Herwig++ News

Keith Hamilton

Università degli Studi di Milano-Bicocca

Herwig++

Bähr Gieseke Gigg Grellscheid Latunde-Dada Plätzer Richardson Seymour Tully Webber

Main News: NLO POWHEG simulations

A method for including NLO corrections in a parton shower simulation

[Nason - 2004, Frixione, Nason, Oleari - 2007].

A POWHEG code just outputs weight +1 events, according to NLO+[N]LL resummed distributions, as Les Houches event files.



POWHEG

Hardest emission:

Always generated according to the NLO real ME

[in A.O. shower ≠ 1st radiation hardest]

Shower independent:

Easy to make [we did it!] Easier for attempting complex NLO processes Insensitive to shower problems / development / fixes [kinematics reconstruction of HW++ has changed 3 times] Opens the door for the NLO community

Only weight +1 events:

Conventional statistical analysis Easier to attack complex analysis [NN's, DA's, ...]

> QCD coherence can be implemented, approximately, in shower MC's by ordering emissions by their angle.

For an A.O. shower the first emission isn't a priori the hardest:



QCD coherence can be implemented, approximately, in shower MC's by ordering emissions by their angle.

For an A.O. shower the first emission isn't a priori the hardest:



- When POWHEG gives a real emission according to the NLO calculation it's supposed to be the hardest.
- So in general you just veto emissions from the shower with $p_T > p_{T,POWHEG}$



[Nason 2004, Implementation KH, Richardson, Tully 2008]



So in general you just veto emissions from the shower with p_T > p_{T,POWHEG}



[Nason 2004, Implementation KH, Richardson, Tully 2008]

- When POWHEG gives a real emission according to the NLO calculation it's supposed to be the hardest.
- So in general you just veto emissions from the shower with p_T > p_{T,POWHEG}



But if the shower is A.O. then the shower should also try to include 'earlier' soft wide angle emissions

[Nason 2004, Implementation KH, Richardson, Tully 2008]

- When POWHEG gives a real emission according to the NLO calculation it's supposed to be the hardest.
- So in general you just veto emissions from the shower with p_T > p_{T,POWHEG}



But if the shower is A.O. then the shower should also try to include 'earlier' soft wide angle emissions

[Nason 2004, Implementation KH, Richardson, Tully 2008]

Shipping with the current release:

- ▷ hh → γ / Z / W / H / ZH / WH [KH, Richardson, Tully]
- Spin correlations in decays [also for real emissions]
 Truncated showers

Also available in Contrib:

POWHEG top and Higgs decays [Latunde-Dada]

Herwig++ in-house POWHEGs

Shipping in a few months:

VBF [Under validation - D'Errico, Richardson]

 $hh \rightarrow WZ / ZZ / WW [Validated - KH]$

Longer term:

Enhance compatibility with POWHEG-BOX [see Re's talk]

Colour coherent CKKW [JHEP 0911:038,2009]

Herwig++ POWHEG validation

1. MCFM:

- Total cross sections agree at O(0.1%)
- O(20) IR safe distributions checked; 'Born variables' are scrutinized and show excellent agreement.
- 2. MC@NLO + 3. Matrix Element Corrections:
- Shift focus to resummation & radiation generation i.e. here we wanted to check the shapes is normalized to 1.
- Good agreement in all 3 approaches. POWHEG & ME Corr tend to have harder spectrum [and we understand why].

4. Data:

▶ For e+e- and DY we compared to the data.

Results

Drell-Yan vector boson production



[KH, Richardson, Tully]

Results: Drell-Yan vector boson production

W boson p_T spectrum compared to DO run I data



Solid line: NLO Herwig++ POWHEG Blue dashes: MC@NLO Red dashes: Herwig++ with ME corrections

Results: Drell-Yan vector boson production

Z boson p_T spectrum compared to DO run II data



Solid line: NLO Herwig++ POWHEG Blue dashes: MC@NLO Red dashes: Herwig++ with ME corrections



Higgs production via gluon fusion



[KH, Richardson, Tully]

Higgs boson rapidities compared to fixed order NLO calculations



Red: NLO Herwig++ POWHEG Magenta: LO Herwig++ POWHEG Black: NLO MCFM fixed order Dashes: LO MCFM fixed order





Higgs-strahlung



[KH, Richardson, Tully]

Results: Higgs-strahlung $(q\overline{q}\rightarrow HZ/W)$

Polar angle of electron vs fixed order NLO calculations: Tevatron



Red: NLO Herwig++ POWHEG Magenta: LO Herwig++ POWHEG Black: NLO MCFM fixed order Dashes: LO MCFM fixed order

