



## Laboratoire de Physique Théorique-Orsay

### Leptogenesis

- 👉 Motivation, Requirements, Framework
- 👉 Neutrino mass → new physics, explains also baryogenesis?
- 👉 Seesaw mechanisms
- 👉 Leptogenesis and Consequences : some phenomenology
- 👉 Leptogenesis at low scales, how difficult?
- 👉 How to Disentangle between seesaws? Collider Physics?

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NuFlavour , 8-10 June 2009, Cosener's House

# Observations : Neutrino masses and BAU

👉 Need dynamical mechanism to generate Baryon Asymmetry of the Universe  $\eta_B$ :

Primordial abundances of light elements + CMB Anisotropies sensitive to  $\eta_B$

$$\eta_B^{\text{BBN}} = \frac{n_B}{n_\gamma} \Big|_{\text{BBN}} \approx (4.7 - 6.5) \times 10^{-10}, \quad \eta_B^{\text{CMB}} = \frac{n_B}{n_\gamma} \Big|_{\text{CMB}} \approx (6.21 \pm 0.16) \times 10^{-10}$$

## Baryo(Lepto)genesis requirements

Sakharov '67 <sup>a</sup>

1. Baryon (Lepton) number violation
2. C & CP Violation
3. Out of equilibrium processes

$$n_{B(L)} \neq n_{\bar{B}(\bar{L})}$$

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<sup>a</sup> 4th constraint :  $B - L$  must be violated if baryogenesis occurs  $>$  EW scale

👉 Need a high scale to generate small neutrino masses and mixings:

Recent data: atmosphere, Solar, Terrestrial expts.

Atmospheric $\nu$ and Reactor Data			Solar $\nu$	
$\Delta m_{\text{atm}}^2$ (eV $^2$ )	$\theta_{\text{atm}}$	$\theta_{13}$	$\Delta m_{\text{sol}}^2$ (eV $^2$ )	$\theta_{\text{sol}}$
$2.6 \pm 0.2 (0.6) \times 10^{-3}$	$43.3^{+4.3}_{-3.8}$	$0^{+5.2}_{-0.0}$	$7.9^{+0.27}_{-0.28} \times 10^{-5}$	$33.7 \pm 1.3$

Absolute mass scale $\nu$	
Oscillations	$m_\nu > \sqrt{\Delta m_{\text{atm}}^2} \sim 5 \times 10^{-2}$ eV
Cosmology	$\sum_i m_{\nu_i} < \sim 1$ eV
Tritium	$m_{\nu_e} \lesssim 2.2$ eV

- Majorana nature? **Lepton Number Violation** →  $(0\nu\beta\beta)$
- Low energy mixing and **Additional  $CP$  phases?**
  - Dirac  $CP$  Phase  $\delta$  → first need to know  $\theta_{13}$  with good accuracy
  - Majorana phases  $\alpha_{1,2}$ : no direct information →  $(0\nu\beta\beta)?$

# Little history

- 👉 1985: Kuzmin, Rubakov & Shaposhnikov, following the discovery of fast  $B + L$  violation at  $T > T_{EW}$   
“On the anomalous electroweak baryon number nonconservation in the early universe”
- 👉 1986: Fukugita & Yanagida, basic idea of leptogenesis  
“Baryogenesis without grand unification”
- 👉 1992-1996: a few remarkable works opened the way to quantitative leptogenesis, e.g.  
Luty ('92), “Baryogenesis via leptogenesis”,  
Covi, Roulet, Vissani ('96), “CP violating decays in leptogenesis scenarios”
- 👉 1993: Farrar, Shaposhnikov, EW baryogenesis?  
“Baryon asymmetry of the Universe in the MSM”
- 👉 1994: Gavela, Hernandez, Orloff, Pene **IMPOSSIBLE!**  
“Standard model CP violation and baryon asymmetry”
- 👉 > SK (2000): Evidence for neutrino oscillations + EW baryogenesis fails in SM and MSSM, fuel the interest on leptogenesis : **a flourishing study begins**

- Need a high scale to generate small neutrino masses and mixings
- Need dynamical mechanism to generate BAU



**AIM: find models that are simultaneously consistent with  
Neutrino masses and BAU**

- Favourite option: new physics at high scale  $M$
- What is this new physics? How does it manifest?
- At which scale? Which implications for low energy phenomenology?

$m_\nu \neq 0 \Rightarrow$  New Physics

- Lepton symmetry is accidental  $\implies$  Non-renormalisable operators dim 5, 6 ...

