General Introduction into 4th Family Model & Phenomenology

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Outline

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Motivation [1]

- Seen 3; why not 4?
- Enhanced prospects for DEWSB
- [Holdom('86); Bardeen et al ('90); Hung & Isidori('97); Hung & Xiong('11)].....
- Enhanced prospects for baryogenesis
 [Jarlskog & Stora('88); Branco et al('98); Hou('08)]

Motivation [II]

- NFC: built in.....Recall this is a very serious problem for almost all BSMs
- Cannot be just simple SM4: DM link? [See e.g. Volovik('03); Lee, Liu & AS ('11)]
- Readily accounts for B-CP-anomalies [A.S et al ('08; '10); Buras et al ('10)]
- Accounts also for AFB(tt) [Davoudiasl,McElmurry & A S]

May be the best reason

• My license plate: udcstbgz (since '89 NY)

• Used to be OSCILL8 (before '89 in CA)

"SM4": a revisit

- Inspired by classic works on "SSB w/o fundamental scalars" by Cornwall and Norton ['73], studied with Norton and students Carpenter & Siegemund-Broka, role of heavy quarks in DEWSB in [87-90]
- potential of B-physics forSM4 studied extensively with George Hou~86-88

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Editors DAVID B. CLINE • AMARJIT SONI



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Parameters & Constraints on 4G

- In the simplest version (SM4) there are quarks t', b' and leptons L₄, v₄
- KM "theorem" tells us that every linearly independent set of 3 generations has a phase; so there will be 2 new CP-odd phases (total 3).
- The 4X4 CKM-mixing matrix will also have 3 new real parameters (total 6)
- As far as masses go, LEPII told us that
 4 G charged lepton, neutrino, quarks have to be at least around 100 GeV

VCKM4

Bearing in mind that rare B-decays are very sensitive to 4 G effects and the hierarchical structure of the Wolfenstein representation, we extend to 4X4 case
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$$V_{us} = \lambda, \qquad V_{cb} = A\lambda^2, \qquad V_{ub} = A\lambda^3 C e^{-i\delta_{ub}},$$
$$V_{t's} = -Q\lambda^2 e^{i\delta_{t's}}, \quad V_{t'd} = -P\lambda^3 e^{i\delta_{t'd}}, \quad V_{t'b} = -r\lambda, \quad (1)$$

$$\begin{pmatrix} 1 - \frac{\lambda^2}{2} + \mathcal{O}(\lambda^4) & \lambda & A\lambda^3 C e^{-i\delta_{ub}} & P\lambda^3 e^{-i\delta_{t'a}} \\ + Q\lambda^3 e^{-i\delta_{t'a}} + ACr\lambda^4 e^{-i\delta_{ub}} \\ -P \frac{\lambda^5}{2} e^{-i\delta_{t'a}} + \mathcal{O}(\lambda^7) \\ -\lambda + \mathcal{O}(\lambda^5) & 1 - \frac{\lambda^2}{2} + O(\lambda^4) & A\lambda^2 & Q\lambda^2 e^{-i\delta_{t'a}} \\ + A\lambda^3 r - P\lambda^4 e^{-i\delta_{t'a}} \\ -\frac{Q}{2}\lambda^4 e^{-i\delta_{t'a}} + \mathcal{O}(\lambda^6) \\ -Pr\lambda^4 e^{i\delta_{t'a}} & +A\lambda^4 \left(\frac{1}{2} - Ce^{i\delta_{ub}}\right) \\ + \frac{1}{2}AC\lambda^5 e^{i\delta_{ub}} + \mathcal{O}(\lambda^7) & + \mathcal{O}(\lambda^6) \\ -P\lambda^3 e^{i\delta_{t'a}} & -Q\lambda^2 e^{i\delta_{t'a}} & -r\lambda & 1 - \frac{r^2\lambda^2}{2} + \mathcal{O}(\lambda^4) \end{pmatrix}$$

