

Neutral meson mixing (Beauty)

IPPP workshop on
Flavour and the 4th family

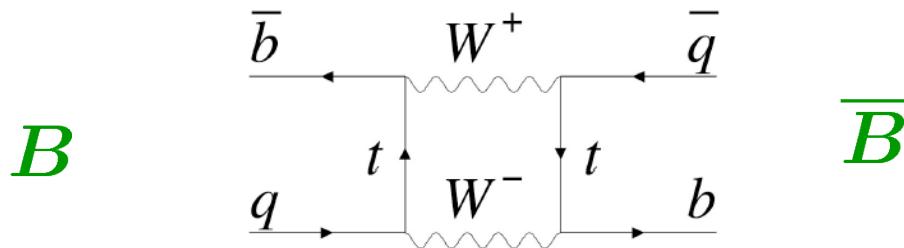


Wouter Hulsbergen



Mixing formalism in a nutshell

- neutral meson mixing: matter changes spontaneously into anti-matter!
- occurs in SM via this 2nd order diagram



- mixing and decay generically described by Schrodinger-like equation

$$i \frac{d}{dt} \begin{pmatrix} \langle B^0 | B(t) \rangle \\ \langle \bar{B}^0 | B(t) \rangle \end{pmatrix} = \begin{pmatrix} M_{11} - \frac{i}{2}\Gamma_{11} & M_{12} - \frac{i}{2}\Gamma_{12} \\ M_{21} - \frac{i}{2}\Gamma_{21} & M_{22} - \frac{i}{2}\Gamma_{22} \end{pmatrix} \begin{pmatrix} \langle B^0 | B(t) \rangle \\ \langle \bar{B}^0 | B(t) \rangle \end{pmatrix}$$

- M and Γ are hermitian: $M_{ij} = M_{ji}^*$ $\Gamma_{ij} = \Gamma_{ij}^*$
 - CPT invariance: $M_{11} = M_{22}$ $\Gamma_{11} = \Gamma_{22}$
 - phase difference of B^0 and \bar{B}^0 is arbitrary

--> 5 physical parameters: M Γ $|M_{12}|$ $|\Gamma_{12}|$ $\phi = \arg\left(-\frac{M_{12}}{\Gamma_{12}}\right)$

From parameters to observables

- time-evolution --> mass eigenstates

$$|B_L\rangle = p|B^0\rangle + q|\overline{B}^0\rangle$$

$$|B_H\rangle = p|B^0\rangle - q|\overline{B}^0\rangle$$

with eigenvalues

$$m_L + \frac{i}{2}\Gamma_L$$

$$m_H + \frac{i}{2}\Gamma_H$$

- usually characterised by observables

$$M \equiv \frac{m_H + m_L}{2}$$

$$\Delta M \equiv m_H - m_L$$

$$\Gamma \equiv \frac{\Gamma_H + \Gamma_L}{2}$$

$$\Delta\Gamma \equiv \Gamma_L - \Gamma_H$$

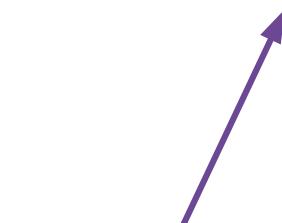
$$\left| \frac{q}{p} \right|$$

- that can (within reasonable approximation) be expressed as

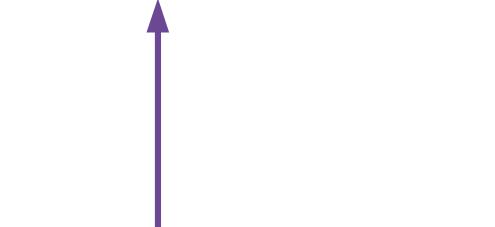
$$\Delta M \simeq 2|M_{12}|$$

$$\Delta\Gamma \simeq 2|\Gamma_{12}| \cos\phi$$

$$1 - \left| \frac{q}{p} \right|^2 \simeq \left| \frac{\Gamma_{12}}{M_{12}} \right| \sin\phi$$