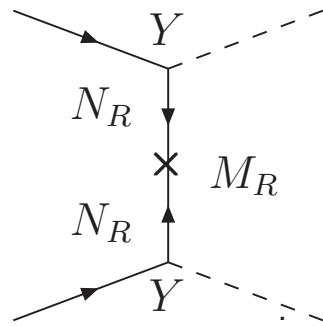
SM  $\equiv$  Effective theory of a larger one at a scale  $\Lambda$

$$O_5 = \frac{C_{d=5}}{\Lambda} \left\{ (\phi\ell)^T (\phi\ell) \right\} \xrightarrow{\langle\phi\rangle=v} m_\nu \sim C_{d=5}^2 v^2 / \Lambda \quad \Rightarrow \quad \Lambda \in [\text{TeV}, M_{\text{GUT}}]$$

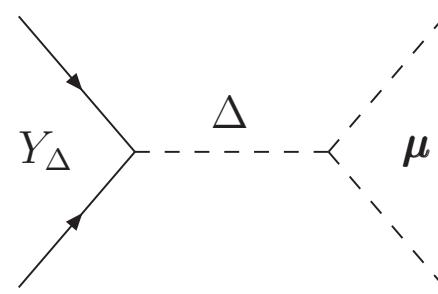
Typically 3 possible ways to generate  $m_\nu \neq 0$ :

- See-saw mechanism: towards GUT  $\rightarrow SO(10), \dots$  can be achieved via
  1. type I with RH neutrino exchange
  2. type II with scalar triplet exchange
  3. type III with fermionic triplet exchange
- Radiative corrections  $\rightarrow$  MSSM or various extensions  $+ R_p$ , Zee model...
- Extra dimensions  $\rightarrow$  alternative to the see-saw

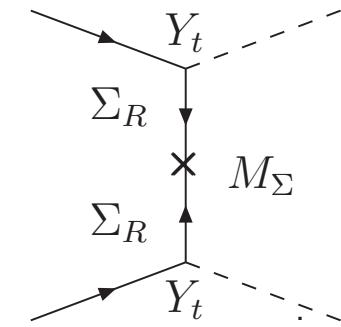
# 3 ways to generate tree level $\mathcal{O}^{d=5}$ : Seesaw I, II, III



type I (fermionic singlet)  
 $\mathbf{m}_\nu = -\frac{v^2}{2} Y_N^T \frac{1}{M_N} Y_N$

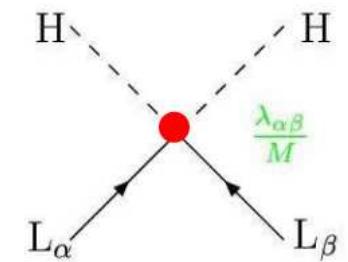


type II (scalar triplet)  
 $\mathbf{m}_\nu = -2v^2 Y_\Delta \frac{\mu_\Delta}{M_\Delta^2}$



type III (fermionic triplet)  
 $\mathbf{m}_\nu = -\frac{v^2}{2} Y_\Sigma^T \frac{1}{M_\Sigma} Y_\Sigma$

👉  $\mathcal{O}^{d=5}$  is common to all models of Majorana neutrinos



# Leptogenesis

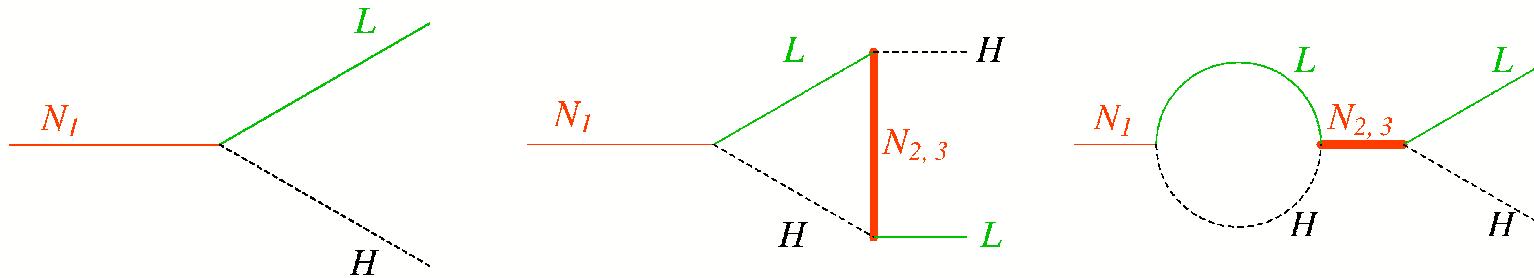
## Basic Leptogenesis Mechanism (Seesaw type I)

Fukugita and Yanagida '86

$$\mathcal{L} = \mathcal{L}_{SM} + \lambda_{Jk}^\nu \bar{L}_k \nu_{R_J} H - \frac{1}{2} \bar{\nu}_{R_J} M_{R_J} \nu_{R_J}^c + h_\alpha H^c \bar{e}_{R_\alpha} \ell_\alpha$$

1. CP violating decay of a heavy particle through an  $L$ -violating interaction can produce a lepton asymmetry
2. Lepton asymmetry transformed into  $BAU$  through sphaleron interactions:  
 $B + L$ -current is anomalous but  $B - L$  and  $1/3B - L_\alpha$  currents are not

$$\eta_B = - \left( \frac{24+4n_H}{42+9n_H} \right) \eta_L$$



1)  $L$  violation  $\leftarrow N_i$  Majorana

$$2) C \& CP \text{ asymmetry: } \epsilon_1 = \sum_{\alpha} \epsilon_{\alpha\alpha} = \frac{\sum_{\alpha} (\Gamma(N_1 \rightarrow H\ell_{\alpha}) - \Gamma(\bar{N}_1 \rightarrow \bar{H}\bar{\ell}_{\alpha}))}{\sum_{\alpha} (\Gamma(N_1 \rightarrow H\ell_{\alpha}) + \Gamma(\bar{N}_1 \rightarrow \bar{H}\bar{\ell}_{\alpha}))}$$

