

2 jets +missing ET
signature and spin of
new particles

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with Jing Shu arXive 1102.0293

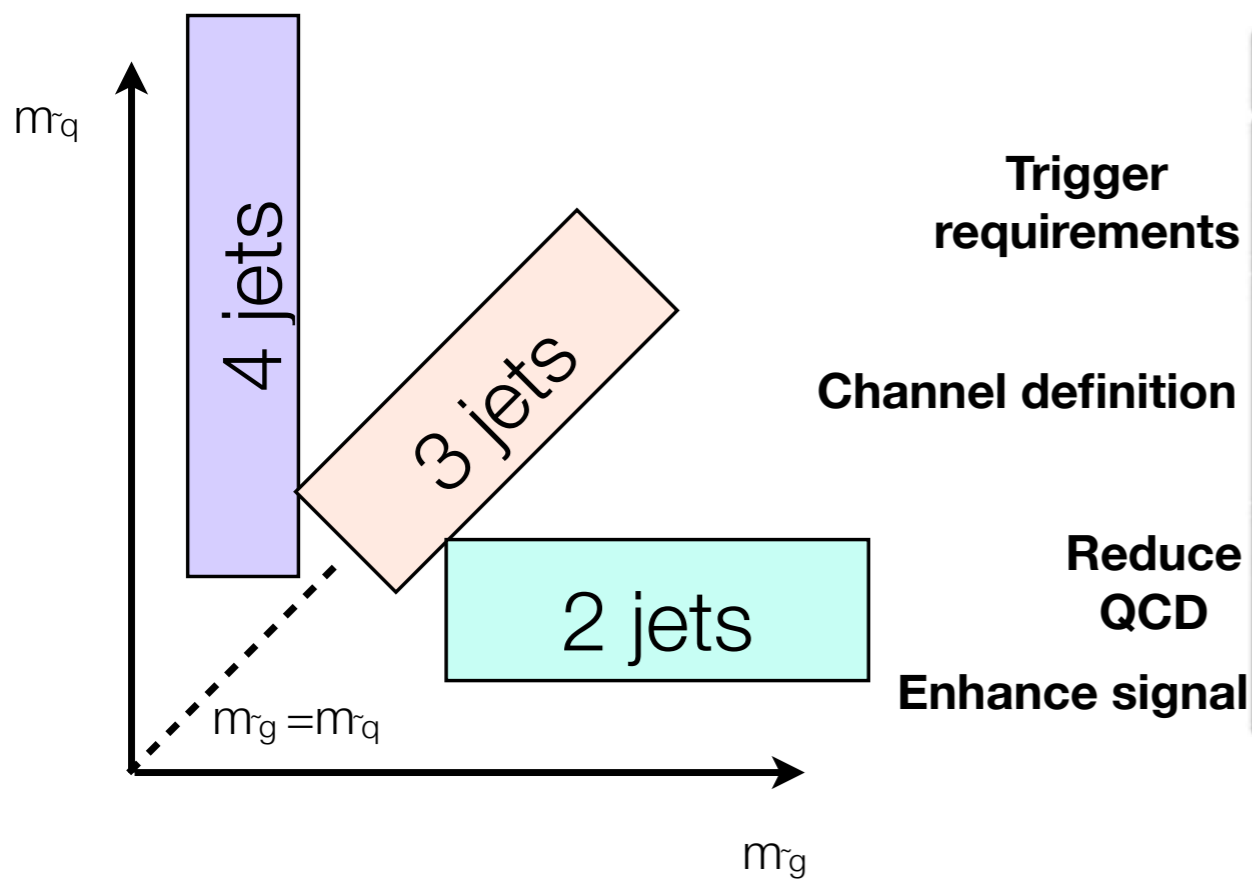
Why Study SUSY models

- Beautiful Symmetry
- Gauge coupling Unification
- New particles that can be searched for at LHC
- Dark matter candidate with R parity
- Can be consistent with low energy measurements.
- Signature missing energy(dark matter) with lots of jets and leptons in the final state.

SUSY search and
measurement
Now and future

Event selection

- Depending on the SUSY mass hierarchy, **different production processes favoured** ($\tilde{g}\tilde{g}$, $\tilde{g}\tilde{q}$, $\tilde{q}\tilde{q}$)
- Signal regions optimised to **maximise sensitivity** to different production processes



Trigger requirements

Channel definition

Reduce QCD

Enhance signal

Signal Region	≥ 2 jets	≥ 3 jets	≥ 4 jets	High mass
E_T^{miss}	> 130	> 130	> 130	> 130
Leading jet p_T	> 130	> 130	> 130	> 130
Second jet p_T	> 40	> 40	> 40	> 80
Third jet p_T	–	> 40	> 40	> 80
Fourth jet p_T	–	–	> 40	> 80
$\Delta\phi(\text{jet}, E_T^{\text{miss}})_{\text{min}}$	> 0.4	> 0.4	> 0.4	> 0.4
$E_T^{\text{miss}}/m_{\text{eff}}$	> 0.3	> 0.25	> 0.25	> 0.2
m_{eff} [GeV]	> 1000	> 1000	$> 500/1000$	> 1100

$$m_{\text{eff}} = \sum_{i=1}^n |\vec{p}_T^{\text{jet } i}| + E_T^{\text{miss}}$$

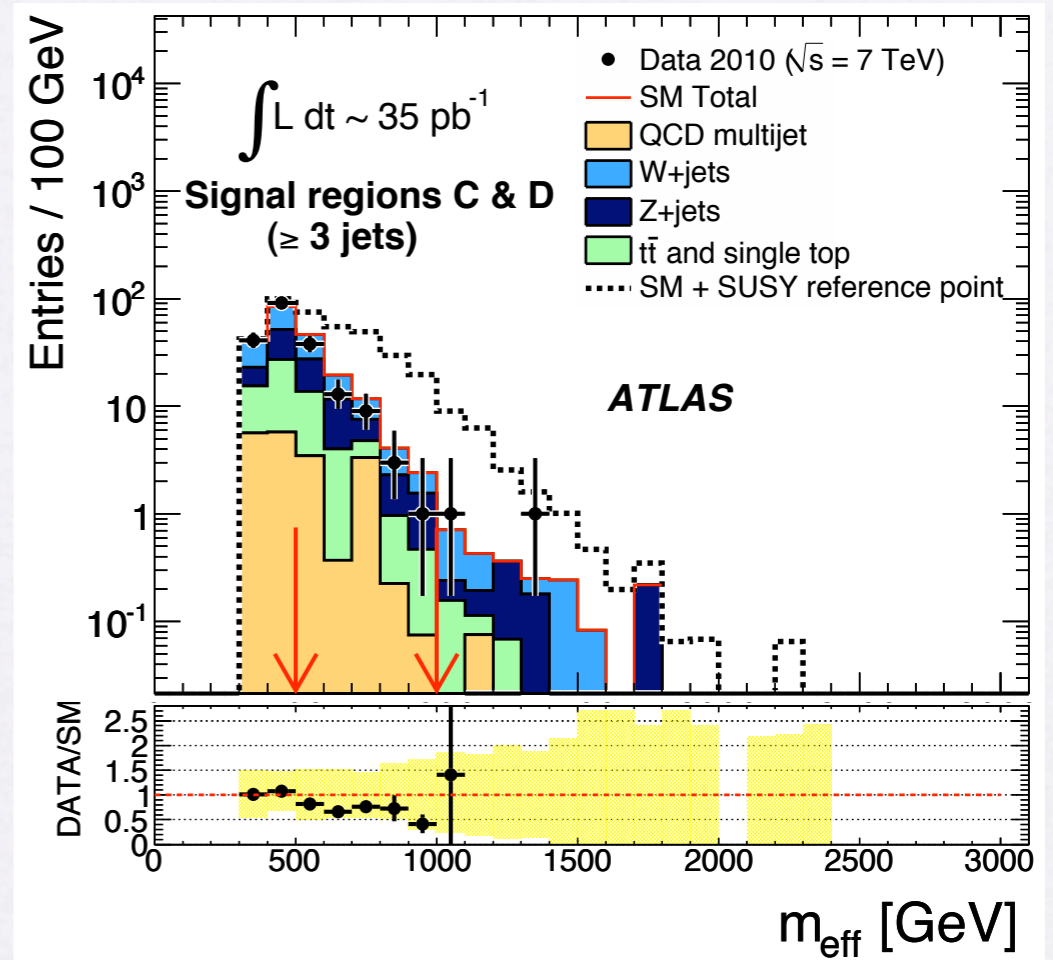
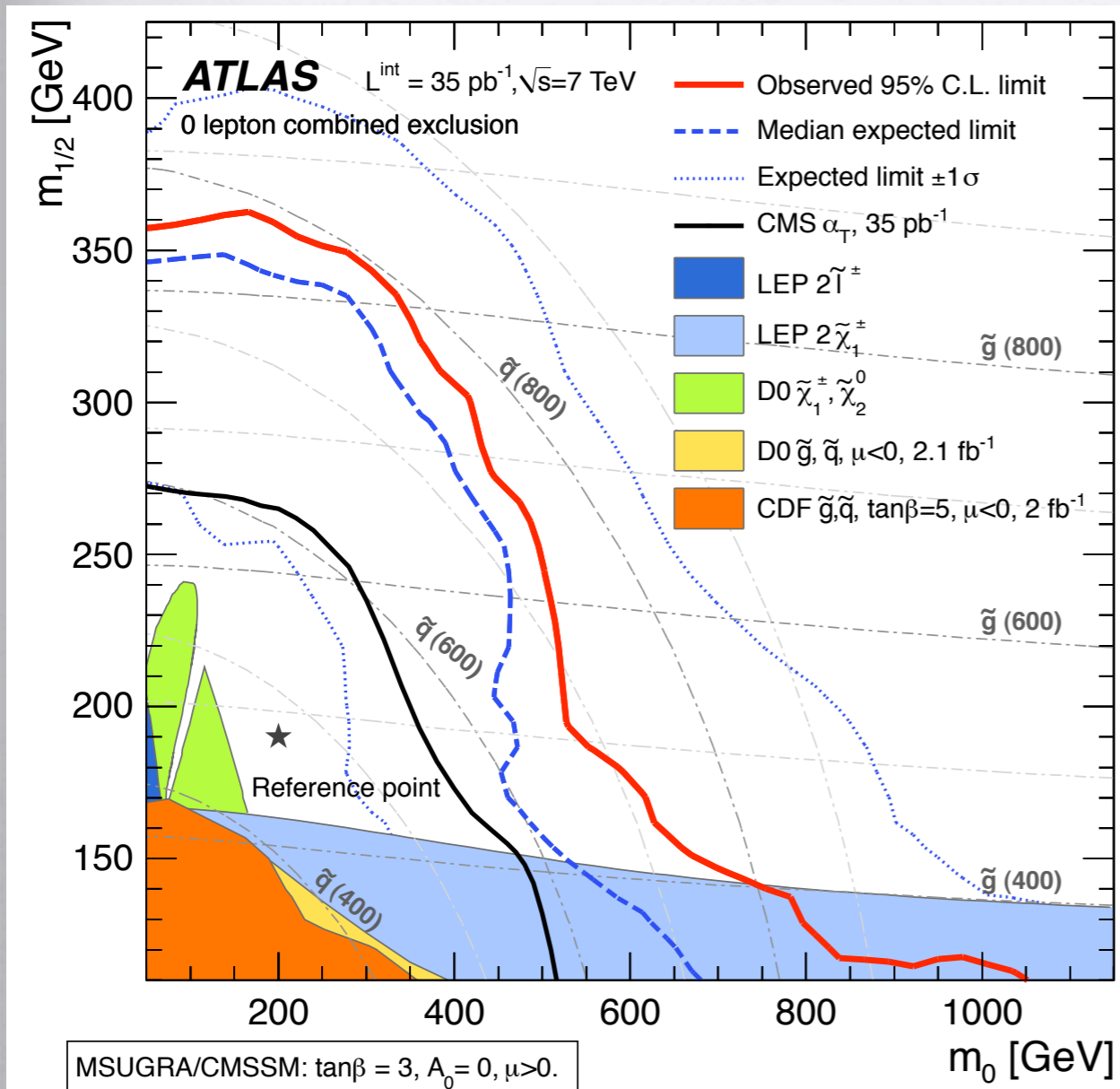
In my view, this is THE BEST way to presenting data

