Rare Kaon decays

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K→ πνν $R_K = K_{e2} / K\mu 2$ Search for ν_H

But no recent updates or plans for $K_L \rightarrow \pi II$

$K \rightarrow \pi \upsilon \upsilon$: Standard Model

q_{2/3}

<u>۸۸۸</u>

e,μ,τ

Smx -

- FCNC loop processes
- Short distance dynamics dominated
- One semileptonic operator
- Hadronic Matrix Element related to measured quantities in semileptonic K decay

$$\lambda = V_{us} \\ \lambda_c = V_{cs}^* V_{cd} \\ \lambda_t = V_{ts}^* V_{td} \\ B(K_L^0 \to \pi^0 v \bar{v}) = \kappa_L \cdot \left(\frac{\operatorname{Im} \lambda_t}{\lambda^5} X(x_t) \right)^2 + \left(\frac{\operatorname{Re} \lambda_t}{\lambda^5} X(x_t) + \frac{\operatorname{Re} \lambda_c}{\lambda} P_c(X) \right)^2 \right] \\ B(K_L^0 \to \pi^0 v \bar{v}) = \kappa_L \cdot \left(\frac{\operatorname{Im} \lambda_t}{\lambda^5} X(x_t) \right)^2 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \right) \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ = r_{K^+} \cdot \frac{3\alpha^2 Br(K^+ \to \pi^0 e^+ v)}{2\pi^2 \sin^4 \theta_W} \cdot \lambda^8 \\ \kappa_+ =$$

$K \rightarrow \pi \ell \ell decays$

Quantitative tests of the SM

Tiny BRs can be

computed to very high

(few percent) precision



	SM prediction	Exp. Limit	Signal(Backg)
$K_L \rightarrow \pi^0 e^+ e^-$	(3.54–1.56)·10 ⁻¹¹ (CPV _{dir} 3·10 ⁻¹²)	< 2.8 ·10 ⁻¹⁰ (FNAL KTeV)	3 ev. (2.05 bkg)
$K_L \rightarrow \pi^0 \mu^+ \mu^-$	$(1.41-0.95) \cdot 10^{-11}$ (CPV _{dir} 1·10 ⁻¹²)	< 3.8 ·10 ⁻¹⁰ (FNAL KTeV)	2 ev. (0.87 bkg)
$K^+ \rightarrow \pi^+ \nu \nu$	(8.5±0.7)·10 ⁻¹¹	1.73 ^{+1.15} -1.05 · 10 ⁻¹⁰ (BNL E787+E949)	7 evt. (bkg. 1.38)
$K_L \rightarrow \pi^0 \nu \nu$	(2.8±0.6)·10 ⁻¹¹	< 6.7 ·10 ⁻⁸ (KEK E391a)	



Differences between SM4, RSc, LHT enhancements

SM4 predictions

Plots from: Buras, Duling, Feldmann, Heidsieck, Promberger, Recksiegel: 1002.2126 Plot density is NOT a probability density; correlations represented by enveloping curves

For general input discussion see S.Nandi talk and references

But mixing in lepton sector assumed to be $W_{I4} \sim O(\lambda)$ and so neglected here (in presence of substantial mixing with 4G leptons, charm contribution in $\pi v v$ will be affected)

Also recent SM4 1 σ limit (Soni, Alok, Giri, Mohanta, Nandi 1002.0595): Br(K+ $\rightarrow \pi + \nu \nu$) = (4.0 \rightarrow 13.0) 10⁻¹¹ m_{t'} = 400-600 GeV Br(K_L $\rightarrow \pi 0 \nu \nu$) = (1.0 \rightarrow 6.2) 10⁻¹¹ m_{t'} = 400-600 GeV



For red points (BS3) now disfavoured by LHC result; Also disfavoured if ϵ'/ϵ taken into account

SM4 correlations 1.2×10^{-9} $4. \times 10^{-10}$ $1. \times 10^{-9}$ $3. imes 10^{-10}$ $\operatorname{Br}(K_L \to \pi^0 \nu \overline{\nu})$ $8. \times 10^{-10}$ $2. \times 10^{-10}$ $6. \times 10^{-10}$ $4. \times 10^{-10}$ $1. imes 10^{-10}$ $2. \times 10^{-10}$ $0^{-1.0}$ -0.50.5 0.0 1.0 -1.0 -0.50.0 0.5 1.0 $S_{\psi\phi}$ $S_{\psi\phi}$ $4. \times 10^{-10}$ 1.2×10^{-9} $1. \times 10^{-9}$ $3. \times 10^{-10}$ $8. \times 10^{-10}$ $\operatorname{Br}(K_L \to \pi^0 \nu \overline{\nu})$ $2. \times 10^{-10}$ $6. \times 10^{-10}$ $4. \times 10^{-10}$ $1. \times 10^{-10}$ $2. \times 10^{-10}$

