

New Methods for Feynman Integrals

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- Feynman integrals: basic notation, definitions and properties. A review of methods of evaluating Feynman integrals

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Feynman Integrals Calculus (Springer 2006)

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Slides and Mathematica files are available at

<http://science.sander.su>

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- Dimensional regularization. Properties of dimensionally regularized Feynman integrals
- A classification of methods of evaluating Feynman integrals

Perturbation theory. Feynman rules. A graph $\Gamma = \{\mathcal{V}, \mathcal{L}, \pi_{\pm}\}$ with vertices and lines (edges)

A given Feynman graph $\Gamma \rightarrow$ tensor reduction \rightarrow various scalar Feynman integrals that have the same structure of the integrand with various distributions of powers of propagators.

$$F_{\Gamma}(a_1, a_2, \dots) = \int \dots \int \frac{d^4 k_1 d^4 k_2 \dots}{(p_1^2 - m_1^2)^{a_1} (p_2^2 - m_2^2)^{a_2} \dots}$$

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Dimensional regularization: $d = 4 - 2\epsilon$; $d^4 k \rightarrow d^d k$

$k = (k_0, \vec{k}) = (k_0, k_1, k_2, k_3)$

k_1, k_2, \dots are loop momenta;

p_1, p_2, \dots are momenta of the lines; they are linear combinations of k_1, k_2, \dots and external momenta q_1, q_2, \dots

The propagator as a building block

$$\frac{1}{k^2 - m^2 + i0} = \lim_{\delta \rightarrow 0} \frac{1}{k^2 - m^2 + i\delta} ,$$
$$k^2 = k_0^2 - \vec{k}^2 = k_0^2 - k_1^2 - k_2^2 - k_3^2$$

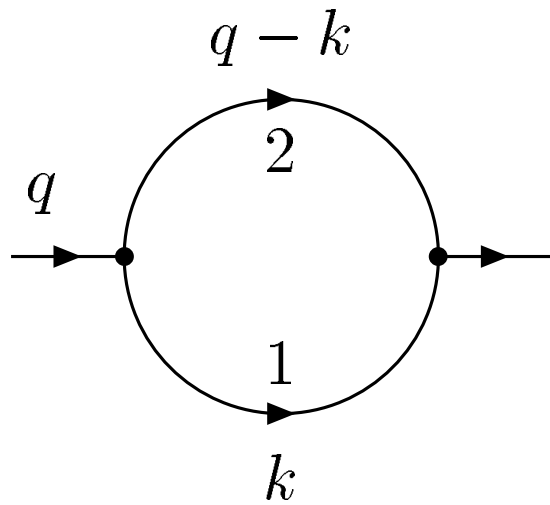
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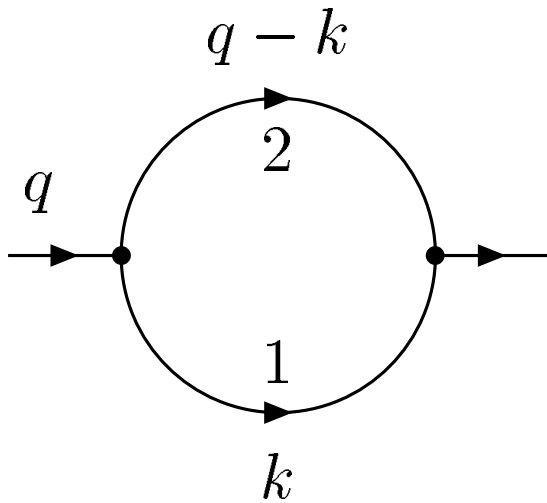
HQET, NRQCD, ... \rightarrow other types of propagators, e.g.

