

Kaon Phenomenology

Lattice Meets Phenomenology
IPPP, University of Durham
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Kaons and Fundamental Physics

- ➔ Determination of fundamental parameters
 - lepton universality
 - CKM unitarity
 - mass determination
- ➔ Test suppression of top-dominated FCNCs
 - rare decays
 - CP violation
- ➔ Interplay between SD & LD is essential

Leptonic and Semileptonic

Observables: $K(\pi) \rightarrow l \bar{\nu}_l$ & $K \rightarrow \pi l \bar{\nu}_l$

$$\frac{\Gamma(K_{\ell 2(\gamma)}^{\pm})}{\Gamma(\pi_{\ell 2(\gamma)}^{\pm})} = \left| \frac{V_{us}}{V_{ud}} \right|^2 \frac{f_K^2 m_K}{f_{\pi}^2 m_{\pi}} \left(\frac{1 - m_{\ell}^2/m_K^2}{1 - m_{\ell}^2/m_{\pi}^2} \right)^2 \times (1 + \delta_{\text{em}})$$

Lattice

$$\Gamma(K_{\ell 3(\gamma)}) = \frac{G_F^2 m_K^5}{192 \pi^3} C_K S_{\text{ew}} |V_{us}|^2 f_+(0)^2 I_K^{\ell}(\lambda_{+,0}) \left(1 + \delta_{SU(2)}^K + \delta_{\text{em}}^{K\ell} \right)^2$$

CKM Unitarity

$$\Gamma(K_{l3}) \quad |V_{us}|f_+(0) = 0.21661(47)$$

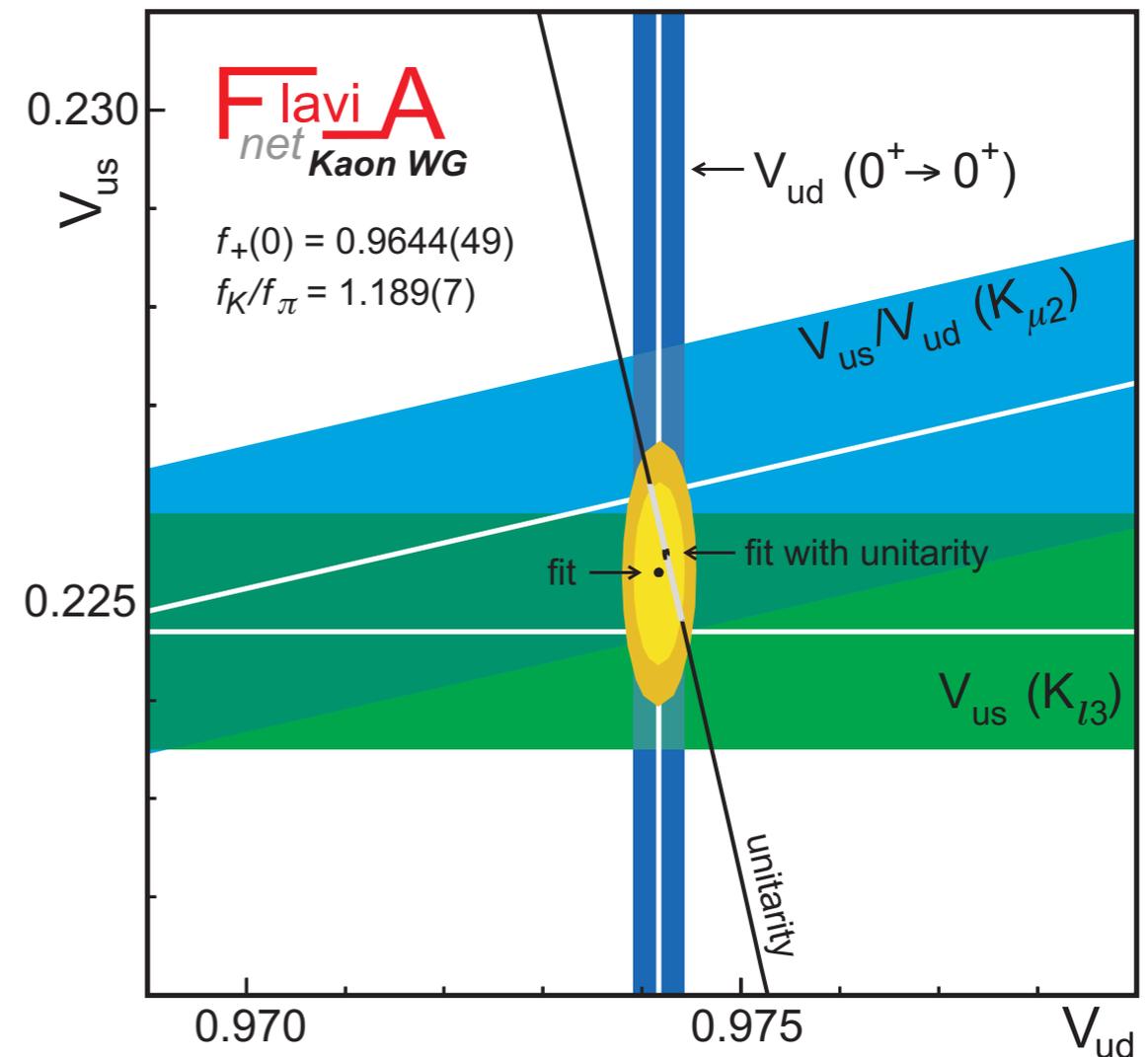
$$\frac{\Gamma(K_{l2})}{\Gamma(\pi_{l2})} \quad \frac{|V_{us}|f_K}{|V_{ud}|f_\pi} = 0.27599(59)$$

[Flavianet]

and nuclear β decay

$$V_{ud} = 0.97425(22)$$

[Hardy, Towner '08]



$$\Delta_{CKM} = |V_{ud}^2| + |V_{us}^2| + |V_{ub}^2| - 1$$

$$= (0.1 \pm 0.6) \times 10^{-3}$$

CKM Unitarity (Model Independent)

[Cirigliano et. al. '09]

$\Lambda_{NP} \gg M_W$ Neglect $\mathcal{O}\left(\frac{M_W}{\Lambda_{NP}}\right)$ corrections

Use $SU(2) \otimes U(1)$ invariant operators [Buchmüller-Wyler '06]
(plus $U(3)^5$ flavour symmetry)

$$O_{lq}^{(3)} = (\bar{l}\gamma^\mu \sigma^a l)(\bar{q}\gamma_\mu \sigma^a q) \quad O_{ll}^{(3)} = \frac{1}{2}(\bar{l}\gamma^\mu \sigma^a l)(\bar{l}\gamma_\mu \sigma^a l)$$

Constrained from EW precision data [Han, Skiba '05]

Redefine

$$G_F(\mu \rightarrow e \nu \bar{\nu}) \rightarrow G_F(1 - 2\bar{\alpha}_{ll}^{(3)}) \longrightarrow G_F^\mu$$
$$G_F(d \rightarrow u e \bar{\nu}) \rightarrow G_F(1 - 2\bar{\alpha}_{lq}^{(3)}) \longrightarrow G_F^{SL}$$

CKM Unitarity (Model Independent)

$$V_{udi}^{\text{PDG}} = \frac{G_F^{\text{SL}}}{G_F^\mu} V_{udi} \longrightarrow \Delta_{\text{CKM}} = 4 \left(\bar{\alpha}_{ll}^{(3)} - \bar{\alpha}_{lq}^{(3)} + \dots \right)$$

$$O_{lq}^{(3)} = (\bar{l}\gamma^\mu \sigma^a l)(\bar{q}\gamma_\mu \sigma^a q)$$

from HEP

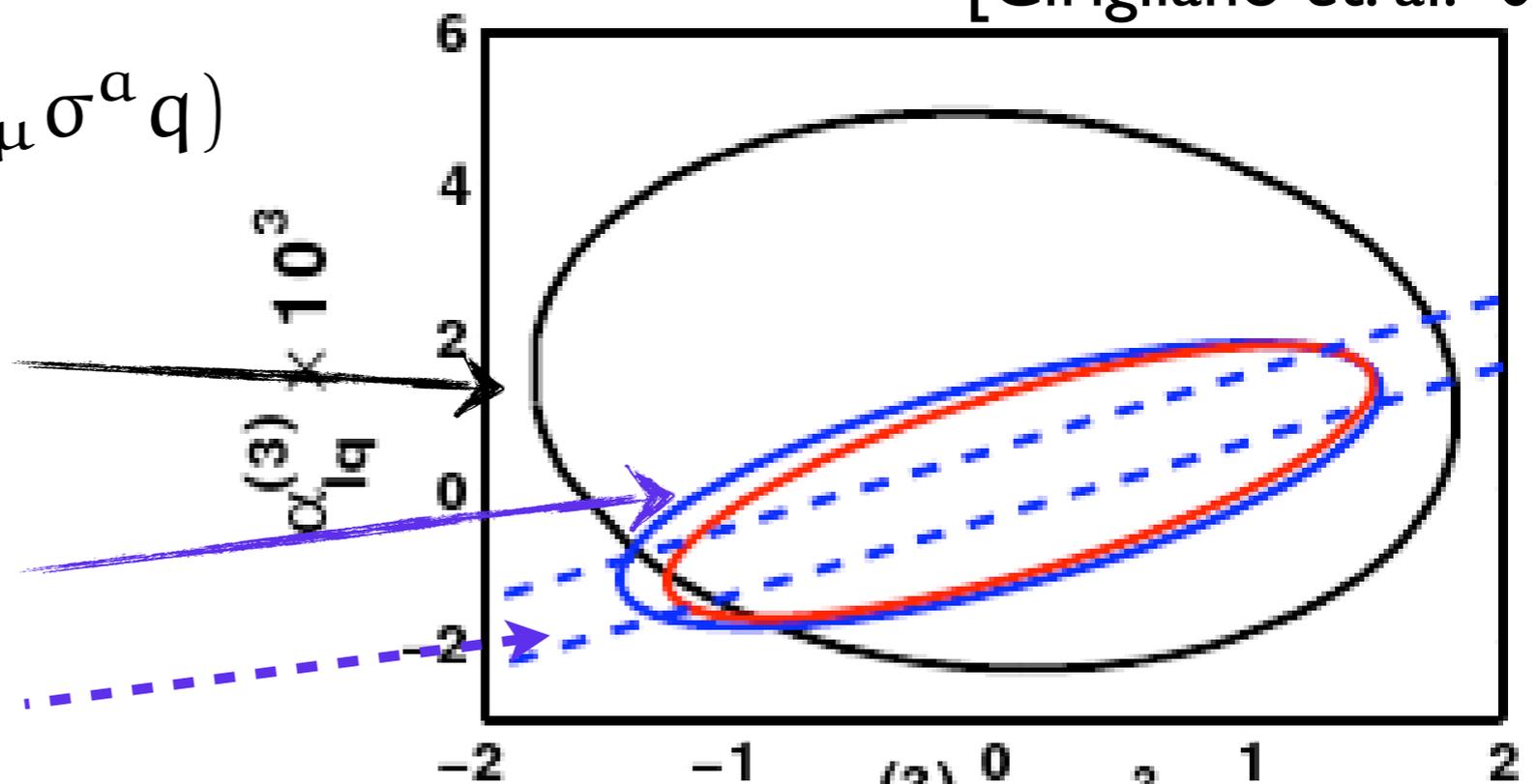
HEP + CKM

CKM

$$\Lambda_{\text{NP}} > 10\text{TeV}$$

$$O_{ll}^{(3)} = \frac{1}{2} (\bar{l}\gamma^\mu \sigma^a l)(\bar{l}\gamma_\mu \sigma^a l)$$

[Cirigliano et. al. '09]



Leptonic and Semileptonic

$$K(\pi) \rightarrow l \bar{\nu}_l \quad \& \quad K \rightarrow \pi l \bar{\nu}_l$$

Observables

$$R_K = \frac{\Gamma(K \rightarrow e \bar{\nu})}{\Gamma(K \rightarrow \mu \bar{\nu})}$$

$$R_K^{SM} = 2.477(1) \times 10^{-5}$$

[Cirigliano, Rosell '07]