 $\operatorname{Br}(K^+ \to \pi^+ \nu \overline{\nu})$

 $\operatorname{Br}(K^+ \to \pi^+ \nu \overline{\nu})$

0

0

 $2.\times10^{-9}$ $4.\times10^{-9}$ $6.\times10^{-9}$ $8.\times10^{-9}$ $1.\times10^{-8}$

 $Br(B_s \rightarrow \mu^+ \mu^-)$

Red points now disfavoured by LHC Even for small effects in B-> $\mu\mu$ and S $\phi\phi$, large effects possible in K 8

0.0

0.5

1.0

 $\operatorname{Br}(B \to X_s \nu \overline{\nu}) / \operatorname{Br}(B \to X_s \nu \overline{\nu})_{SM}$

1.5

2.0



Measurement of K⁺ \rightarrow π⁺υυ with new decay in-flight technique Intense un-separated (6% K⁺) 75 GeV/*c* hadron beam: 5 ·10¹² ppp High-energy: high yield, large decay volume, more powerful vetoing Track incoming K⁺ in 800MHz beam, particle ID, photon vetoing

5 10¹² K⁺ decays/year
55 SM events/(<Snowmass) year

Expected precision: 10%

NA62 : principle of experiment

Decay-in-flight very different from stopped kaon technique

BR(SM) = 8×10^{-11} ~ 10^{12} K⁺ decays Acceptance= 10%



- K decays in flight
- Intense beam of protons from SPS
- High energy K (P_K = 75 GeV/c)
- Cherenkov K ID: CEDAR

Kinematic rejection



- Kaon: beam tracking
- Pion: spectrometer
- Excellent timing for K-π association
- γ/μ : calorimeter
- Charge Veto : spectrometer
- **\pi/\mu separation : RICH**





First dedicated pilot experiment to search for $K_L \rightarrow \pi^0 \upsilon \upsilon$ at the KEK-PS Improve over KTeV (Dalitz) limit: BR < 5.9 ·10⁻⁷

- High intensity: 2 ·10¹² ppp 12 GeV/*c* (50% DC)
- •"Pencil" beam as transverse constraint: ~ $2 \text{ GeV}/c \text{ K}_{\text{L}}$ at 4° and 11m
- Photon veto hermeticity down to 1-2 MeV: Pb/scint in high vacuum
- Good EM calorimetry: ~500 pure CsI 7x7 cm², with central hole

Three runs (2004-2005): 12 month total

BR(K_L→π⁰υυ)<6.7 · 10⁻⁸ (90% CL)







Higher beam intensity, acceptance Lower DC, yield (angle): Statistics: 3000 x E391a

Halo n/K: 240x E391a: new beam line

Improved **background** control: new EM calorimeter (> granularity, longer), new backside charged veto, new beam-hole γ veto (25x Pb/aerogel)

Step 1: SES = 2.7 SM events (3 Snowmass years) with 2.2 background

Step 2 upgrade: **100 SM events** (dedicated, smaller targeting angle beam line, larger detector)

66 people, 16 institutions (Japan, Korea, USA, Russia, Taiwan) Stage 2 approval, beam line commissioned, in preparation

First physics run in 2013(?)

$K \rightarrow \pi \upsilon \upsilon$ beyond SM



Purple bands: foreseen 10% precision

$R_{K} = K_{e2}/K_{\mu 2}$ beyond the SM

<u>2HDM – tree level</u> (including SUSY) K_{12} can proceed via exchange of charged Higgs H[±] (in place of W[±]) \rightarrow Does not affect the ratio R_K

<u>2HDM – one-loop level</u>

Dominant contribution to ΔR_{K} : H[±] mediated LFV (rather than LFC) with emission of v_{τ} $\rightarrow R_{K}$ enhancement can be experimentally accessible

$$\mathbf{R}_{\mathbf{K}}^{\text{LFV}} \approx \mathbf{R}_{\mathbf{K}}^{\text{SM}} \left[1 + \left(\frac{\mathbf{m}_{\mathbf{K}}^4}{\mathbf{M}_{\mathbf{H}^{\pm}}^4} \right) \left(\frac{\mathbf{m}_{\tau}^2}{\mathbf{M}_{\mathbf{e}}^2} \right) | \boldsymbol{\Delta}_{\mathbf{13}} |^2 \text{tan}^6 \, \beta \right]$$

JHEP 0811 (2008) 042 \vec{s} \vec{h} \vec{h}

PRD 74 (2006) 011701.