$$\frac{1}{v \cdot k \pm i0} , \quad v = (1, \vec{0})$$

For example,

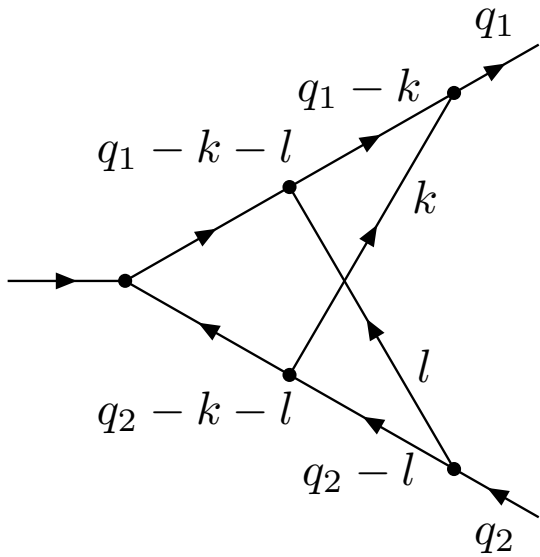


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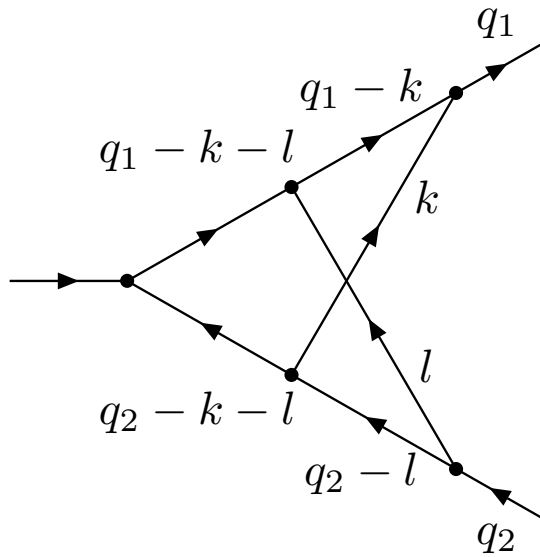


$$F_{\Gamma}(a_1, a_2, d) = \int \frac{\mathbf{d}^d k}{(m_1^2 - k^2 - i0)^{a_1} (m_2^2 - (q - k)^2 - i0)^{a_2}}$$

$$p_1^2 = p_2^2 = 0, \quad Q^2 = -(p_1 - p_2)^2 = 2p_1 \cdot p_2$$



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$$F_{\Gamma}(Q^2; a_1, \dots, a_6, d) = \int \int \frac{\mathbf{d}^d k \mathbf{d}^d l}{1 \left[(k+l)^2 - 2q_1 \cdot (k+l) \right]^{a_1}}$$

$$\times \frac{1}{\left[(k+l)^2 - 2q_2 \cdot (k+l) \right]^{a_2} (k^2 - 2q_1 \cdot k)^{a_3} (l^2 - 2q_2 \cdot l)^{a_4} (k^2)^{a_5} (l^2)^{a_6}}$$

UV, IR and collinear divergences \rightarrow a regularization

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[G. 't Hooft and M. Veltman'72]

[C.G. Bollini and J.J. Giambiagi'72; P. Breitenlohner and D. Maison'77]

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Informally, use alpha parameters

$$\frac{1}{(-k^2 + m^2 - i0)^a} = \frac{e^{i\pi a}}{\Gamma(a)} \int_0^{\infty} \alpha^{a-1} e^{i(k^2 - m^2)\alpha} d\alpha$$

$$\frac{1}{(-v \cdot k - i0)^a} = \frac{e^{i\pi a}}{\Gamma(a)} \int_0^{\infty} \alpha^{a-1} e^{i(v \cdot k)\alpha} d\alpha$$

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\int_{-\infty}^{+\infty} \mathbf{d}^1 k e^{\pm i\alpha k^2} &= \sqrt{\frac{\pi}{\alpha}} e^{\pm i\pi/4}, \\
\int \mathbf{d}^4 k e^{i\alpha k^2} &= \int \int \int \int \mathbf{d}k_0 \mathbf{d}k_1 \mathbf{d}k_2 \mathbf{d}k_3 e^{i\alpha k_0^2 - i\alpha k_1^2 - i\alpha k_2^2 - i\alpha k_3^2} \\
&= \left(\frac{\pi}{\alpha}\right)^2 e^{i\pi/4 - 3i\pi/4} = -i \left(\frac{\pi}{\alpha}\right)^2,
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\int \mathbf{d}^d k e^{i\alpha k^2} &= \left(\frac{\pi}{\alpha}\right)^{d/2} e^{i\pi/4 - (d-1)i\pi/4} = e^{i\frac{\pi}{2}(1-d/2)} \left(\frac{\pi}{\alpha}\right)^d,
\end{aligned}$$