Use K,B, D ...experimental info such as b-> s gamma, b-> s II, Delta Ms(Md), ϵ K, K⁺ -> π + $\nu\nu$[Nandi and A. S, arXiv: 1011.6091] SEE SOUMITRA NANDI's talk

| Parameter | Allowed range | Parameter | Allowed range |
|-----------|---------------------------|-----------------------|---------------|
| λ | 0.2205 ± 0.0018 | $ V_{t'b} $ | < 0.12 |
| С | $0.32 \rightarrow 0.42$ | $ V_{t'd} $ | < 0.05 |
| Α | $0.825 \rightarrow 0.865$ | $ V_{t's} $ | < 0.11 |
| γ | $(73 \pm 13)^{\circ}$ | $ V_{ub'} $ | < 0.05 |
| r | < 0.5 | $ V_{cb'} $ | < 0.11 |
| Р | <5.0 | $\lambda_{db}^{t'}$ | < 0.002 |
| Q | <2.5 | $ \lambda_{sb}^{t'} $ | < 0.01 |
| | | $ \lambda_{uc}^{b'} $ | < 0.0025 |

TABLE III. Allowed ranges of the CKM4 parameters obtained from our analysis.

Electroweak precision tests[EWPT]

- Provide important constraints on SM4 (see talk by Juergen Rohrwild for details).
- Here we'll simply note that due EWPT in SM4 mass difference of t', b' needs to be less than ~ 70 GeV.

That prejudice suggests in SM4 prominent modes are t'-> b W; b' -> t W

HINTS of CKM breaking

Based on Enrico Lunghi+AS 0707.0212; 0803.4340; 0903.5059;09⊉200002 JHEL 1010.6069 JHEL SRL

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Important to Examine only DeltaF=2 observables:Leave out Vub sin 2 β = 0.87+-.09{Lunghi+AS,hep-ph/08034340} (became possible only due significantly reduced error in B_{κ})



requires careful follow-up

Continuing saga of Vub

- For past many years exclusive & inclusive show discrepancy (Latest; gotten worse) lattice $B \rightarrow X l Y$ C.g. $B \rightarrow T f y$
- Exc ~ (29.7 +-3.1)X10⁻⁴
- Inc ~ $(40.1+-2.7+-4.0)X10^{-4}$

-> Let's try NOT use Vub: initiated in '08 (EL&AS'08)... Not just for the above reason DNLY BECAME VIABLE DUE TO BETTER SIGNIFICANT BETTER

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UPDATE 2010







GOLD"



BRB->T)XIO Deviation INPUT $\sin(2\beta), \epsilon_K, \Delta M_q, |V_{cb}|, \gamma$ 0.768±0.099 (2.7\sigma) Sina $\sin(2\beta), \epsilon_K, \Delta M_q, |V_{cb}|$ 0.776±0.10 (2.7\sigma) 7,2.76 $\sin(2\beta), \epsilon_K, \Delta M_g, |V_{cb}|, \gamma, |V_{ub}|$ 0.762±0.096 (2.8) $\sin(2\beta), \epsilon_K, \Delta M_q, \gamma$ 0.761±0.099 (2.8\sigma) $\sin(2\beta), \Delta M_q, |V_{cb}|, \gamma$ 0.752±0.099 (2.8\sigma) ND SIM2B < 0.96 $\epsilon_K, \Delta M_q, |V_{cb}|, \gamma$ 1.29±0.40 (0.86) 0.87±0.26 (2.0) fBat, Vub tot $1.4 \pm 0.20 (< 1\sigma)$ $f_{R}^{\text{lat}}, |V_{\text{ub}}|_{\text{incl}}$ 0.74±0.14 (2.8) $f_B^{\text{lat}}, |V_{\text{ub}}|_{\text{excl}}$ $BT(B \rightarrow \tau \nu)_{exp}$ 1.68 ± 0.31 -20 2 4 Flavor & 4G; IPPP DURHAM Α. 19 Soni

In a nutshell

- Bulk of NP effects is in Bd,Bs mixing & in sin2β {CONFIRMS our 2008 findings}
- Bulk of NP NOT in B->τν, or in ε_K [Presence of subdominant effects therein certainly possible]
- Many, many checks for robustness of the conclusions
- DIFFICULT to RECONCILE RESULTS with CKM-SM

A BRIFF N25 years history of Br



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Average the four 2+1 flavor calculations presented

Except for BMW, all are preliminary, although all groups have recently published B_K results from earlier datasets, so preliminary work should be fairly reliable.