Analogous SUSY effect in pion decay is suppressed by a factor $(M_{\pi}/M_{K})^{4} \approx 6 \times 10^{-3}$ (see also PRD76 (007) 095017)

Up to ~1% effect : slepton mixing $\Delta_{13}=5\times10^{-4}$, tan $\beta=40$, M_H=500 GeV/c² lead to R_K^{MSSM} = R_KSM(1+0.013)



Backgrounds and Uncertainty

Backgrounds

Source	B/(S+B)
Κ _{μ2}	(5.64±0.20)%
K _{μ2} (μ→e)	(0.26±0.03)%
$\overline{\mathrm{K}_{\mathrm{e}2\gamma}}\mathrm{(SD^+)}$	(2.60±0.11)%
Beam halo	(2.11±0.09)%
K _{e3(D)}	(0.18±0.09)%
Κ _{2π(D)}	(0.12±0.06)%
Wrong sign K	(0.04±0.02)%
Total	(10.95±0.27)%

Lepton momentum bins are differently affected by backgrounds and thus the systematic uncertainties.

Uncertainties

Source	$\delta R_{K} \times 10^{5}$
Statistical	0.007
K _{u2}	0.004
BR(K _{e2v} SD+)	0.002
$K^{\pm} \rightarrow \pi^{0} e^{\pm} v, K^{\pm} \rightarrow \pi^{\pm} \pi^{0}$	0.003
Beam halo	0.002
Helium purity	0.003
Acceptance	0.002
DCH alignment	0.001
Positron ID	0.001
Lkr readout inef	0.001
1-track trigger	0.001
Total	0.010

(40% data set: PLB 698 (2011), 105)



2.4 2.4 20 30 40 50 60 Data sample Lepton momentum, GeV/c Independent measurements 40 measurements (4 data samples x 10 momentum bins) in lepton momentum bins including correlations: $\chi^2/ndf = 47/39$

(systematic errors included, partially correlated)

R_K: world average

tanβ 06 tanβ July 2011 average **PDG 2008** 2HDM-II R_K=2.488(9)×10⁻⁵ Clark et al. (1972) BJTV K-JHVIK 80 Direct searches (LEP) Heard et al. (1975) 70 60 Heintze et al. (1976) 50 40 KLOE (2009) = PDG 2010 30 R_k: 95% CL exclusion NA62 (2011) 20 $\Delta_{\rm 13} = {\rm 1} \times {\rm 10}^{-3}$ full data set $\Delta_{13} = 5 \times 10^{-4}$ 10 $\Delta_{13}^{13} = 1 \times 10^{-1}$ SM 0<u>,</u> 2.8 R_K×10⁵ 2.5 2.7 2.3 2.4 2.6 0.4 0.5 0.6 0.7 0.8 0.9 0.1 0.2 0.3 H^{\pm} mass, TeV/c²

Other limits on 2HDM-II: PRD 82 (2010) 073012. SM with 4 generations: JHEP 1007 (2010) 006. **18**

 $(M_H, \tan\beta)$ 95% exclusion limits

4th generation constraints

Lacker, Menzel arXiv:1003.4532v2

Radiative and leptonic μ , τ Kl3, π l2 constraints

4th gen. PMNS matrix unitary:

$$\sum_{j=1,2,3} |U_{\mu j}|^2 = 1 - |U_{\mu 4}|^2$$
$$\sum_{i=1,2,3} |U_{ei}|^2 = 1 - |U_{e4}|^2$$



$$R_K = (1 - |U_{e4}|^2) / (1 - |U_{\mu4}|^2)$$

...cont.



Old plot shows deviation of Ue4 from zero Since then:

- 1) Kl3 results changed, lowering the effect
- 2) RK added, and pulls in opposite direction

Good agreement with Universality and improved Ue4 limit

Prospects with new NA62

• At present Kµ2/Ke2 can not be separated kinematically at high lepton momentum: particle identification is essential. Model-dependent MC correction required to evaluate the muon mis-ID probability.