$$\int_{-\infty}^{+\infty} \mathbf{d}^1 k e^{\pm i \alpha k^2} = \sqrt{\frac{\pi}{\alpha}} e^{\pm i \pi / 4},$$

$$\begin{aligned} \int \mathbf{d}^4 k e^{i \alpha k^2} &= \int \int \int \int \mathbf{d}k_0 \mathbf{d}k_1 \mathbf{d}k_2 \mathbf{d}k_3 e^{i \alpha k_0^2 - i \alpha k_1^2 - i \alpha k_2^2 - i \alpha k_3^2} \\ &= \left(\frac{\pi}{\alpha}\right)^2 e^{i \pi / 4 - 3 i \pi / 4} = -i \left(\frac{\pi}{\alpha}\right)^2, \end{aligned}$$

$$\int \mathbf{d}^d k e^{i \alpha k^2} = \left(\frac{\pi}{\alpha}\right)^{d/2} e^{i \pi / 4 - (d-1) i \pi / 4} = e^{i \frac{\pi}{2} (1-d/2)} \left(\frac{\pi}{\alpha}\right)^d,$$

$$\int \mathbf{d}^d k e^{i(\alpha k^2 - 2q \cdot k)} = e^{i \frac{\pi}{2} (1-d/2)} \left(\frac{\pi}{\alpha}\right)^d e^{-i q^2 / \alpha}$$

Dimensional regularization:

when deriving alpha representations, apply this rule with
 $d = 4 - 2\epsilon$

$$\int \mathbf{d}^4 k e^{i(\alpha k^2 - 2q \cdot k)} = -i\pi^2 \alpha^{-2} e^{-iq^2/\alpha}$$

→

$$\int \mathbf{d}^d k e^{i(\alpha k^2 - 2q \cdot k)} = e^{i\pi(1-d/2)/2} \pi^{d/2} \alpha^{-d/2} e^{-iq^2/\alpha}$$

$$\begin{aligned}
\int \frac{\mathbf{d}^d k}{(-k^2)^{a_1} (-(q-k)^2)^{a_2}} &= \frac{e^{i\pi(a_1+a_2)}}{\Gamma(a_1)\Gamma(a_2)} \int_0^\infty \int_0^\infty \mathbf{d}\alpha_1 \mathbf{d}\alpha_2 \alpha_1^{a_1-1} \alpha_2^{a_2-1} \\
&\quad \times \int \mathbf{d}^d k e^{i[\alpha_1 k^2 + \alpha_2 (k^2 + 2q \cdot k + q^2)]} \\
&= \frac{e^{i\pi(a_1+a_2+1-d/2)/2} \pi^{d/2}}{\Gamma(a_1)\Gamma(a_2)} \int_0^\infty \int_0^\infty \mathbf{d}\alpha_1 \mathbf{d}\alpha_2 \frac{\alpha_1^{a_1-1} \alpha_2^{a_2-1}}{(\alpha_1 + \alpha_2)^{d/2}} e^{i\alpha_1 \alpha_2 q^2 / (\alpha_1 + \alpha_2)}
\end{aligned}$$

$\alpha_1 = \eta\xi$, $\alpha_2 = \eta(1 - \xi)$, with the Jacobian η , integrate over η and ξ

$$\int \frac{\mathbf{d}^d k}{(-k^2)^{a_1} [-(q-k)^2]^{a_2}} = i\pi^{d/2} \frac{G(a_1, a_2)}{(-q^2)^{a_1+a_2+\epsilon-2}} ,$$

$$G(a_1, a_2) = \frac{\Gamma(a_1 + a_2 + \epsilon - 2)\Gamma(2 - \epsilon - a_1)\Gamma(2 - \epsilon - a_2)}{\Gamma(a_1)\Gamma(a_2)\Gamma(4 - a_1 - a_2 - 2\epsilon)}$$