*CP from the complexity of the  $\lambda^{\nu}$  Yukawa matrix*

3) Out-of-equilibrium decay of  $N_1 \leftarrow \frac{\Gamma_{N_1}}{H(T=M_1)} \ll 1$

$\eta_L \propto K \epsilon_1$  is produced

4) Sphalerons act:  $\Delta L \rightarrow \Delta B$

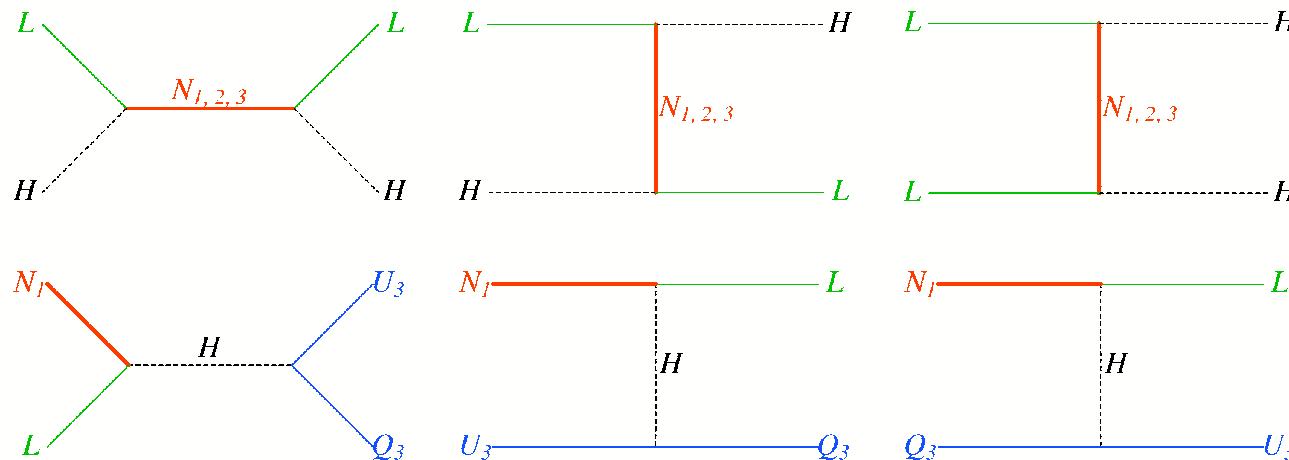
To see if it works , Solve B.E

## Boltzmann Equations 1 Flavour Approx

Lightest  $N_1$  produced in the thermal bath after inflation

$$\begin{aligned}\frac{dY_{N_1}}{dz} &= -\frac{z}{sH(M_1)} \left( \frac{Y_{N_1}}{Y_N^{eq}} - 1 \right) (\gamma_{D_1} + \gamma_{S_1}) \\ \frac{dY_L}{dz} &= -\frac{z}{sH(M_1)} \left[ \epsilon_1 \gamma_{D_1} \left( \frac{Y_{N_1}}{Y_N^{eq}} - 1 \right) - \gamma_w \frac{Y_L}{Y_L^{eq}} \right]\end{aligned}$$

- \*  $Y_{N_1} \equiv \frac{n_{N_1}}{s}$  is  $N_1$  abundance,  $Y_L \equiv \frac{n_\ell - n_{\bar{\ell}}}{s} = \sum_\alpha Y_L^{\alpha\bar{\alpha}}$  is total lepton asymmetry
- \*  $z = \frac{M_1}{T}$ ,  $s$ : entropy density,  $\epsilon_1$ : CP-asymmetry
- \*  $\gamma_{D_1}, \gamma_{S_1}$ : interaction rates for the decay and scattering  $\Delta L = 1$
- \*  $\gamma_w$  = function ( $\gamma_{D_1}, \gamma_{S_1}, \Delta L = 2$ ) washout factor