Analysis without electron identification allows to avoid the uncertainty. Possibility: future NA62 analysis at p<40GeV/c without electron identification.

• For NA62 conditions (75 GeV/c beam, ~100 m decay volume): $P(K_{\mu 2}, \mu \rightarrow e \text{ decay})/R_{K} \sim 10$ but they are naturally suppressed by the muon polarisation Irreducible background: ~0.2% however known to excellent precision

...cont.

R_K is inclusive of IB radiation by definition. SD radiation is a background. INT is negligible.



• Ke2γ background:

KLOE measurement (arXiv:0907:3594): precision improved by a factor of 4; NA62 (2007 data): analysis in progress; expect precision similar to KLOE; future NA62 hermetic veto (large-angle and small-angle veto counters) will strongly decrease the background. SD background will not be relevant for a future NA62 precision RK measurement

• beam halo: future beam spectrometer (beam tracker plus beam Cherenkov) will allow time correlation between incoming kaons and decay products Expect background to be reduced to negligible level.

Larger statistics collectable in shorter time

To explain MiniBooNE results

- There is a 4th sterile neutrino, but it mixes with v_{μ} and not with v_{e} ;
- All LSND and MiniBooNE electron distributions can be explained assuming that this 4th neutrino is heavy (mass = 40 80 MeV) and decays to an ordinary neutrino + a photon: $v_H \rightarrow v + \gamma$ (LSND and MiniBooNE do not distinguish electrons from photons)



Gninenko's hypothesis requires $|U_{\mu H}|^2 = (2 - 8) \times 10^{-3}$; $\tau(v_H) \le 10^{-9} \text{ s}$ $\Gamma(K \rightarrow \mu v_H) = \Gamma(K \rightarrow \mu v_{\mu}) |U_{\mu H}|^2 \rho(m_{v_H})$ (in the rest frame) Ratio of phase space

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- The missing mass resolution of the NA62 detector in the 2007 configuration is not good enough to identify $v_{\rm H}$ from the MM² distribution.

- In future, search for events containing a μ^+ , one photon and NOTHING ELSE exploiting the full coverage of photon veto (LAV, LKr, IRC, SAC) 24

Summary/Discussion

- Kaon decays can give several constraints on 4th generation
- BR(K+,L $\rightarrow \pi v v$) about to be measured
- RK reached record precision, and will improve in the future
- \bullet search for ν_{H} just started
- Possible to have strong effects in K and B at the same time
- Even for <u>SM-like</u> S $\phi\phi$ and Bs $\rightarrow\mu\mu$, possible large effects and correlations in K
- Interplay between rare decays and RK when 4G lepton mixing is not negligible
- If 4G lepton mixing is large, but RM is SM (Ue4=U μ 4), how do rare decays change
- Nature, correlation and variation of inputs; theory errors
- How $\pi\nu\nu$ change for various 4G quark mixing matrix scenarios
- If $\pi\nu\nu$ deviates from SM, how important for SM4 is to measure π II

Spares

Leptonic meson decays: $P^+ \rightarrow I^+ v$



Subject to hadronic uncertainties



$R_{K} = K_{e2}/K_{\mu 2}$ in the SM

Observable sensitive to lepton flavour and its SM expectation:



- <u>SM prediction</u>: excellent <u>sub-per mille</u> accuracy due to cancellation of hadronic uncertainties.
- Measurements of R_K and R_π have long been considered as tests of lepton universality.
- Recently understood: helicity suppression of R_K might enhance sensitivity to non-SM effects to an experimentally accessible level.

$$K^{+} \qquad W^{+} \qquad e^{+}, \mu^{+}$$

$$R_{K}^{SM} = (2.477 \pm 0.001) \times 10^{-5}$$
$$R_{\pi}^{SM} = (12.352 \pm 0.001) \times 10^{-5}$$

Phys. Rev. Lett. 99 (2007) 231801

CERN NA48/NA62



Measurement strategy

- (1) $K_{e2}/K_{\mu 2}$ candidates are collected <u>concurrently</u>:
 - analysis does not rely on kaon flux measurement;
 - several systematic effects cancel at first order (e.g. reconstruction/trigger efficiencies, time-dependent effects).