Results

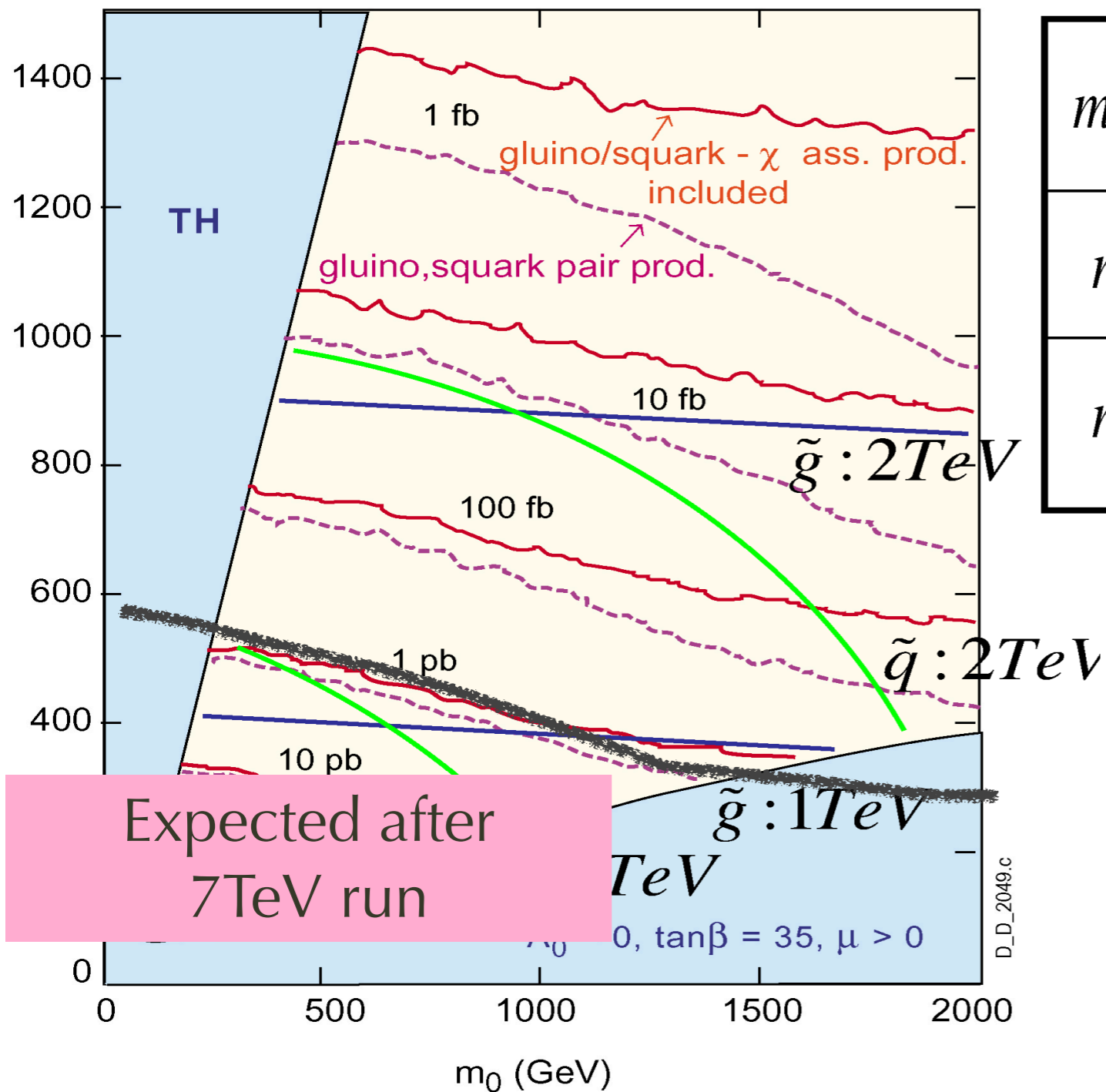
Process	Signal Region				
	≥ 2 -jet	≥ 3 -jet	≥ 4 -jet, $m_{\text{eff}} > 500$ GeV	≥ 4 -jet, $m_{\text{eff}} > 1000$ GeV	High mass
Z/ γ +jets	32.5 \pm 2.6 \pm 6.8	25.8 \pm 2.6 \pm 4.9	208 \pm 9 \pm 37	16.2 \pm 2.1 \pm 3.6	3.3 \pm 1.0 \pm 1.3
W+jets	26.2 \pm 3.9 \pm 6.7	22.7 \pm 3.5 \pm 5.8	367 \pm 30 \pm 126	12.7 \pm 2.1 \pm 4.7	2.2 \pm 0.9 \pm 1.2
$t\bar{t}$ + single top	3.4 \pm 1.5 \pm 1.6	5.6 \pm 2.0 \pm 2.2	375 \pm 37 \pm 74	3.7 \pm 1.2 \pm 2.0	5.6 \pm 1.7 \pm 2.1
QCD jets	0.22 \pm 0.06 \pm 0.24	0.92 \pm 0.12 \pm 0.46	34 \pm 2 \pm 29	0.74 \pm 0.14 \pm 0.51	2.10 \pm 0.37 \pm 0.83
Total	62.3 \pm 4.3 \pm 9.2	55 \pm 3.8 \pm 7.3	984 \pm 39 \pm 145	33.4 \pm 2.9 \pm 6.3	13.2 \pm 1.9 \pm 2.6
Data	58	59	1118	40	18
excluded σ_{acc} (fb)	24	30	477	32	17

- **No discrepancy** with respect to SM predictions. upper limit of each search channel
- The result is interpreted as a **95% CL exclusion limit** on effective cross sections using a profile likelihood ratio approach following the CLs prescriptions.
- Analysis giving best expected limit used in each point.

LHC SUSY search



14TeV projection



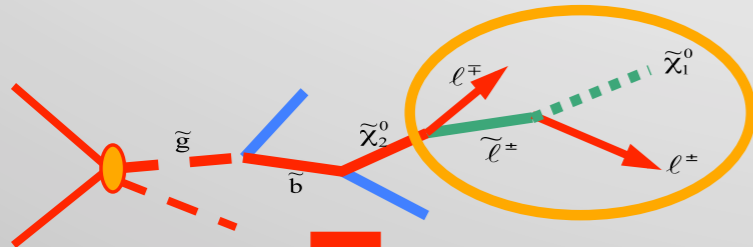
$m(\tilde{q}) = m(\tilde{g}) = 0.5TeV$	$\sigma \sim 100pb$ $\tilde{g}\tilde{g}$ が main
$m(\tilde{q}) = m(\tilde{g}) = 1TeV$	$\sigma \sim 3pb$
$m(\tilde{q}) = m(\tilde{g}) = 2TeV$	$\sigma \sim 20fb$ $\tilde{u}\tilde{u}, \tilde{u}\tilde{d}$ が main

- 7TeV run excluded significant parameter space
- production at 14TeV would be 1 pb or less. significantly limits statistics at 14TeV run already.

Sparticle Detection & Reconstruction

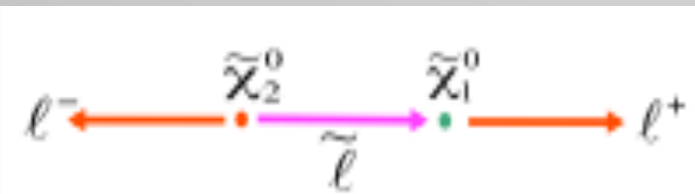
Mass precision for a favorable benchmark point at the LHC
LCC1 ~ SPS1a ~ point B'

$m_0 = 100$ GeV
 $m_{1/2} = 250$ GeV
 $A_0 = -100$
 $\tan\beta = 10$
 $\text{sign}(\mu) = +$

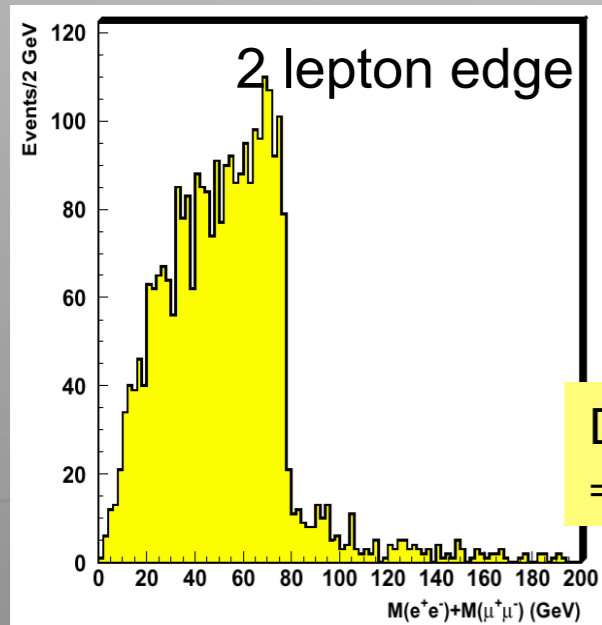


hep-ph/0508198
100 fb⁻¹, 14 TeV

Lightest neutralino → Dark Matter?
Fit SUSY model parameters to the measured SUSY particle masses to extract $\Omega_\chi h^2 \Rightarrow O(10\%)$ for LCC1