See also recent summary by FLAG working group of FLAVIANET (arXiv:1011.4408) They quedator $\mathbb{A}_{4}^{G; IRPP/DURHAM}$ for $N_{f} = 2+1$ 22



 $11 ~ 14.8 \pm 2.80/0$ LUNGHI + AS 09 2.8±2.8 Numerous tests to check stableity under lage variation of imputs.

Delta K pi

Possibilties for DEWSB

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Bundman & Rold arXiV: 0710.0623 WARPED SPACE



Figure 1: Two contributions to four-fermion interactions of the up-type fourth-generation quark: (a) from the interactions with a KK gluon; (b) from the four-fermion interactions induced by the bulk operators of (2.11).

0600-700GeV $m_H \sim 2m_U$

HEAVY fourth family cannot co-exist with light Higgs

Higgs description loses meaning

- $m_{t',b'} \approx 600$ GeV close to the unitarity upper bound
- Goldstone bosons of electroweak symmetry breaking couple strongly to t', b', so strong interactions unitarize WW scattering

Higgs is no longer needed

- strong interactions can also be responsible for the Goldstone bosons
- a fermion condensate replaces the Higgs vev
- so where do quark and lepton masses come from?

$$\langle \overline{\Psi}\Psi \rangle = \Lambda_{ew}^3 \qquad \frac{1}{\Lambda_{fl}^2} \overline{\Psi}\Psi\overline{\psi}\psi \qquad m_{\psi} = \frac{\Lambda_{ew}^3}{\Lambda_{fl}^2}$$

Heavy 4th Generation's scenario What can a heavy 4th generation do? A heavy Higgs-4th Yukawa system can give rise to condensates and bound states of 4th generation fermions \Rightarrow Implications on the vacuum structure of the SM and the number of Higgs doublets (a mixture of fundamental and composites) The appearance of quasi-fixed points in the

 The appearance of quasi-fixed points in the 2-loop approximation at a scale Λ_{FP} ~ O(TeV) hints at a possible scale-invariant theory above Λ_{FP}.

CAN NOT BE JUST SM4

- Neutrino4 must be very heavy [LEP, Fermilab]
- D0,CDF mt', mb' > 350GeV [bounds in SM4]
- 4G fermions very likely play important role IN DEWSB -> likely a strongly coupled 4G
- Generic 2HDMmodels [see Bar-Shalom, Nandi and A. S '11] suggest significant differences in search strategies from SM4

and constraints of oblique parameters

Baryogenesis

• For SM3, there is unique CP invariance [Jarlskog '87]

$$J = \operatorname{Im} \operatorname{det}[M_{u}M_{u}^{\dagger}M_{d}M_{d}^{\dagger}]$$

$$= 2(m_{t}^{2} - m_{u}^{2})(m_{t}^{2} - m_{c}^{2})(m_{c}^{2} - m_{u}^{2})$$

$$(m_{b}^{2} - m_{d}^{2})(m_{b}^{2} - m_{s}^{2})(m_{s}^{2} - m_{d}^{2})A$$

$$= \operatorname{exceding} \operatorname{ly} \operatorname{cnal} \operatorname{first} \operatorname{small} \operatorname{cnas} \operatorname{small} \operatorname{$$

In 4G models

$$J_{234} = 2(m_{t'}^2 - m_t^2)(m_{t'}^2 - m_c^2)(m_t^2 - m_c^2)$$

$$(m_{b'}^2 - m_b^2)(m_{b'}^2 - m_s^2)(m_b^2 - m_s^2)A_{234}$$

$$\int \frac{16}{334} \int 0^{16} \int \frac{16}{344} \int \frac{16}{344$$

1

DM possibilities in 4G models

- Lee,Liu, AS arXiv: 1105.3490 [c also Volovik hepph/0310006]
- Link heaviness of v4 with new abelian gauge interaction: B – 4 L4 distinguishing it from SM3 neutrinos
- Just a heavy v4 in SM4 is NOT a viable DMcandidate ...conflicts with direct detection bounds[CDMS & XENON} both for Dirac (spin independent) or Majorana (spin dependent) cases [XENON 10]