(2) counting experiment, measured independently in <u>10 lepton momentum bins</u> (owing to strong momentum dependence of backgrounds and event topology)

 $\mathsf{R}_{\mathsf{K}} = \frac{1}{\mathsf{D}} \cdot \frac{\mathsf{N}(\mathsf{K}_{e2}) - \mathsf{N}_{\mathsf{B}}(\mathsf{K}_{e2})}{\mathsf{N}(\mathsf{K}_{\mu2}) - \mathsf{N}_{\mathsf{B}}(\mathsf{K}_{\mu2})} \cdot \frac{\mathsf{A}(\mathsf{K}_{\mu2}) \times \mathsf{f}_{\mu} \times \varepsilon(\mathsf{K}_{\mu2})}{\mathsf{A}(\mathsf{K}_{e2}) \times \mathsf{f}_{e} \times \varepsilon(\mathsf{K}_{e2})} \cdot \frac{1}{\mathsf{f}_{\mathsf{LKr}}}$

 $\begin{array}{lll} N(K_{e2}), N(K_{\mu 2}): & \text{numbers of selected } K_{l2} \text{ candidates;} \\ N_B(K_{e2}), N_B(K_{\mu 2}): & \text{numbers of background events;} & & & & \\ N_B(K_{e2}), N_B(K_{\mu 2}): & \text{numbers of background events;} & & & & \\ A(K_{e2}), A(K_{\mu 2}): & MC \text{ geometric acceptances (no ID);} \\ f_{e'}, f_{\mu}: & \text{directly measured particle ID efficiencies;} \\ \epsilon(K_{e2})/\epsilon(K_{\mu 2}) > 99.9\%: & E_{LKr} \text{ trigger condition efficiency;} \\ f_{LKr} = 0.9980(3): & global LKr \text{ readout efficiency;} \\ D = 50-150 & & \text{downscaling factor of the } K_{\mu 2} \text{ trigger.} \end{array}$

(3) MC simulations use minimised:

- Geometrical part of the acceptance correction comes from simulation;
- PID, trigger, readout efficiencies are measured directly.

K_{e2} vs $K_{\mu 2}$ selection



$K_{\mu 2}$ background in K_{e2} sample

Main background source

Muon "catastrophic" energy loss in LKr by emission of energetic bremsstrahlung photons. $P_{ue} \sim 3 \times 10^{-6}$ (and momentum-dependent).

 $P_{\mu e}$ / $R_K \sim 10\%$: K_{u2} decays represent a major background

Direct measurement of P_{µe}

Pb wall (9.2X₀) in front of LKr: suppression of $\sim 10^{-4}$ positron contamination due to $\mu \rightarrow e$ decay.

 $K_{\mu 2}$ candidates, track traversing Pb, p>30GeV/c, E/p>0.95: positron contamination <10⁻⁸.



- \rightarrow ionization losses in Pb (low p);
- \rightarrow bremsstrahlung in Pb (high p).



The correction $f_{Pb} = P_{\mu e} / P_{\mu e}^{Pb}$ is evaluated with a dedicated Geant4-based simulation

[Muon bremsttranlung: Phys. Atom. Nucl. 60 (1997) 576] **33**

Muon mis-identification



$K_{\mu 2}$ with $\mu \rightarrow e$ decay in flight

For NA62 conditions (74 GeV/c beam, ~100 m decay volume), $N(K_{\mu 2}, \mu \rightarrow e \text{ decay})/N(K_{e2}) \sim 10$ $K_{\mu 2} (\mu \rightarrow e)$ naïvely seems a huge background

Muons from $K_{\mu 2}$ decay are fully polarized: Michel electron distribution

 $d^2\Gamma/dxd(\cos\Theta) \sim x^2[(3-2x) - \cos\Theta(1-2x)]$

 $x = E_e/E_{max} \approx 2E_e/M_{\mu'}$

 Θ is the angle between p_e and the muon spin (all quantities are defined in muon rest frame).

Result: $B/(S+B) = (0.26 \pm 0.03)\%$

Important but not dominant background





10⁻⁵

0.05

0.1

0.15

0.2

0.25

E_v, MeV

INT= interference; IB=inner bremsstrahlung

Beam halo background

Electrons produced by beam halo muons via $\mu \rightarrow e$ decay can be kinematically and geometrically compatible to genuine K_{e2} decays

Background measurement:

- Halo background much higher for K_{e2}^{-} (~20%) than for K_{e2}^{+} (~1%).
- Halo background in the $K_{\mu 2}$ sample is considerably lower.
- ~66% of the data sample is K⁺ only, ~7% is K⁻ only, ~27% has both
- K^+ halo component is measured directly with the K^- sample and vice versa.