Graph $\Gamma \rightarrow$ dimensionally regularized Feynman integral

$$F_{\Gamma}(a_1 \dots, a_L; d) = \frac{e^{i\pi(a+h(1-d/2))/2} \pi^{hd/2}}{\prod_l \Gamma(a_l)} \\ \times \int_0^{\infty} \mathbf{d}\alpha_1 \dots \int_0^{\infty} \mathbf{d}\alpha_L \prod_l \alpha_l^{a_l-1} \mathcal{U}^{-d/2} e^{i\mathcal{V}/\mathcal{U} - i \sum m_i^2 \alpha_i},$$

where $a = \sum a_i$

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For a Feynman integral with $1/(m^2 - k^2 - i0)^{a_l}$ propagators,

$$\mathcal{U} = \sum_{\text{trees } T} \prod_{l \notin T} \alpha_l,$$

$$\mathcal{V} = \sum_{\text{2-trees } T} \prod_{l \notin T} \alpha_l (q^T)^2.$$

Alpha representation \rightarrow

- Mathematical proofs. (for Feynman integrals at Euclidean external momenta, $(\sum q_i)^2 < 0$)
Analysis of convergence.

[K. Hepp'66; P. Breitenlohner and D. Maison'77; E. Speer'68,'77]

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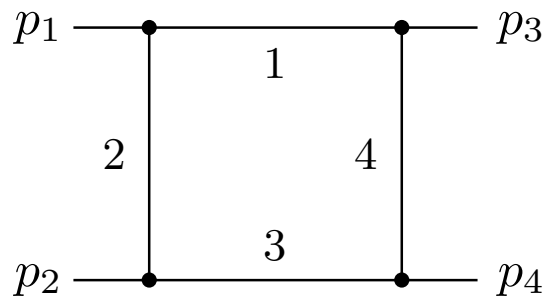
[E. Speer'77]

- A tool to evaluate Feynman integrals analytically.
- A tool to evaluate Feynman integrals numerically.
Sector decompositions

[T. Binoth and G. Heinrich'00; C. Bogner & S. Weinzierl'07; A.V. Smirnov & M.N. Tentyukov'08]

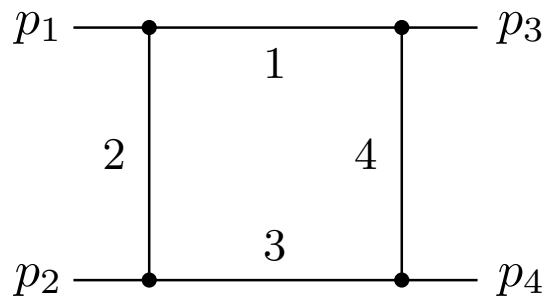
The massless box,

$$p_i^2 = 0, \quad i = 1, 2, 3, 4, \quad s = (p_1 + p_2)^2, \quad t = (p_1 + p_3)^2$$

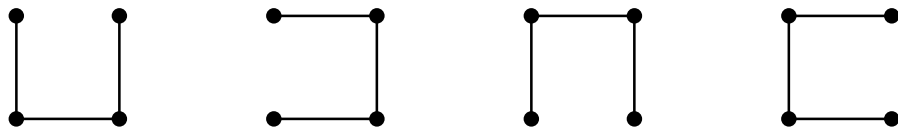


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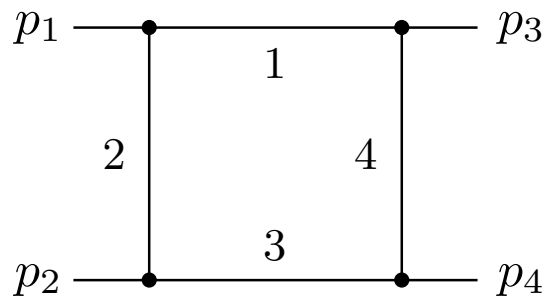


trees

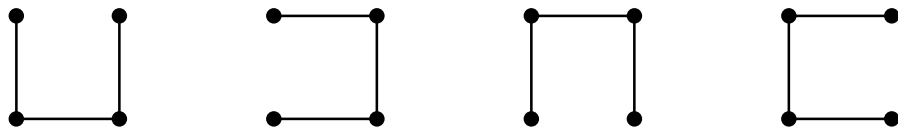


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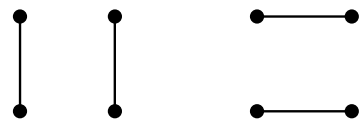
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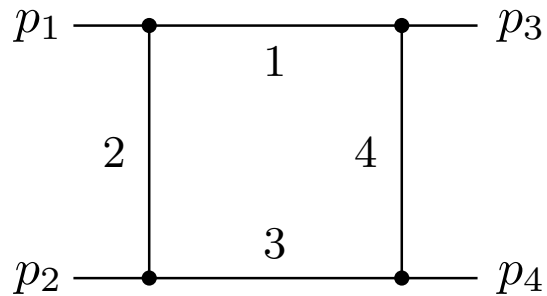