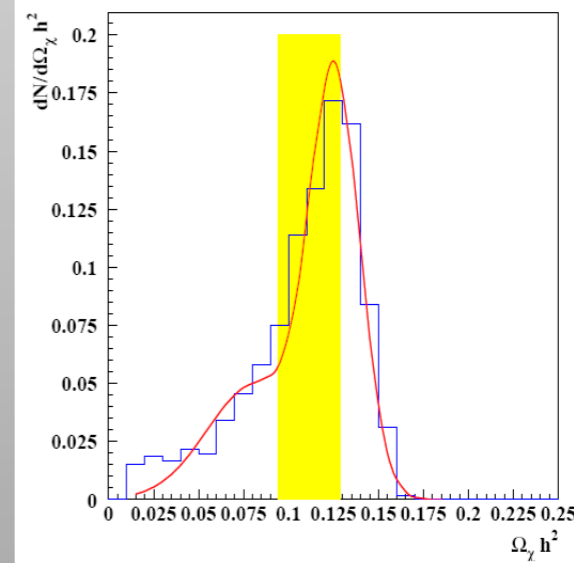


GeV	LHC
$\Delta m_{\tilde{\chi}_1^0}$	4.8
$\Delta m_{\tilde{\chi}_2^0}$	4.7
$\Delta m_{\tilde{\chi}_4^0}$	5.1
$\Delta m_{\tilde{l}_R}$	4.8
$\Delta m_{\tilde{l}_L}$	5.0
Δm_{τ_1}	5-8
$\Delta m_{\tilde{q}_L}$	8.7
$\Delta m_{\tilde{q}_R}$	7-12
$\Delta m_{\tilde{b}_1}$	7.5
$\Delta m_{\tilde{b}_2}$	7.9
$\Delta m_{\tilde{g}}$	8.0



D. Miller et al
⇒ Use shapes

25



This point and much more of the CMSSM space is ruled out
What can LHC still say on DM?

SUSY mass determination using jets+ 2 lepton channel

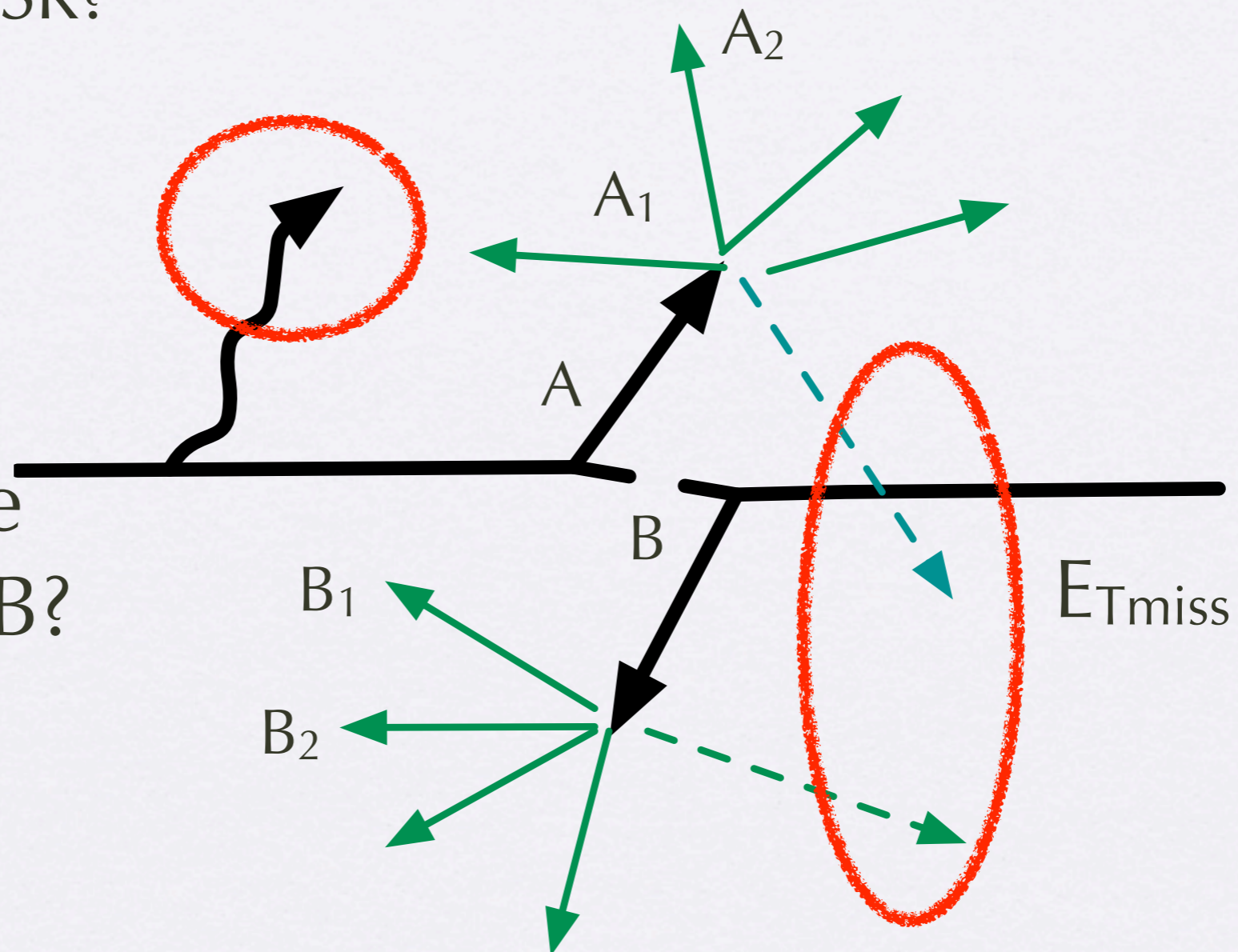
- production cross section is determined by squark gluino mass
- Branching ratio into the second lightest neutralino 30% , lepton branch 6~20% → total 2~6%.
- $30\text{fb}^{-1} \times 1\text{pb} = 30000 \rightarrow$ **600 events(2% branch) are not enough** to determine EW SUSY particles masses precisely
- **Need full use of hadronic channels to determine SUSY scale when it is discovered.**

Combinatorial background in hadronic channel

2) ISR; Which jet comes from ISR?

USE MT_2 ~ mass of parents

1) jets are from A or B?



Combinatorial background in hadronic channel

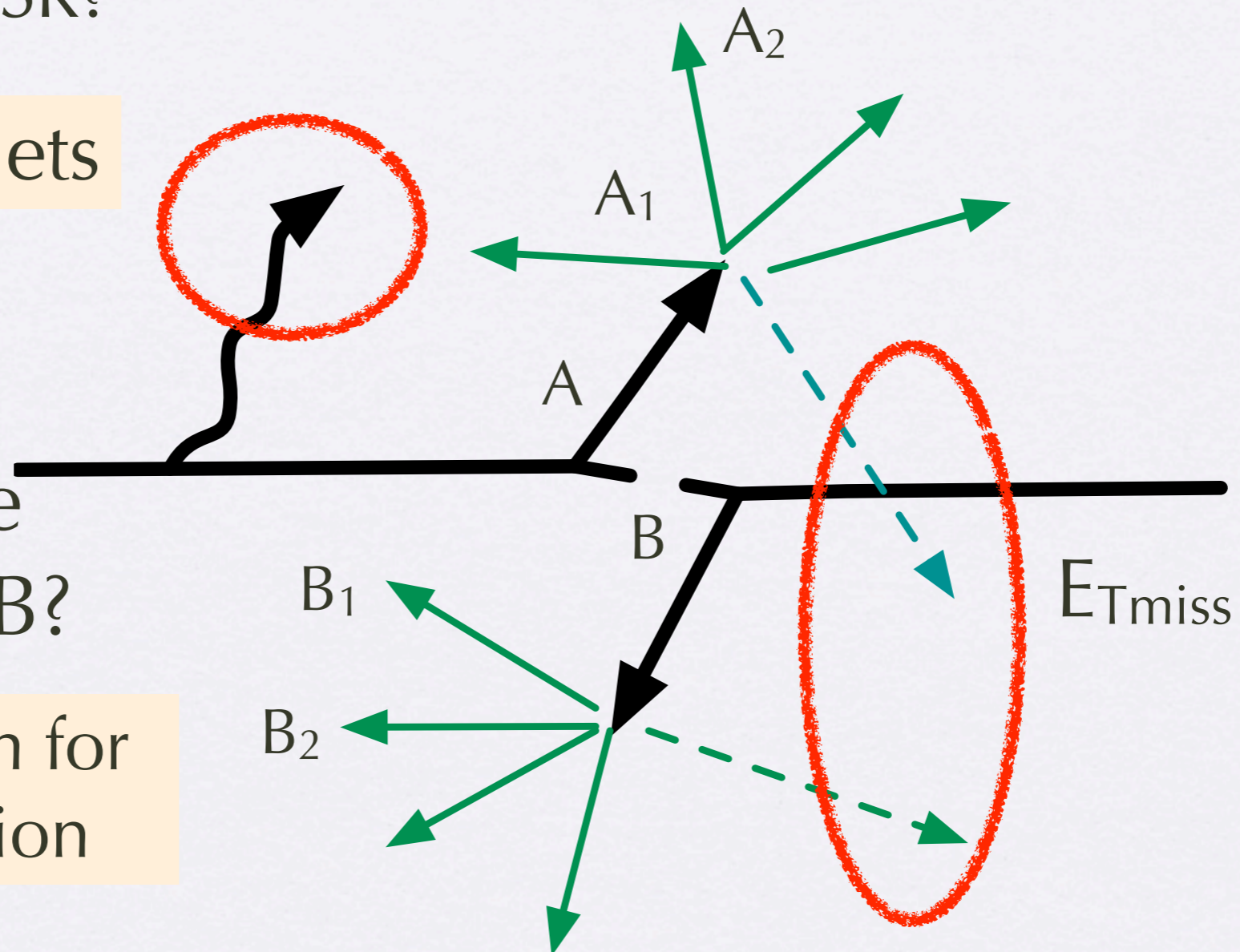
2) ISR; Which jet comes from ISR?

USE MT_2 ~ mass of parents

use part of the jets

1) jets are from A or B?

