



# Phenomenology and Models for Minimum Bias and the Underlying Event

Mike Seymour University of Manchester

SM@LHC 11th-14th April 2011, IPPP Durham





# Recent Progress in Models for Minimum Bias and the Underlying Event

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# Recent Progress in Models for Soft Inclusive and the Underlying Event

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Minimum bias = experimental statement Models = zero bias? i.e. inclusive sample of all inelastic (non-diffractive?) events

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### General-purpose event generators for LHC physics

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### Abstract

We review the physics basis, main features and use of general-purpose Monte Carlo event generators for the simulation of proton-proton collisions at the Large Hadron Collider. Topics included are: the generation of hardscattering matrix elements for processes of interest, at both leading and nextto-leading QCD perturbative order; their matching to approximate treatments of higher orders based on the showering approximation; the parton and dipole shower formulations; parton distribution functions for event generators; non-perturbative aspects such as soft QCD collisions, the underlying event and diffractive processes; the string and cluster models for hadron formation; the treatment of hadron and tau decays; the inclusion of QED radiation and beyond-Standard-Model processes. We describe the principal features of the ARIADNE, Herwig++, PYTHIA 8 and SHERPA generators, together with the Rivet and Professor validation and tuning tools, and discuss the physics philosophy behind the proper use of these generators and tools. This review is aimed at phenomenologists wishing to understand better how parton-level predictions are translated into hadron-level events as well as experimentalists wanting a deeper insight into the tools available for signal and background simulation at the LHC.



How similar are they?

Fluctuations and correlations play crucial role



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### Fluctuations and correlations

 $\log \sigma$ **p**<sub>t</sub> Steep distribution ⇒ small sideways shift = large vertical

Rare fluctuations can have a huge influence

 $1/p_t^n \rightarrow n$ th moment

 $\Rightarrow$  corrections depend on physics process





### The Basics: Multiparton Interaction Model For small p<sub>t min</sub> and high energy inclusive parton—parton cross section is larger than total proton—proton cross section.

→ More than one parton—parton scatter per proton—proton



Need a model of spatial distribution within proton

 $\rightarrow$  Perturbation theory gives you n-scatter distributions

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- The University of Manchester Usually assume x and b factorize ( $\rightarrow$  see later)  $n_i(x, b; \mu^2, s) = f_i(x; \mu^2) G(b, s)$ 
  - and *n*-parton distributions are independent (→ see soon)  $n_{i,i}(x_i, x_i, b_i, b_i) = n_i(x_i, b_i) n_i(x_i, b_i)$
  - $\Rightarrow$  scatters Poissonian at fixed impact parameter  $\sigma_n = \int d^2b \, \frac{(A(b)\sigma^{inc})^n}{n!} \exp(-A(b)\sigma^{inc})$  $A(b) = \int d^2b_1 G(b_1) d^2b_2 G(b_2) \,\delta(b - b_1 + b_2)$



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### Flavour conservation

- The University of Manchester Primary interactions "use up" valence partons
  - additional scatters are sea quarks and gluons









# The Herwig++ Model (formerly known as Jimmy+Ivan) The Universit of Mancheste

Take eikonal+partonic scattering seriously

$$\sigma_{tot} = 2 \int d^2 b \left( 1 - e^{-\frac{1}{2}A(b)\sigma_{inc}} \right)$$
$$B = \left[ \frac{d}{dt} \left( \ln \frac{d\sigma_{el}}{dt} \right) \right]_{t=0} = \frac{1}{\sigma_{tot}} \int d^2 b \, b^2 \left( 1 - e^{-\frac{1}{2}A(b)\sigma_{inc}} \right)$$

• given form of matter distribution  $\Rightarrow$  size and  $\sigma_{inc}$ 

Bähr, Butterworth & MHS, JHEP 0901:067, 2009

• too restrictive  $\Rightarrow$ 

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$$\sigma_{tot} = 2 \int d^2 b \left( 1 - e^{-\frac{1}{2} (A_{\text{soft}}(b)\sigma_{\text{soft,inc}} + A_{\text{hard}}(b)\sigma_{\text{hard,inc}})} \right)$$

•  $\Rightarrow$  two free parameters

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### **Final state implementation**

