Introduction to Jet Finding and Jetography (2)

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A full set of IRC-safe jet algorithms

Generalise inclusive-type sequential recombination with

$$d_{ij} = \min(k_{ti}^{2\mathbf{p}}, k_{tj}^{2\mathbf{p}}) \Delta R_{ij}^2 / R^2$$
 $d_{iB} = k_{ti}^{2\mathbf{p}}$

	Alg. name	Comment	time
<i>p</i> = 1	k _t	Hierarchical in rel. k_t	
	CDOSTW '91-93; ES '93		NIn N exp.
<i>p</i> = 0	Cambridge/Aachen	Hierarchical in angle	
	Dok, Leder, Moretti, Webber '97	Scan multiple <i>R</i> at once	N In N
	Wengler, Wobisch '98	$\leftrightarrow QCD \text{ angular ordering}$	
<i>p</i> = -1	anti- k_t Cacciari, GPS, Soyez '08	Hierarchy meaningless, jets	
	\sim reverse- k_t Delsart	like CMS cone (IC-PR)	$N^{3/2}$
SC-SM	SISCone	Replaces JetClu, ATLAS	
	GPS Soyez '07 + Tevatron run II '00	MidPoint (xC-SM) cones	$N^2 \ln N \exp$.

Compromise between having a limited set of algs. and a good range of complementary properties

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Towards an understanding of jets

How a jet is and isn't like a parton — quantitatively

And how this relationship is affected by the jet radius

Small jet radius

single parton @ LO: jet radius irrelevant

Small jet radius

Large jet radius



Small jet radius





perturbative fragmentation: large jet radius better (it captures more)

Small jet radius



Large jet radius



non-perturbative fragmentation: large jet radius better (it captures more)



underlying ev. & pileup "noise": **small jet radius better** (it captures less)

Small jet radius



Large jet radius



multi-hard-parton events: **small jet radius better** (it resolves partons more effectively)

Parton pt v. jet pt

3 physical effects:

Gluon radiation from the parton
 Hadronisation
 Underlying Event

One important consideration:

Whether the parton is a quark or a gluon [quarks radiate with colour factor $C_F = 4/3$ gluons radiate with colour factor $C_A = 3$]

e.g. de Florian & Vogelsang '07

The question's dangerous: a "parton" is an ambiguous concept

Three limits can help you:

Threshold limit

[Understanding jets]

 \lfloor [Parton p_t v. jet p_t]

- ► Parton from color-neutral object decay (Z')
- Small-R (radius) limit for jet

One simple result (small-*R* limit)

$$\frac{\langle p_{t,jet} - p_{t,parton} \rangle}{p_t} = \frac{\alpha_s}{\pi} \ln R \times \begin{cases} 1.01 C_F & quarks\\ 0.94 C_A + 0.07 n_f & gluons \end{cases} + \mathcal{O}(\alpha_s)$$

only $\mathcal{O}(\alpha_s)$ depends on algorithm & process cf. Dasgupta, Magnea & GPS '07

Jet p_t v. parton p_t : hadronisation?

Hadronisation: the "parton-shower" \rightarrow hadrons transition

Method:

[Understanding jets]

 \lfloor [Parton p_t v. jet p_t]

- "infrared finite α_s"
- **prediction** based on e^+e^- event shape data
- could have been deduced from old work

à la Dokshitzer & Webber '95

Korchemsky & Sterman '95 Seymour '97

Main result

$$\langle p_{t,jet} - p_{t,parton-shower} \rangle \simeq - rac{0.4 \text{ GeV}}{R} imes \left\{ egin{array}{c} C_F & quarks \ C_A & gluons \end{array}
ight.$$

cf. Dasgupta, Magnea & GPS '07

coefficient holds for anti- k_t ; see Dasgupta & Delenda '09 for k_t alg.

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"Naive" prediction (UE \simeq colour dipole between *pp*): $\Delta p_t \simeq 0.4 \text{ GeV} \times \frac{R^2}{2} \times \begin{cases} C_F & q\bar{q} \text{ dipole} \\ C_A & \text{gluon dipole} \end{cases}$

lodern Monte Carlo tunes tell you (
$$\sqrt{s}=$$
 7 TeV) $\Delta
ho_t \simeq {f 8} \,\, {f GeV} imes {R^2 \over 2} \simeq 1.2 \,\, {f GeV} imes (\pi R^2)$

This big coefficient motivates special effort to understand interplay between jet algorithm and UE: "jet areas" How does coefficient depend on algorithm? How does it depend on jet p_t ? How does it fluctuate? cf. Cacciari, GPS & Soyez '08 "Naive" prediction (UE \simeq colour dipole between *pp*): $\Delta p_t \simeq 0.4 \text{ GeV} \times \frac{R^2}{2} \times \begin{cases} C_F & q\bar{q} \text{ dipole} \\ C_A & \text{gluon dipole} \end{cases}$

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Ν

A qualitative example: top reconstruction



 $\frac{\text{Game: measure top mass to 1 GeV}}{\text{example for Tevatron}}$ $m_t = 175 \text{ GeV}$

 Small R: lose 6 GeV to PT radiation and hadronisation, UE and pileup irrelevant

 Large R: hadronisation and PT radiation leave mass at ~ 175 GeV, UE adds 2 – 4 GeV.