Take minimum for jet combination

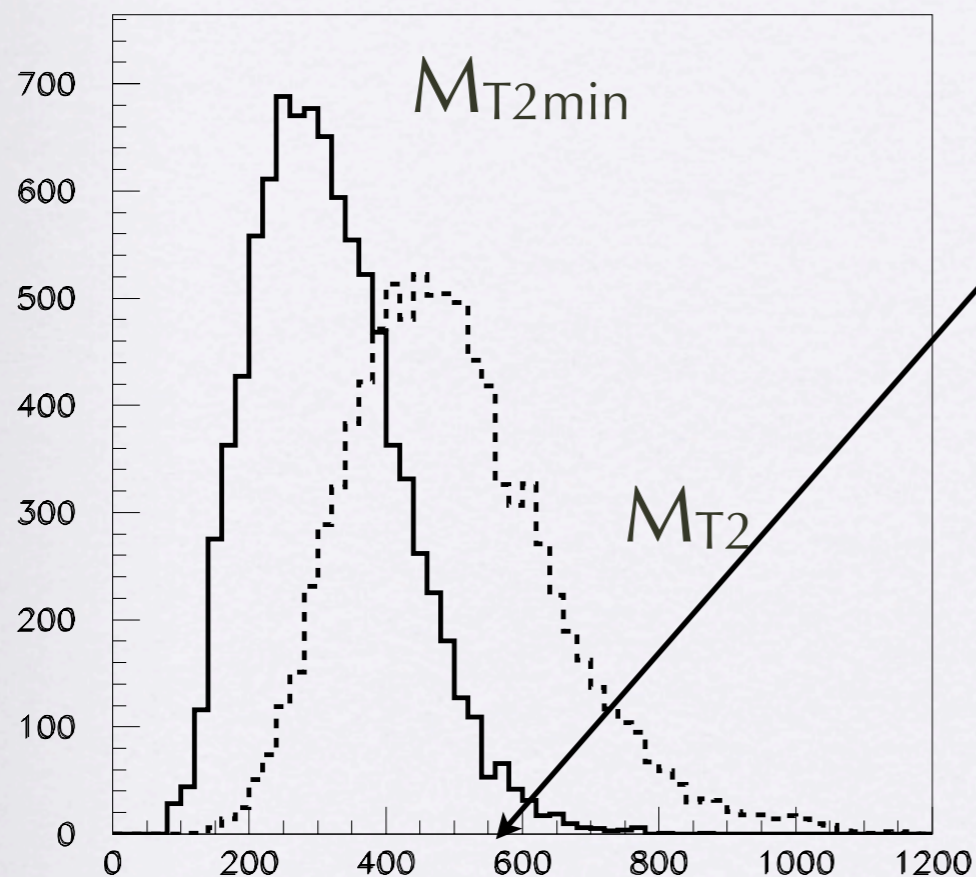


MT2 and mass reconstruction

$$m_{T2}(\mathbf{p}_T^{vis(1)}, m_{vis}^{(1)}, \mathbf{p}_T^{vis(2)}, m_{vis}^{(2)}, m_\chi) \equiv \min_{\{\mathbf{p}_T^{\chi(1)} + \mathbf{p}_T^{\chi(2)} = -\mathbf{p}_T^{vis(1)} - \mathbf{p}_T^{vis(2)}\}} \left[\max\{m_T^{(1)}, m_T^{(2)}\} \right],$$

7TeV 100fb⁻¹

m_{gl}=558GeV m_{ul}=825 GeV



input gluino mass

M_{T2} for multijet final state = minimization for all jet combination

M_{T2min} =ISR removal ~remove one jet from the minimization (among 5 leading jets)

Nojiri Sakurai 2010

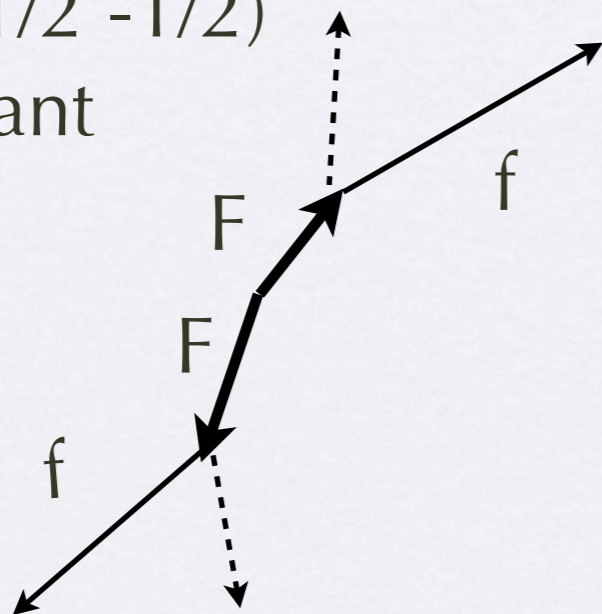
Reconstruction of (squark /gluino mass -LSP mass) may be possible

How about spin measurements

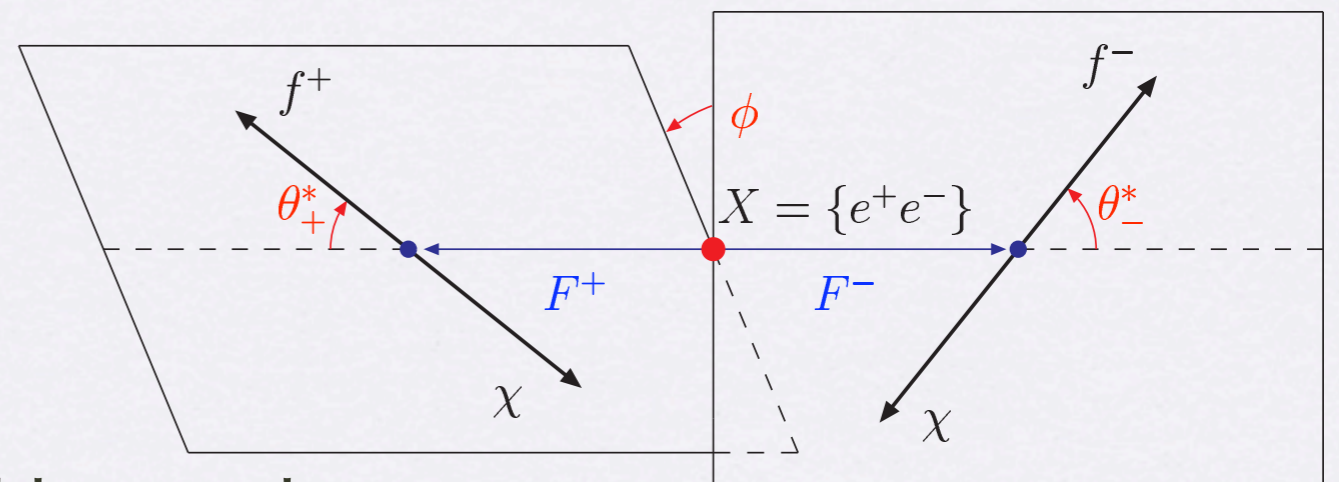
- in Jet +2 lepton channel spin effect in the inv. mass distribution, able to distinguish SUSY vs "Same spin partner" models (such as LHT, UED)
- jet channel: there are jet ID problem, but jets from two body decay of quark partner is easy to identify because of the PT
- If the interaction of quark partner is chiral, there are visible spin effect

polarization

$(1/2, 1/2)$ **or** $(-1/2, -1/2)$
is dominant

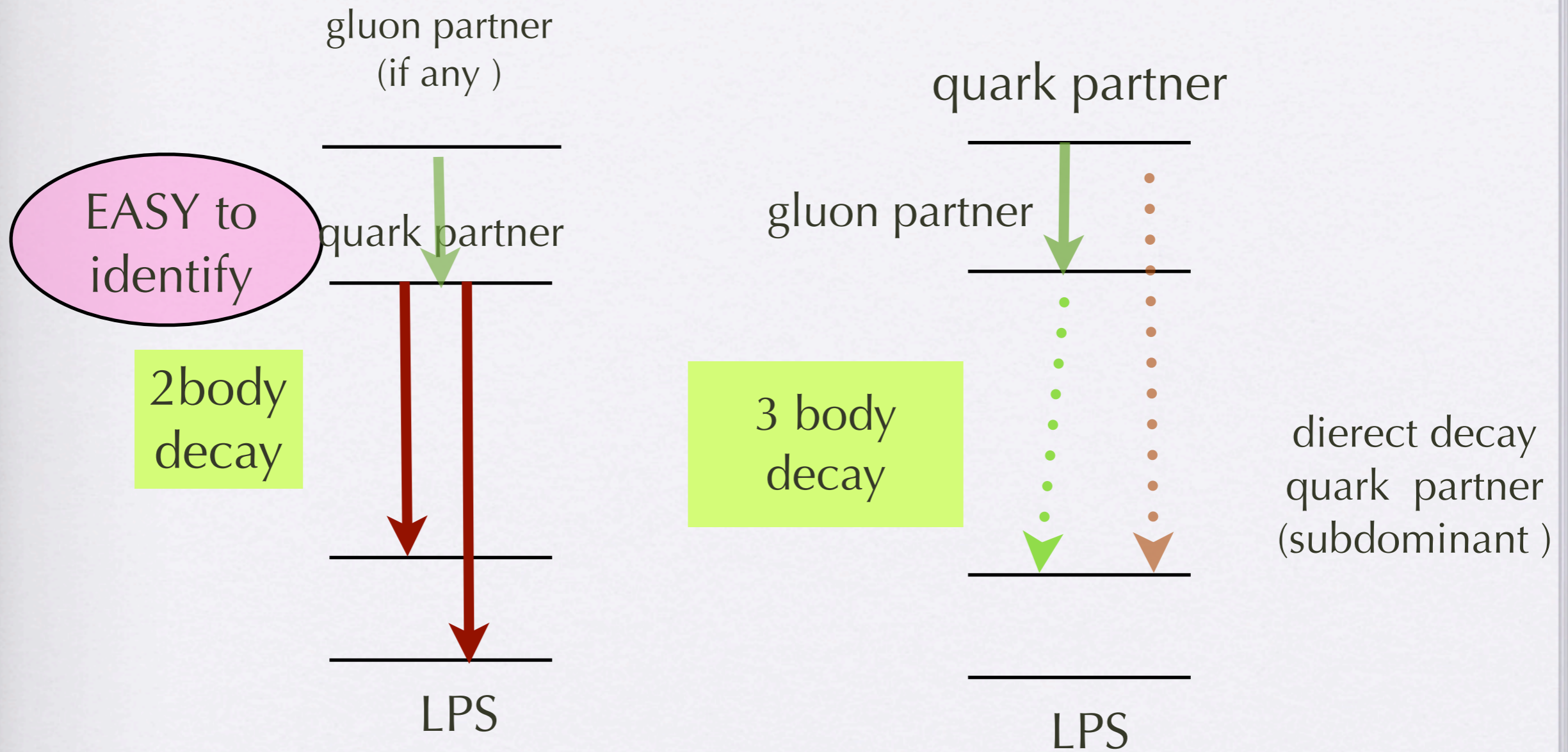


Azimuthal angle correlation
 $(1/2, 1/2)$ **and** $(-1/2, -1/2)$
amplitude is same order