- The Universit of Mancheste Pure independent perturbative scatters above PTMIN
  - Gluonic scattering below PTMIN with total  $\sigma_{soft,inc}$ and Gaussian distribution in p<sub>t</sub>
  - $d\sigma/dp_t$  continuous at PTMIN



### Colour reconnection model

- Röhr, Siodmok and Gieseke have implemented a new model based on momentum structure
- Refit LEP-I and LEP-II data
- Conclusion: hadronization parameters correlated with reconnection probability, but good fit can be obtained for any value of  $p_{reco}$



### Retrospective: particle flow in $WW \rightarrow 4j$ at LEP



- small effects here
- marginal improvement (if at all)

### data from [DELPHI Collaboration, Eur. Phys. J. C51 (2007) 249-269]

Colour Reconnection in Cluster Hadronisation, 8th MCnet Meeting, Cambridge, 22-24 Sept 2010

7/14







### Parameter tuning

- Procedure: • Procedure:
  - fix parton shower and hadronization parameters to LEP data, as a function of colour reconnection p<sub>reco</sub>
  - choose a total cross section and elastic slope parameter  $\Rightarrow$   $A_{soft,inc}(b)$  and  $\sigma_{tot,inc}$
  - fit  $A_{hard,inc}(b)$ ,  $p_{t,min} (\Rightarrow \sigma_{hard,inc} \text{ and } \sigma_{soft,inc})$  and  $p_{reco}$  to minimum bias and underlying event data







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Nchg

16

14

 $(x_{\min}) + R(F(x_{\max}) - F(x_{\min}))$ 

### Underlying event at 1800 GeV

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The University of Manchester N<sub>ch</sub> (away) for min-bias (transverse) for min-bias  $N_{ch}$ sum / GeV CDF data CDF data mu2 1.1 mu2 1.1 mu2 0.9 mu2 0.9 Prof. - Prof. З 2 1 0 0 1.4 1.4 MC/data MC/data 1.2 1.2 0.8 0.8 0.6 Nch (transverse) for min-bias 0.6 0 5 Nch  $p_{\perp}^{\text{lead}}$  / GeV 10 15 0 5 CDF data mu2 1.1 mu2 0.9 Prof. 0 1.4 MC/data 1.2 0.8 0.6  $p_{\perp}^{\rm lead}$  / GeV 10 0 5 15 SM@LHC Monte Carlo MCnet net  $(x_{\min}) + R(F(x_{\max}) - F(x_{\min}))$ 

### The University of Manchester Underlying event at 7000 GeV Away $N_{chg}$ density vs. $p_{\perp}^{trk_1}$ , $\sqrt{s} = 7$ TeV Transverse $\sum p_{\perp}$ density vs. $p_{\perp}^{\rm trki}$ , $\sqrt{s} = 7$ TeV 2 ---------. . . . . . . . . . 1.4 🖂 $(d^2 N_{chg}/d\eta d\phi)$ $\langle d^2 \sum p_{\perp} / d\eta d\phi \rangle$ [GeV] 1.2 1.5 τ. \*\*\*\* 0.8 1 0.6 ATLAS data ATLAS data mu2 1.1 mu2 1.1 0.4 0.5 mu2 0.9 mu2 0.9 Prof. Prof. 0.2 0 1.4 0 1.4 MC/data 1.2 MC/data 1.2 1 1 0.8 0.8 0.6 Transverse $N_{chg}$ density vs. $p_{\perp}^{trk_1}$ , $\sqrt{s} = 7$ TeV 0.6 Transverse $(p_{\perp})$ vs. $N_{chg}$ , $\sqrt{s} = 7$ TeV 6 $(d^2 N_{chg}/d\eta d\phi)$ (p\_) [GeV] 2 4 6 2 4 1.2 1 0.8 o.8 0.6 0.6 ATLAS data ATLAS data 0.4 0.4 mu2 1.1 mu2 1.1 mu2 0.9 mu2 0.9 0.2 0.2 Prof. Prof. 0 0 1.4 1.4 MC/data 1.2 MC/data 1.2 1 0.8 **o.8** 0.6 0.6 ----10 12 16 18 2 4 6 8 14 20 5 10 15 20 25 $p_{\perp}$ (leading track) [GeV] Nchg Monte