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6 partons v. 6 jets?



6 partons v. 6 jets?





Using our understanding to help discover a dijet resonance, $q\bar{q} \rightarrow X \rightarrow q\bar{q}$.

E.g. to reconstruct $m_X \sim (p_{tq} + p_{t\bar{q}})$ $\frac{\text{PT radiation:}}{q: \quad \langle \Delta p_t \rangle \simeq \frac{\alpha_{\text{s}} C_{\text{F}}}{\pi} p_t \ln R}$ q Hadronisation: $\overline{q:\quad \langle \Delta p_t
angle \simeq - rac{C_F}{R} \cdot 0.4 \; {
m GeV}}$ q q р р **Underlying event:** $\overline{q,g:} \ \langle \Delta p_t
angle \simeq rac{R^2}{2} \cdot 2.5 - 15 \; {
m GeV}$ a

Minimise fluctuations in p_t

Use crude approximation:

 $\langle \Delta p_t^2
angle \simeq \langle \Delta p_t
angle^2$

in small-*R* limit (!) NB: full calc, correct fluct: Soyez '10



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Dijet mass: scan over R [Pythia 6.4]



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Dijet mass: scan over *R* [Pythia 6.4]











After scanning, summarise "quality" v. R. Minimum \equiv BEST picture not so different from crude analytical estimate

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Best R is at minimum of curve

 Best *R* depends strongly on mass of system
 Increases with mass can reproduce this anayltically Soyez '10

Message received by CMS: they combine all R = 0.5 jets ($p_t >$ 10 GeV) within $\Delta R = 1.1$ of two hardest to improve resolution. ATLAS '11 still just use R = 0.6



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http://quality.fastjet.fr/



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Fat jets boosted massive hadronically decaying objects

E.g. when a known particle, W, Z or a top \rightarrow a single jet or a new particle, Higgs, gluino, neutralino \rightarrow a single jet

This will be common for electroweak-scale objects at LHC: $m_W, \, m_t \ll 14 \,\, {\rm TeV}$

 $[1 \text{ jet} \gtrsim 2 \text{ partons}]$

E.g. $X \rightarrow t\bar{t}$ resonances of varying difficulty



RS KK resonances $\rightarrow t\bar{t}$, from Frederix & Maltoni, 0712.2355

NB: QCD dijet spectrum is $\sim 10^3$ times $t\bar{t}$

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 $[1 \ {
m jet} \gtrsim 2 \ {
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Boosted massive particles, e.g.: EW bosons

Hadronically decaying EW boson at high $p_t \neq two$ jets



Rules of thumb:

 $m = 100 \text{ GeV}, \ p_t = 500 \text{ GeV}$

$$R < \frac{2m}{p_t}$$
: always resolve two jets $R < 0.4$ $R \gtrsim \frac{3m}{p_t}$: resolve one jet in ~75% of cases $(\frac{1}{8} < z < \frac{7}{8})$ $R \gtrsim 0.6$

Boosted ID strategies







Select on the jet mass with one large (cone) jet Can be subject to large bkgds [high- p_t jets have significant masses]

Choose a small jet size (*R*) so as to resolve two jets Easier to reject background if you actually see substructure [NB: must manually put in "right" radius]

Take a large jet and split it in two Let jet algorithm establish correct division



Jet mass gives clear sign of massive particles inside the jet;



Jet mass gives clear sign of massive particles inside the jet;



Jet mass gives clear sign of massive particles inside the jet; but QCD jets are massive too — must learn to reject them

QCD principle: soft divergence



Background



Splitting probability for Higgs:

 $P(z) \propto 1$

Splitting probability for quark:

 $P(z) \propto rac{1+z^2}{1-z}$

1/(1-z) divergence enhances background

Remove divergence in bkdg with cut on z Can choose cut analytically so as to maximise S/\sqrt{B}

Originally: cut on (related) k_t -distance

Butterworth, Cox & Forshaw '02



QCD jet mass distribution has the approximate

$$\frac{dN}{d\ln m} \sim \alpha_{\rm s} \ln \frac{p_t R}{m} \times {\rm Sudakov}$$

Work from '80s and '90s + Almeida et al '08

The logarithm comes from integral over soft divergence of QCD:

 $\int_{\frac{m^2}{p_t^2 R^2}}^{\frac{1}{2}} \frac{dz}{z}$

A hard cut on z reduces QCD background & simplifies its shape



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How well can an algorithm identify the "blobs" of energy inside a jet that come from different partons?

This is crucial for identifying the kinematic variables of the partons in the jet (e.g. z).

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Its last step is to merge two hard pieces. Easily undone to identify underlying kinematics



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This meant it was the first algorithm to be used for jet substructure.