Buckley et al 2008

leading objects and new particle decays



SUSY quark partner=spin 0
 UED/LHT=Spin 1/2

Interactions of the same spin partner model

$$L_{int} = \int dx^4 [g_s G_H^{\mu a} \bar{Q} T^a \gamma_\mu q + g W_H^{\mu a} \bar{Q}_L T^a \gamma_\mu P_L q_L + g' B_H^\mu (Y_L \bar{Q}_L \gamma_\mu P_L q_L + Y_L \bar{Q}_R \gamma_\mu P_R q_R) + (\text{Lepton part}) + h.c.]$$

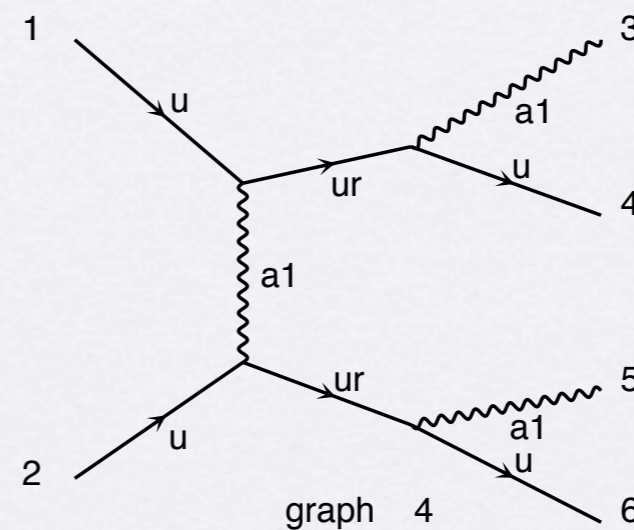
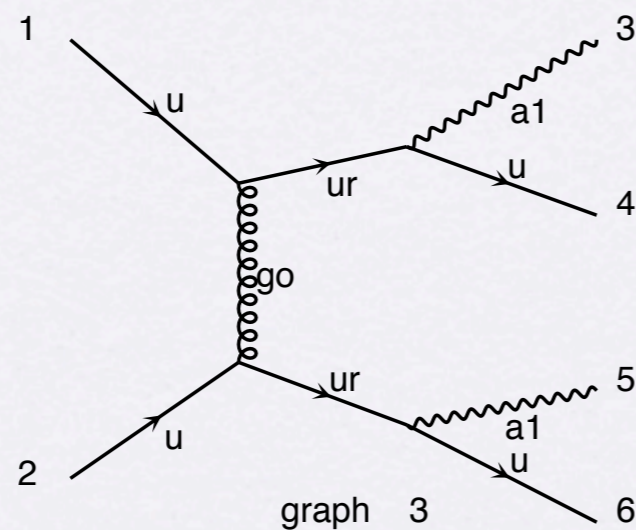
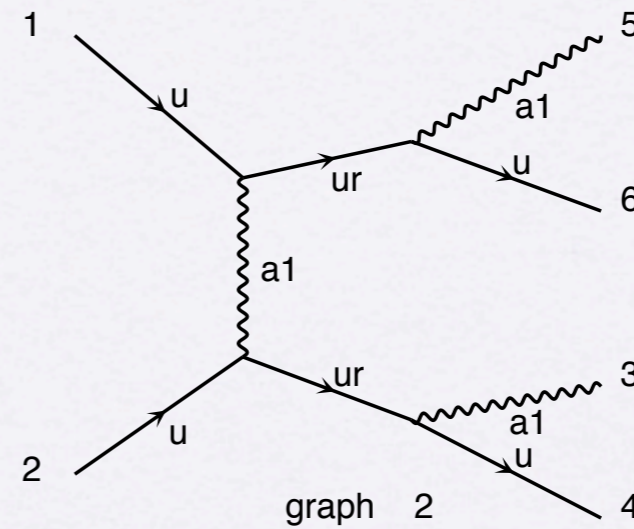
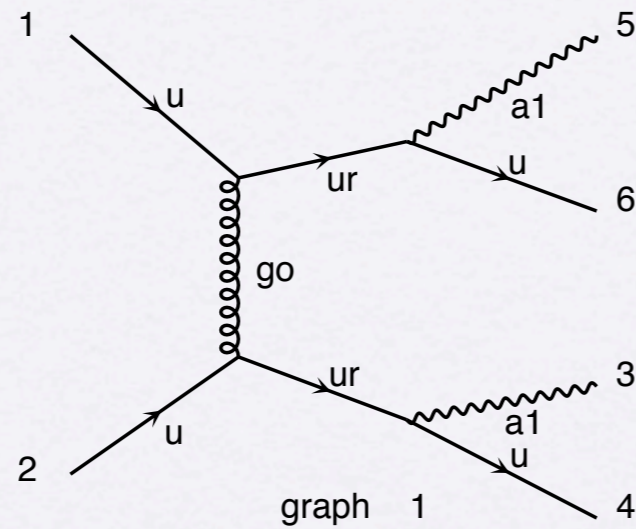
- Heavy gauge boson (spin 1) and quark partners (spin 1/2) have $Z_2 = -1$
- Haven't specify Higgs sector: Gauge invariance of the amplitude must be carefully checked.
 - in UED like model: 5th component of gauge boson is the goldstone boson.
 - split UED/ a three site model allow mass splitting of the partners
 - big difference in the distributions

Production process at LHC and decay

$pp \rightarrow UR UR \rightarrow u u BH BH$

Diagrams by MadGraph

$u u \rightarrow a1 u a1 u$



Amplitude Tips

- Production and chiral interaction

$$\begin{aligned}
 i\mathcal{M}(uu \rightarrow U_R U_R) &= \frac{ig'^2 Y_u^2 \delta_{aa'} \delta_{bb'} \left(-g^{\mu\nu} + \frac{q^\mu q^\nu}{m_{BH}^2} \right)}{q^2 - m_B^2} \bar{u}_{h_f}(p_f) \gamma_\mu P_R u_{h_i}(p_i) \bar{u}_{h'_f}(p'_f) \gamma_\nu P_R u_{h'_i}(p'_i) \\
 &\quad + \text{cross diagram} + \text{gluon exchange contribution} \tag{2}
 \end{aligned}$$

$\frac{-ig'^2 \delta_{aa'} \delta_{bb'} m_Q^2}{(q^2 - m_B^2) m_B^2} \bar{u}(p'_f) P_R u(p_i) \bar{u}(p'_f) P_R u(p'_i)$

helicity conserving in the $\beta \rightarrow \infty$ limit
 helicity (1/2, 1/2)

chirality flip
 helicity (-1/2, -1/2)

Decay and polarization

- Polarized particle decay non-spherically

$$iM \propto \epsilon^{*\mu} \epsilon^\nu \text{Tr} \left[\gamma_\mu \overset{\text{chiral vertex}}{\circlearrowleft} P_R \not{p}_f \gamma_\nu \overset{\text{chiral vertex}}{\circlearrowleft} P_R \frac{1 + \not{h}\gamma_5}{2} (\not{p}_i + m_Q) \frac{1 + \not{h}\gamma_5}{2} \right] = \frac{2k_B \cdot p_f m_Q}{m_B^2} (E_B - k_{B//})$$

chiral vertex

projection to the helicity state

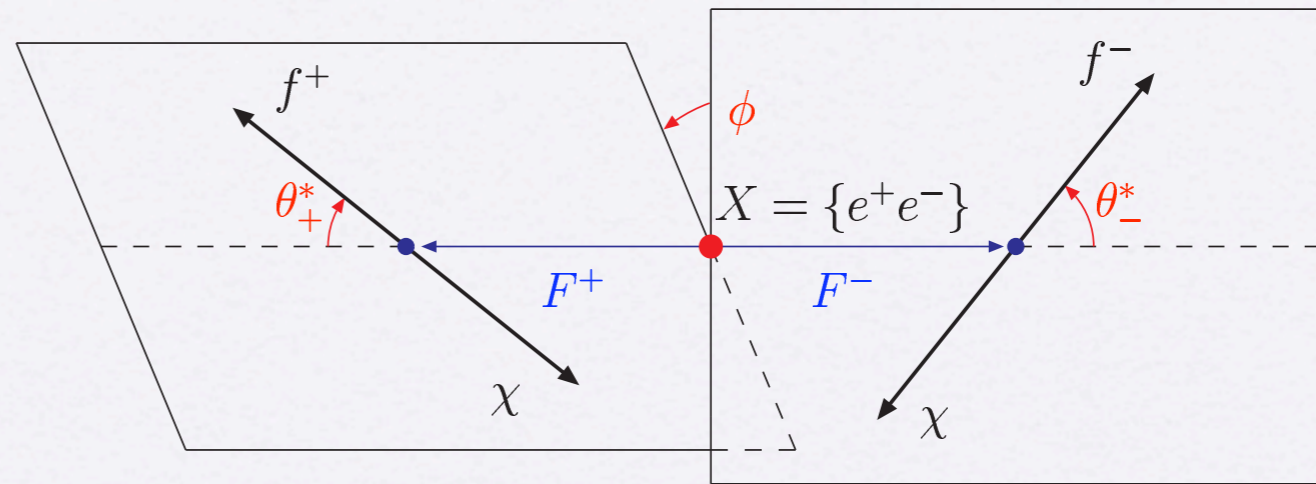
$h=1/2$ particle + chiral vertex \rightarrow quark distribution is $\propto 1 + \cos \theta$

(The distribution is same for Q_L decay into $q_L \chi$)

If number of particle is larger than number of antiparticle(LHC) effect remains.

Note : $h=0$ massive gauge boson dominates in the decay.