mixing frequency



decay width difference



CP violation in mixing

From parameters to observables

- observables sensitive to NP in mixing

$$\Delta M \simeq 2|M_{12}|$$

 mixing frequency

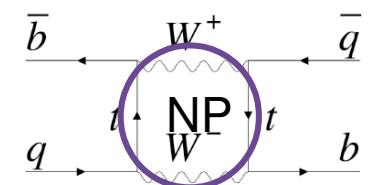
$$\Delta\Gamma \simeq 2|\Gamma_{12}| \cos \phi$$

 decay width difference

$$1 - \left| \frac{q}{p} \right|^2 \simeq \left| \frac{\Gamma_{12}}{M_{12}} \right| \sin \phi$$

 CP violation in mixing

- in the SM, the phase ϕ is small both in Bd and in Bs
- NP models usually only consider contributions to M_{12}



$$M_{12} = M_{12}^{\text{SM}} r^{\text{NP}} e^{i\phi^{\text{NP}}}$$

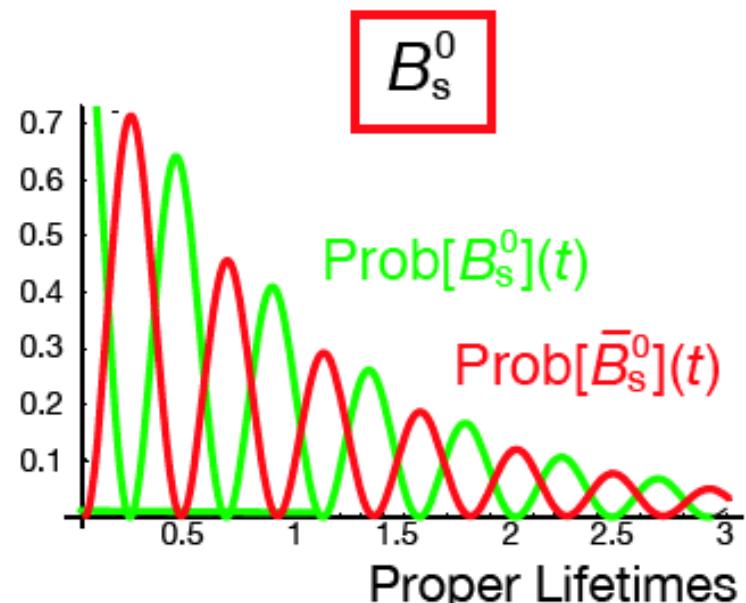
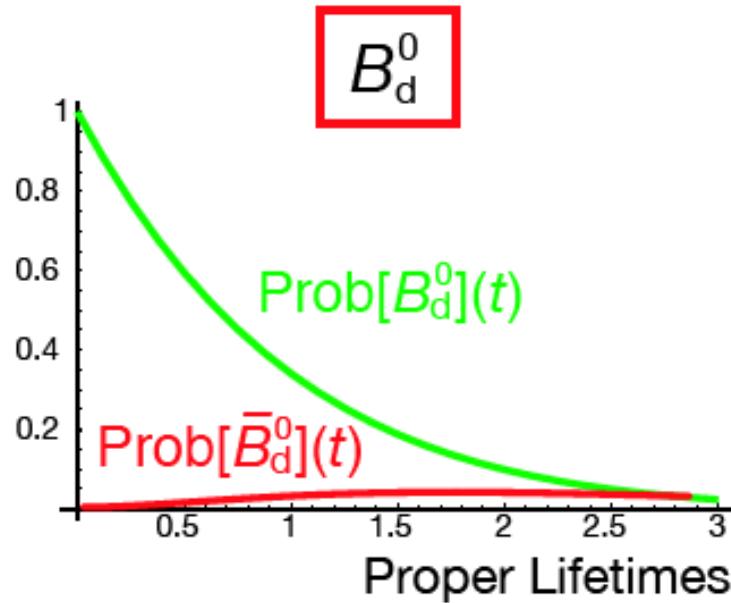
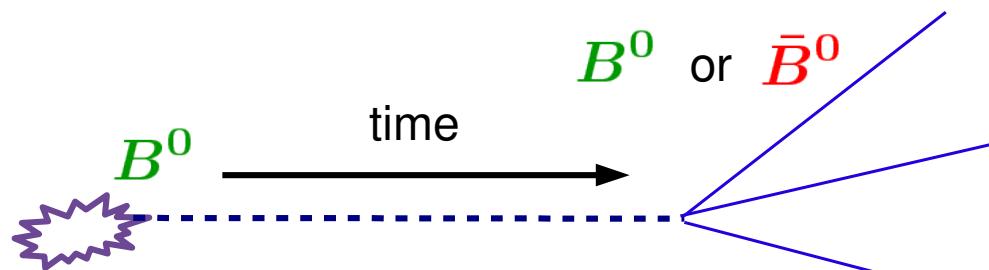
 constrained by mixing frequency