Results: Higgs-strahlung $(q\bar{q}\rightarrow HZ)$



Results: Higgs-strahlung $(q\bar{q}\rightarrow HW)$



Results: Higgs-strahlung $(q\bar{q}\rightarrow HW)$



Results: Diboson production





Results: Diboson production at LHC



LO total inclusive cross section

PS has half as many 80 GeV jet events [LL approx]

Direction of hard jets not like the NLO real correction

Results: Diboson production at LHC



 p_{T,μ^+} [GeV]

Monday, 6 September 2010

 $p_{T,\mu^{-}}$ [GeV]

Results: Diboson production at LHC



Really "soft" W emission from hard quark $qg \rightarrow qZ \rightarrow qWZ$

Also big enhancement due to incident gluon flux [Nason et al]

Results



[D'Errico, Richardson]

Results: VBF



[D'Errico, Richardson]

Results: VBF



[D'Errico, Richardson]



Multiple Interactions

Colour Coherent Formulation of CKKW

Multiple Parton Interations

- Original Herwig: UA1 model of the soft UE but option to link to JIMMY model of MPIs
- Herwig++:built in MPI model building on JIMMY technology.



[Bahr, Butterworth, Gieseke, Seymour]

Colour Coherent CKKW: LEP jet rates



[KH, Richardson, Tully]

Colour Coherent CKKW: fit to LEP data



CKKW tune at $y_{merge} = 10^{-2}$ gives much better fit to LEP data than standard tune [HW++ with ME corrections]:

Observable	Hw+ME χ^2 /d.o.f	CKKW χ^2 /d.o.f
Thrust	23.48	10.62
Sphericity	5.638	0.580
Oblateness	2.450	0.339
Planarity	1.249	1.211
y_{23}	2.400	0.867
y_{34}	1.887	1.026
y_{45}	4.571	2.018
$\cos lpha_{34}$	0.569	3.301
$\cos\chi_{BZ}$	1.002	0.775
$\cos \Phi_{KSW}$	1.469	1.337
$\cos heta_{NR}$	4.509	0.702

[KH, Richardson, Tully]

Colour Coherent CKKW: CDF jet E_T spectra



[Richardson, Tully]

Main news summary

Herwig++ ships with the following in-house POWHEG codes:

hh $\rightarrow \gamma / Z / W / H / ZH / WH$

- All processes underwent lengthy validations via comparison with MCFM, MC@NLO, ME Corrections and real data.
- The only sizeable discrepancies found were in p_T spectra and rapidity correlations w.r.t. MC@NLO for gg \rightarrow H.
- Our studies conclude that these discrepancies stem from the dependency of MC@NLO on the underlying shower MC, in particular its phase space partitioning [should not occur in MC@NLO+PYTHIA].
- > VBF is being validated and WW/ZZ/WZ production is validated: HW++ will ship with these in the coming months.

Bonus material

Results: Drell-Yan vector boson production

Z boson rapidity compared to DO run II data



Solid line: NLO Herwig++ POWHEG Red dashes: MC@NLO Blue dashes: Herwig++

Radiative phase space variables: x & y

NLO real emission correction: $g+g \rightarrow H+g$



The Dead Zone

In angular ordered parton showers like Herwig there is a region of phase space that the shower can't emit into: the dead zone.



The dead zone is in the region of phase space corresponding to wide angle, high p_T emission of the first radiated parton.





$$x = 1 - E_{\text{emitted}} / E_{\text{incoming}}$$

 $y = \cos \theta$





$$x = 1 - E_{emitted}/E_{incoming}$$

 $y = \cos \theta$





$$x = 1 - E_{emitted}/E_{incoming}$$

 $y = \cos \theta$



The Dead Zone

• Now we superimpose contours for $p_T = m_H \text{ GeV}, 80 \text{ GeV}, 40 \text{ GeV}, 10 \text{ GeV}$ for a 160 GeV Higgs at the Tevatron and a 115 GeV Higgs at the LHC.



The dead zone is in the region of phase space corresponding to wide angle, high p_T emission of the first radiated parton.