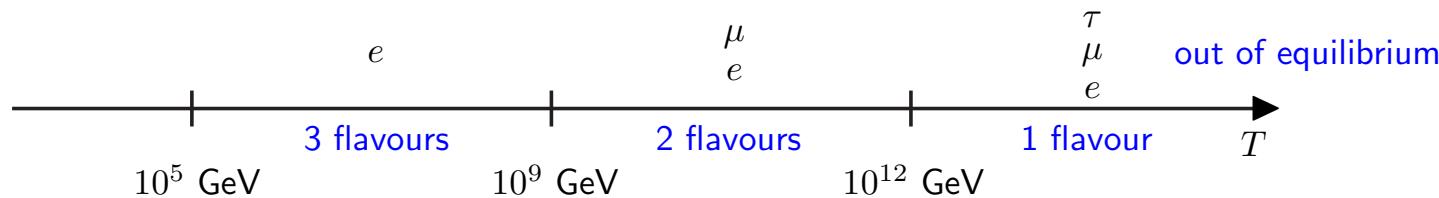


## First results: Constraints in 1 Flavour Approx

- $M_1 \ll M_2 \ll M_3 \rightarrow |\epsilon_1| \leq \frac{3\sqrt{2}}{16\pi} M_1 \frac{\sqrt{m_1^2 + m_2^2 + m_3^2}}{v^2}$ 
  - If **degenerate** light  $\nu \rightarrow \epsilon_1 = 0$
  - If **hierarchical** light  $\nu \rightarrow |\epsilon_1| \leq \frac{3\sqrt{2}}{16\pi} M_1 \frac{\sqrt{\Delta m_{\text{atm}}^2}}{v^2}$  [Davidson, Ibarra]
  - $\lambda \sim \mathcal{O}(1)$  &  $\kappa \sim 1$  and  $\eta_B^{\text{WMAP}} \rightarrow M_1 \gtrsim 6.4 \times 10^8$  GeV
  - $\eta_{B-L}^{\max} \rightarrow m \lesssim 0.15$  eV [Buchmuller, Di Bari, Plumacher]
- Quasi-degenerate  $M_J^2 - M_I^2 \ll M_J^2 \simeq M_I^2 \rightarrow |\epsilon_1| \neq 0 \rightarrow m \lesssim 1$  eV [Hambye et al]
- $M_1 = M_2 = M_3 \rightarrow \epsilon_1 = 0$

## Why should we worry about flavour?

- Thermal leptogenesis at  $T \simeq M_1$  when  $\Gamma_{N_1} < H(M_1)$
- One flavour approximation: correct when the interactions mediated by charged lepton Yukawa are out of equilibrium
- $\Gamma_\alpha \simeq 5 \times 10^{-3} h_\alpha^2 T$



- One flavour approx valid for  $M_1 > 10^{12}$  GeV,  
below, solve FLAVOURED Boltzmann eqs to generate  $Y_{L_{\alpha\beta}}$
- diagonal terms ( $\alpha\alpha$ ): lepton asymmetries stored in each flavour
- off-diagonal terms encode the quantum correlations between the flavour asymmetries

## One flavour

$$\frac{dY_L^{\alpha\alpha}}{dz} = -\frac{z}{sH(M_1)} \left[ \epsilon_{\alpha\alpha} \gamma_{D_1} \left( \frac{Y_{N_1}}{Y_N^{eq}} - 1 \right) - \gamma_W^{\alpha\alpha} \frac{Y_L^{\alpha\alpha}}{Y_L^{eq}} \right],$$

$$\epsilon_{\alpha\alpha} = \frac{\Gamma(N_1 \rightarrow H\ell_\alpha) - \Gamma(\bar{N}_1 \rightarrow \bar{H}\bar{\ell}_\alpha)}{\Gamma(N_1 \rightarrow H\ell_\alpha) + \Gamma(\bar{N}_1 \rightarrow \bar{H}\bar{\ell}_\alpha)}$$

## Sum over flavour

$$\begin{aligned} \sum_{\alpha} \frac{dY_L^{\alpha\alpha}}{dz} &= -\frac{z}{sH(M_1)} \left[ \underbrace{(\epsilon_{ee} + \epsilon_{\mu\mu} + \epsilon_{\tau\tau})}_{\epsilon_1} \gamma_{D_1} \left( \frac{Y_{N_1}}{Y_N^{eq}} - 1 \right) \right. \\ &\quad \left. - \frac{1}{Y_L^{eq}} \underbrace{(\gamma_W^{ee} Y_L^{ee} + \gamma_W^{\mu\mu} Y_L^{\mu\mu} + \gamma_W^{\tau\tau} Y_L^{\tau\tau})}_{\neq (\gamma_W^{ee} + \gamma_W^{\mu\mu} + \gamma_W^{\tau\tau}) Y_L} \right], \\ &\neq -\frac{z}{sH(M_1)} \left[ \epsilon_1 \gamma_{D_1} \left( \frac{Y_{N_1}}{Y_N^{eq}} - 1 \right) - \gamma_W \frac{Y_L}{Y_L^{eq}} \right] \text{ (One Flav Approx)} \end{aligned}$$

## Flavoured B.E

$Y_{\Delta\alpha} \Rightarrow B/3 - L_\alpha$  asymmetry in lepton flavour  $\alpha$

Baryonic asymmetry:  $Y_B \simeq \frac{12}{37} \sum_\alpha Y_{\Delta\alpha} \simeq -1.26 \times 10^{-3} \sum_\alpha \epsilon_\alpha \eta_\alpha$

$N_1$  Abundance:  $Y'_{N_1}(z) = -\kappa (D(z) + S(z)) (Y_{N_1}(z) - Y_{N_1}^{eq}(z)), \quad z = M_1/T$

$Y_{\Delta\alpha}$  asymmetry:

$$Y_{\Delta\alpha}'(z) = -\epsilon_\alpha \kappa (D(z) + S(z)) (Y_{N_1}(z) - Y_{N_1}^{eq}(z)) - \kappa_\alpha W(z) \sum_\beta A_{\alpha\beta} Y_{\Delta\beta}(z)$$

$Y_L$  Abundance:  $Y_{L_\alpha} = A_{\alpha\beta} Y_{\Delta\beta}; \quad A = - \begin{pmatrix} -151/179 & 20/179 & 20/179 \\ 25/358 & -344/537 & 14/537 \\ 25/358 & 14/537 & -344/537 \end{pmatrix}$

