

# Late time stau dilution by Q balls

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Ref. SK & F.Takahashi, arXiv:0708nnnn

**Aug. 22, 2007, at COSMO-07, Sussex, UK**

# §1. Introduction

Gauge-med. ~~SUSY~~  $\implies$  No FCNC & ~~CP~~ problems.

$\hookrightarrow m_{3/2} \ll M_W \implies$  Gravitino LSP = good DM

If  $\tau_{\text{stau}}$  is NLSP  $\implies$  Long lifetime

$\implies M_p$  could be measurable for  $m_{3/2} \sim O(10)$  GeV.

Buchmüller et al. (04)

Collider experiments (LHC, ILC,...)

$\implies$  Affects BBN

- Energetic decay destroys light elements.

- Bound state ( ${}^4\text{He-stau}^-$ ) catalyze nucl. react.

${}^6\text{Li}$  is over-produced very much.

Pospelov (07)

Kohri, Takayama (07)

Kaplinghat, Rajaraman (06)  
and others...

Cosmological disaster

$\implies \tau_{\text{stau}} < 10^3 \text{ sec}$  or  $\Delta \sim 10^3$

Hamaguchi et al. (07)

# §1. Introduction (cont'd)

Late-time dilution is not so easy.

$T_D \gtrsim$  a few MeV Kawasaki et al.(99,00), Hannestad (04), Ichikawa et al. (05)

$$T_D < T_{fo} \sim 10 \text{ GeV} \left( \frac{m_{\tilde{\tau}}}{200 \text{ GeV}} \right)$$

$$T_D < T_{eq}$$

Example

$X$ : even R-parity

$$\begin{cases} m_X < m_{\tilde{\tau}} < m_{\tilde{X}} \\ X \text{ should couple to SM more strongly than stau.} \end{cases}$$

$\implies$  Very difficult, if not impossible.

Large Q ball

Long lifetime

$$m_{eff} \sim \omega_Q < m_{\tilde{\tau}}$$

Decay products are SM particles.

## §2. Q balls in GMSB

A Q ball is a kind of **non-topological soliton**, the energy min. configuration of the scalar field with **non-zero charge Q**. Coleman (85)

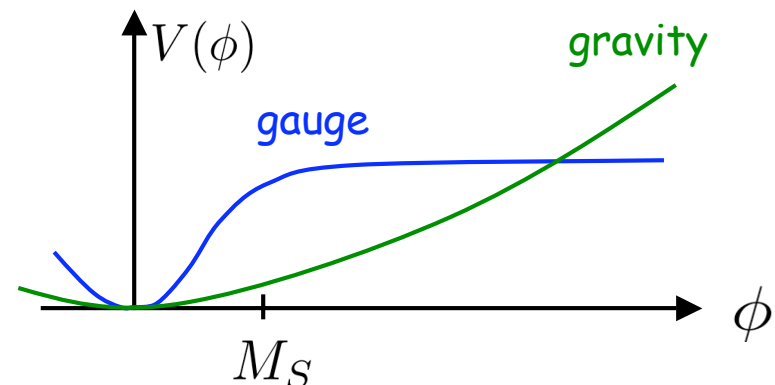
$\Phi$  : MSSM flat directions (e.g.,  $LH_u$ ,  $udd$ ,  $LLe$ , ...) Dine et al. (96)  
Gherghetta et al. (96)

$$V(\Phi) \simeq M_F^4 \left( \log \frac{|\Phi|^2}{M_S^2} \right)^2 + c_g m_{3/2}^2 \left[ 1 + K \log \frac{|\Phi|^2}{M^2} \right] |\Phi|^2 + \frac{\lambda^2 |\Phi|^{2(n-1)}}{M_P^{2(n-3)}} - c_H H^2 |\Phi|^2$$