Matrix Element Corrections

- Before the days of MC@NLO and POWHEG we used Matrix Element Corrections (MECs) to fill the dead zone and correct the shower.
- > With MECs you get an emission in the dead zone with probability:

$$\mathcal{P}_{\text{dead}}^{\text{HW}}\left(\Phi_{B}\right) = \int_{\text{dead}} \mathrm{d}\Phi_{R} \, \frac{\widehat{R}\left(\Phi_{B}, \Phi_{R}\right)}{B\left(\Phi_{B}\right)}$$

- If an emission occurs in the dead zone it's distributed according to the real emission matrix element with PDFs etc: $\hat{R}(\Phi_B, \Phi_R)$
- All this does is correct the shape of distributions sensitive to extra emissions:
 - > all normalizations and scale dependencies are as at LO
 - no virtual effects are included at all
- N.B. integrating $\mathcal{P}_{dead}^{HW}(\Phi_B)$ over the Born variables gives the fraction that the dead zone contributes to the NLO x-section, up to terms of order α_8^2 .





Hardest jet rapidity – Higgs rapidity ($p_T > 10 \text{ GeV}$) Hardest jet rapidity – Higgs rapidity ($p_T > 40 \text{ GeV}$)

0.3 0.3 ∕ơdơ∕d(y_{jet}−y_H) 0.2 0.2 0.1 0.1 0.0 -10 0.0 -10 -5 0 5 10 $^{-5}$ 0 5 10 $y_{jet} - y_H$ $y_{jet} - y_H$ Hardest jet rapidity - Higgs rapidity ($p_T > 80$ GeV) $1/\bar{s}_{max}$ Х shower a [TVT m_H=160 GeV] 0.3 $\sigma d\sigma/d(y_{jet}-y_{H})$ 0.2 dead zone У 0.1 shower b 0.0 -10 0 5 10 -5 $y_{jet} - y_H$

 $1/\sigma d\sigma/d(y_{jet}-y_H)$

Matrix Element Corrections

- In comparing our POWHEG with MC@NLO and the MEC method for this process we also used a modified version of the MEC.
- Only the emission rate (not the shape) in the dead zone was altered:

$$\mathcal{P}_{\text{dead}}^{\text{HW}}\left(\Phi_{B}\right) \to \mathcal{P}_{\text{dead}}^{\text{NLO}}\left(\Phi_{B}\right) = \int_{\text{dead}} \mathrm{d}\Phi_{R} \; \frac{\widehat{R}\left(\Phi_{B}, \Phi_{R}\right)}{\overline{B}\left(\Phi_{B}\right)}$$

- Reminder: $\overline{B}(\Phi_B)$ is the NLO differential x-sec integrated over the radiative phase space (x and y).
- Integrating $\mathcal{P}_{dead}^{NLO}(\Phi_B)$ over the remaining Born variables gives the fraction that the dead zone contributes to the NLO x-section exactly.
- This means it should put the same fraction of events in the dead zone as an exact NLO calculation would i.e. the same fraction as MC@NLO.

The phase space maps showed that for $p_T > m_n$ all emissions are in the dead zone. So from the MEC they all occur with probability:

$$\mathcal{P}_{m_n}^{\mathrm{HW}}\left(\Phi_B\right) = \int_{m_n} \mathrm{d}\Phi_{R_1} \, \frac{\widehat{R}_1\left(\Phi_B, \Phi_{R_1}\right)}{B\left(\Phi_B\right)}$$

• MC@NLO reproduces NLO results so the fraction of emissions it has with $p_T > m_n$ will be the corresponding fraction of the NLO x-sec:

$$\mathcal{P}_{m_n}^{\text{NLO}}\left(\Phi_B\right) = \int_{m_n} \mathrm{d}\Phi_{R_1} \ \frac{\widehat{R}_1\left(\Phi_B, \Phi_{R_1}\right)}{\overline{B}\left(\Phi_B\right)}$$

In POWHEG, the probability that an emission occurs with $p_T > m_n$ is one minus the probability that no emission occurs with $p_T > m_n$:

$$\mathcal{P}_{m_n}^{\text{PH}} = 1 - \Delta_{\hat{R}} \left(m_n \right) = \int_{m_n} \mathrm{d}\Phi_{R_1} \, \frac{\widehat{R}_1 \left(\Phi_B, \Phi_{R_1} \right)}{B \left(\Phi_B \right)}$$

Hence: $\mathcal{P}_{m_n}^{\mathrm{HW}} \approx \mathcal{P}_{m_n}^{\mathrm{PH}} \approx \mathcal{K} \mathcal{P}_{m_n}^{\mathrm{NLO}}$

Results: a word about p_T

What would this rate (fraction) be if we had an NNLO calculation?