See also [Marciano, Sirlin '93]

$$R_K^{NA62} = 2.500(16) \times 10^{-5}$$

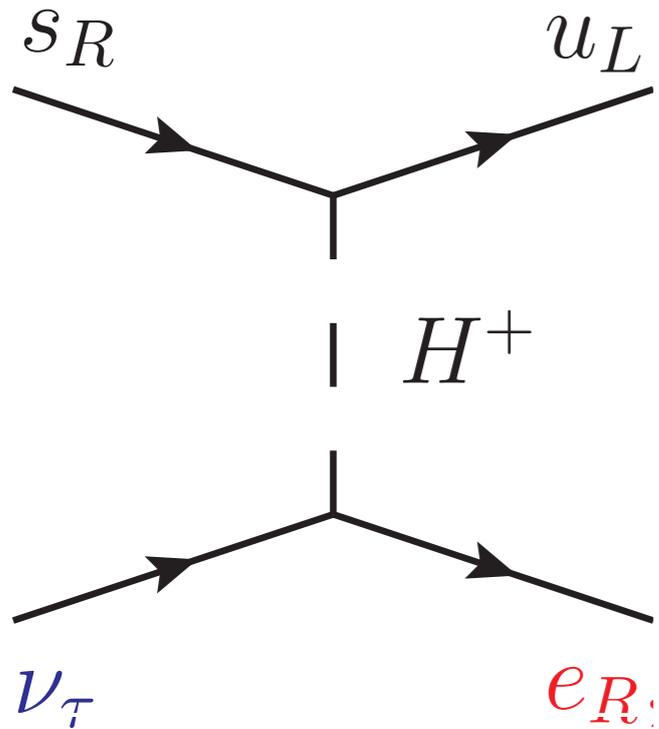
[numbers from KAON09]

$$R_K^{KLOE} = 2.493(25)(19) \times 10^{-5}$$

**New numbers
from NA62 in
agreement with SM**

**Test of lepton universality violation
driven by experimental precision**

Lepton Universality in the MSSM



LF Conserving: \sim lepton mass

Lepton Flavour Violation: $\Delta_R^{31} \sim \frac{g_2^2}{16\pi^2} \delta_{RR}^{31}$
 [Masiero, Paradisi, Petronzio '08]

$$R_K^{\text{LFV}} = \frac{\Gamma_{SM}(K \rightarrow e \nu_e) + \Gamma_{SM}(K \rightarrow e \nu_\tau)}{\Gamma_{SM}(K \rightarrow \mu \nu_\mu)}$$

$$\Delta r_K \sim \frac{m_K^4}{m_{H^+}^4} \frac{m_\tau}{m_e} |\Delta_R^{31}|^2 \tan^6 \beta \longrightarrow \text{can reach } 10^{-2}$$

But: finetuning of m_e necessary [Girrbach et. al. '09]

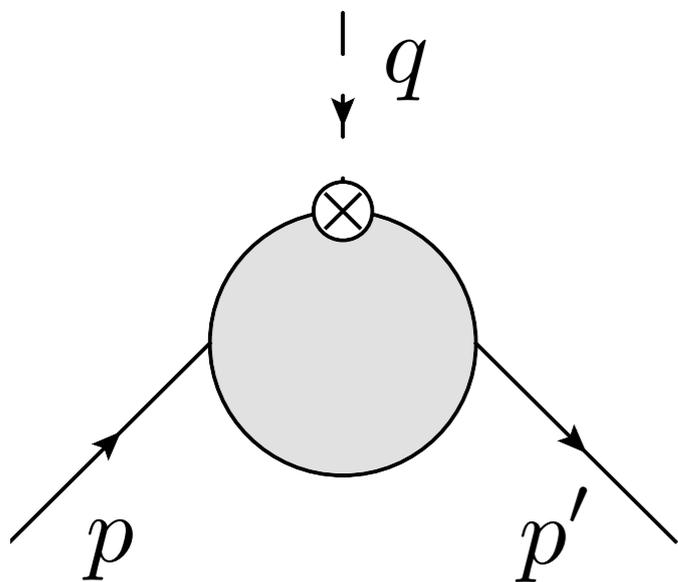
Model independent MLFV and GUT analysis [Isidori et. al. '09]

Lattice & Continuum: Masses

How do lattice results correspond to the continuum?

- Lattice perturbation theory
- Non-perturbative renormalisation

Transform to $\overline{\text{MS}}$ scheme using perturbation theory



$$S_B(p') \Lambda_{P,B}(p, p') S_B(p) =$$

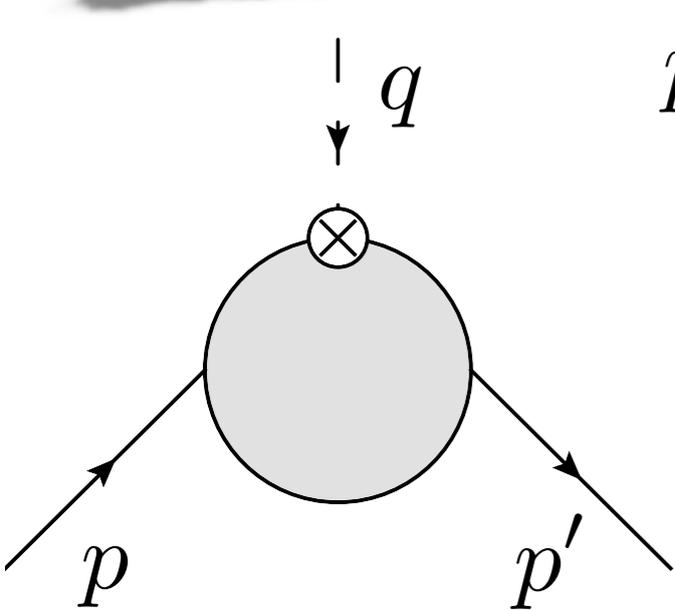
$$\int d^4x d^4y e^{ip'x} e^{-ipy} \langle T([i\bar{u}_B \gamma_5 s_B](0) u_B(x) \bar{s}_B(y)) \rangle$$

condition:

$$Z_q^{-1} Z_m^{-1} \text{tr}[\Lambda_{P,B}(p, p') \gamma_5] = 12$$

momentum configuration plus 2 point function:
fix mass and field renormalisation

Conversion RI/SMOM \rightarrow \overline{MS}



$$p^2 = p'^2 = -\mu^2$$

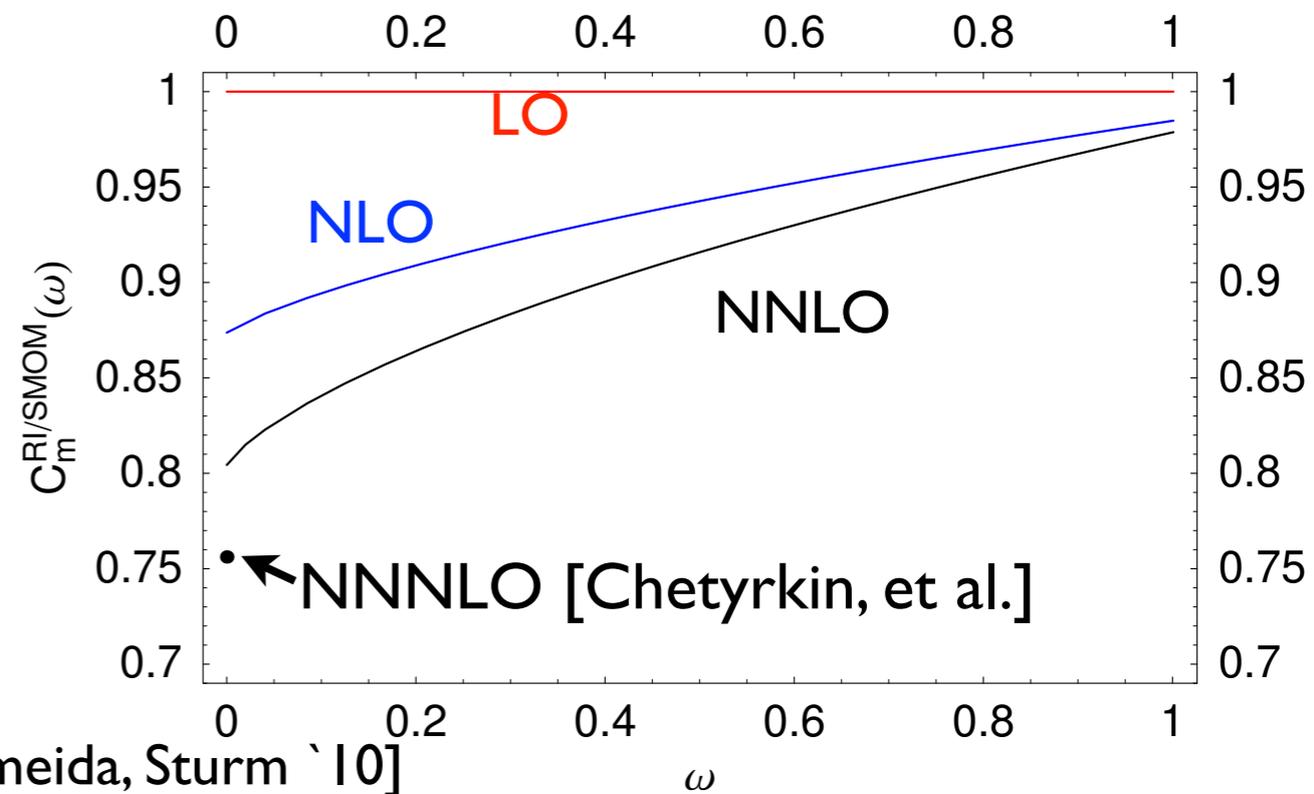
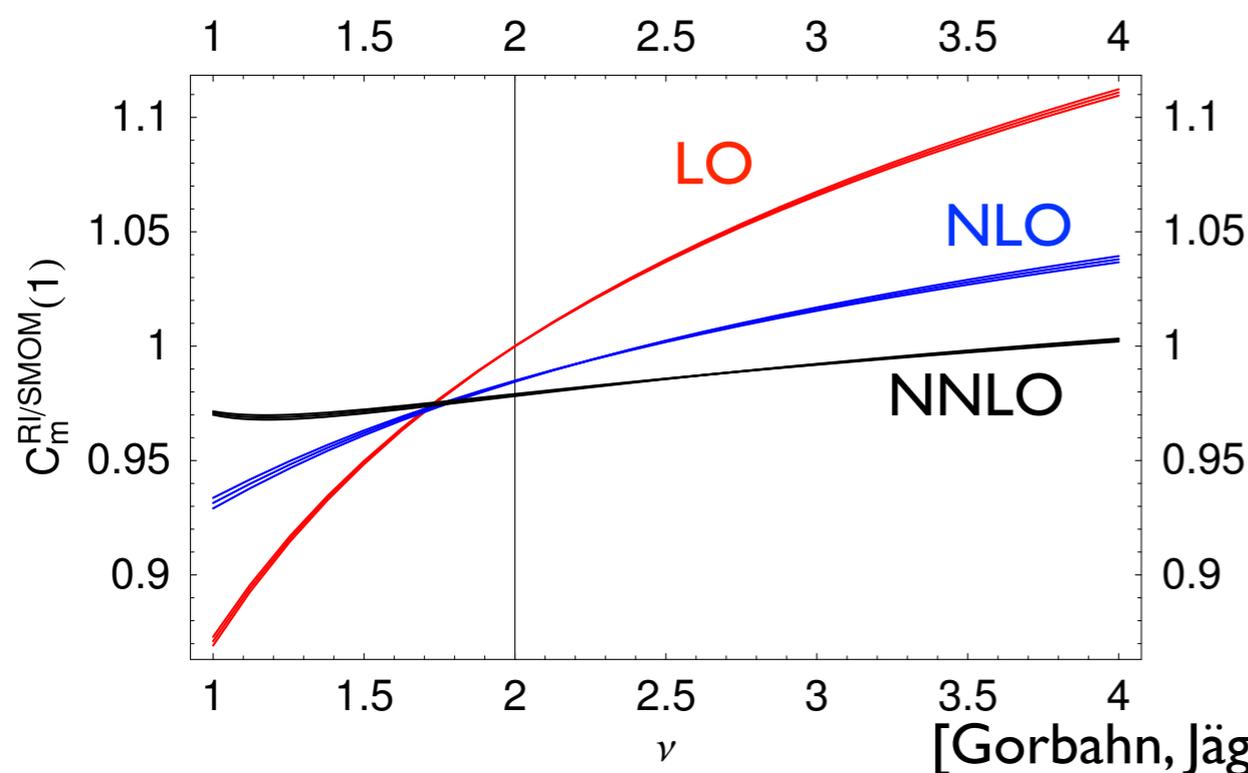
$$q^2 = 0$$

[Martinelli et. al. '93-'95]

corresponds to: RI/MOM scheme
 matching@NNLO: poor convergence

find a scheme good for lattice & loops
 $p^2 = p'^2 = q^2 = -\mu^2$ RI/SMOM [Sturm et. al. '09]

off shell: Λ^2/q^2 suppression and good convergence



Rare Kaon Decays

FCNCs which are dominated by top-quark loops:

$$\begin{array}{lll} b \rightarrow s : & b \rightarrow d : & s \rightarrow d : \\ |V_{tb}^* V_{ts}| \propto \lambda^2 & |V_{tb}^* V_{td}| \propto \lambda^3 & |V_{ts}^* V_{td}| \propto \lambda^5 \end{array}$$

CKM suppression: enhanced sensitivity to NP

$$V_{ts}^* V_{td} + V_{cs}^* V_{cd} = -V_{us}^* V_{ud}$$

λ^5 λ λ

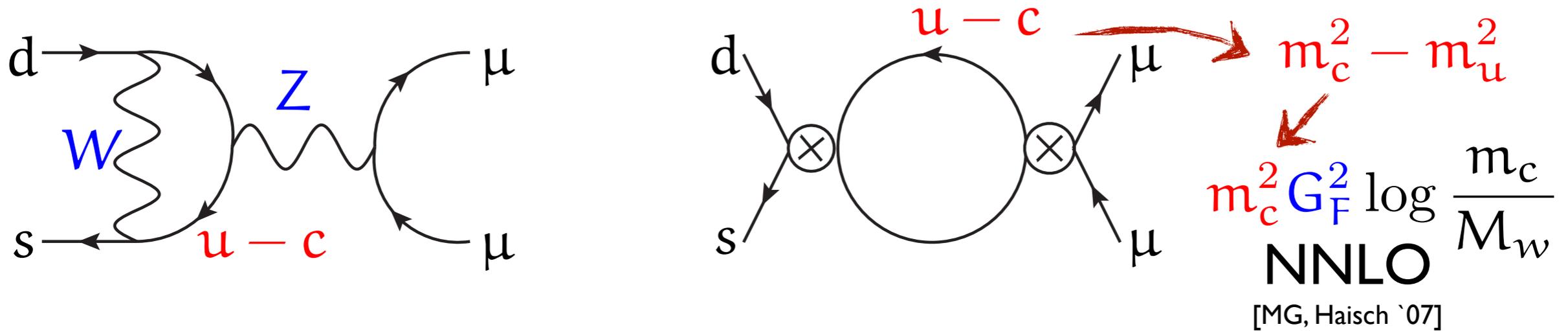
how can we suppress the light quark contribution?