Washout parameters:  $\kappa_\alpha \equiv \frac{\Gamma_{N_1 \ell_\alpha}}{H(M_{N_1})} = \lambda_{1\alpha} \lambda_{1\alpha}^* \frac{v^2}{M_1 m^*} \equiv \frac{\tilde{m}_\alpha}{m^*}$

$$\kappa = \sum_\alpha \kappa_\alpha \equiv \frac{\tilde{m}}{m^*}, \quad m^* \simeq 1.08 \times 10^{-3} \text{ eV}$$

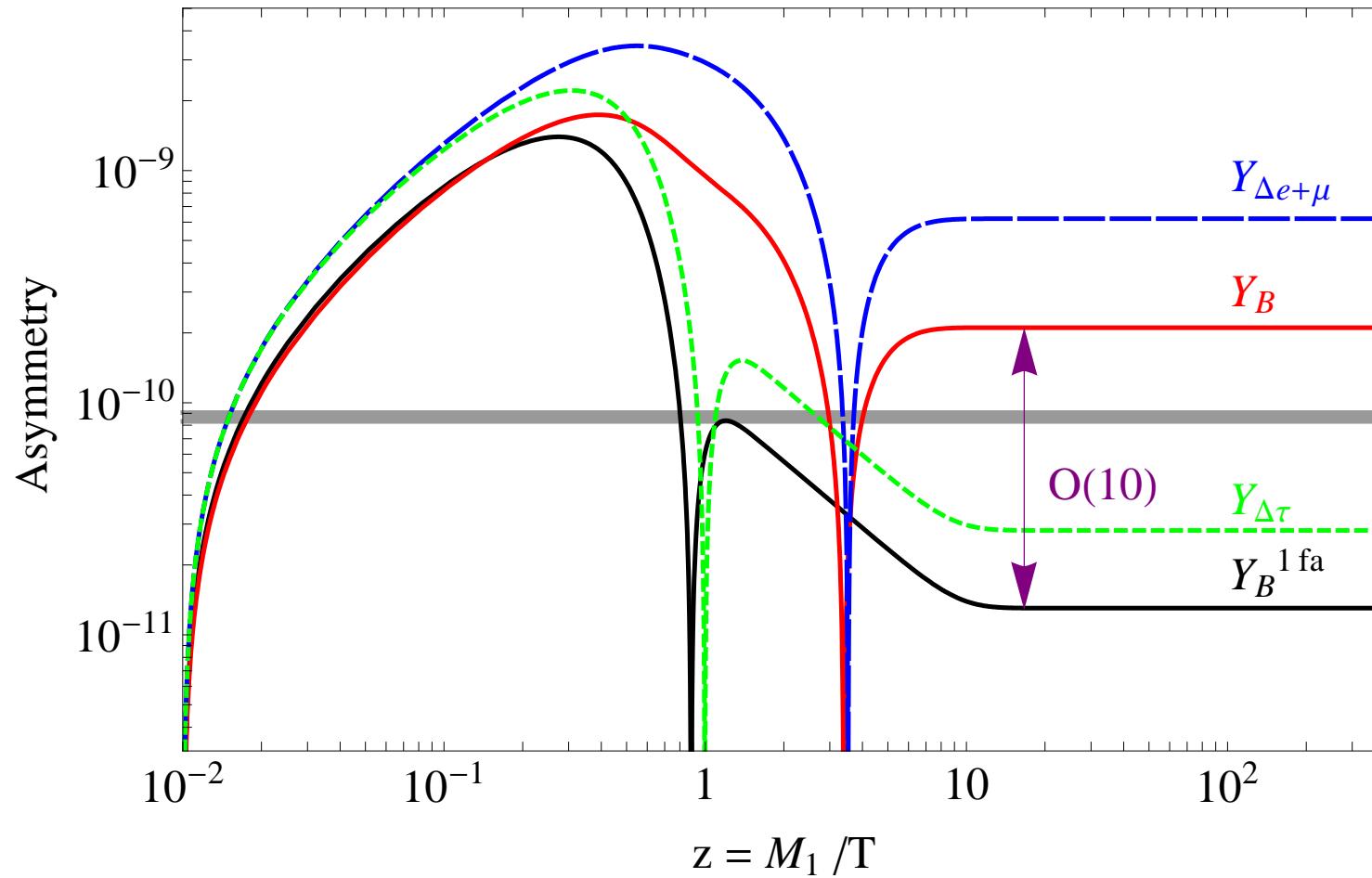
The CP-asymmetry:

$$\epsilon_\alpha = \frac{\Gamma_{N_1 \ell_\alpha} - \Gamma_{N_1 \bar{\ell}_\alpha}}{\sum_\alpha (\Gamma_{N_1 \ell_\alpha} + \Gamma_{N_1 \bar{\ell}_\alpha})} = \frac{1}{(8\pi)} \frac{1}{[\lambda \lambda^\dagger]_{11}} \sum_j \text{Im} \left\{ (\lambda_{1\alpha})(\lambda \lambda^\dagger)_{1j} (\lambda_{j\alpha}^*) \right\} g \left( \frac{M_j^2}{M_1^2} \right)$$