New B-xiL gauge symmetry for 4G

- [its similar to earlier 3G studies: Ma arXiv:9709.474; Lee & Ma, 1001.0768]
- Relevant effective new interaction is:

quarks

NO

Als

$$\mathcal{L}_{int}^{Z'} = g_X Z'_{\mu} \left[\frac{1}{6} \bar{q} \gamma^{\mu} q - 2 \bar{e}_4 \gamma^{\mu} e_4 - 2 \bar{\nu}_{4L} \gamma^{\mu} \gamma_5 \nu_{4L} - 2 \bar{\nu}_{4R} \gamma^{\mu} \gamma_5 \nu_{4R} \right]$$

TE: (almost) leptophobic Z'
o couples rather weakly to

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RELIC DENSITY VS DM MASS



L'HC Seanches NOT Yet Sensitive



Implications for High Intensity Experiments

- Effects on CP-conserving observables are expected to be small.[e.g b s gamma; b s l l..]
- A very nice exception is Bs->mu mu
- CP-violating observables are more susceptible due to the 2 new phases.
- Using null observables of SM3 is a good idea..
 [see e.g. Gershon & AS hep-ph/0607230]
- Here are some candidates: Bs-> psi phi,
 B(Bs) -> X(Xs) | nu, D0->phi pi0, eta' rho, ks omega; KL-> pi0 nu nu, epsilon'/epsilon
- May also appear as deviation from sin2beta


Predicted range of S(psiKs) in SM4 (with mt'=400 GeV) is (shown in red) compared with the experimentally measured value via the psi Ks mode (1 sigma error) and with the SM (1 sigma)



SM4 seems to predict sin2 beta around 0.70 with an error of about 0.06; S. Nandi and A.S, 1011.6091

Recent D0 result is vertical axis and combined D0, CDF each For SM4 error on Delta_Gamma_s is increased by a factor of two resulting in ~50% increase in a_sl^s



DO update 11





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(Sin2beta)tree vs (sin2beta)penguin

- The other hint is that b-> s penguin transitions have a non-vanishing BSM phase.
- Cleanest modes (eta' Ks, phi Ks) seem to show smaller value ~0.59 – 0.56
- Central value for several other modes is smaller than 0.68 [and of course practically all are well below the theory prediction ~0.85] Note sin 42 deg = 0.67 sin 35 deg = 0.59 NEED TO TARGET such small deviation MMM MMM MMM MMM MMM MMM MMM

"....but we have only just got started", G Raven at LP11

A sentiment I enthusiastically share

@ colliders







O. S 6 W 1.2DSlMET 2.2 SSl MET <u>ب</u> ۲ 3.11, 26 MET • A. Soni

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Detection of 4G [SM4] quarks [Atwood, Gupta, AS'11,arXiv:1104.3874]

- An important feature: To a rather good approximation, pair production is near threshold. Using this feature showed robust signals:
- 1) Single lepton
- 2) Opposite sign leptons
- 3) Same sign dileptons

All 3 viable

Search for 4th generation at hadron colliders

- Large production rates at the LHC
 - b'→tW gives same sign leptons, plus jets, E^{miss}
 - Small Standard Model background
 - Can reconstruct b' mass



Experimental bounds

- CDF, D0 and CMS have looked for t', b' assuming simple SM4 decay patterns.
- CDF: mb' > 338 GeV; mt' > 358 GeV
- D0:
- CMS (LHC7, 34/pb): mb' between 255 GeV & 361 GeV excluded @ 95%CL.
- ATLAS [UCI group actively involved] talk by Michael Werth

Theory bounds (prejudice)

 Chanowitz, Furman, Hinchliffe (`78-79) showed perturbative unitarity requires quark mass < ~ 550 GeV; leptons < ~ 1 TeV

Important Caveat: Nature cares less above perturbation theory than many of us [e.g. QCD]

[Infact strong coupling could lead to DEWSB]

Cannot be simple SM4

- Even if a 4th generation exists it is unlikely to be a simple replica of SM3:
- Neutrino mass provides a strong clue
- DM possibility and baryogenesis both strongly suggest 4G not in SM4
- Highly implausible that heavy quarks will not be used for DEWSB.....mH tends to be heavy in these models
- Simple SM4 light higgs seems strongly disfavored by data



Other exotic fermions are still alive and interesting, but the sequential 4th generation is in deep troupble!