Azimuthal Angle Correlations



amplitude till decay

$$W(E_{\text{cm}}; \Theta; \theta_{\pm}^*, \phi_{\pm}^*) = \sum_{\lambda_{\pm}, \lambda'_{\pm} = -j}^j \mathcal{P}_{\lambda'_{-}\lambda'_{+}}^{\lambda_{-}\lambda_{+}}(E_{\text{cm}}, \Theta) \mathcal{D}_{\lambda_{-}\lambda'_{-}}^{-}(\theta_{-}^*, \phi_{-}^*) \mathcal{D}_{\lambda_{+}\lambda'_{+}}^{+}(\theta_{+}^*, \phi_{+}^*),$$

decay matrix

spin dependent phase

$$\mathcal{P}_{\lambda'_{-}\lambda'_{+}}^{\lambda_{-}\lambda_{+}} = \sum_{\sigma_{\pm} = \pm 1/2} \mathcal{T}_{\sigma_{-}\sigma_{+}; \lambda_{-}\lambda_{+}} \mathcal{T}_{\sigma_{-}\sigma_{+}; \lambda'_{-}\lambda'_{+}}^*, \quad \mathcal{D}_{\lambda_{\pm}\lambda'_{\pm}}^{\pm}(\theta_{\pm}^*, \phi_{\pm}^*) = D_{\lambda_{\pm}\lambda'_{\pm}}^{\pm}(\theta_{\pm}^*) e^{\mp i(\lambda_{\pm} - \lambda'_{\pm})\phi_{\pm}^*},$$

If spin=0 no azimuthal angle correlation correlation.
 If a helicity state dominates no azimuthal angle correlation
 → independent information

- Amplitude of spin j particle pair production and decay

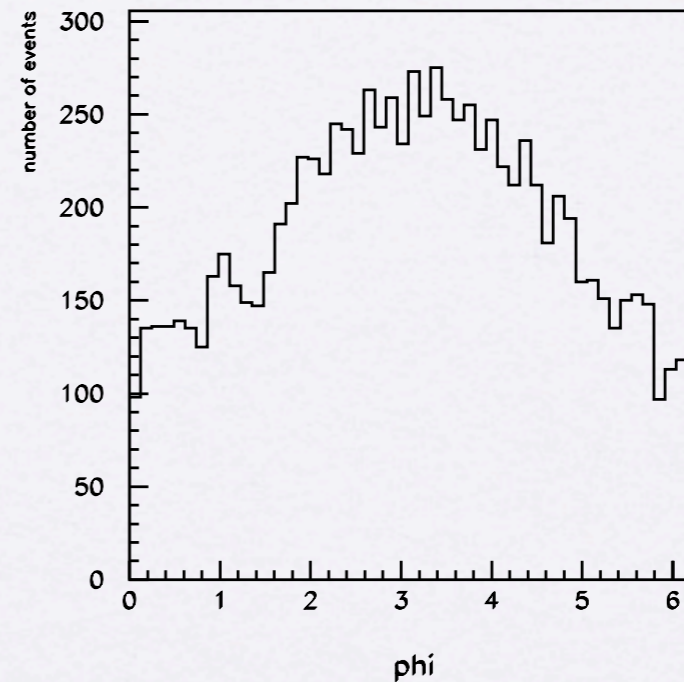
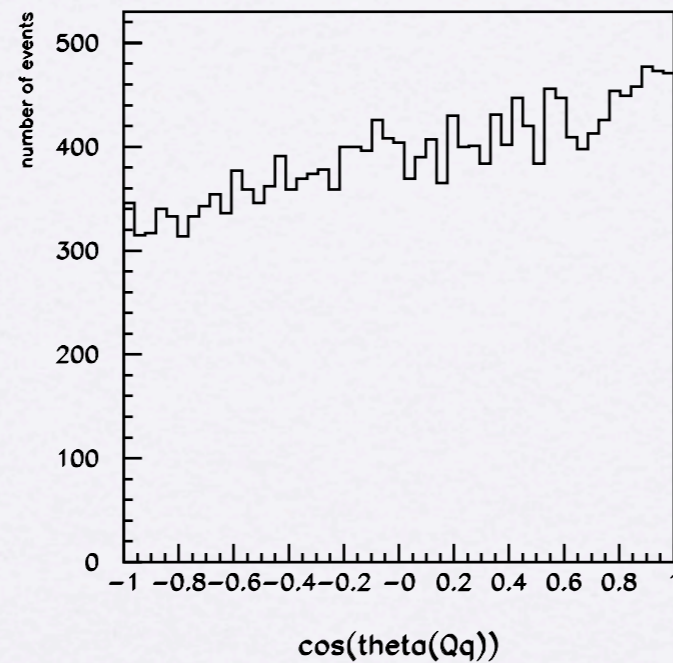
$$\frac{1}{\mathcal{C}} \frac{d\mathcal{C}}{d\phi} = \frac{1}{2\pi} [1 + A_1 \cos(\phi) + \dots + A_{2j} \cos(2j\phi)].$$

parton level distribution

- $m_{BH}=100 \text{ GeV}/200\text{GeV}$ $m_Q=600\text{GeV}$ $m_G=700\text{GeV}$

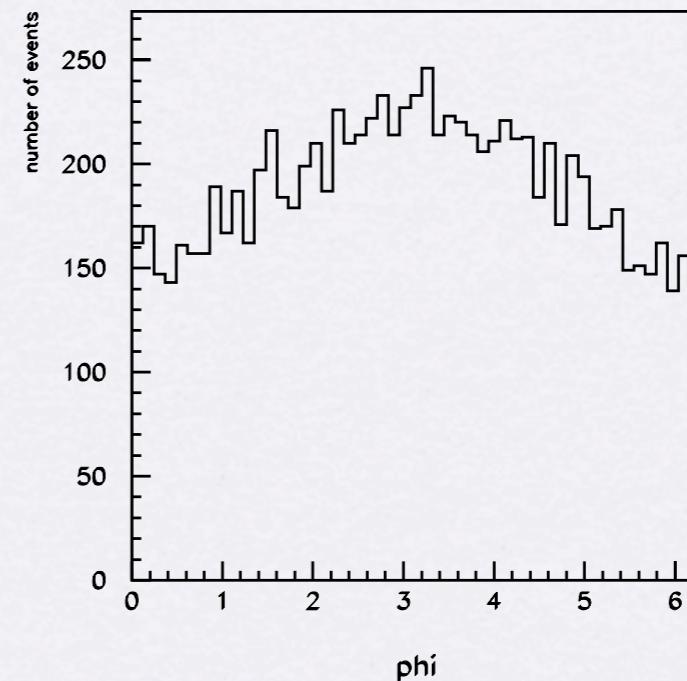
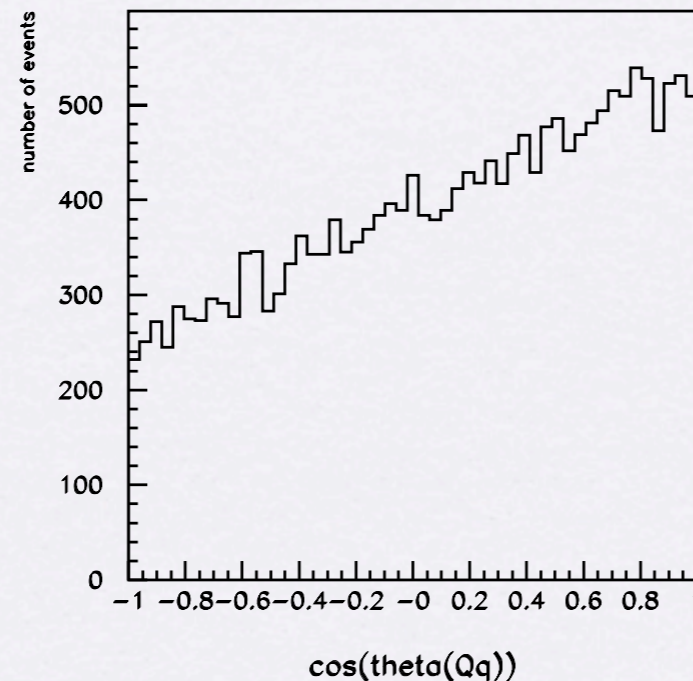
$m_{BH}=100 \text{ GeV}$

both $(1/2, 1/2)$
 $(-1/2, -1/2)$ are
large



$m_{BH}=200 \text{ GeV}$

$(1/2, 1/2)$
dominates



How to see it

schematic figure of distributions

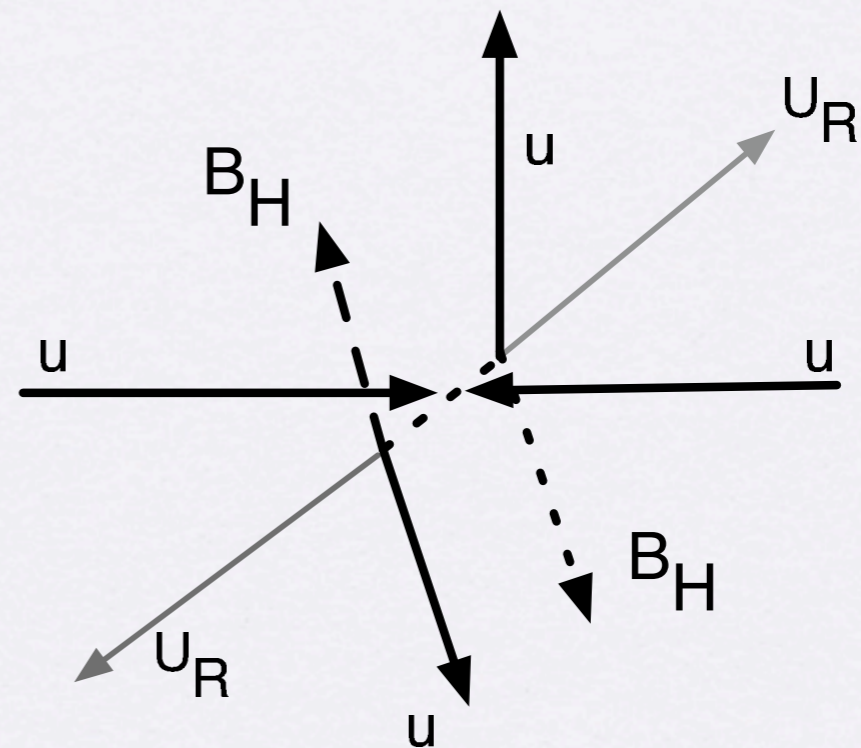
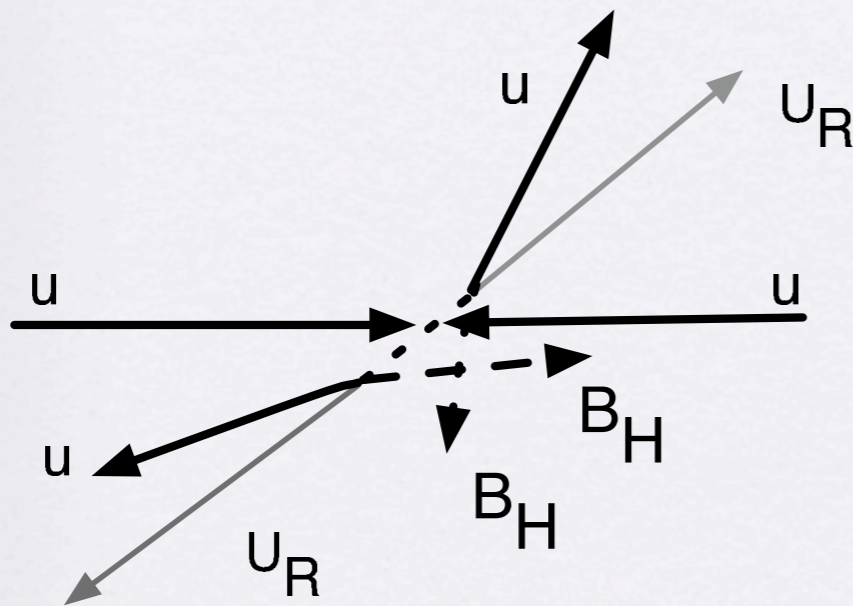
$m_{BH} = 200 \text{ GeV}$