$$\mathcal{P}_{m_n}^{\text{NNLO}}\left(\Phi_B\right) = \int_{m_n} \mathrm{d}\Phi_{R_1} \left[\widehat{R}_1\left(\Phi_B, \Phi_{R_1}\right) + R_{1+1}\left(\Phi_B, \Phi_{R_1}\right) + \int \mathrm{d}\Phi_{R_2} R_2\left(\Phi_B, \Phi_{R_1}, \Phi_{R_2}\right) \right]$$

$$\div \mathrm{d}\sigma_{\text{NNLO}}\left(\Phi_B\right)$$

$$= \int_{m_n} \mathrm{d}\Phi_{R_1} \, \frac{\widehat{R}_1\left(\Phi_B, \Phi_{R_1}\right)}{B\left(\Phi_B\right)} \left[1 - \frac{\overline{B}\left(\Phi_B\right)}{B\left(\Phi_B\right)} + \frac{\overline{R}_1\left(\Phi_B, \Phi_{R_1}\right)}{\widehat{R}_1\left(\Phi_B, \Phi_{R_1}\right)}\right]$$

where

$$\overline{R}_{1}(\Phi_{B},\Phi_{R_{1}}) = \widehat{R}_{1}(\Phi_{B},\Phi_{R_{1}}) + R_{1+1}(\Phi_{B},\Phi_{R_{1}}) + \int d\Phi_{R_{2}}R_{2}(\Phi_{B},\Phi_{R_{1}},\Phi_{R_{2}})$$

Now $\frac{\overline{B}(\Phi_B)}{B(\Phi_B)}$ and $\frac{\overline{R}_1(\Phi_B, \Phi_{R_1})}{\widehat{R}_1(\Phi_B, \Phi_{R_1})}$ are basically differential K-factors for e.g. in the case of gluon fusion, gg \rightarrow H and gg \rightarrow H+jet respectively

Results: a word about p_T

Now it turns out that the NLO K-factors for gg→H and gg→H+jet are very similar ≈ 1.6/1.7:

Grazzini, Kunszt, De Florian PRL 82 [1999]



And they are fairly insensitive to the kinematics too! This is because the large contribution to the cross section are due to soft emissions and these don't alter the LO / NLO kinematics too much.

Results: a word about p_T

So one should really expect the final two terms in the NNLO emission rate to basically cancel each other out!

$$\mathcal{P}_{m_{n}}^{\text{NNLO}}(\Phi_{B}) = \int_{m_{n}} d\Phi_{R_{1}} \left[\hat{R}_{1}(\Phi_{B}, \Phi_{R_{1}}) + R_{1+1}(\Phi_{B}, \Phi_{R_{1}}) + \int d\Phi_{R_{2}} R_{2}(\Phi_{B}, \Phi_{R_{1}}, \Phi_{R_{2}}) \right]$$

$$\div d\sigma_{\text{NNLO}}(\Phi_{B})$$

$$= \int_{m_{n}} d\Phi_{R_{1}} \frac{\hat{R}_{1}(\Phi_{B}, \Phi_{R_{1}})}{B(\Phi_{B})} \left[1 - \frac{\overline{B}(\Phi_{B})}{B(\Phi_{B})} + \frac{\overline{R}_{1}(\Phi_{B}, \Phi_{R_{1}})}{\widehat{R}_{1}(\Phi_{B}, \Phi_{R_{1}})} \right]$$

where

$$\overline{R}_{1}(\Phi_{B},\Phi_{R_{1}}) = \widehat{R}_{1}(\Phi_{B},\Phi_{R_{1}}) + R_{1+1}(\Phi_{B},\Phi_{R_{1}}) + \int d\Phi_{R_{2}}R_{2}(\Phi_{B},\Phi_{R_{1}},\Phi_{R_{2}})$$

And that then makes:

$$\mathcal{P}_{m_n}^{\text{NNLO}}\left(\Phi_B\right) = \int_{m_n} \mathrm{d}\Phi_{R_1} \, \frac{\widehat{R}_1\left(\Phi_B, \Phi_{R_1}\right)}{B\left(\Phi_B\right)}$$
$$\approx \mathcal{P}_{m_n}^{\text{HW}} \approx \mathcal{P}_{m_n}^{\text{PH}} \approx \mathcal{K} \, \mathcal{P}_{m_n}^{\text{NLO}}$$

Results: QED radiation in the decay [Sophty]

Total photon energy radiated in Z decays



Each line corresponds to a different pair of charged leptons. The lowest / innermost lines are for some fictional 'heavy leptons'.