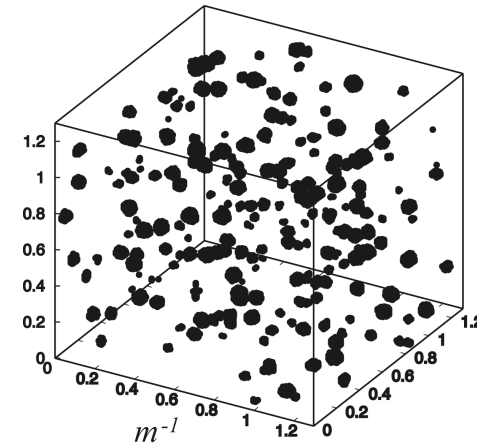
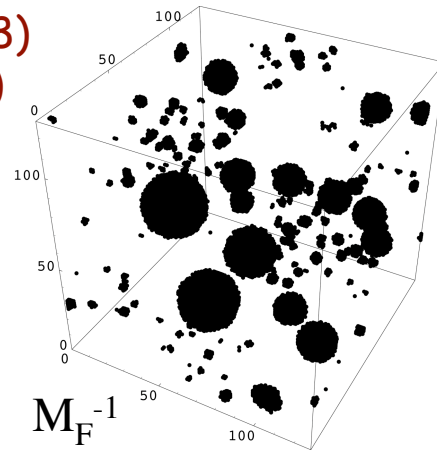
Gauge-med.
Gravity-med.
NR op.
H-induced

Q-ball formation Kusenko, Shaposhnikov (98)  
Enqvist, McDonald (98,99)  
SK, Kawasaki (00,01)

- (1) Large VEV during inflation, trapped by H-ind. & NR op.
- (2)  $\Phi$  starts rotation when  $H \sim m_{\phi, \text{eff}}$ .
- (3)  $\Phi$  feels spatial instabilities.
- (4) Instabilities go non-linear, and  $\Phi$  deforms into Q balls.



Kusenko, Shaposhnikov (98)  
 Enqvist, McDonald (98,99)  
 SK, Kawasaki (00,01)



Gauge-med. type

Gravity-med. (or New) type

Charge

$$Q = \beta \left( \frac{\phi_{osc}}{M_F} \right)^4$$

$$Q = \beta' \left( \frac{\phi_{osc}}{m_{3/2}} \right)^2$$

Size

$$R_Q \sim M_F$$

$\omega_Q < m_{\text{stau}}$  for large Q

$$|K|^{-1/2} m_{3/2}^{-1}$$

Mass

$$M_Q \sim M_F$$

( $m_{3/2} \sim 10 \text{ GeV}$ )

$$\sim m_{3/2} Q$$

Mass/Charge

$$\omega_Q \sim M_F Q^{-1/4}$$

$$\omega_Q \sim m_{3/2}$$

Decay rate

Cohen et al. (86)

$$\Gamma_Q \sim \frac{M_F}{48\pi} Q^{-5/4}$$

$$\Gamma_Q \sim \frac{m_{3/2}}{48\pi |K|} Q^{-1}$$

Long lifetime for large Q

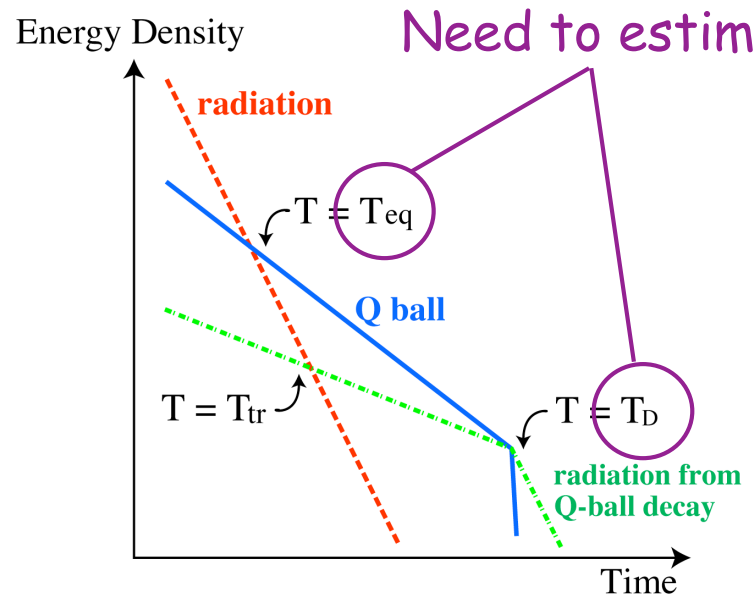
# §3. Dilution factor

$$\frac{n_{\tilde{\tau}}}{s} = \frac{1}{\Delta} \left( \frac{n_{\tilde{\tau}}}{s} \right)_{\text{thermal}}$$