q goes direction of Q

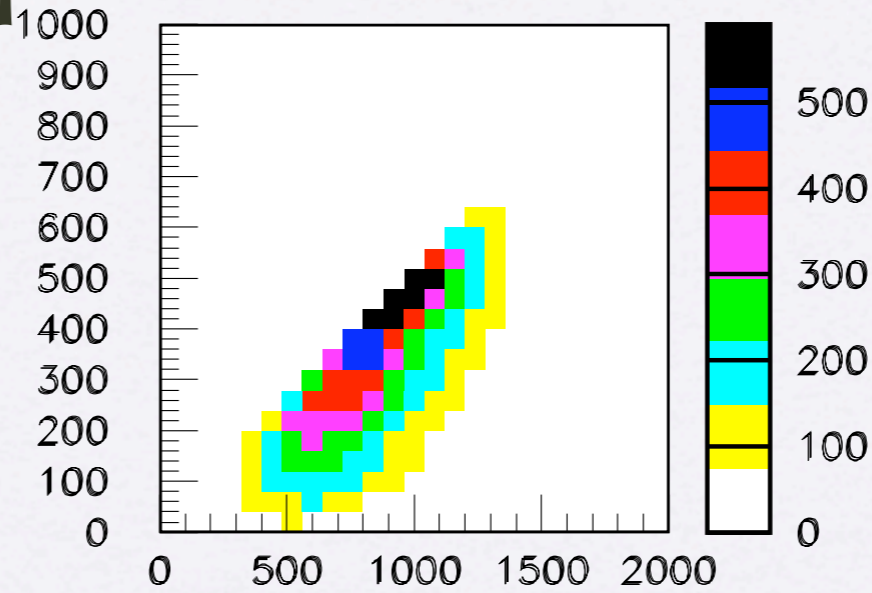
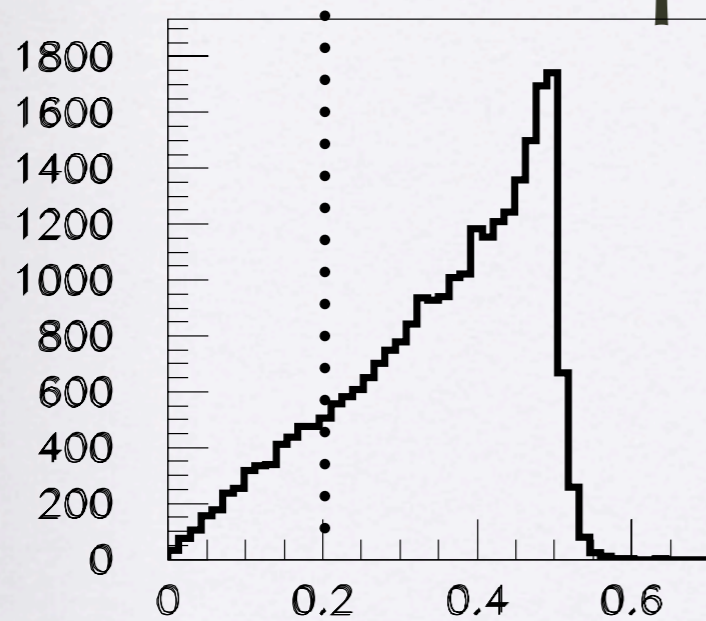
or

$m_{BH} = 100 \text{ GeV}$

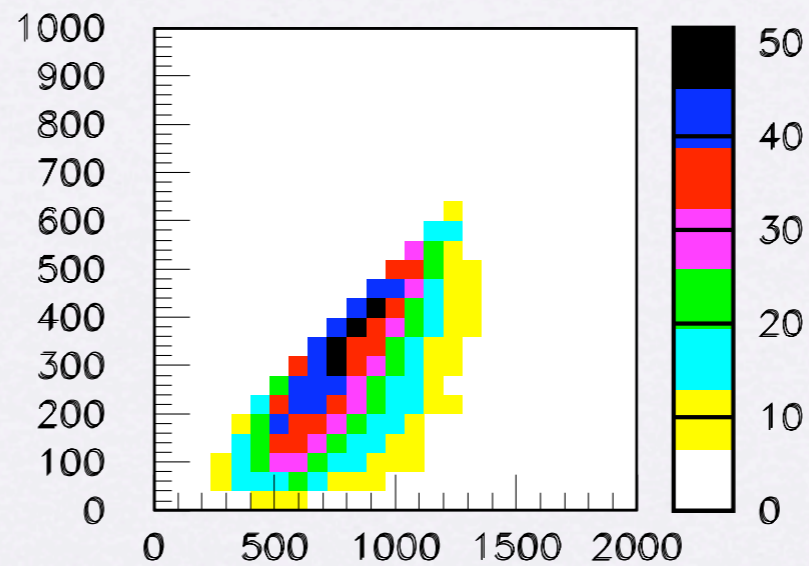
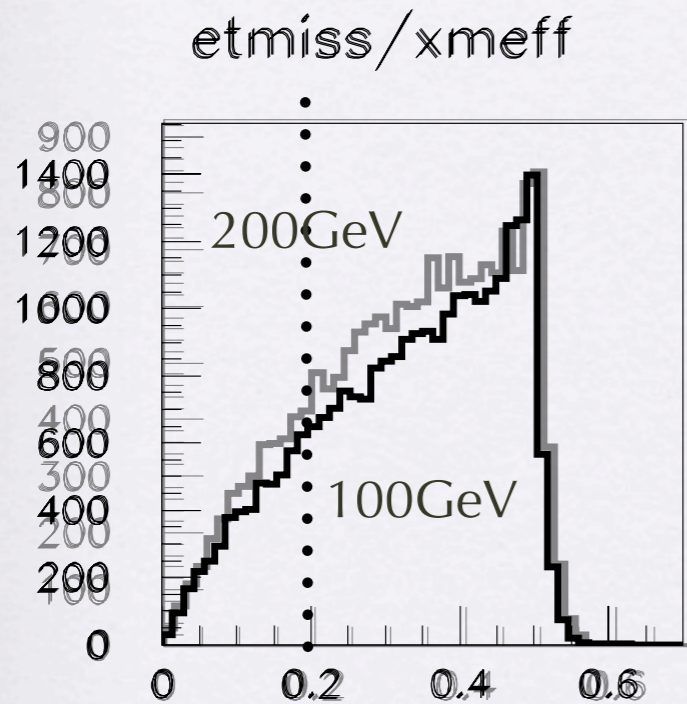
q goes to back to back



The jet level distribution for $pp \rightarrow U_R U_R$ channel



No spin
(Madgraph 2 by 2
→ pythia/bridge)



Madgraph
2 by 4

Nojiri, J. Shu to appear

etmiss/xmeff

M_{T2} and reconstruction of a-angle in jet level

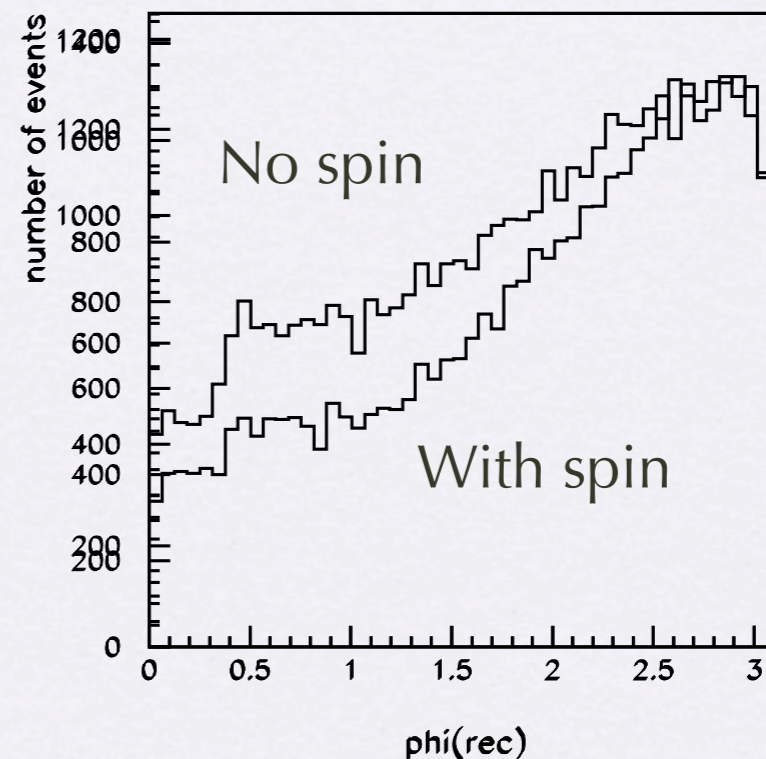
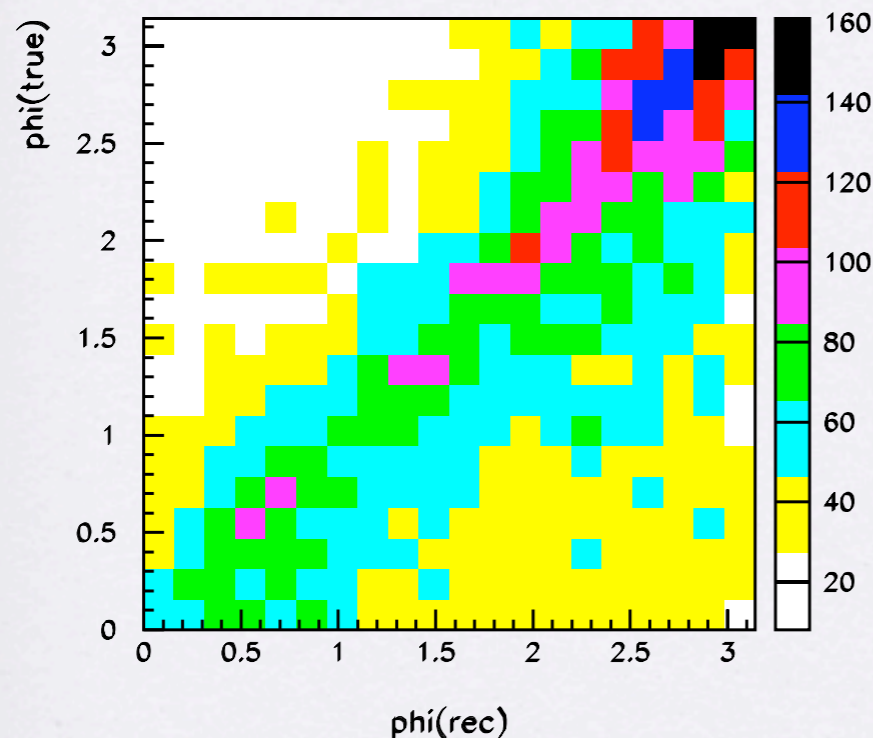
- M_{T2}
$$M_{T2} = \min_{p_1^T + p_2^T = p_{T\text{miss}}} \left[\max \left(m_{T2}(p_{\text{vis}}^{(1)}, p_1^T), m_{T2}(p_{\text{vis}}^{(2)}, p_2^T) \right) \right]$$