Quadratic GIM: $\lambda \frac{m_c^2}{M_W^2}$

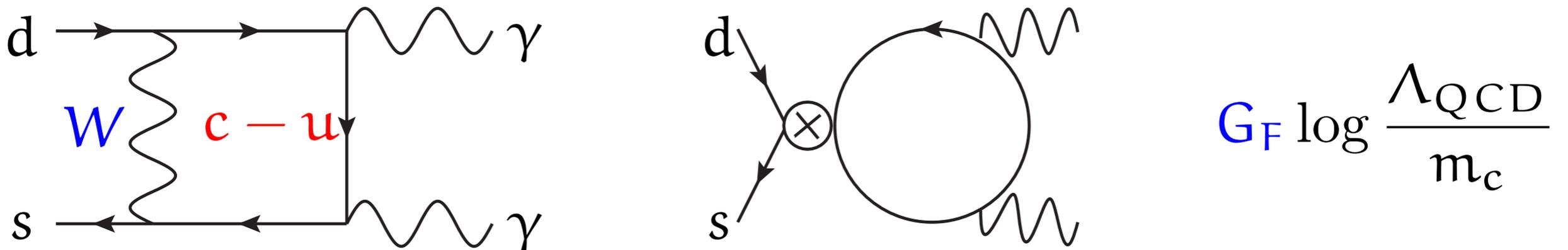
CP violation: $\text{Im}(V_{cs}^* V_{cd})$

GIMnastics

Quadratic GIM suppresses light quark contribution

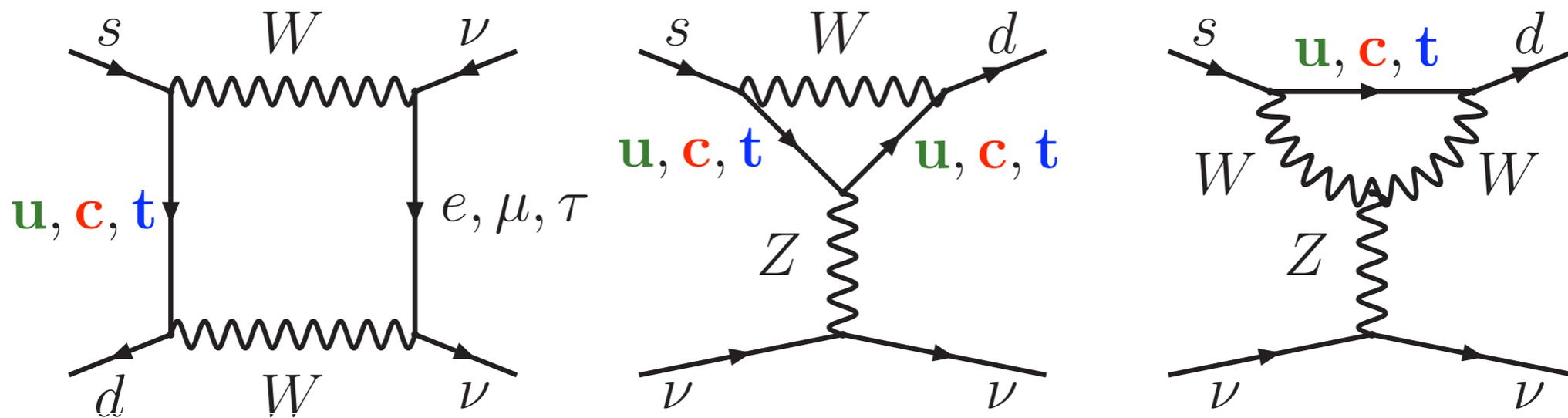


No quadratic suppression for $K_L \rightarrow \gamma\gamma$



$\frac{\alpha}{4\pi} \times K_L \rightarrow \gamma\gamma$ also contributes to: $K_L \rightarrow \mu^+ \mu^-$

No couplings to γ s: $K \rightarrow \pi \nu \bar{\nu}$



- Dominant Operator: $Q_\nu = (\bar{s}_L \gamma_\mu d_L)(\bar{\nu}_L \gamma^\mu \nu_L)$

$$\sum_i V_{is}^* V_{id} F(x_i) = V_{ts}^* V_{td} (F(x_t) - F(x_u)) + V_{cs}^* V_{cd} (F(x_c) - F(x_u))$$

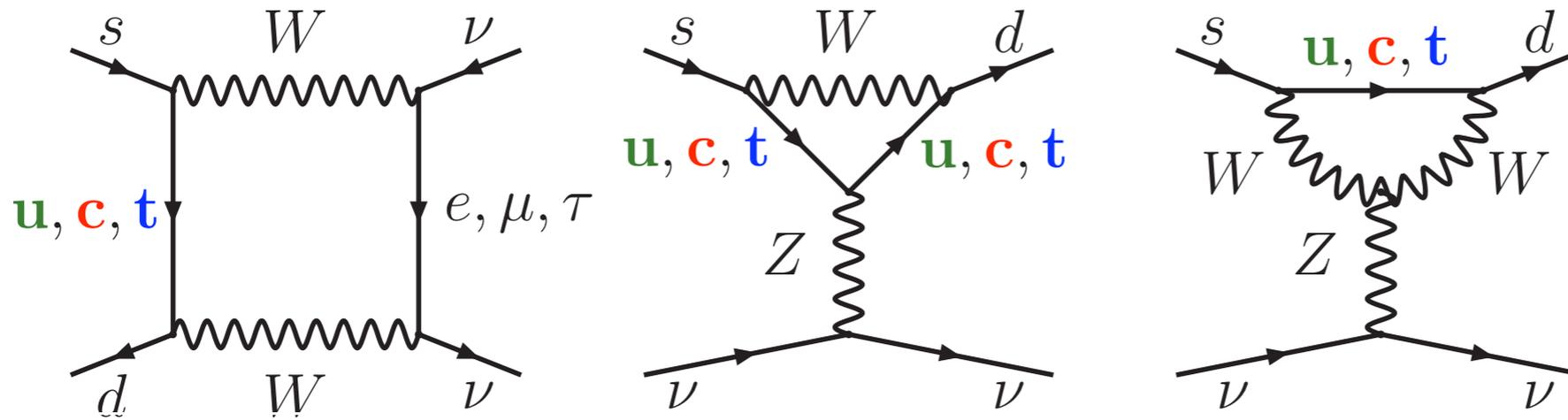
Quadratic GIM:

$$\lambda^5 \frac{m_t^2}{M_W^2}$$

$$\lambda \frac{m_c^2}{M_W^2} \ln \frac{M_W}{m_c}$$

$$\lambda \frac{\Lambda^2}{M_W^2}$$

$K_L \rightarrow \pi^0 \bar{\nu} \nu$: Effective Hamiltonian



CP violating

including
NLO EW

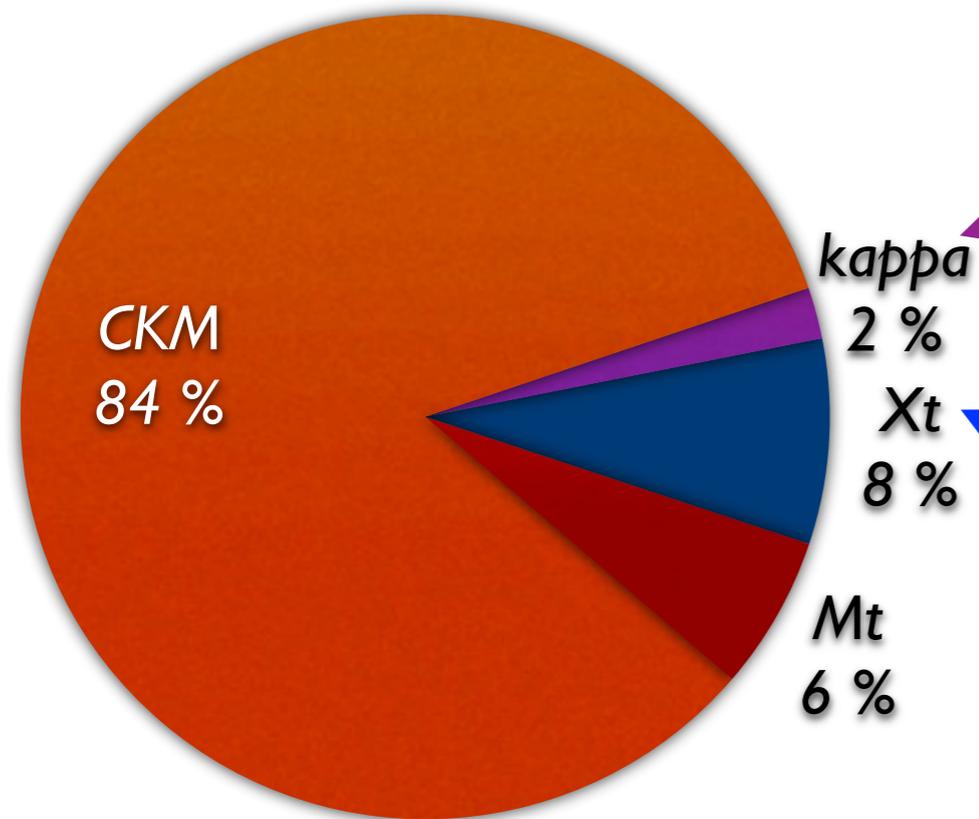
[Bord, Gorbahn, Stamou '10]

Only top quark contributes: $H_{\text{eff}} = \frac{4G_F}{\sqrt{2}} \frac{\alpha V_{ts}^* V_{td}}{2\pi \sin^2 \Theta_W} X(x_t) Q_\nu$

Use isospin symmetry and normalise to: $K^+ \rightarrow \pi^0 e^+ \nu$

$$\text{Br}(K_L \rightarrow \pi^0 \bar{\nu} \nu) = \kappa_L \left(\frac{\text{Im}(V_{ts}^* V_{td})}{\lambda^5} X(x_t) \right)^2$$

$K_L \rightarrow \pi^0 \nu \bar{\nu}$: Theoretical Status



Matrix element extracted from K_{13} decays. $N^{\frac{3}{2}}$ LO χ PT
[Mescia, Smith '07; Bijens, Ghorbani '07]

$X(x_t)$: Full NLO electroweak corrections
[Brod, Gorbahn, Stamou '10]
Reduce theory uncertainty by factor of 2

$$Br_{K_L} = 2.57(37)(4) \times 10^{-11}$$

Experiment: $< 6.7 \times 10^{-8}$ [E391a '08]

=> KOTO

$$K^+ \rightarrow \pi^+ \bar{\nu} \nu \text{ and } K_L \rightarrow \pi^0 \bar{\nu} \nu$$

Different from $K_L \rightarrow \pi^0 \bar{\nu} \nu$

- CP conserving: **Top** & **charm** contribute

$$\text{Br}(K^+ \rightarrow \pi^+ \nu \bar{\nu}(\gamma)) = \kappa_+ (1 + \Delta_{\text{EM}})$$

