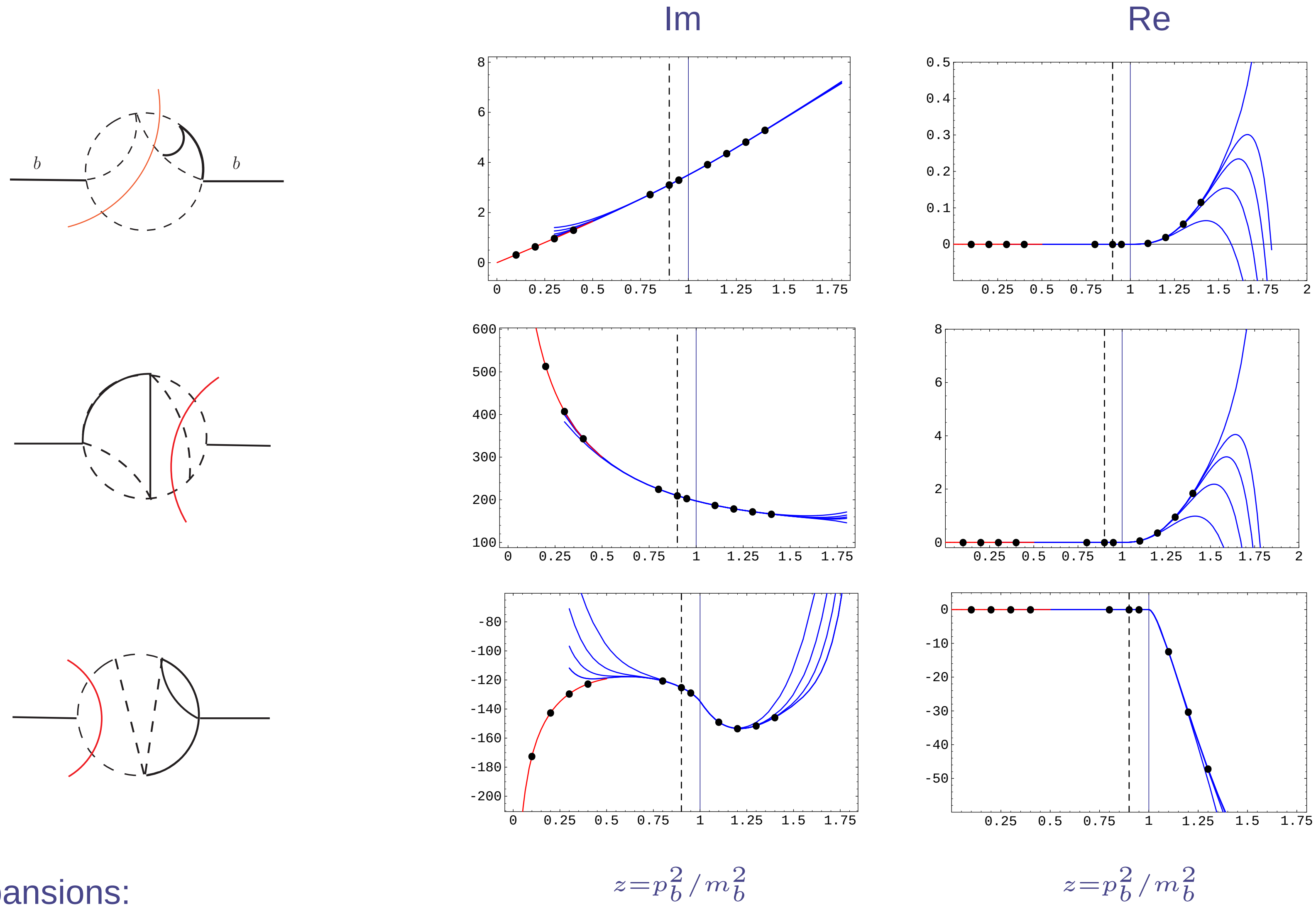
- solve recursively for c^0_{inmk} up to high powers in z
- boundary conditions:
  - Mellin Barnes $\&$ diagrammatic large-mass expansions for $p_b^2 \ll m_b^2$
    $\Rightarrow$ high precision values for $z \approx 0$
- use these values as starting point for numerical integration (in complex plane) up to $z \approx 1$ (zvode)



- perform another power logarithmic expansion around $z \to 1$ and solve coefficients c^1_{inmk} recursively
- use numerical integration to fix the remaining c^1_{inmk}

Some Results for 2- and 3-particle cuts

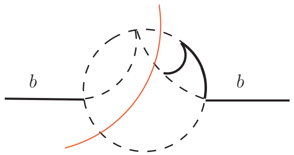
Preliminary results: sample masters with 2- and 3-particle cuts



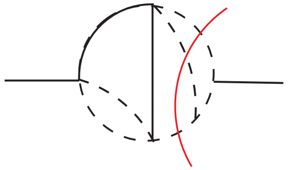
- Expansions:
  - $z \longrightarrow 0$: up to z^{18}
  - $z \longrightarrow 1$: up to $(1-z)^{12}$
- Numerical integration: starts at $z_0 = 0.02$
- Matching: done at $z_1 = 0.9$

Some Results for 2- and 3-particle cuts

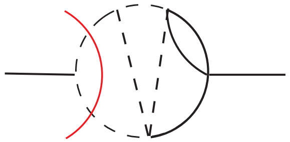
- Preliminary results at $z = 1$: sample masters with 2- and 3-particle cuts



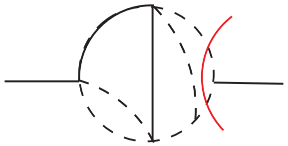
$$= \frac{3.10453\, i}{\varepsilon} + \mathcal{O}\left(1\right)$$



$$= \frac{3.14159\, i}{\varepsilon^3} + \frac{20.0142\, i}{\varepsilon^2} + \frac{77.1378\, i}{\varepsilon} + 209.713\, i + \mathcal{O}\left(\varepsilon\right)$$