4G2HDM:LEET for DEWSB

[Bar-Shalom,Nandi, A.S '11, arXiv:1105.6095] see also He & Valencia: arXiv:1108:0222

- 2 Higgs doublet model for the 4th; 2nd doublet with v2/v1 >> 1 couples only to 4th gen
- FCNC only among 3rd and 4th gen
- Significant modifications to t', b' decays
- t'->t h, b h+ very large BR, similarly for b'
- Requires modification to search strategies
- Can account for tt FB asymmetry

t' decay pattern



b' decay pattern



 $m_{b'}$ =500 GeV, m_{h} =220 GeV, m_{H+} =300 GeV, $tan\beta$ =1, ϵ_t = $m_t/m_{t'}$

4G2HDM: significantly alleviates EWPC on 4G



FIG. 13: Allowed regions in the $\Delta m_{q'} - \Delta m_{\ell'}$ plane within the 95%CL contour in the S-T plane, for the 4G2HDMs (left) and for the SM4 (right). The data points are varied as in Fig. 11.

Top pair FB asymmetry from strongly coupled 4G quarks

- [Davoudiasl, McElmurry and AS arXiv:1108.1173]
- CDF has been reporting ~2 sigma deviation from SM in the integrated asymm. And in the high invariant top pair mass about 3.5 sigma deviation from SM:
- 0.158 +- 0.075 vs 0.058+-0.009
- For high (>450 GeV), 0.475+-0.114 vs 0.088+-0.013
- NEW (EPS'11) D0 finds similar integrated asymmetery (a bit more significant than CDF); in the high mass region they see some increase but not so pronounced

We consider a simple toy 4G model with [FC] scalars

 $\mathcal{L} \supset \lambda_{ut'} \phi^0 \bar{u}t' + \lambda_{ub'} \phi^+ \bar{u}b' + \lambda_{dt'} \phi^- \bar{d}t'$ $+ \lambda_{db'} \phi^0 \bar{d}b' + \lambda_{tt'} \phi^0 \bar{t}t' + \lambda_{tb'} \phi^+ \bar{t}b' + \text{H.C.},$ $\text{NOTE } t' \rightarrow t \phi \text{Dom} \text{INANT} \text{Not} t' \rightarrow 6 \text{W}$ ¢ $\overline{\varrho}$ \overline{o} ASYMMETRY from interce with





FIG. 4: Asymmetry as a function of the invariant mass $M_{t\bar{t}}$ for $(m_Q, m_{\phi}, \lambda_u, \lambda_d, \lambda_t) = (380 \,\text{GeV}, 140 \,\text{GeV}, 0, 3.5, 4.5)$ (solid), $(380 \,\text{GeV}, 140 \,\text{GeV}, 1, 2, 6)$ (dashed), and (440 $\,\text{GeV}, 120 \,\text{GeV}, 0, 3.5, 4.5$) (dotted). DIAGNOSTIC Power 6f ASYM VS MASS Flavor & 4G IPPP DURHAM A. Soni

Summary

- 4G extensions of SM3 should be seriously considered
- It is highly unlikely to be just SM4
- Heavier 4G quarks can trigger DEWSB and consequently can be a simpler solution to hierarchy problem w/o the need for multitude of new parameters
- It facilitates significantly baryogenesis over SM3
- Opens up new avenues for DM
- Unlike almost all BSM models, 4G extension need not cause flavor or CP problem
- (Sin2beta)tree less than SM3 and (sin2beta)penguin less than both, may be hints of 4G
- AFB (tt) may be a hint of a strongly coupled 4G
- Experimental searches for t' -> t h, b'->t h+- should be given a high priority [in addition to conventional modes]

