harm physics from Lattice QCD Christine Davies University of Glasgow

Lattice QCD meets phenomenology Durham, Sept.2010

QCD is a key part of the Standard Model but quark confinement complicates things.



QCD only tested to 5-10% level at high energies from comparison of e.g. jet phenomena to pert.th.

But properties of hadrons calculable from QCD if fully nonperturbative calc. is done can test QCD and determine parameters very accurately (1%).





Compare to exptl rate gives $V_{qq'}$ accurately



charm physics important component of this as sits between light and bottom physics

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ \pi \rightarrow l\nu & K \rightarrow l\nu & B \rightarrow \pi l\nu \\ K \rightarrow \pi l\nu & V_{cd} & V_{cs} & V_{cb} \\ D \rightarrow l\nu & D_s \rightarrow l\nu & B \rightarrow D l\nu \\ D \rightarrow \pi l\nu & D \rightarrow K l\nu & D \\ V_{td} & V_{ts} & V_{tb} \\ \langle B_d | \overline{B}_d \rangle & \langle B_s | \overline{B}_s \rangle \end{pmatrix}$$

Also, lots of interesting new hadrons being discovered in charmonium and charm-light sectors raising QCD issues ...

Lattice QCD: solving a path integral

Path integral over gluon and quark fields on a 4-d spacetime lattice - quarks anticommute so do by hand.

 $\mathcal{L}_{QCD} = \frac{1}{2} Tr F_{\mu\nu}^2 + \overline{\psi} (\gamma \cdot D + m) \psi$ \$\sim = a huge matrix, M $\int \mathcal{D}U\mathcal{D}\psi \mathcal{D}\overline{\psi}O(\psi,\overline{\psi})e^{-S_{QCD}} \rightarrow$ Integral over gluon $\int \mathcal{D}UO(M^{-1})e^{-(S_g - \ln(\det M))}$ complicated prob, $\int e^{-i\mu ons} - in$ valence quarks distn for gluons - inc. inc. in operator effects of sea quarks Fit as fn of t to $> = \langle \dot{H}(t)H^{\dagger}(0) \rangle = \sum A_n e^{-E_n t}$ get hadron mass. A_n gives e.g. ensemble average decay constant





Lattice QCD = fully nonperturbative QCD calculation

RECIPE

- Generate sets of gluon fields for Monte Carlo integrn of Path Integral (inc effect of u, d and s sea quarks)
 - Calculate averaged "hadron correlators" from valence q props.
 - Fit for masses and simple matrix elements
 - Fix m_q and determine a to get results in physical units.
 - extrapolate to $a = 0, m_{u,d} = phys$ for real world

Lattice results need to be extrapolated to the real world where a=0 and $m_{u/d} = small$. To do this



The gold-plated meson spectrum - HPQCD 2009

inc u, d and s sea quarks critical for agreement with expt. across the board



Wednesday, 15 September 2010

Charm quarks in lattice QCD - heavy or light?

Advantages of relativistic quark formalism:

• $E_{sim} = m$

- PCAC relation (if enough chiral symmetry) gives Z = 1
 same action as for u, d, s, so cancellation in ratios

Key issue is discretisation errors:

 $E = E_{a=0}(1 + A(m_c a)^2 + B(m_c a)^4 + \dots)$ $m_c a \approx 0.4, (m_c a)^2 \approx 0.2, \alpha_s (m_c a)^2 \approx 0.06, (m_c a)^4 \approx 0.04$ for $a \approx 0.1 fm$

Need to remove *all* of these errors for precision results Highly Improved Staggered Quarks (HISQ) formalism does this, twisted mass formalism has a² errors. HPOCD

ETMC

Previous method: heavy Fermilab formalism Fermilab/MILC

Results with HISQ charm quarks on MILC 2+1 configs



Results at 5 values of a, several sea $m_{u/d}$ and m_s and two volumes CTHD, Eduardo Follana, Craig McNeile et al, HPQCD, 1008.4018



Compare to recent experimental results (lower than 2008) - no f_{Ds} puzzle remains!?



Can use to set limits on charged Higgs in some models





Xin,CKM2010



Semileptonic form factors $D \rightarrow K l \nu$



Comparison to expt gives direct determination of V_{cs} assuming Standard Model



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Harder calculations in progress



Radiative transition rates from exotic charmonia



Future - use relativistic formalisms such as HISQ or twisted mass for heavier quarks up to b, i.e. all quarks same method

• Exactly normalised operators, but need very fine lattices



Full error budgets now available for lattice calcs

| Error | m_{D_s} | f_{D_s} |
|--------------------------------|-----------|-----------|
| statistical/valence tuning | 0.094% | 0.57% |
| r_1/a | 0.025% | 0.15% |
| r_1 | 0.051% | 0.57% |
| a^2 extrapoln | 0.044% | 0.40% |
| $m_{q,sea}$ extrapoln | 0.048% | 0.34% |
| finite volume | 0% | 0.10% |
| m_{η_s} | 0.056% | 0.13% |
| em effects in D_s | 0.036% | 0.10% |
| em and annihln in m_{η_c} | 0.076% | 0.00% |
| em effects in η_c | _ | _ |
| missing c in sea | 0.01% | 0% |
| Total | 0.16% | 1.0% |
| | | |

→ will tell you what is possible in future e.g. is error from disc. errors, m_{u,d} extrapoln, stats ...

Conclusions

• relativistic formalisms with absolutely normalised currents are good for gold-plated charm physics - HISQ currently most accurate.

• HPQCD updates give m_{Ds} to 3 MeV and f_{Ds} to 1% result moved closer to expt - discrepancy now 2 σ $D \rightarrow K$ semileptonic form factors now to 2.5% gives V_{cs} in agreement with CKM unitarity

• Updates to Fermilab and twisted mass results in progress. Also harder charm calculations ...

• Relativistic techniques have allowed precision for c physics - now applying to b physics - watch this space ...

Future

• sets of 'next generation' gluon configs will have $m_{u,d}$ at physical value (so no extrapoln) or

a down to 0.03fm (so b quarks are 'light') *or much* higher statistics (for harder hadrons) also can include charm in the sea now.

• Pushing errors down to 1% level will mean em corrections and $m_u \neq m_d$ must be understood.

• some harder calculations (flavor singlet, excited states, nuclear physics) will also become possible