2-trees

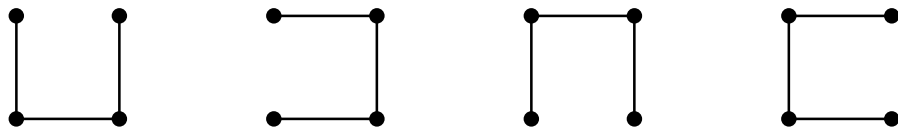


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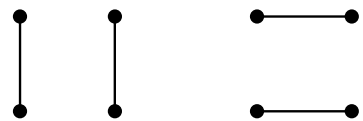
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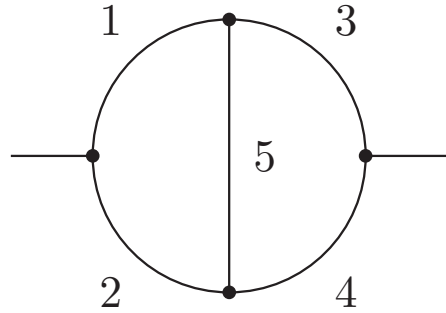
trees



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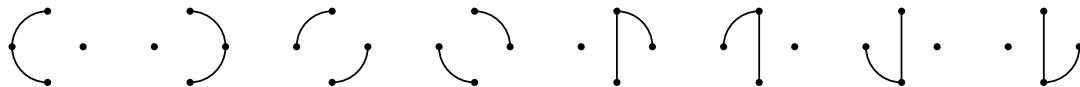
$$\mathcal{U} = \alpha_1 + \alpha_2 + \alpha_3 + \alpha_4, \quad \mathcal{V} = s\alpha_1\alpha_3 + t\alpha_2\alpha_4.$$



trees



2-trees



$$\mathcal{U} = (\alpha_1 + \alpha_2 + \alpha_3 + \alpha_4)\alpha_5 + (\alpha_1 + \alpha_2)(\alpha_3 + \alpha_4) ,$$

$$\begin{aligned} \mathcal{V} &= [(\alpha_1 + \alpha_2)\alpha_3\alpha_4 + \alpha_1\alpha_2(\alpha_3 + \alpha_4) + (\alpha_1 + \alpha_3)(\alpha_2 + \alpha_4)\alpha_5]q^2 \\ &\equiv \overline{\mathcal{V}}q^2 \end{aligned}$$

$\alpha = \eta \alpha'_l$, $l = 1, 2, \dots, L - 1$, $\eta = \sum_{l=1}^L \alpha_l$, integrate over η ,
introduce $\alpha'_L = 1 - \sum_{l=1}^{L-1} \alpha'_l$ by inserting an integration over
 α'_L with $\delta \left(\sum_{l=1}^L \alpha_l - 1 \right)$, replace α'_l by α_l :

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$$F_{\Gamma}(q_1, \dots, q_n; d; a_1 \dots, a_L) = \frac{\left(i\pi^{d/2}\right)^h \Gamma(a - hd/2)}{\prod_l \Gamma(a_l)}$$

$$\times \int_0^\infty \dots \int_0^\infty \delta\left(\sum_{l=1}^L \alpha_l - 1\right) \frac{\prod_l \alpha_l^{a_l-1} \mathcal{U}^{a-(h+1)d/2}}{\left(-\mathcal{V} + \mathcal{U} \sum m_l^2 \alpha_l\right)^{a-hd/2}} \mathbf{d}\alpha_1 \dots \mathbf{d}\alpha_L$$