- M_{T2} assisted reconstruction: The process give transverse test LSP momentum $p_{(1)}, p_{(2)}$ of that gives M_{T2} . calculate p_z momentum that gives correct m_Q , and m_χ

$$(p_i + p_{\text{vis}}^{(i)})^2 = m_Q^2, \quad (p^{(i)})^2 = m_{BH}^2$$

- calculate φ for the momentum.

Small phi tend to fail

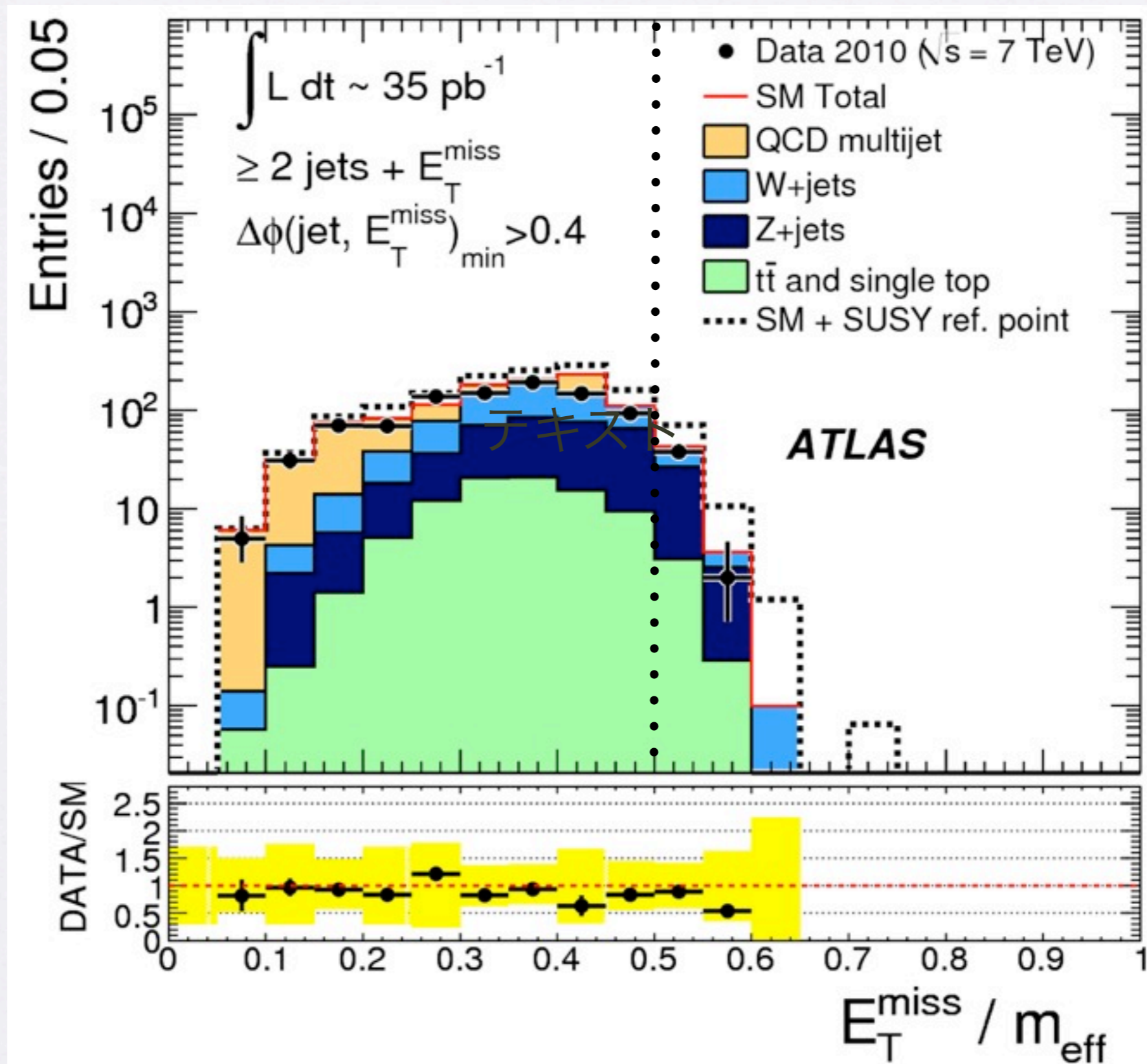


Lessons and some info

- “phase space decay” for leading objects fails to reproduce physics processes even in such simple case. (Some phase space generators are not useful)
 - Moortgat-Pick et al : distribution of production cross section in forward region is also different
- “Consistent treatment”: Production in T channel and decay are correlated.
 - Full amplitude generator (Madgrap and Herwig++) work. No Pythia.
- Working on little higgs model with T parity (no gluon partner, heavy wino like object)

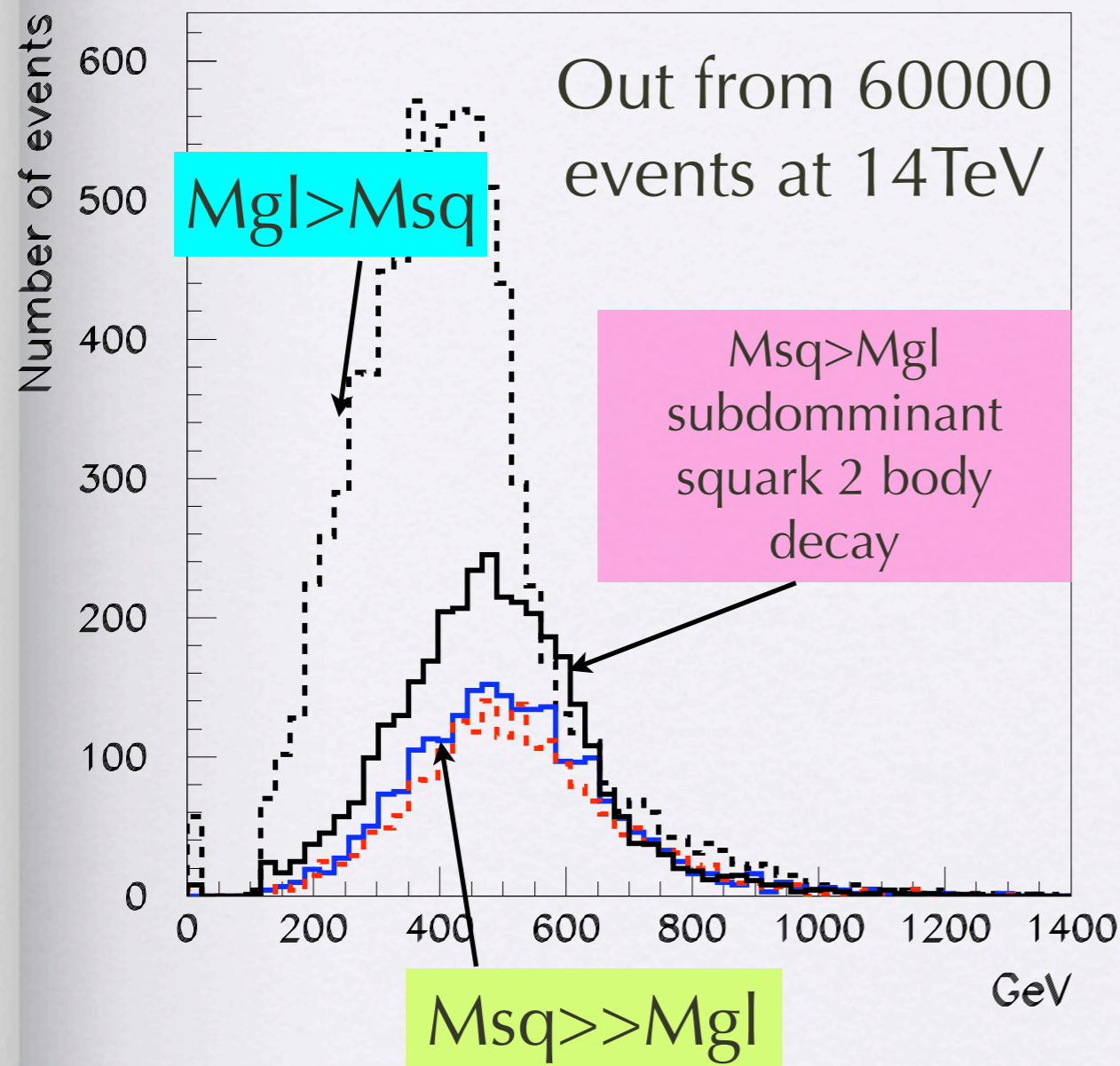
おまけ(ATLAS distribution)

$M_{\text{eff}} =$
 $E_{\text{Tmiss}} + \text{two highest } p_{\text{T}} \text{ jet}$



予備

two high pT jets in SUSY events



Nojiri, Sakurai

- Inclusive MT_2 distribution for $M_{gl} \sim 600$ GeV
 - divide events into two using Lund distance and calculate MT_2 from two visible system
- Selection: Events at least 2 jets with $p_T > 200$ GeV
- $m_{sq} < m_{gl}$: large branch sharp edge. The mode with 2 high p_T jet stands!

under mixed SUSY production

- In SUGRA like mass spectrum ($m_{sq}, m_{gl} \gg m_\chi$) jets from $sq \rightarrow q \chi$ is prominent.

(pT of the 3rd jet) / pT of the 1st jet