XTRAS



Random scans of free-parameters:

 $0 \le \alpha \le 2\pi \quad \tan\beta \le 30, \quad \theta_{34} \le 0.3;$

 $100 \text{ GeV} \le m_h \le 1 \text{ TeV}$ $100 \text{ GeV} \le m_A \le 1 \text{ TeV}$ $m_h \le m_H \le 1.5 \text{ TeV}$

400 GeV $\leq m_{t'}, m_{b'} \leq 600$ GeV 100 GeV $\leq m_{\nu'}, m_{\tau'} \leq 1.2$ TeV



Order of magnitude more points within 68%CL contour in 4G2HDMs

S-T & the $M_{\text{H+}}$ - tan β plane



Favors: low $tan\beta \sim O(few)$

S-T & quark/lepton mass splitting:



$\Delta m_{q}, \Delta m_{l} \rightarrow 0$ allowed also for $\theta_{34} \sim O(0.1)$

⇒ Isospin breaking required (to compensate for effects of extra fermions and Higgs particles) provided by Higgs sector; more room for a light Higgs ...

Special/interesting points that survive combined constraints from EWPD & flavor physics

e.g., for tan β =1:





C also: Novikov, Rozanov, Vysotsky, 0904.4570 & earlier works

4th family is not inconsistent with LEP EWPC
See also M. Chanowitz, arXiv:0903.3570;
1007.0043; Erler abd Langacker 1003.3211

TABLE I. Examples of the total contributions to ΔS and ΔT from a fourth generation. The lepton masses are fixed to $m_{\nu_4} = 100 \text{ GeV}$ and $m_{\ell_4} = 155 \text{ GeV}$, giving $\Delta S_{\nu\ell} = 0.00$ and $\Delta T_{\nu\ell} = 0.05$. The best fit to data is (S, T) = (0.06, 0.11) [35]. The standard model is normalized to (0, 0) for $m_t = 170.9 \text{ GeV}$ and $m_H = 115 \text{ GeV}$. All points are within the 68% C.L. contour defined by the LEP EWWG [35].

| Parameter set | m_{u_4} | m_{d_4} | m_H | $\Delta S_{\rm tot}$ | $\Delta T_{\rm tot}$ |
|---------------|-----------|-----------|-------|----------------------|----------------------|
| (a) | 310 | 260 | 115 | 0.15 | 0.19 |
| (b) | 320 | 260 | 200 | 0.19 | 0.20 |
| (c) | 330 | 260 | 300 | 0.21 | 0.22 |
| (d) | 400 | 350 | 115 | 0.15 | 0.19 |
| (e) | 400 | 340 | 200 | 0.19 | 0.20 |
| (f) | 400 | 325 | 300 | 0.21 | 0.25 |






Recently D0 (V M Abazov et al, arXiv: 1005.2757) reported a_sl^s = -0.0146 +- 0.0075)