Efficiency  $\eta_\alpha$ : need to solve B.E

## Enhancing $\eta_B$

$\alpha = 1 \equiv ("e + \mu") \rightarrow$  strong washout  $\alpha = 2 \equiv (\tau) \rightarrow$  weak washout



Significant enhancement of BAU compared to One Flavour Approximation

# Consequences of treatment of flavours

With **Yukawa couplings**:  $\lambda = \frac{1}{v} \sqrt{M_N} \ R \ \sqrt{m_\nu} \ U^\dagger$  [Casas-Ibarra]

**CP asymmetry**

$$\begin{aligned}\epsilon_1 &= -\frac{3 m_{N_1}}{8 \pi v_2^2} \frac{\text{Im} \left[ \sum_\beta m_{\nu_\beta}^2 |R_{1\beta}|^2 \right]}{\sum_\delta m_{\nu_\delta} |R_{1\delta}|^2} \\ \epsilon_{1,\alpha} &= -\frac{3 m_{N_1}}{8 \pi v_2^2} \frac{\text{Im} \left[ \sum_{\beta\rho} m_{\nu_\beta}^{1/2} m_{\nu_\rho}^{3/2} (U_{\text{MNS}})_{\alpha\beta}^* (U_{\text{MNS}})_{\alpha\rho} R_{1\beta} R_{1\rho} \right]}{\sum_\delta m_{\nu_\delta} |R_{1\delta}|^2}\end{aligned}$$

**1 flav approx:** if light  $\nu$  are **degenerate** ( $m_i \simeq \bar{m}$ )  $\rightarrow \epsilon_1 = 0$

**flavoured:**  $\epsilon_1 = \sum_\alpha \epsilon_{1,\alpha} = 0$ ; individual  $\epsilon_{1,\alpha} \leq \frac{3 M_1 \bar{m}}{8 \pi v^2}$  grows with absolute scale  $\bar{m}$

$\epsilon_\alpha$  have opposite sign but are weighted by different washout factors

$\rightarrow$  Can have successful leptogenesis in the case of degenerate light neutrino spectrum

$\rightarrow$  Additional sources of *CP* violation

## Sources of $CP$ violation

$R$  and  $U_{\text{MNS}}$ !

Crucial:  $R \neq 1$

Input parameters - **Yukawa couplings**  $\lambda$ :

$$\lambda = \frac{1}{v} \sqrt{M_N} R \sqrt{m_\nu} U^\dagger \quad [\text{Casas-Ibarra}]$$

BAU from **flavoured** leptogenesis:  $(R(\theta_i), M_1, m_1, \theta_{13}, \delta, \varphi_1, \varphi_2)$

**Flavoured BAU** ( $3 \nu_R$ ): not sensitive to phases in  $U$  [Davidson, Garayoa, Palorini, Rius]

$\eta_B$  accounted for with  $R$  phases (any value of  $\theta_{13}, \delta, \varphi_1, \varphi_2$ )  
even if  $U$  phases measured,  $\eta_B$  can have any value

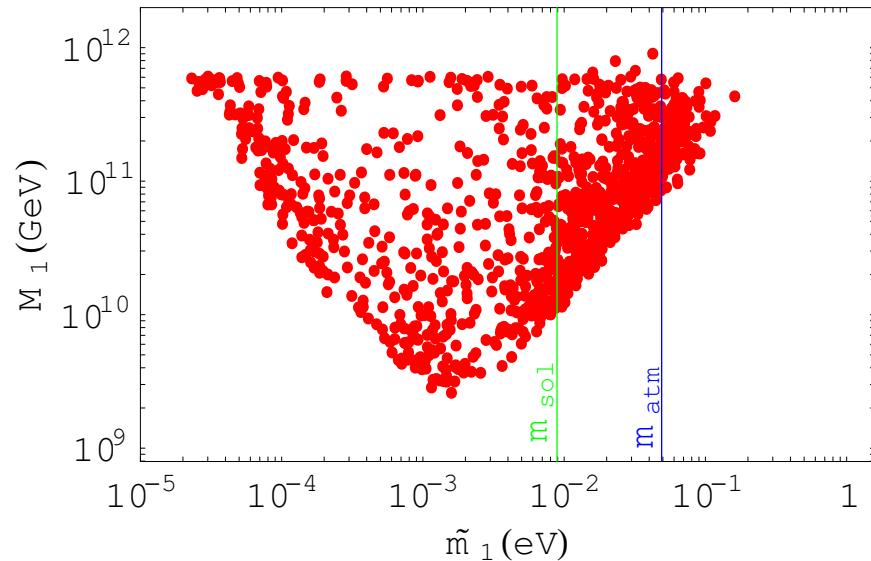
**BAU with Real  $R$**

SM: feasible [Pascoli, Petcov, Riotto, Branco, Gonzalez-Felipe, Joaquim...]

Disfavoured in SUSY seesaw ( $T_{RH}$ ) [Antusch, Teixeira ...]