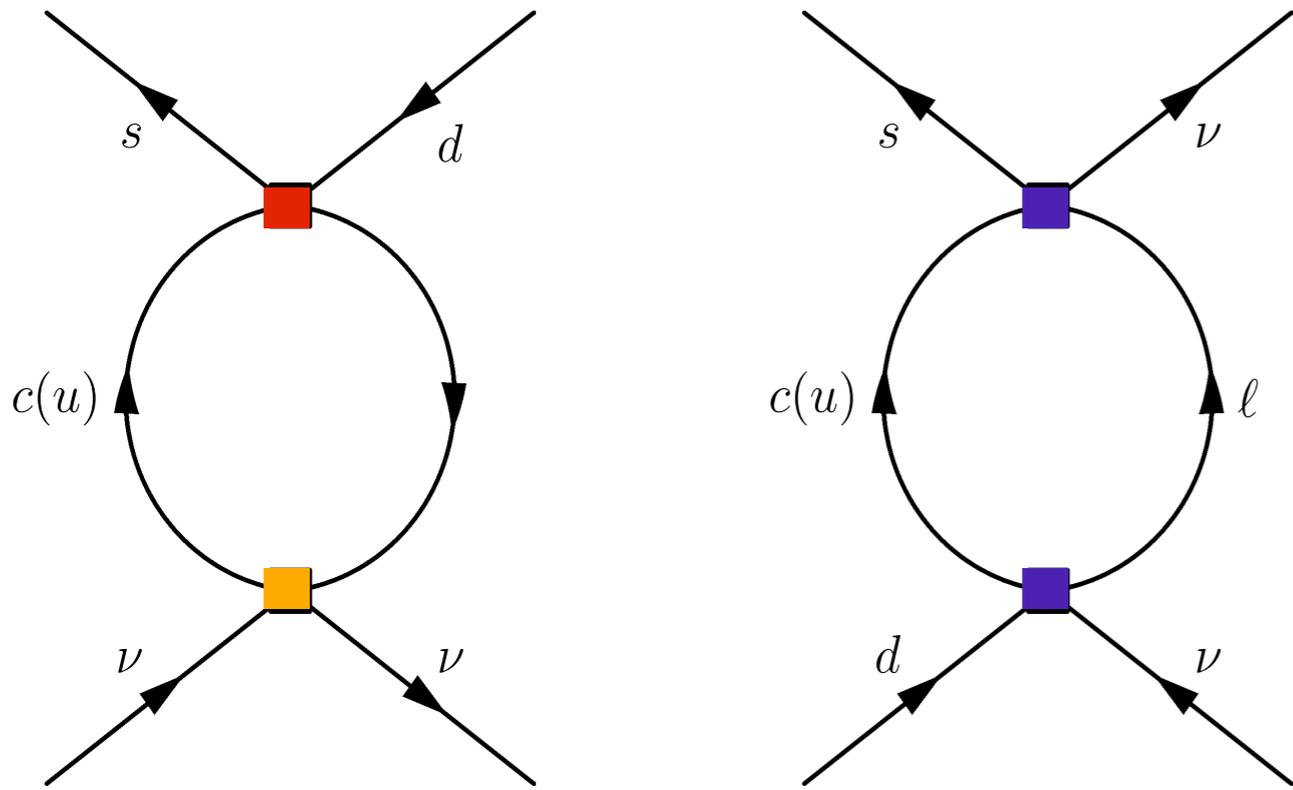
$$\times \left| \frac{V_{ts}^* V_{td} X_t(m_t^2) + \lambda^4 \text{Re} V_{cs}^* V_{cd} (P_c(m_c^2) + \delta P_{c,u})}{\lambda^5} \right|^2 \cdot$$

$\frac{m_c^2}{M_W^2}$ suppression lifted by $\log\left(\frac{m_c}{M_W}\right) \frac{1}{\lambda^4}$

Like in $K_L \rightarrow \pi^0 \bar{\nu} \nu$

- Only Q_ν : Quadratic GIM & Isospin symmetry
- Top quark contribution like in $K_L \rightarrow \pi^0 \bar{\nu} \nu$

$K^+ \rightarrow \pi^+ \bar{\nu} \nu$ charm contribution



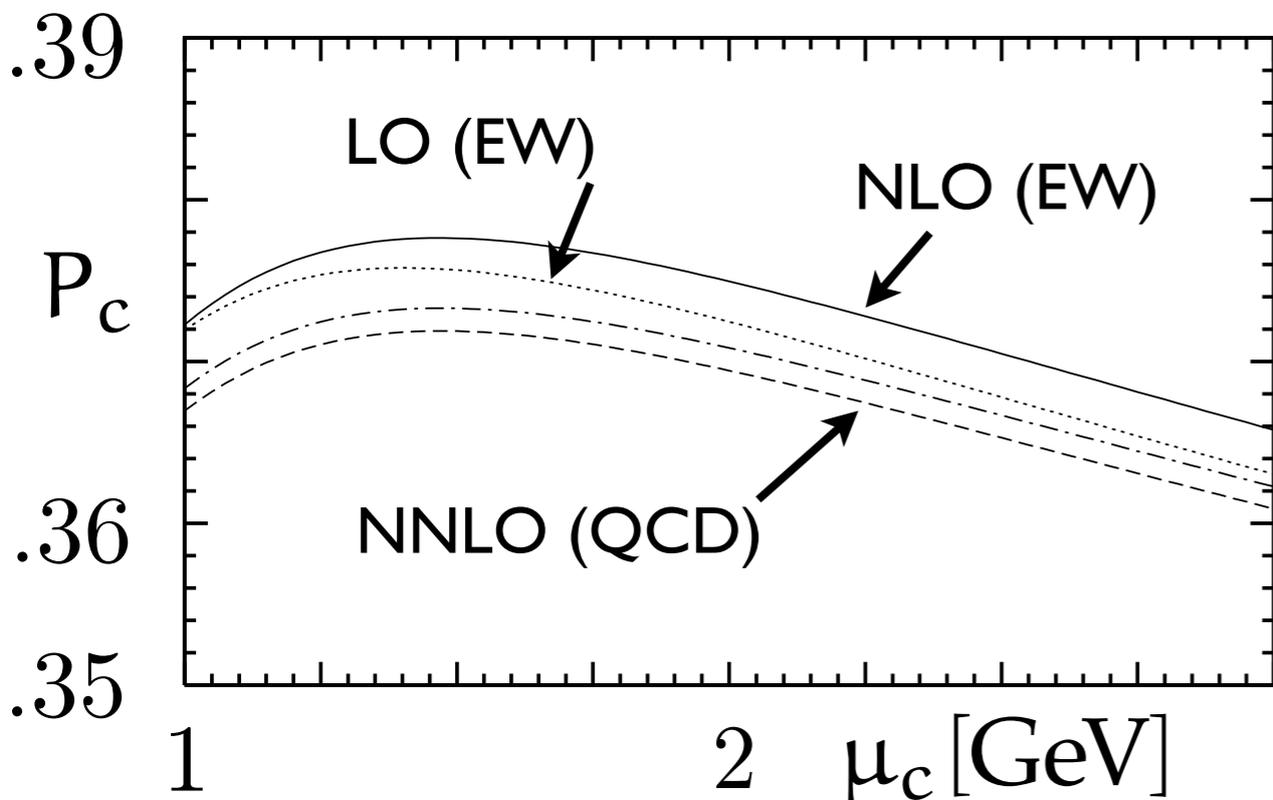
$$\bar{s}_i \gamma^\mu (1 - \gamma_5) q_j \bar{q}_j \gamma_\mu (1 - \gamma_5) d_i$$

$$\bar{s}_i \gamma^\mu (1 - \gamma_5) q_i \bar{q}_j \gamma_\mu (1 - \gamma_5) d_j$$

$$\bar{c} \gamma^\mu \gamma_5 c \bar{\nu} \gamma^\mu (1 - \gamma_5) \nu$$

$$\bar{s} \gamma^\mu (1 - \gamma_5) q \bar{\nu}_l \gamma_\mu (1 - \gamma_5) l$$

$$\bar{l} \gamma^\mu (1 - \gamma_5) \nu_l \bar{q} \gamma_\mu (1 - \gamma_5) d$$



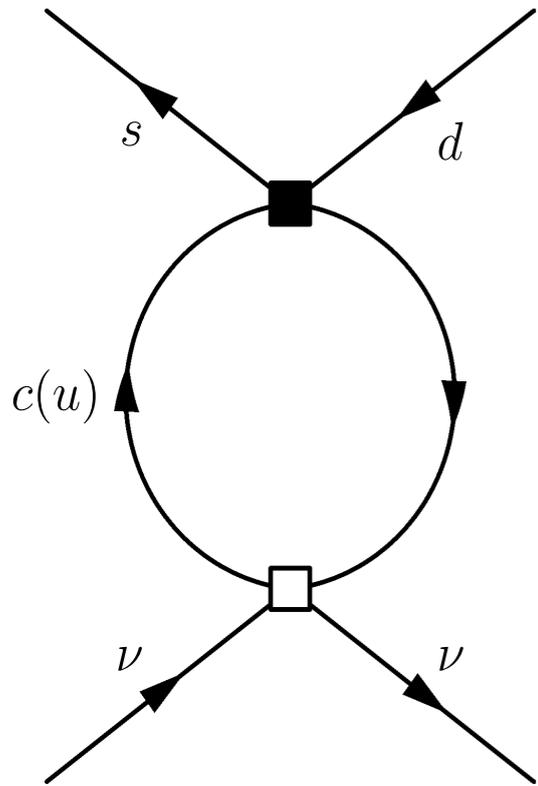
- Resum $\log \frac{m_c}{M_W}$ in P_c

P_c at NNLO: $\pm 2.5\%$ (theory)

[Buras, Gorbahn, Haisch, Nierste '06]

NLO EW [Brod, Gorbahn '08]

Long Distance Contribution



No GIM below the charm quark mass scale

q^2/m_c^2 higher dimensional operators
UV scale dependent

One loop CHPT calculation approximately
cancels this scale dependence [Isidori, Mescia, Smith '05]

Also: box-type diagrams considered
(from two semileptonic operator insertions)
cancelation is more complicated

$$\delta P_{c,u} = 0.04 \pm 0.02 \quad [\text{Isidori, Mescia, Smith '05}]$$

One Current & One Operator

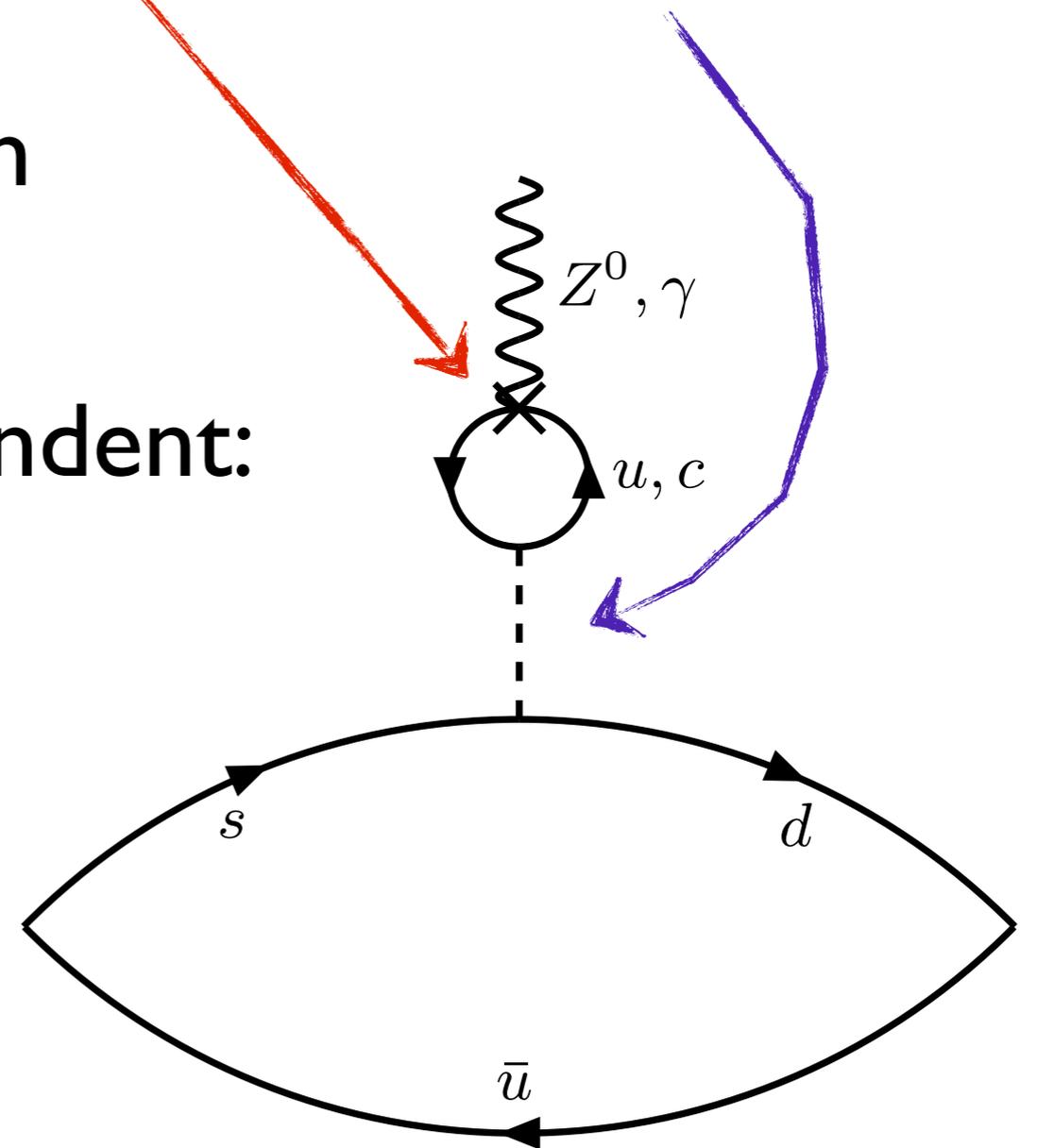
$$(\mathcal{T}_i^j)_{\text{em}}^\mu(q^2) = -i \int d^4x e^{-iq \cdot x} \langle \pi^j(p) | T \{ J_{\text{em}}^\mu(x) [Q_i^u(0) - Q_i^c(0)] \} | K^j(k) \rangle$$

Current and operator insertion

[Isidori, Martinelli, Turchetti '06]

$\mathcal{O} \left(\frac{1}{a^2} \right)$ divergence mass independent:
cancelled by GIM

$\mathcal{O} \left(\frac{1}{a} \right)$ appear \rightarrow maximally
twisted fermions



also: no semileptonic operators discussed

$K^+ \rightarrow \pi^+ \bar{\nu} \nu$ Long distance

- Matrix element extracted from K_{l3} decays

[Mescia, Smith '07]