MCnet

Carlo

net

 $min) + R(F(x_{max}) - F(x_{min}))$ 

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### Tuning conclusion

- The University of Manchester Not possible to fit with energy-independent parameters
  - Possible to fit with energy-dependent p<sub>t.min</sub> and all else energy-independent

| • e.q. | For µ <sup>2</sup> = | 1.1 GeV <sup>2</sup>                                 | 0.9 GeV <sup>2</sup>  |                |
|--------|----------------------|--|---|----------------|
| Ŭ      | and √s=              | PTMIN=   | PTMIN=  |                |
|        | 900                  | 2.34   | 2.17  |                |
|        | 1800                 | 3.09   | 2.80  |                |
|        | 2760                 | 3.31   | 2.92 -  | Prediction/    |
|        | 7000                 | 4.02   | 3.36  | Recommendation |
| SM@LHC | M                    | Cnet $\frac{d\sigma}{dx_1 dx_2} = \frac{2\pi}{2\pi}$ | Monte<br>Carlo<br>net<br>$x_{min} + R(F(r_{max}) - F(r_{optin}))$ | Mike Seymour   |



### Tuning conclusion

- The University of Manchester Not possible to fit with energy-independent parameters
  - Possible to fit with energy-dependent p<sub>t.min</sub> and all else energy-independent
  - e.g.



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### **Tuning conclusion**

The University of Manchester Public version: 2.5.0, released 8<sup>th</sup> Feb 2011

- does not come with latest tuned values
  - check

http://projects.hepforge.org/herwig/trac/wiki/MB\_UE\_tunes for updates





- Most existing models use factorization of x and b
  - or (Herwig++) crude separation into hard and soft components (simple hot-spot model)
- R.Corke and T.Sjöstrand, arXiv:1101.5953 consider Gaussian matter distribution with width



Figure 1: (a) The rise of the total and non-diffractive pp cross section with energy, and (b) the ratio  $a_0(E_{\rm CM})/a_0(200 \,{\rm GeV})$ , over the same energy range, for a set of different  $a_1$  values

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- Most existing models use factorization of x and b
  - or (Herwig++) crude separation into hard and soft components (simple hot-spot model)
- R.Corke and T.Sjöstrand, arXiv:1101.5953 consider Gaussian matter distribution with width  $a(x) = a_0 \left(1 + a_1 \ln \frac{1}{x}\right)$
- for  $a_1 \approx 0.15$ , matter distribution can be E-indep





Figure 11: Tune 4C, using the log profile, and with a raised  $p_{\perp 0}$  in the MPI framework, compared against an overlap profile with p = 1.6, also with a raised  $p_{\perp 0}$ , and LHC data



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- (My) conclusion: for soft inclusive and jet underlying event data compatible with data but not required, but sheds interesting light on energy dependence
- Interesting correlation with hardness of hard scatter, e.g. less underlying event in 1 TeV Z' events than in Z events





## Conclusions

Despite ~25 year history, multi-parton interaction models are still in their infancy

- LHC experiments'
  - step up in energy
  - high efficiency, purity and phase space coverage
  - emphasis on physical definition of observables
  - have given us a huge amount of useful data
- existing models describe data well with tuning
- need more understanding of correlations/corners of phase space/relations between different model components

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### Conclusions

The University of Manchester don't forget that jet corrections depend on correlations and high moments of distributions and are physics-process dependent



### Backup slides

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## 'Interesting features' of Herwig++ #1

- The additional scatters are not p<sub>t</sub> ordered, so it can occasionally happen that a high p<sub>t</sub> jet comes from a low p<sub>t</sub> primary scattering event
  - this is a disaster if you generate weighted primary scatters or mix event samples with different  $p_t$  ranges
  - it is safe to remove such events from your sample
    - provided they are a small fraction of the eikonal cross section
    - i.e. provided it is an underlying event not part of a soft inclusive sample





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## 'Interesting features' of Herwig++ #2

- Soft inclusive event generation is built on top of the underlying event machinery: a fictional (totally ineffective) hard process is needed
  - unfortunately ThePEG does not know about this and reports the cross section of the fictional process as the total cross section
    - needs to be fixed! But in the meantime, extract the correct cross section from the log file!



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## 'Interesting features' of Herwig++ #3

- Double- (or more) scattering is built in just
  select a list of the hard processes you want to
  include in each event
  - unfortunately ThePEG does not know about this and reports the cross section of the first process as the total cross section
    - needs to be fixed! But in the meantime, extract the correct cross section from the log file!