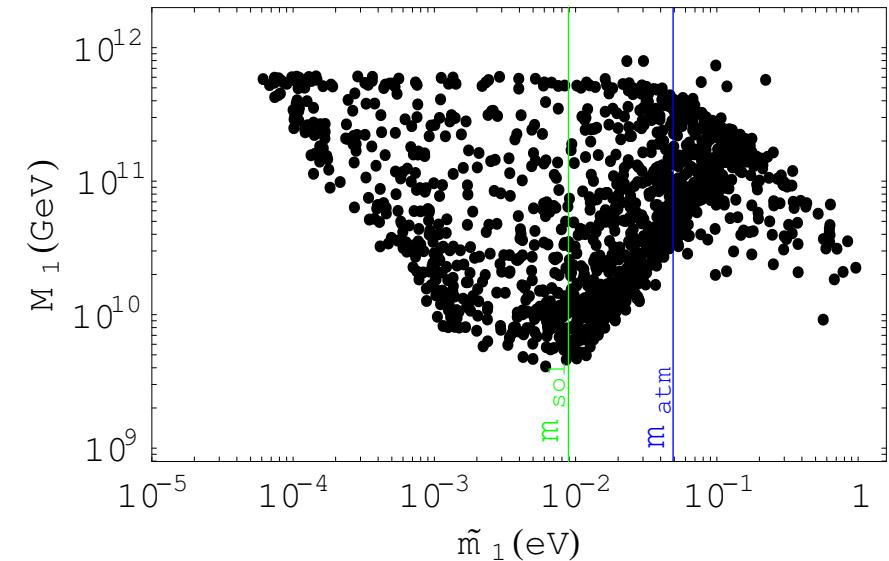
## \* Implication for $M_\nu$ : Lower bound on $N_1$ mass

- Allowed parameter space  $M_1 - \tilde{m} = \sum_\alpha \kappa_\alpha$   
compatible with  $5.2 \times 10^{-10} \lesssim 7.04 \times Y_B \lesssim 7.2 \times 10^{-10}$



One Flavour approx

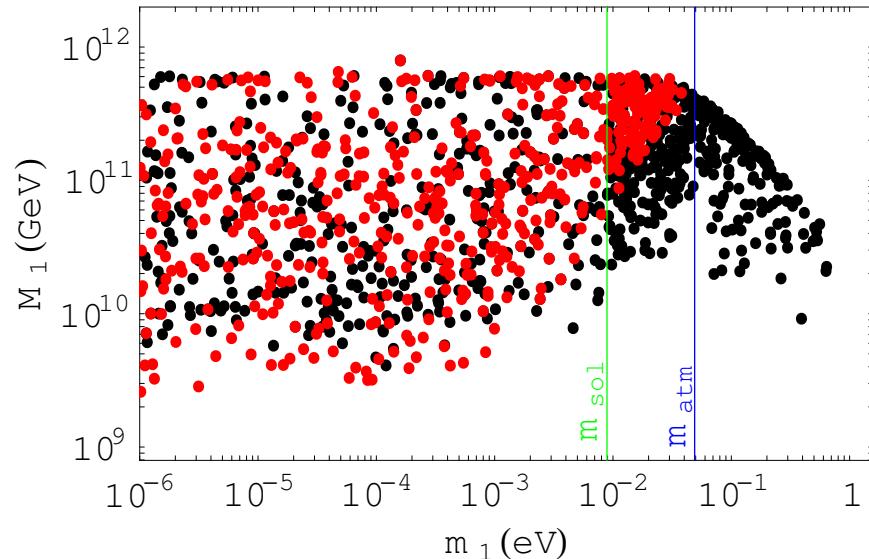
$$M_{N_1} \gtrsim \begin{cases} 6 \times 10^8 \text{ GeV} & \text{in the thermal case} \\ 4 \times 10^9 \text{ GeV} & \text{in the dynamical case} \end{cases}$$



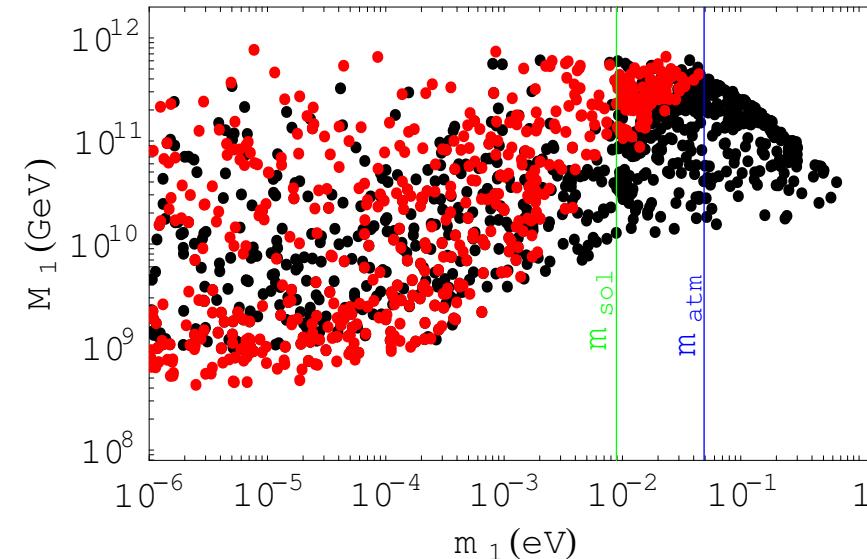
Full flavour treatment

Flavours open a wider range for  $\tilde{m}_1$  (extremely small values allowed)

\* Implication for lightest neutrino mass  $m_1$



Dynamical



Thermal

There is **NO BOUND** on absolute scale of light neutrinos from the requirement of successful leptogenesis!