- $K^+ \rightarrow \pi^+ \nu \bar{\nu}$ is $K^+ \rightarrow \pi^+ \nu \bar{\nu}(\gamma)$

QED radiative corrections included:

$$\Delta_{EM}(E_\gamma < 20\text{MeV}) = -0.003$$

- Uncertainty in $\kappa_+(1 - \Delta_{EM})$ reduced by $\frac{1}{7}$

- Below charm scale: Dimension 8 operators

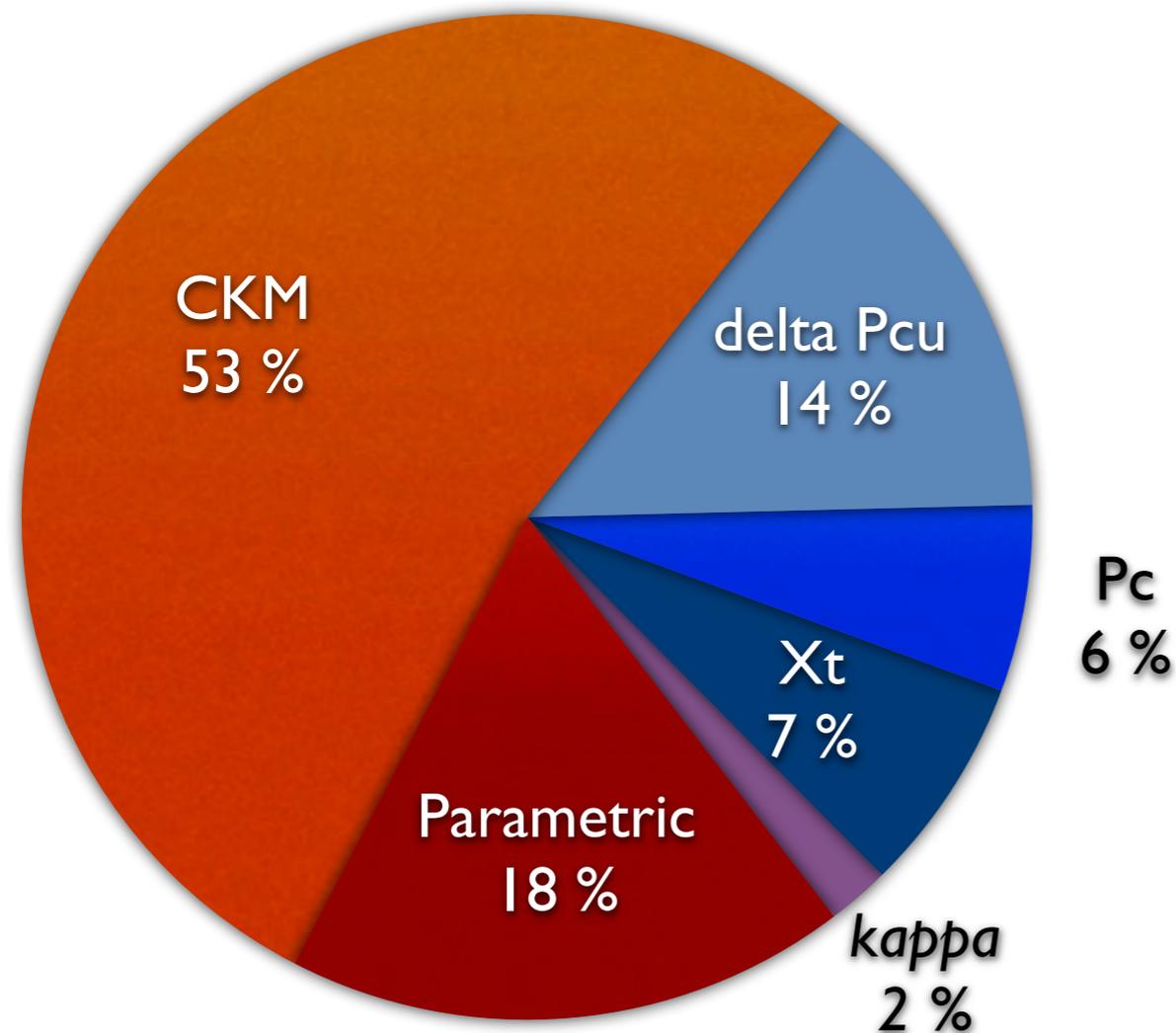
[Falk et. al. '01]

- Together with light quarks: $\delta P_{c,u} = 0.04 \pm 0.02$

[Isidori, Mescia, Smith '05]

$K^+ \rightarrow \pi^+ \bar{\nu} \nu$ Error Budget

Theory error budget $\mathcal{B}_{K^+} = 0.822(69)(29) \times 10^{-10}$

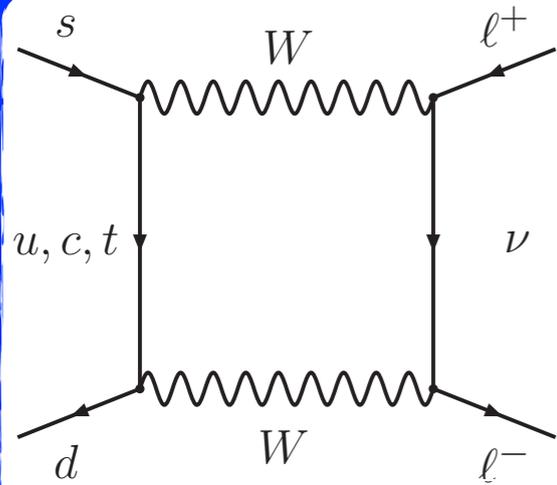


Experiment [E787, E949 '08]

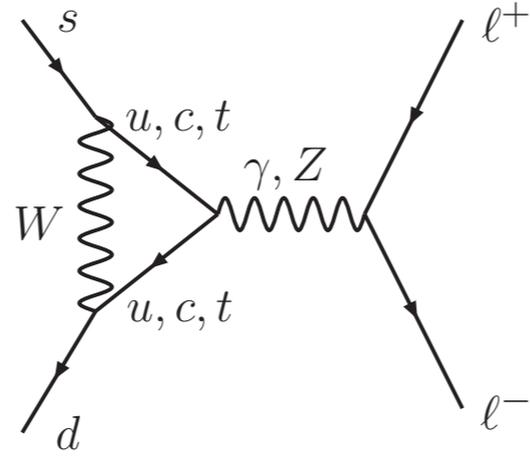
$$\text{Br}_{K^+} = (1.73^{+1.15}_{-1.05}) \times 10^{-10}$$

Talk by Goadzovski
on NA62 => 10%

$K_L \rightarrow \pi^0 l^+ l^-$: Three Contributions



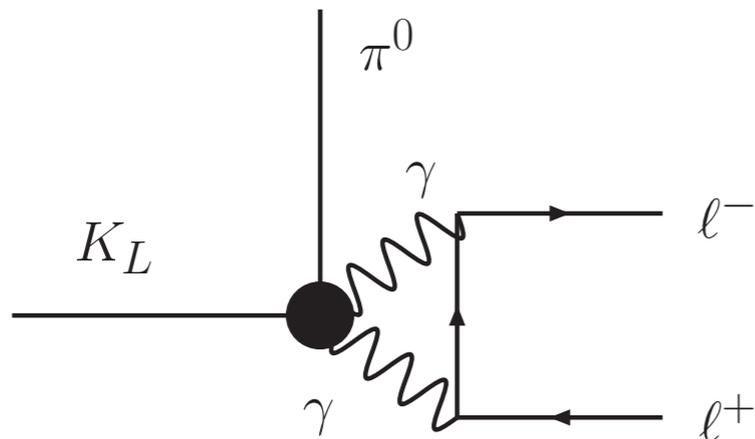
Direct CP Violating



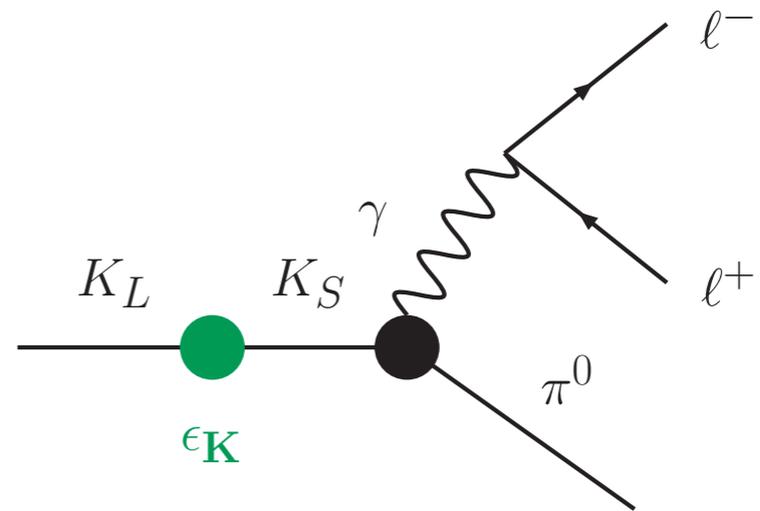
$$Q_{7V} = (\bar{s}_L \gamma_\mu d_L) (\bar{l} \gamma^\mu l) \rightarrow 1^{--}$$

$$Q_{7A} = (\bar{s}_L \gamma_\mu d_L) (\bar{l} \gamma^\mu \gamma_5 l) \rightarrow 1^{++}, 0^{-+}$$

Wilson Coefficients: y_{7V}, y_{7A}
at NLO [Buchalla et al. '96]

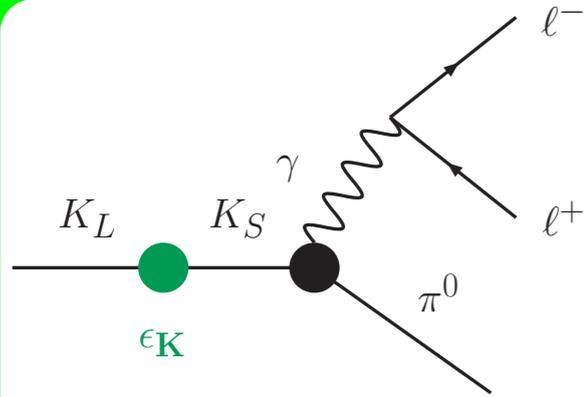


CP Conserving



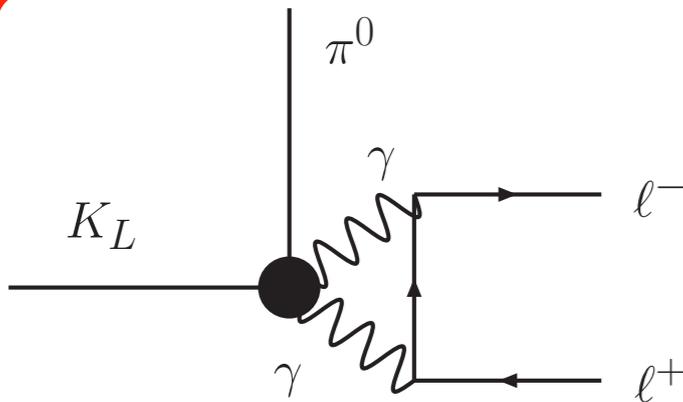
Indirect CP Violating

$K_L \rightarrow \pi^0 l^+ l^-$: Three Contributions



Counterterm $|a_S| = 1.2 \pm 0.2$ from
 [D'Ambrosio et. al. '98, Mescia et. al. '06] $K_S \rightarrow \pi^0 l^+ l^-$

For 1^{--} interference with Q_{7V}
 [Buchalla et. al. '03, Friot et al. '04]