Q-domi. starts after stau freeze-out: **Case A**  $H \sim \frac{T^2}{M_P}$

Q-domi. starts before stau freeze-out.  
 Radiation dominated by relic one: **Case B**  $H \sim \frac{(T^3 T_{eq})^{\frac{1}{2}}}{M_P}$

Radiation dominated by new one: **Case C**  $H \sim \frac{T^4}{T_D^2 M_P}$



$$\Delta \sim \begin{cases} \frac{T_{eq}}{T_D} & \text{(Case A : } T_{eq} < T_{fo}) \\ \frac{(T_{fo} T_{eq})^{\frac{1}{2}}}{T_D} & \text{(Case B : } T_{tr} < T_{fo} < T_{eq}) \\ \left( \frac{T_{fo}}{T_D} \right)^3 & \text{(Case C : } T_D < T_{fo} < T_{tr}) \end{cases}$$

$$T_{fo} \sim \frac{m_{\tilde{\tau}}}{20} \sim 10 \text{ GeV} \left( \frac{m_{\tilde{\tau}}}{200 \text{ GeV}} \right)$$

$$T_{tr} \sim (T_{eq} T_D^4)^{\frac{1}{5}}$$

# §4. Entropy production by the Q ball decay

Large Q, i.e., large  $\phi \implies$  Gravity (New) type Q ball

$\Phi$  starts oscillation when  $H \sim m_{3/2}$ .

$$\text{At } T=T_{RH}, \quad \left. \frac{\rho_Q}{\rho_{inf}} \right|_{RH} \sim \frac{m_{3/2}^2 \phi_{osc}^2}{m_{3/2}^2 M_P^2} \sim \left( \frac{\phi_{osc}}{M_P} \right)^2$$

$$\implies T_{eq} \sim T_{RH} \left( \frac{\phi_{osc}}{M_P} \right)^2$$

$$T_D \sim (\Gamma_Q M_P)^{\frac{1}{2}} \sim \left( \frac{m_{3/2} M_P}{48\pi |K|} \right)^{\frac{1}{2}} \beta^{-\frac{1}{2}} \frac{m_{3/2}}{\phi_{osc}}$$

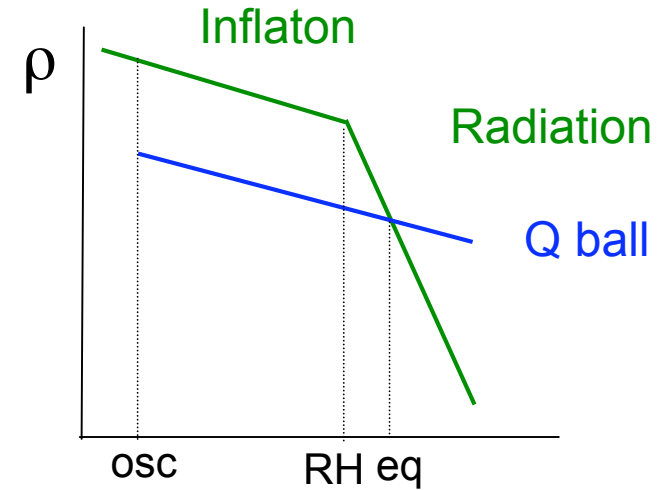
$$\implies \Delta \sim \frac{T_{eq}}{T_D} \sim 10^3 \left( \frac{T_{RH}}{8 \times 10^7 \text{ GeV}} \right) \left( \frac{\phi_{osc}}{5 \times 10^{14} \text{ GeV}} \right)^3$$

Dynamically,

$$\phi_{osc} \sim \left( \frac{H_{osc}}{\lambda M_P} \right)^{\frac{1}{n-2}} M_P$$

$n=7$  dddLL,  $\lambda \sim 10$

$n=6$  LLe or udd,  $\lambda \sim 0.002$



$$Q \sim 10^{24}$$

$$T_{eq} \sim 4 \text{ GeV}$$

$$T_D \sim 4 \text{ MeV}$$

## §5. Summary

Cosmic abundance of a long-lived charged particle such as a stau is tightly constrained by the catalyzed big bang nucleosynthesis.

One of the solutions is to dilute them by a huge entropy production.

We evaluate the dilution factor in the case when the freeze-out temperature is relatively low as in the stau NLSP scenario.

Q balls are the most promising source for this purpose.

Such Q balls are naturally produced in the gauge-med. ~~SUSY~~ scenario.