$$\begin{aligned} |V_{cb}|_{excl} &= (39.0 \pm 1.2)10^{-3} & \eta_1 = 1.51 \pm 0.24 \ [18] \\ |V_{cb}|_{incl} &= (41.31 \pm 0.76)10^{-3} & \eta_2 = 0.5765 \pm 0.0065 \ [19] \\ |V_{cb}|_{tot} &= (40.43 \pm 0.86)10^{-3} & \eta_3 = 0.494 \pm 0.046 \ [20, 21] \\ |V_{ub}|_{excl} &= (29.7 \pm 3.1)10^{-4} & \eta_B = 0.551 \pm 0.007 \ [22] \\ |V_{ub}|_{incl} &= (40.1 \pm 2.7 \pm 4.0)10^{-4} & \xi = 1.23 \pm 0.04 \ [23, 24] \\ |V_{ub}|_{tot} &= (32.7 \pm 4.7)10^{-4} & \lambda = 0.2255 \pm 0.0007 \\ \Delta m_{B_d} &= (0.507 \pm 0.005) \ \text{ps}^{-1} & \alpha = (89.5 \pm 4.3)^{\circ} \\ \Delta m_{B_s} &= (17.77 \pm 0.12) \ \text{ps}^{-1} & \kappa_{\varepsilon} = 0.94 \pm 0.02 \ [25-27] \\ S_{\psi K_S} &= 0.668 \pm 0.023 \ [28] & \gamma = (74 \pm 11)^{\circ} \\ m_c(m_c) &= (1.268 \pm 0.009) \ \text{GeV} & \hat{B}_K = 0.740 \pm 0.025 \\ m_{t,pole} &= (172.4 \pm 1.2) \ \text{GeV} & f_K &= (155.8 \pm 1.7) \ \text{MeV} \\ f_{B_s} \sqrt{\hat{B}_s} &= (276 \pm 19) \ \text{MeV} \ [23] & \varepsilon_K &= (2.229 \pm 0.012)10^{-3} \\ f_B &= (208 \pm 8) \ \text{MeV} \ [23, 24]^{-a} & \hat{B}_d &= 1.26 \pm 0.10 \ [23, 24] \\ \mathcal{B}_{B \to \tau\nu} &= (1.68 \pm 0.31) \times 10^{-4} \ [30-32] \end{aligned}$$

^aOur value of f_B reflects the change in the overall scale (r_1) recently adopted by the Fermilab/MILC and HPQCD collaborations [29] \rightarrow AS in 2008 WC ASERFAG IRPP DURHAM Vuls. Soni

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TABLE I. Inputs that we use in order to constrain the SM4 parameter space, when not explicitly stated, we take the inputs from Particle Data Group [7]; for the lattice inputs see also [6].

| $B_K = 0.740 \pm 0.025 \ [35-37]$ | $R_{bb} = 0.216 \pm 0.001$ |
|---|--|
| $f_{bd}\sqrt{B_{bd}} = 0.224 \pm 0.015 \text{ GeV} [38,39]$ | $ V_{ub} = (32.8 \pm 2.6) \times 10^{-4a}$ |
| $\xi = 1.232 \pm 0.042$ [38,39] | $ V_{cb} = (40.86 \pm 1.0) \times 10^{-3}$ |
| $\eta_c = 1.51 \pm 0.24$ [40] | $\gamma = (73.0 \pm 13.0)^{\circ}$ |
| $\eta_t = 0.5765 \pm 0.0065 \ [41]$ | $\mathcal{BR}(B \rightarrow X_s \gamma) = (3.55 \pm 0.25) \times 10^{-4}$ |
| $\eta_{ct} = 0.494 \pm 0.046$ [42] | $\mathcal{BR}(B \to X_s \ell^+ \ell^-) = (0.44 \pm 0.12) \times 10^{-6}$ |
| $\Delta M_s = (17.77 \pm 0.12) \text{ ps}^{-1}$ | $\mathcal{BR}(K^+ \to \pi^+ \nu \nu) = (0.147^{+0.130}_{-0.089}) \times 10^{-9}$ |
| $\Delta M_d = (0.507 \pm 0.005) \text{ ps}^{-1}$ | $\mathcal{BR}(B \to X_c \ell \nu) = (10.61 \pm 0.17) \times 10^{-2}$ |
| $ \epsilon_k \times 10^3 = 2.32 \pm 0.007$ | $T_4 = 0.11 \pm 0.14$ |
| $\kappa_{\epsilon} = 0.94 \pm 0.02 \ [43]^{\mathrm{b}}$ | $m_t(m_t) = (163.5 \pm 1.7) \text{ GeV}$ |

^aIt is the weighted average of $V_{ub}^{inl} = (40.1 \pm 2.7 \pm 4.0) \times 10^{-4}$ and $V_{ub}^{exl} = (29.7 \pm 3.1) \times 10^{-4}$. In our numerical work to follow, we increase the error on $|V_{ub}|$ by 50% and take the total error to be around 12% because of the appreciable disagreement between the two determinations. ^bWe tacitly assume that κ_{ϵ} in SM4 is approximately the same as in SM3.