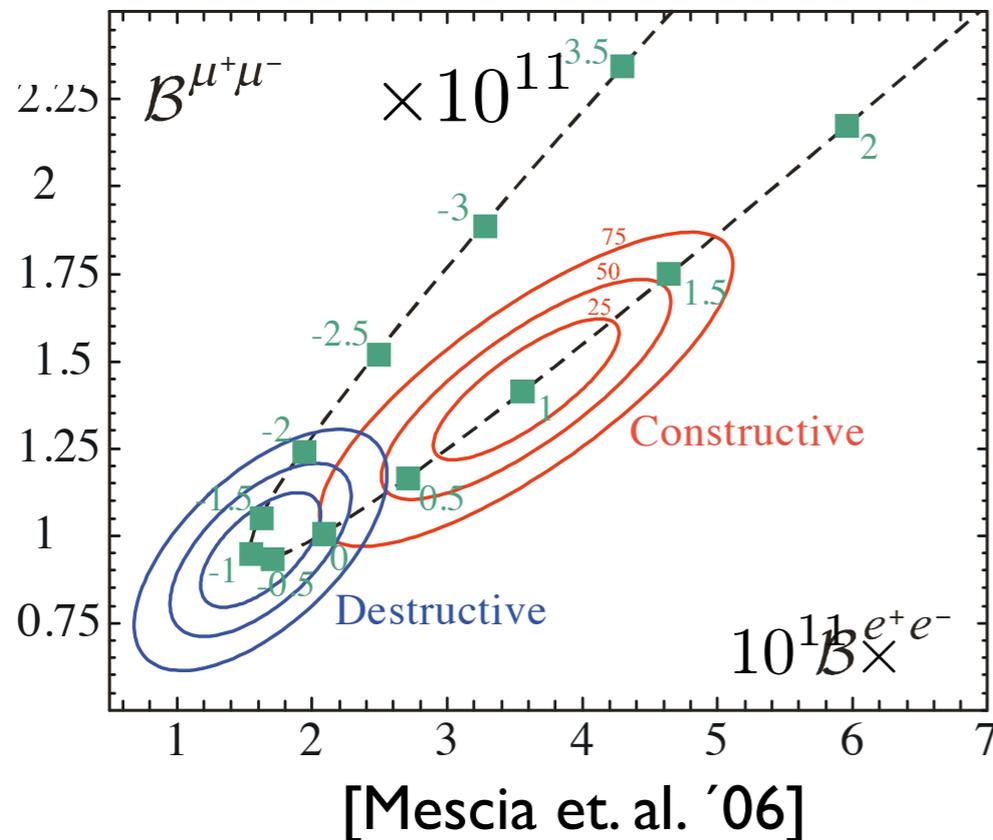
Estimate from $K_L \rightarrow \pi^0 \gamma \gamma$
 [Isidori et. al. '04]

$$\text{Br}(K_L \rightarrow \pi^0 l^+ l^-) = (C_{\text{dir}}^l \pm C_{\text{int}}^l |a_S| + C_{\text{mix}}^l |a_S|^2 + C_{\gamma\gamma}^l) \times 10^{-12}$$

l	C_{dir}^l	C_{int}^l	C_{mix}^l	$C_{\gamma\gamma}^l$
e	$(4.62 \pm 0.24)(y_V^2 + y_A^2)$	$(11.3 \pm 0.3)y_V$	14.5 ± 0.5	≈ 0
μ	$(1.09 \pm 0.05)(y_V^2 + 2.32y_A^2)$	$(2.63 \pm 0.06)y_V$	3.36 ± 0.20	5.2 ± 1.6

$K_L \rightarrow \pi^0 l^+ l^-$: Improvements

- Measure both $\text{Br}_{e^+e^-}$ and $\text{Br}_{\mu^+\mu^-}$: [Mescia et. al. '06]
Disentangle short distance contribution (y_{7V}, y_{7A})
- Dominant theory error in a_s :
Forward backward asymmetry. [Mescia, Smith, Trine '06]
Better measurement of $K_S \rightarrow \pi^0 l^+ l^-$. [Smith '07]



[KTEV '04]

$\text{Br}_{e^+e^-}$

$< 28 \times 10^{-11}$

[KTEV '00]

$\text{Br}_{\mu^+\mu^-}$

$< 38 \times 10^{-11}$

One Current & One Operator

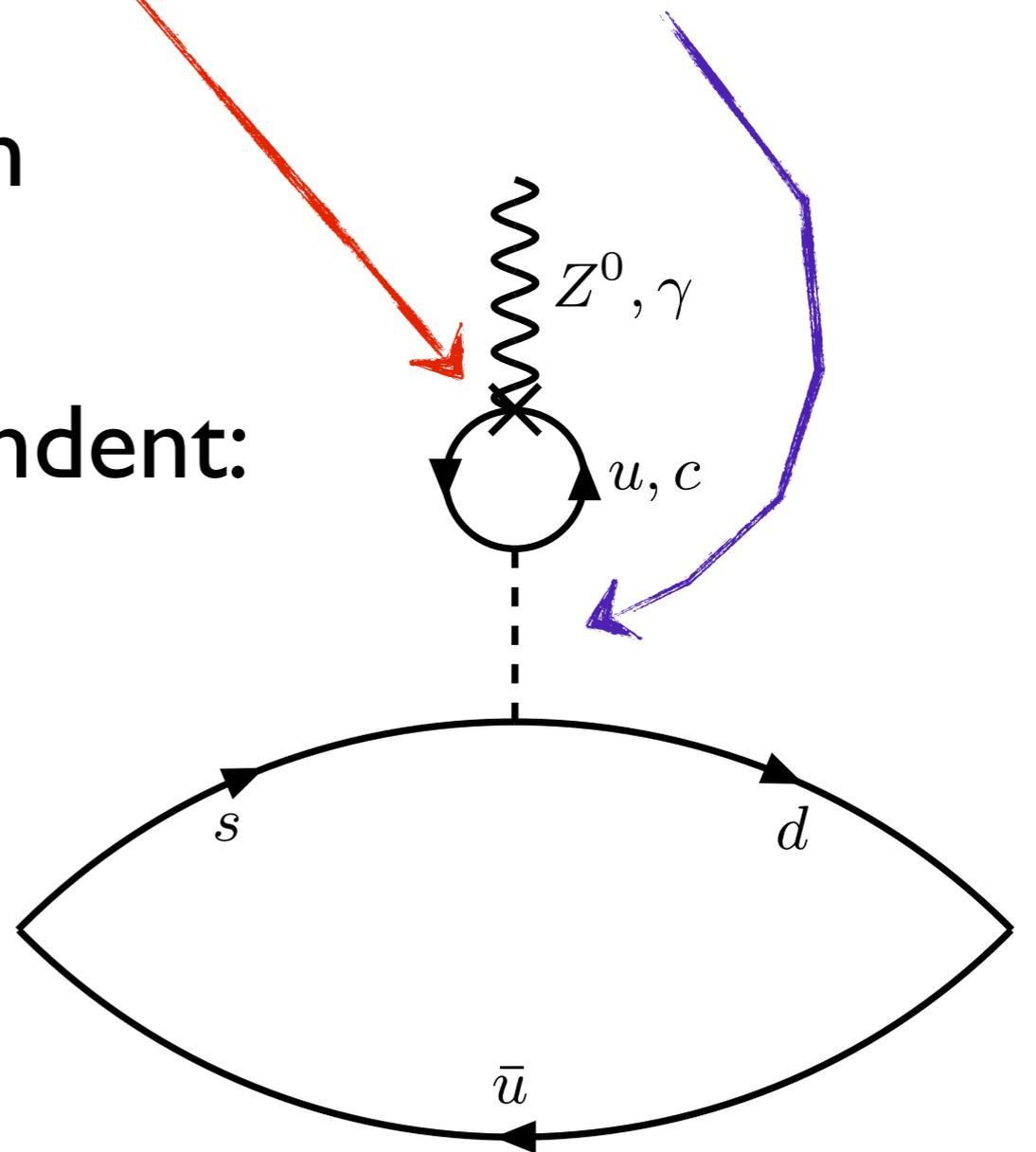
$$(\mathcal{T}_i^j)_{\text{em}}^\mu(q^2) = -i \int d^4x e^{-iq \cdot x} \langle \pi^j(p) | T \{ J_{\text{em}}^\mu(x) [Q_i^u(0) - Q_i^c(0)] \} | K^j(k) \rangle$$

Current and operator insertion

[Isidori, Martinelli, Turchetti '06]

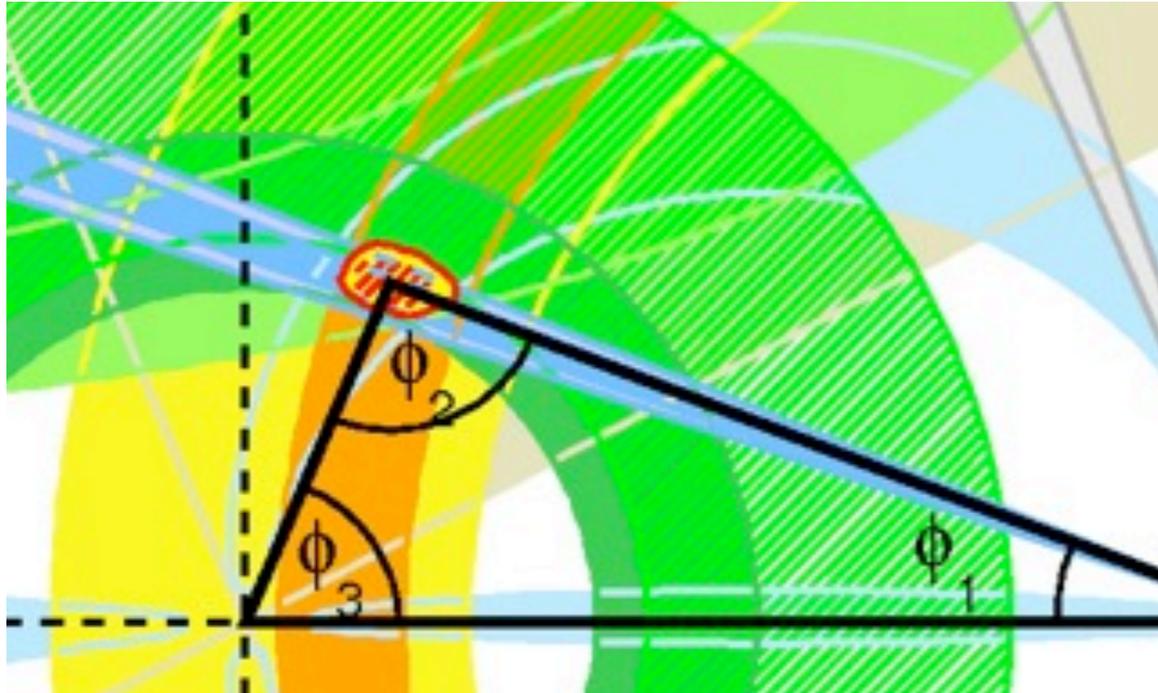
$\mathcal{O} \left(\frac{1}{a^2} \right)$ divergence mass independent:
cancelled by GIM

$\mathcal{O} \left(\frac{1}{a} \right)$ vector current \rightarrow
define conserved current



What about the two photon contribution?

ϵ_K : Indirect CP Violation



$$\epsilon_K \simeq \frac{\langle (\pi\pi)_{I=0} | K_L \rangle}{\langle (\pi\pi)_{I=0} | K_S \rangle} \frac{\Im A_0}{\Re A_0}$$

$$\epsilon_K = e^{i\phi_\epsilon} \sin \phi_\epsilon \left(\frac{\text{Im}(M_{12}^K)}{\Delta M_K} + \xi \right)$$

- In almost all old analysis: $\phi_\epsilon = 45^\circ$ and $\xi = 0$
- In reality: $\xi \neq 0$ $\phi_\epsilon \neq 45^\circ$ [Andriyash et. al.'04]

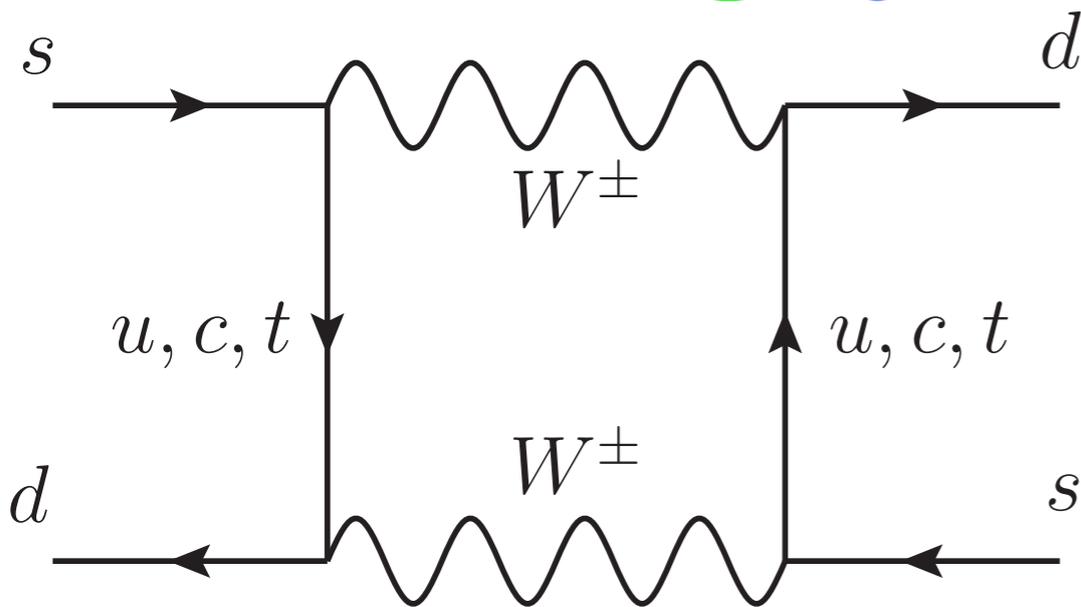
$$|\epsilon_K^{SM}| = \kappa_\epsilon |\epsilon_K| (\phi_\epsilon = 45^\circ, \xi = 0)$$

+ similar contribution as $\delta P_{c,u}$ in ϵ_K

$$\kappa_\epsilon = 0.94 \pm 0.02 \quad [\text{Buras, Guadagnoli, Isidori '10}]$$

Calculation of $M_{12}^K = \langle K^0 | \mathcal{H}_{eff}^{\Delta S=2} | \bar{K}^0 \rangle$