1-7% background in K⁺/ K⁻ Pb/no Pb. The background is measured 0.1-0.2% precision, and strongly depends on decay vertex position and track momentum. The selection criteria (esp. Z_{vertex}) are optimized to minimize the halo background.

 $B/(S+B) = (2.11\pm0.09)\%$

Uncertainty: 1) limited size of control sample; 2) π, K decays upstream vacuum tank.





K₁₃: lepton universality test

Comparison of $|V_{us}|$ determined from K_{e3} vs K_{u3} decays

$$r_{\mu e} = \frac{[|V_{us}|f_{+}(0)]_{\mu 3, \exp}^{2}}{[|V_{us}|f_{+}(0)]_{e3, \exp}^{2}} = \frac{\Gamma_{K\mu 3}}{\Gamma_{Ke3}} \frac{I_{e3} (1 + 2\delta_{\rm EM}^{Ke})}{I_{\mu 3} (1 + 2\delta_{\rm EM}^{K\mu})} = (g_{\mu}/g_{e})^{2} = 1$$
Experimental results
$$K^{\pm}: \quad r_{\mu e} = 0.998(9)$$

$$K^{0}: \quad r_{\mu e} = 1.003(5) \rightarrow r_{\mu e} = 1.002(4)$$
Non-kaon measurements:

 $\begin{array}{ll} \pi \rightarrow |_{V}: & r_{\mu e} = 1.0042(33) & (\text{PRD 76 (2007) 095017}) \\ \tau \rightarrow |_{VV}: & r_{\mu e} = 1.000(4) & (\text{Rev.Mod.Phys. 78 (2006) 1043}) \end{array}$

The sensitivity in kaon sector approaches those obtained in the other fields.

SM

FNAL P996 proposal

5% measurement of BR(K⁺ $\rightarrow \pi^+ \upsilon \upsilon$) with **stopped beam** technique, improving **x100** over BNL E949 by using:

- •9.6 ·10¹³ 150 GeV/c p (kaon yield x7)
- TeVatron as a stretcher ring (95% DC), same detector rates (≈ 8 MHz)
- •Separated 550 MeV/c K+ beam (K/ $\pi \approx 2.5$, 13.5 m long, K⁺ stops x4.5)

Goal: 194^{+89}_{-79} events/year (1 year = 1.8 Snowmass years) with S/N ≈ 4



Kaons at Project-X

High-Intensity frontier path for Fermilab

Expect CD-3: 2014 Start construction: 2015 Complete: 2019

Flux potential for **ultimate** ultra-rare K decay measurements

~500 K⁺ $\rightarrow \pi^+ \upsilon \upsilon$ events/year (S/B ~ 4)

 K_L → π^0 υυ experiment: - Intrinsic high-precision timing: TOF approach (KOPIO) beam microbunching (50ps/40ns) - Round and small beam (acceptance and bkg rejection)

~200 K_L $\rightarrow \pi^0 \upsilon \upsilon$ evts/year (S/B ~ 5-10)

 $K^+ \rightarrow \pi^+ \nu \bar{\nu}$ Experiment





Beam photon veto

Two Protvino projects

SPHINX+GAMS+ISTRA → **OKA** at Protvino: 65-70 GeV 10¹³ ppp at U-70 (38% DC) **12.5 GeV** RF-separated K⁺ beam **5** 10⁶ Kpp (K/π ≈ 4) Commissioning beam and detector with runs started 2009 10-100x improvement on ISTRA Kaon program + spectroscopy



Ongoing R&D for a $K_L \rightarrow \pi^0 \upsilon \upsilon$ experiment **KLOD** Neutral pencil beam extracted @ 35 mrad, 10 GeV/c K⁰ 300 MHz n background: dual-readout spaghetti calorimeter Aim at **1 SM event** (S/B \approx 3) with 10 days of beam



$K \rightarrow \pi \upsilon \upsilon BR predictions$

2.2%

Theoretical improvements make

 $K^+ \rightarrow \pi^+ \upsilon \upsilon$

the errors small

 $K_L \rightarrow \pi^0$ υυ



CKM, parametric

Comparable, unprecedented, *tiny* theoretical errors

K→πυυ beyond SM



K \rightarrow πυυ **remains clean** also beyond SM:

single effective uu operator, calculable Wilson coeff., no long-distance effects

The long quest for $K^+ \rightarrow \pi^+ \upsilon \upsilon$



BR(K⁺ $\rightarrow \pi^+ \nu \nu$) = **1.73** $^{+1.15}_{-1.05} \times$ **10**⁻¹⁰