## Role of $N_2$ and $N_3$

Relaxing the assumption of  $M_3 \gg M_2 \gg M_1$

### ► Impact of $N_2$

$M_{N_{2,3}} < T_{\text{RH}}$ ,  $L$ -asymmetry from  $N_{2,3}$  survives  $N_1$  leptogenesis

- **$N_1$ - decoupling** ( $Y_{N_2} \gg Y_{N_1}$ ) [ Di Bari, Vives, Blanchet ... ]
- **strong  $N_1$ - coupling** ( $N_1$  decoherence effects project  $L$ -asymmetry from  $N_2$  onto direction protected against  $N_1$  washout) [ Barbieri et al, Nardi, Strumia, ... ]

### ► Lowering the seesaw scale: resonant leptogenesis

**Resonant** enhancement of  $CP$  asymmetry if  $M_{N_2} - M_{N_1} = \Gamma_2/2$ :  $\varepsilon_{\text{CP}} \sim \mathcal{O}(1)$

[Pilaftsis, Underwood, Branco, Buras, Sarkar, Paschos, Flanz, Broncano, Plumacher ...]

Viable BAU **independently** of  $M_{N_i} \rightarrow$  low scale leptogenesis

[Ma, Hambye, Frigerio, Sarkar, Albright, Riotto, Babu, Isidori ...]

# SUSY leptogenesis

- **SUSY-Seesaw:** small qualitative and quantitative differences

$$Y_{B-L}^{\text{MSSM}} \approx -\eta \left[ \frac{1}{2}(\varepsilon_1 + \tilde{\varepsilon}_1) Y_{N_1}^{\text{eq}} + \frac{1}{2}(\varepsilon_{\tilde{1}} + \tilde{\varepsilon}_{\tilde{1}}) Y_{\tilde{N}_1}^{\text{eq}} \right]$$

flavours: correct  $T$  by  $(1 + \tan^2 \beta)$  factor in general

but *gravitino problem*: constraints on  $T_{RH}$  → can compromise viable BAU

$T_{RH}$  bounded from above (*gravitino problem*) ⇒ **upper bound on  $m_{N_1}$**   
(due to dramatic **loss of efficiency** when  $m_{N_1} \gg T_{RH}$  [Giudice et al ('03)] )

- “**new mechanisms**”:  $\left\{ \begin{array}{l} \text{- Affleck-Dine (BAU from coherent oscillations of } \\ \text{ } B, L \text{ charged flat directions of } V_{\text{SUSY}}^{\text{scalar}} \text{)} \\ \text{- Soft leptogenesis (unavoidable but subdominant in} \\ \text{the SUSY-seesaw) } \not\perp \text{ and } C/P \text{ from } \mathcal{L}_{\tilde{N}}^{\text{soft}} \\ \text{ } [\text{Grossman et al., D'Ambrosio et al, ...}] \end{array} \right.$

## Leptogenesis beyond type I see-saw

### ► Type II ( $SU(2)_L$ triplet + singlet)

- New sources of lepton asymmetry (and neutrino masses)
- The  $CP$  asymmetry  $\varepsilon_T$  increases with  $m_\nu$
- Less severe lower bound on  $M_T$  and  $M_R$

[Frère, Hambye, Raidal, Strumia, Chun, Akhmedov, Lavignac, Frigerio, Antusch, ...]

### ► Type III (fermionic triplet)

- Decay rate  $\Gamma_{FT} > \Gamma_N$
- The asymmetry  $\varepsilon_{FT} < \varepsilon_N$  (cancellations in the loops)
- Ignoring flavour effects, stricter bound on  $M_{FT}$  ( $\sim 10^{10}$  GeV) [Hambye, Senjanovic, ...]

## Other leptogenesis mechanisms

### ► Dirac leptogenesis [Akhmedov, Rubakov, Smirnov, Dick, Leidner, Ratz, Wright...]

- Dirac neutrinos with  $Y_\nu \sim 10^{-11}$  (from soft symmetry breaking)
- No  $L$ -violation
- Lepton asymmetry: decays of new heavy particles producing  $L$  for  $f_L$  and  $-L$  for  $f_R$
- Charged fermions quickly equilibrate  
 $\nu_L$  and  $\nu_R$  in equilibrium for  $T < T_{EW}$  ( $Y_\nu \sim 10^{-11}$ )
- Sphalerons: part of  $L$  stored in  $\nu_L$  converted into  $B$

**Phenomenological consequences: no  $(0\nu\beta\beta)$  decay!**

## Open Questions

- 👉 Is there a way to disprove leptogenesis?
- 👉 If seesaw leptogenesis, how to disentangle between seesaws? Collider Physics?
- 👉 Is it realistic to envisage a bridge between high and low energy  $CP$  phases?
- 👉 How natural is a real  $R$  matrix?  $CP$  exact symm in the RH sector
- 👉 Do LFV processes provide any additional constraints?