Box diagram
with internal **u, c, t**



$$\lambda_i \lambda_j A(x_i, x_j)$$

$$\lambda_i = V_{is}^* V_{id}$$

plus GIM:

$$\lambda_c + \lambda_t = -\lambda_u$$

Gives three different
contributions for

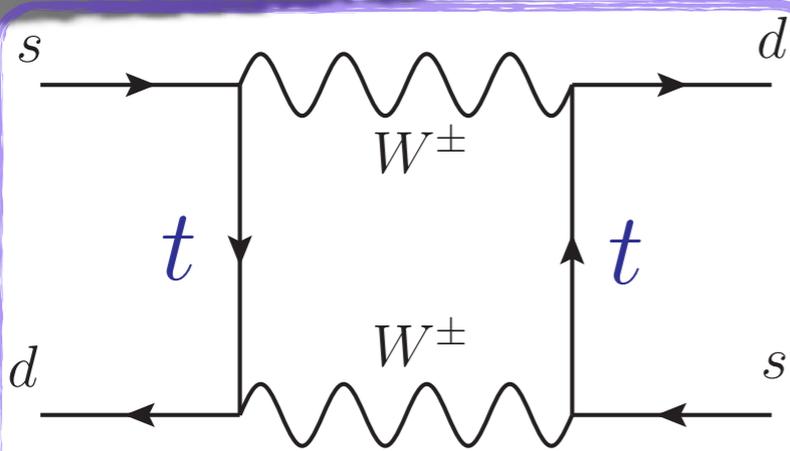
$$M_{12}^K = \langle K^0 | \mathcal{H}_{eff}^{\Delta S=2} | \bar{K}^0 \rangle$$

↑
Caveat: first only SD

$$\begin{aligned} \mathcal{H} \propto & \left[\lambda_t^2 \eta_t S(x_t) \quad \text{top} \right. \\ & + 2\lambda_c \lambda_t \eta_{ct} S(x_c, x_t) \quad \text{charm top} \\ & \left. + \lambda_c^2 \eta_c S(x_c) \right] b(\mu) \tilde{Q} \quad \text{charm} \end{aligned}$$

$$\tilde{Q} = (\bar{s}_L \gamma_\mu d_L) (\bar{s}_L \gamma^\mu d_L)$$

Calculation of $M_{12}^K = \langle K^0 | \mathcal{H}_{eff}^{\Delta S=2} | \bar{K}^0 \rangle$

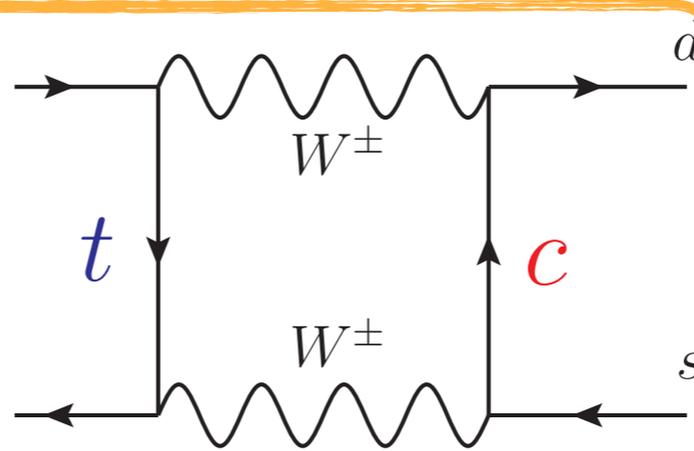


top
 $\log \chi_t$

LO $(\alpha_s \log \chi_c)^n$
NLO $\alpha_s (\alpha_s \log \chi_c)^n$

ϵ_K
scale

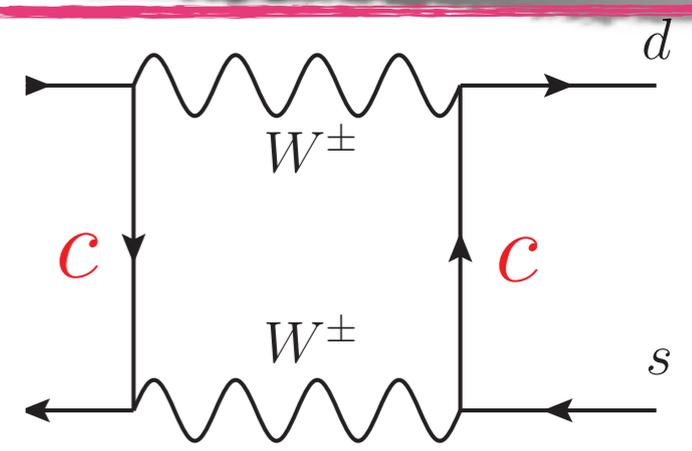
75%
1.8%



charm top
 $\log \chi_c$

LO $(\alpha_s \log \chi_c)^n \log \chi_c$
NLO $(\alpha_s \log \chi_c)^n$

37%
16%

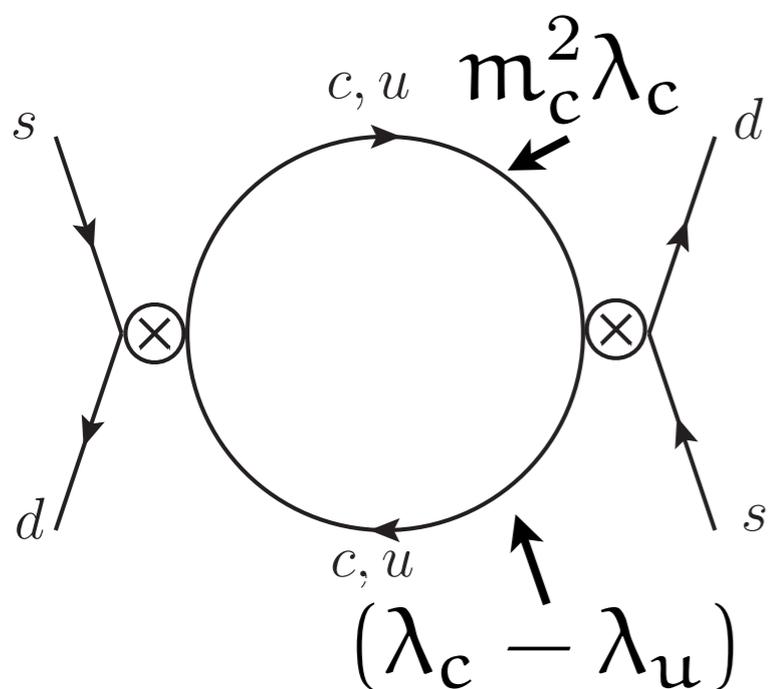
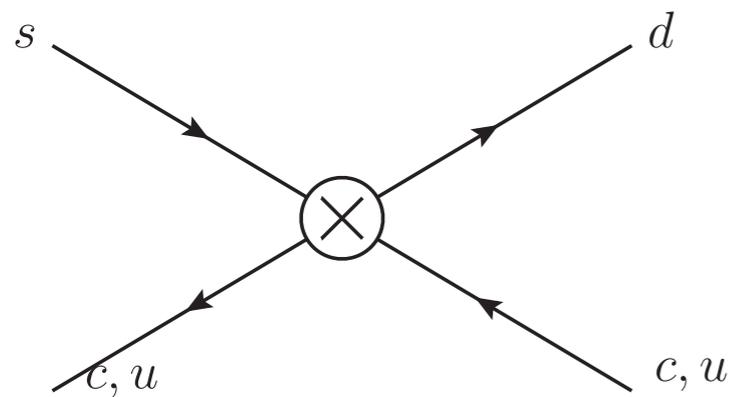
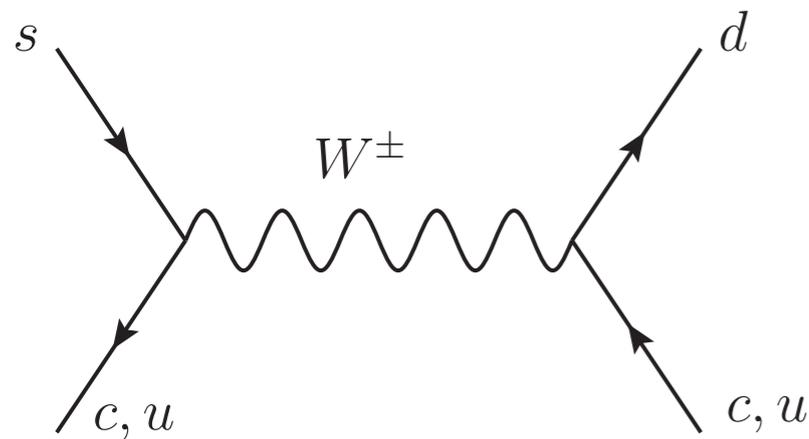


charm
 $(\log \chi_c)^0$ 
hard GIM

LO $(\alpha_s \log \chi_c)^n$
NLO $\alpha_s (\alpha_s \log \chi_c)^n$

-12%
17.7%

η_{ct} : Charm Top at LO

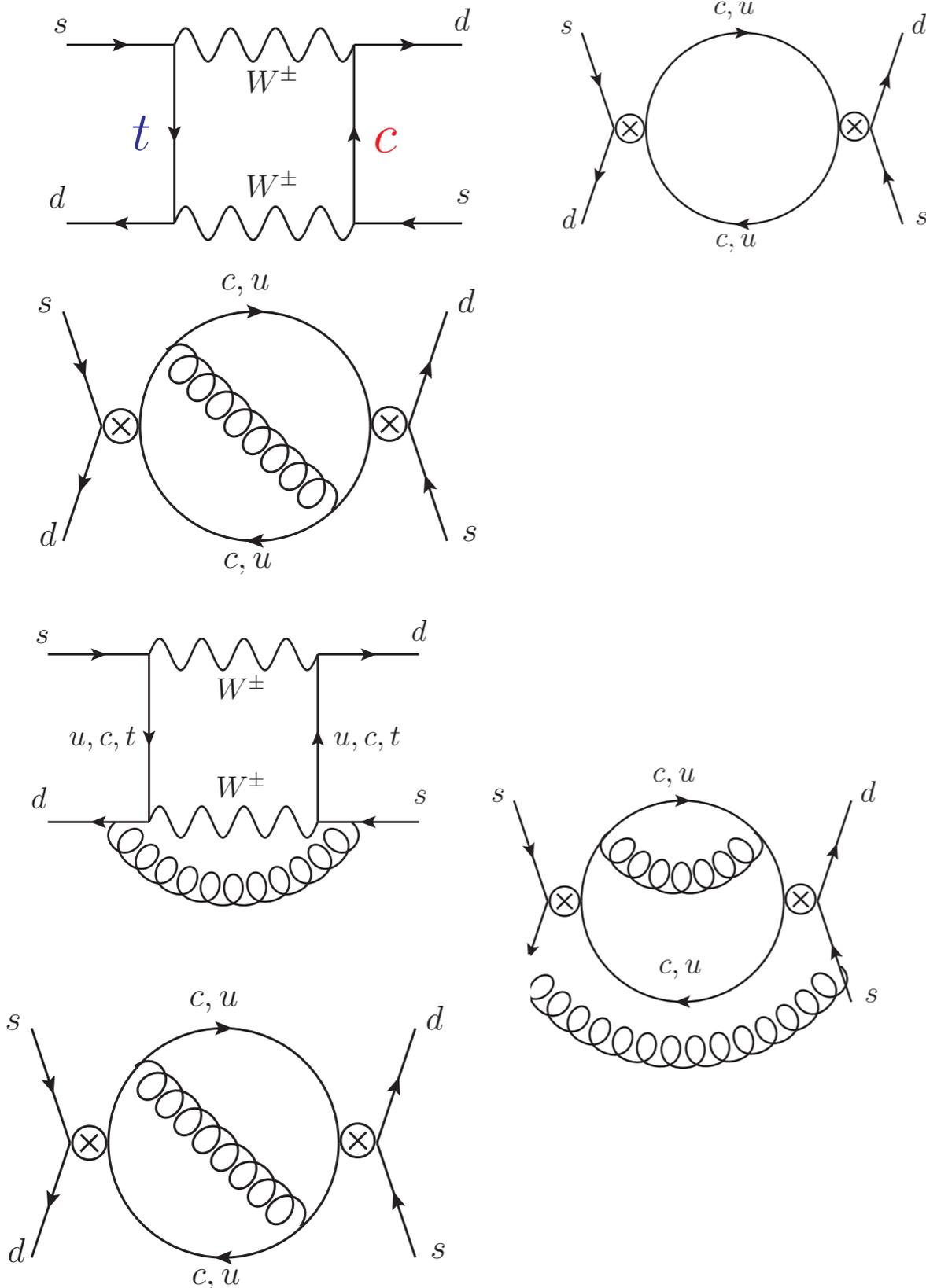


- The Leading Order result
 $(\alpha_s \log x_c)^n \log x_c$
 starts with a $\log x_c$
- Tree level matching
- One-loop Renormalisation Group Equation