$\alpha = \eta \alpha'_l$, $l = 1, 2, \dots, L - 1$, $\eta = \sum_{l=1}^L \alpha_l$, integrate over η , introduce $\alpha'_L = 1 - \sum_{l=1}^{L-1} \alpha'_l$ by inserting an integration over α'_L with $\delta\left(\sum_{l=1}^L \alpha_l - 1\right)$, replace α'_l by α_l :

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Cheng–Wu theorem:

$$\delta\left(\sum_{l=1}^L \alpha_l - 1\right) \rightarrow \delta\left(\sum_{l \in \nu} \alpha_l - 1\right) \rightarrow \delta\left(\alpha_l - 1\right)$$

$\alpha = \eta \alpha'_l$, $l = 1, 2, \dots, L - 1$, $\eta = \sum_{l=1}^L \alpha_l$, integrate over η , introduce $\alpha'_L = 1 - \sum_{l=1}^{L-1} \alpha'_l$ by inserting an integration over α'_L with $\delta\left(\sum_{l=1}^L \alpha_l - 1\right)$, replace α'_l by α_l :

$$F_{\Gamma}(q_1, \dots, q_n; d; a_1 \dots, a_L) = \frac{\left(i\pi^{d/2}\right)^h \Gamma(a - hd/2)}{\prod_l \Gamma(a_l)} \\ \times \int_0^\infty \dots \int_0^\infty \delta\left(\sum_{l=1}^L \alpha_l - 1\right) \frac{\prod_l \alpha_l^{a_l-1} \mathcal{U}^{a-(h+1)d/2}}{\left(-\mathcal{V} + \mathcal{U} \sum m_l^2 \alpha_l\right)^{a-hd/2}} \mathbf{d}\alpha_1 \dots \mathbf{d}\alpha_L$$

Cheng–Wu theorem:

$$\delta\left(\sum_{l=1}^L \alpha_l - 1\right) \rightarrow \delta\left(\sum_{l \in \nu} \alpha_l - 1\right) \rightarrow \delta\left(\alpha_l - 1\right)$$

Proof. Use $\eta = \sum_{l \in \nu} \alpha_l$ instead of $\eta = \sum_{l=1}^L \alpha_l$

Feynman parameters:

$$\frac{1}{(m_1^2 - p_1^2)^{a_1} (m_2^2 - p_2^2)^{a_2}} = \frac{\Gamma(a_1 + a_2)}{\Gamma(a_1)\Gamma(a_2)} \int_0^1 \frac{d\xi \xi^{a_1-1} (1-\xi)^{a_2-1}}{[(m_1^2 - p_1^2)\xi + (m_2^2 - p_2^2)(1-\xi)]^{a_1+a_2}}$$

Feynman parameters:

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$$\frac{1}{\prod A_l^{a_l}} = \frac{\Gamma(\sum a_l)}{\prod \Gamma(a_l)} \int_0^1 d\xi_1 \dots \int_0^1 d\xi_L \prod_l \xi_l^{a_l-1} \frac{\delta(\sum \xi_l - 1)}{(\sum A_l \xi_l)^{\sum a_l}}$$

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They are even better ;-)

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They are even better ;-)

Use integration by parts (**IBP**) and neglect surface terms:

$$\int \dots \int \left[\left(q_i \cdot \frac{\partial}{\partial k_j} \right) \frac{1}{(p_1^2 - m_1^2)^{a_1} (p_2^2 - m_2^2)^{a_2} \dots} \right] d^d k_1 d^d k_2 \dots = 0;$$

$$\int \dots \int \left[\left(k_i \cdot \frac{\partial}{\partial k_j} \right) \frac{1}{(p_1^2 - m_1^2)^{a_1} (p_2^2 - m_2^2)^{a_2} \dots} \right] d^d k_1 d^d k_2 \dots = 0.$$

Methods to evaluate Feynman integrals: analytical, numerical, semianalytical

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An old **straightforward** analytical strategy:

to evaluate, by some methods, every scalar Feynman
integral generated by the given graph.

The **standard** modern strategy:

to derive, without calculation, and then apply IBP identities between the given family of Feynman integrals as **recurrence relations**.

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The whole problem of evaluation→

- constructing a reduction procedure
- evaluating master integrals

Methods to evaluate master integrals:

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- method of differential equations
[A.V. Kotikov'91, E. Remiddi'97, T. Gehrmann & E. Remiddi'00]