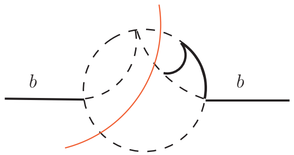
$$= \frac{-2.0944\, i}{\varepsilon^3} - \frac{12.5778\, i}{\varepsilon^2} - \frac{35.6402\, i}{\varepsilon} - 125.153\, i + \mathcal{O}\left(\varepsilon\right)$$



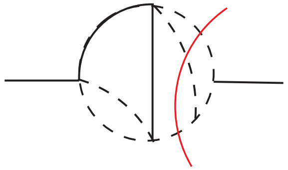
$$= \frac{-2.0944\, i}{\varepsilon^3} - \frac{4.91208\, i}{\varepsilon^2} - \frac{30.5699\, i}{\varepsilon} - 40.7068\, i + \mathcal{O}\left(\varepsilon\right)$$

Some Results for 2- and 3-particle cuts

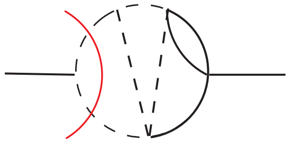
- Preliminary results at $z = 1$: sample masters with 2- and 3-particle cuts



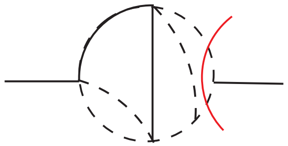
$$= \frac{3.10453\, i}{\varepsilon} + \mathcal{O}\left(1\right)$$



$$= \frac{3.14159\, i}{\varepsilon^3} + \frac{20.0142\, i}{\varepsilon^2} + \frac{77.1378\, i}{\varepsilon} + 209.713\, i + \mathcal{O}\left(\varepsilon\right)$$



$$= \frac{-2.0944\, i}{\varepsilon^3} - \frac{12.5778\, i}{\varepsilon^2} - \frac{35.6402\, i}{\varepsilon} - 125.153\, i + \mathcal{O}\left(\varepsilon\right)$$

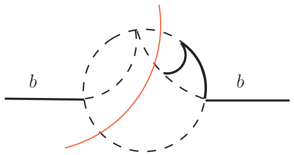


$$= \frac{-2.0944\, i}{\varepsilon^3} - \frac{4.91208\, i}{\varepsilon^2} - \frac{30.5699\, i}{\varepsilon} - 40.7068\, i + \mathcal{O}\left(\varepsilon\right)$$

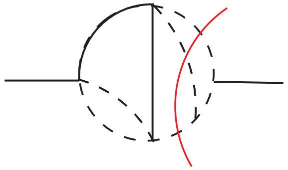
masters with 2-particle cuts are obtained with two independent calculations
$\longrightarrow$ cross checks will be done soon

Some Results for 2- and 3-particle cuts

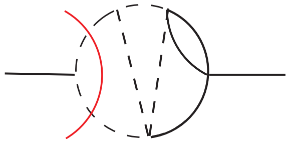
- Preliminary results at $z = 1$: sample masters with 2- and 3-particle cuts



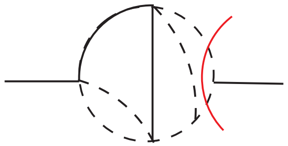
$$= \frac{3.10453\, i}{\varepsilon} + \mathcal{O}\left(1\right)$$



$$= \frac{3.14159\, i}{\varepsilon^3} + \frac{20.0142\, i}{\varepsilon^2} + \frac{77.1378\, i}{\varepsilon} + 209.713\, i + \mathcal{O}\left(\varepsilon\right)$$



$$= \frac{-2.0944\, i}{\varepsilon^3} - \frac{12.5778\, i}{\varepsilon^2} - \frac{35.6402\, i}{\varepsilon} - 125.153\, i + \mathcal{O}\left(\varepsilon\right)$$



$$= \frac{-2.0944\, i}{\varepsilon^3} - \frac{4.91208\, i}{\varepsilon^2} - \frac{30.5699\, i}{\varepsilon} - 40.7068\, i + \mathcal{O}\left(\varepsilon\right)$$

masters with 2-particle cuts are obtained with two independent calculations
$\longrightarrow$ cross checks will be done soon

- what we have:
  - masters with massless internal lines:
    all 2- and 3-particle cuts
    all 4-particle cuts but one
  - masters with b-quark internal lines:
    2- and 3-particle cuts are almost there

- still to be calculated: masters with 4-particle cuts and internal b-lines

Summary

- Matching current and future experimental precision for $\bar{B} \longrightarrow X_s\gamma$ decay necessitates NNLO corrections on the theory side
  crucial missing piece: $O(\alpha_s^2)$ correction to $\langle s\gamma|O_{1,2}|b\rangle$

- Reducing the interpolation uncertainty: needs $O(\alpha_s^2)$ correction to $\langle s\gamma|O_{1,2}|b\rangle$ at $m_c = 0$
  $\rightarrow$ 70% of the project is completed

- Removing the interpolation uncertainty: needs $O(\alpha_s^2)$ correction to $\langle s\gamma|O_{1,2}|b\rangle$ at physical m_c
  $\longrightarrow$ completed the fermionic contribution
  - $\rightarrow$ massless case: calculated in two ways and confirmed the findings of [Bieri, Greub, Steinhauser 03]
  - $\rightarrow$ massive case: impact on the branching ratio $+1.1\%$ for $\mu_b = 2.5$GeV [R. B, Czakon, Schutzmeier 07]

  $\longrightarrow$ bosonic contribution: work in progress