$$m_c^2 \lambda_c (\lambda_c - \lambda_u) \log \frac{m_c}{M_W}$$

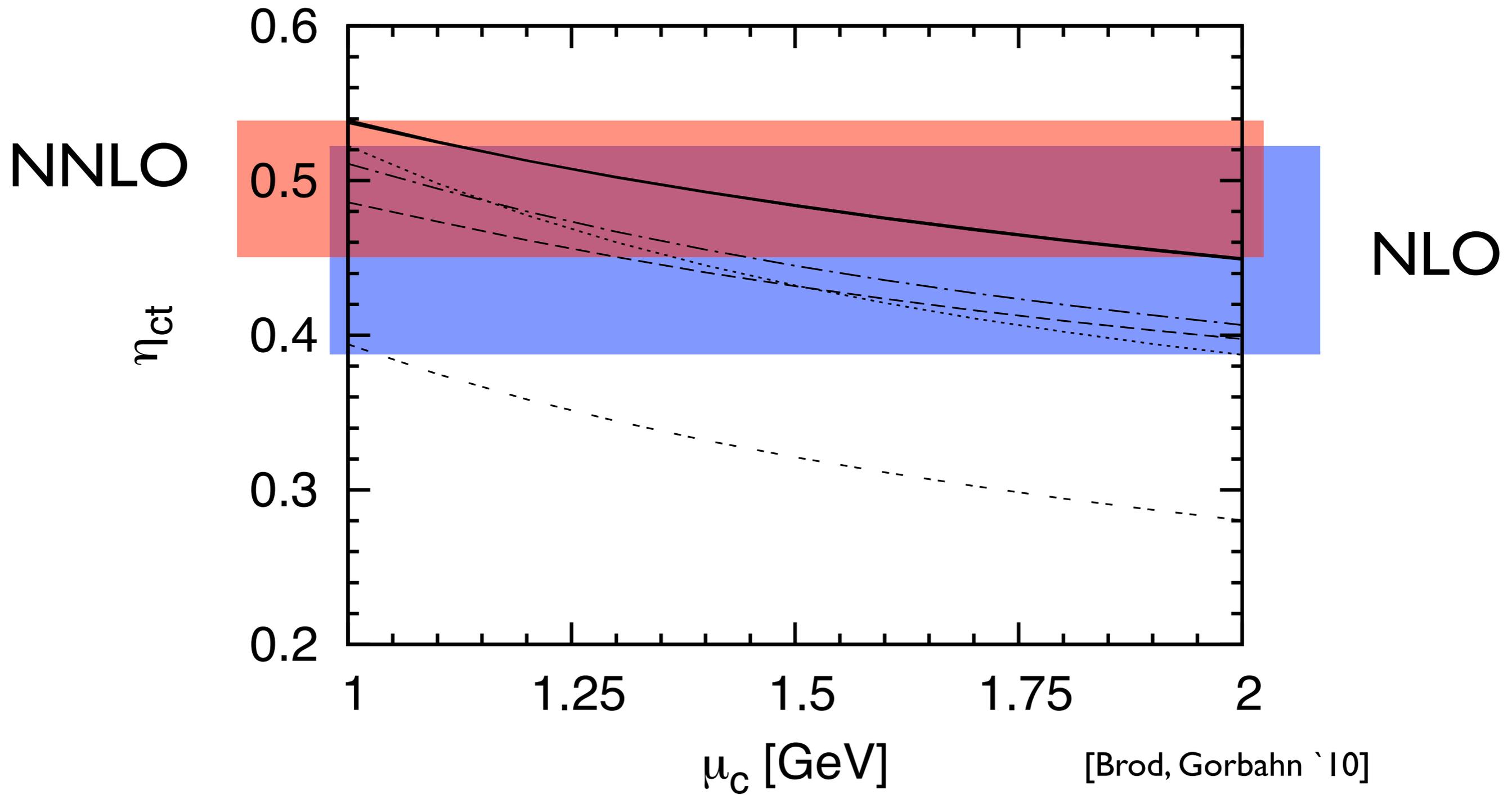
$$\rightarrow m_c^2 \lambda_c \lambda_t \tilde{Q} \log x_c$$

η_{ct} : Charm Top beyond LO



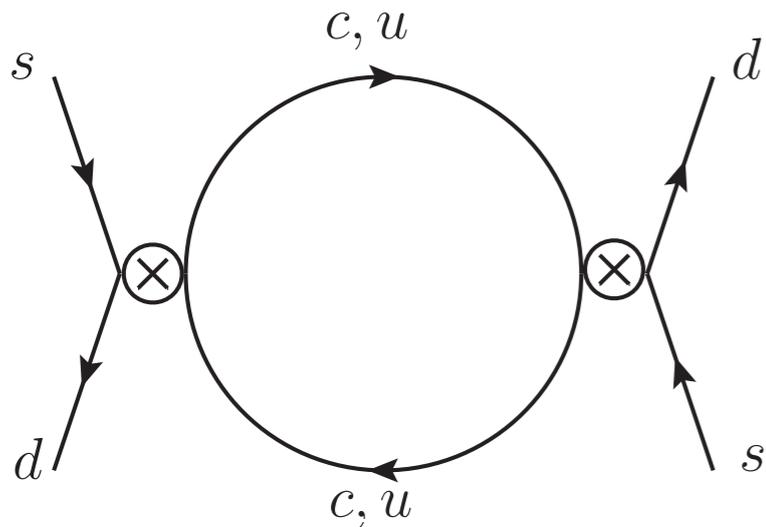
- One-loop matching at μ_t
- One-loop matching at μ_c
- Two-loop RG running
- Plus $d=6$ operators NLO [Herrlich, Nierste]
- NNLO: RGE and matching for $d=6$ operators RGE: [MG, Haisch '04], Matching: [Bobeth, et. al. '00]
- $O(10000)$ diagrams were calculated [Brod, Gorbahn '10]

η_{ct} at NNLO



Long Distance Contribution

ϵ_K LD of the matrix element is known precisely



$$\int d^4x \langle K^0 | H^{|\Delta S|=1}(x) H^{|\Delta S|=1}(0) | \bar{K}^0 \rangle$$

dispersive
part

absorptive
part

$$\epsilon_K = e^{i\phi_\epsilon} \sin \phi_\epsilon \left(\frac{\text{Im}(M_{12}^K)}{\Delta M_K} + \xi \right)$$

estimated form ϵ'

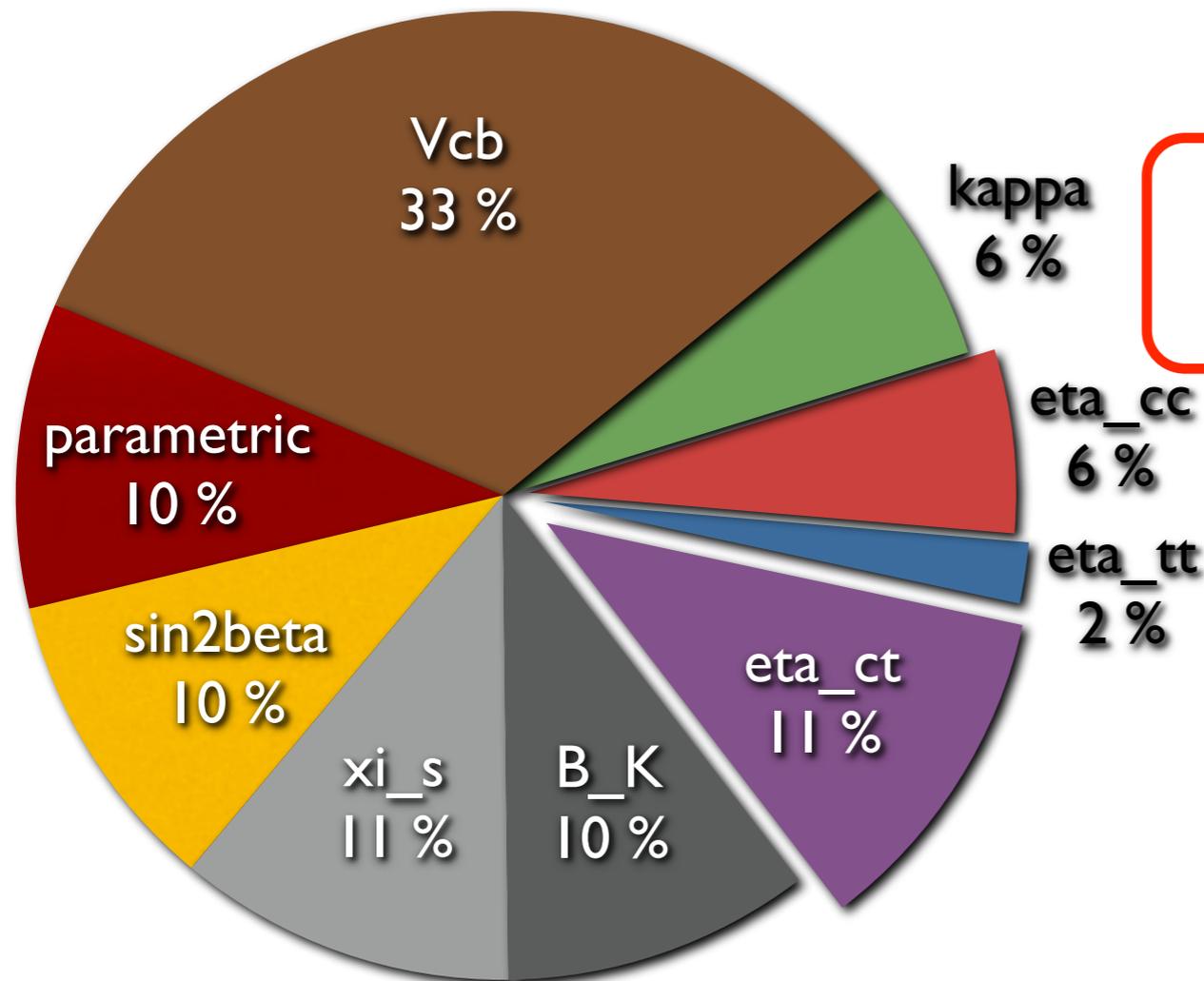
dispersive part estimated in CHPT

no higher dimensional operators and scale cancellation

put everything in: $\kappa_\epsilon = 0.94 \pm 0.02$

[Buras, Isidori, Guadagnoli '10]

$|\epsilon_K|$ and Error Budget



New input [PDG `10]

$$|\epsilon_K| = 1.89(27) \times 10^{-3}$$

using

$$\eta_{ct} = 0.494 \pm 0.046$$

$$|V_{cb}| = 406(13) \times 10^{-4}$$

Experiment [PDG `10]:

$$|\epsilon_K|^{\text{exp.}} = 2.229(12) \times 10^{-3}$$

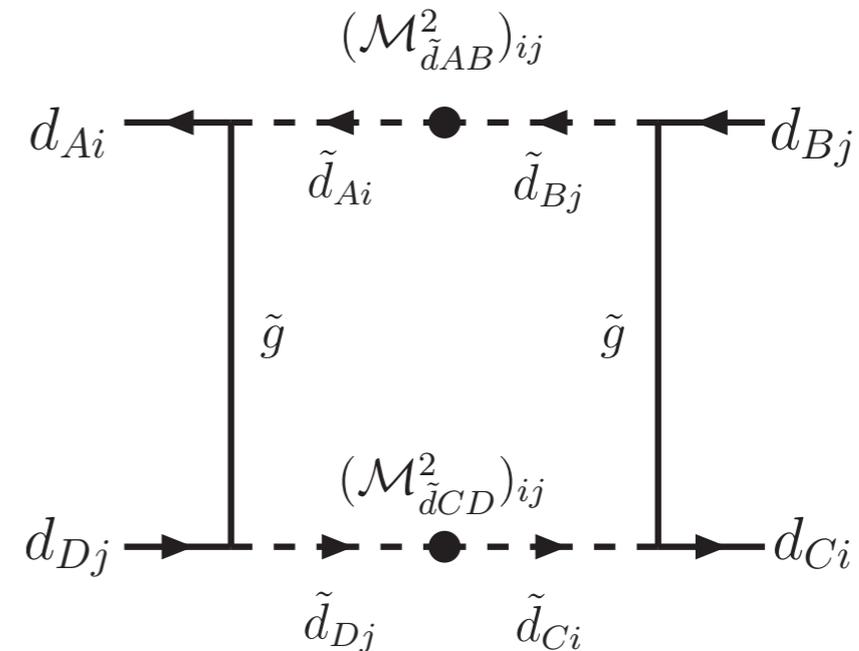
New Physics & New Operators

So far we mostly discussed
the SM background

New operators are
generated by your
favourite NP model

Talks by Petrov and Jäger

4 more operators



$$Q_1^{\text{VLL}} = (\bar{s}^\alpha \gamma_\mu P_L d^\alpha) (\bar{s}^\beta \gamma^\mu P_L d^\beta),$$

$$Q_1^{\text{LR}} = (\bar{s}^\alpha \gamma_\mu P_L d^\alpha) (\bar{s}^\beta \gamma^\mu P_R d^\beta),$$

$$Q_2^{\text{LR}} = (\bar{s}^\alpha P_L d^\alpha) (\bar{s}^\beta P_R d^\beta),$$

$$Q_1^{\text{SLL}} = (\bar{s}^\alpha P_L d^\alpha) (\bar{s}^\beta P_L d^\beta),$$

$$Q_2^{\text{SLL}} = (\bar{s}^\alpha \sigma_{\mu\nu} P_L d^\alpha) (\bar{s}^\beta \sigma^{\mu\nu} P_L d^\beta)$$

Conclusions

$K \rightarrow \pi \nu \bar{\nu}$: very clean and sensitive to short distances
Lattice could clarify the long distance contribution to K^+

the same for ϵ_K (thanks to the improvement by Lattice)
Matrix elements for NP?

Closer contact of the
perturbative and lattice
community could be
beneficial (quark masses)

Many more exciting things
in kaon physics:
 ϵ' , unitarity, lepton univ.