Seymour '93

Butterworth, Cox & Forshaw '02










Identifying jet substructure: Cam/Aachen



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C/A identifies two hard blobs with limited soft contamination, joins them, and then adds in remaining soft junk

The interesting substructure is buried inside the clustering sequence — it's less contamined by soft junk, but needs to be pulled out with special techniques

Butterworth, Davison, Rubin & GPS '08 Kaplan, Schwartz, Reherman & Tweedie '08 Butterworth, Ellis, Rubin & GPS '09 Ellis, Vermilion & Walsh '09



Example improvement from boosted regime

Search for main decay of light Higgs boson in W/Z+H, H \rightarrow $b\bar{b}$



restricting search to $p_{tH} > 200 \text{ GeV}$ using the method from Butterworth, Davison, Rubin & GPS '08

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Closing

LHC events will cover 2 orders of magnitude in jet p_t

Flexibility in the choice of jet definitions has potential to bring significant gains [there is no unique best definition; anti- k_t with R = 0.5 or 0.6 will sometimes be far from optimal]

EW-scale particles are "light" relative to the TeV scale Using the full power of jet algorithms & their substructure helps pull out signals that might otherwise be missed [currently a very active research field]

EXTRAS

quality: 5 algorithms, 3 processes



 $[1 \text{ jet} \gtrsim 2 \text{ partons}]$

└[Quality]

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└[Quality]

Other work \rightarrow improving the methods

- Using matrix-element methods for the substructure Done analytically Soper & Spannowsky '11 Most "physically interesting"
- Using jet shapes. E.g. subjettiness: break a jet into subjets 1, 2, ... N

$$S_N = \frac{1}{p_t} \sum_i p_{ti} \min(\delta R_{i1}, \dots \delta R_{iN})$$

J-H Kim '10; Thaler & Van Tilburg '10

Using boosted decision trees
 Cui, Han & Schwartz '10; seems powerful



Biggest improvements are to be had at moderate signal efficiencies

Conclusion from Boost 2010 comparison study of top taggers The method to be adopted depends on the signal efficiency you want

Pileup

high $p_t \rightarrow$ requires high lumi \rightarrow high pileup

28/03/2011	LHC 8:30 meeting		
2011 Records			3 5 TeV
Items in red are records set in the past week			0.0 10 1
Peak Stable Luminosity Delivered	2.49x10 ³²	Fill 1645	11/03/22, 17:12
Maximum Peak Events per Bunch Crossing	13.08	Fill 1644	11/03/22, 02:20
Maximum Average Events per Bunch Crossing	8.93	Fill 1644	11/03/22, 02:20

 \gtrsim 10 events per bunch crossing $\mathcal{O}(10 \text{ GeV})$ of extra p_t per jet, with large fluctuations

$$p_{t,jet}^{\text{subtracted}} = p_{t,jet} - \rho \times A_{jet}$$

Cacciari, GPS & Soyez '08

$$egin{aligned} & A_{jet} = {
m jet} ext{ area} \ &
ho = eta_t ext{ per unit area from pileup} \ & ({
m or ``background''}) \end{aligned}$$

This procedure is intended to be common to pp ($\rho \sim 1-2$ GeV), pp with pileup ($\rho \sim 2-15$ GeV) and Heavy-Ion collisions ($\rho \sim 100-300$ GeV)

As proposed so far: jet-by-jet area determination, event-by-event ρ determination

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Compare FastJet median ρ to Monte Carlo truth (ρ_{Direct})



Comparing pileup estimation methods



A non-trivial issue: rapidity dependence

The original method assumed rapidity dependence was small

- \blacktriangleright In some sense it is, $\lesssim 1.5~{
 m GeV}$
- Measure ρ globally, and include a rapidity-dependent rescaling

 $p_t^{sub} = p_t - f(y)\rho A$

determine f(y) from min-bias • Measure ρ "locally" in strips of $|\Delta y| < 1.5$



Conclusion: global ρ determination with fixed rapidity-dependent rescaling is probably the most effective choice

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Conclusion: global ρ determination with fixed rapidity-dependent rescaling is probably the most effective choice

Jets lecture 2 (Gavin Salam)

A non-trivial issue: rapidity dependence

The original method assumed rapidity dependence was small

- \blacktriangleright In some sense it is, $\lesssim 1.5~{\rm GeV}$
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But lower number of total jets more biased by hard jets (e.g. $t\bar{t}$)



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Hints from charged tracks

Dispersion of offset gives another measure of the subtraction "quality"

- several GeV without subtraction
- only partially reduced with FJ subtraction
- alternative: use PF to remove PU charged tracks in each jet if PU is in-time
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Direct knowledge of PU from tracks can be beneficial

Detector impact harder to judge

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Jet masses etc.?

Fat-jet studies need more than just the jet p_t . E.g. jet mass

There are methods to limit PU sensitivity of jet masses.

> Filtering: Butterworth et al '08 Pruning: Ellis et al '09 Trimming: Thaler et al '09

4-vector subtraction can also help

 $p_{\mu}^{(sub)} = p_{\mu} - f(y)
ho A_{\mu}$

"Automatically" corrects mass as long as hadron masses set to zero

> Many more things can be corrected for PU beyond jet p_t Tests are still in v. early stages / drawing board

 $[1 \text{ jet} \gtrsim 2 \text{ partons}]$ [Pileup]

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