$$\phi = \phi^{\text{SM}} + \phi^{\text{NP}}$$

 constrained by CPV and $\Delta\Gamma$

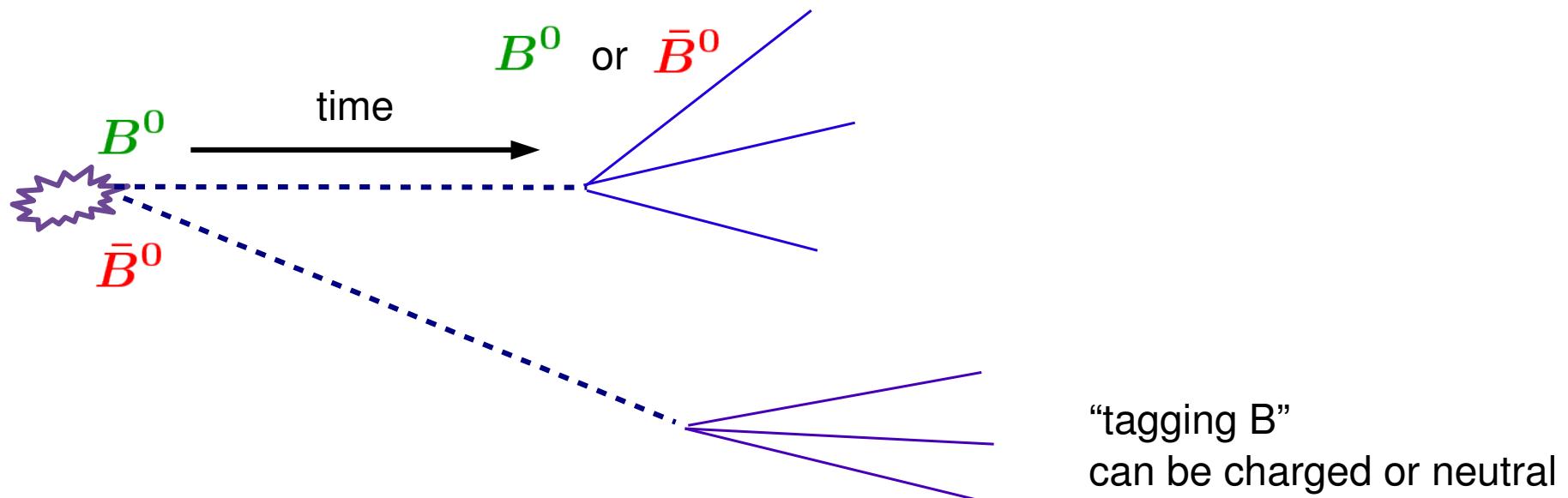
Experimental observables: mixing frequency

- decay time allows to observe time-evolution

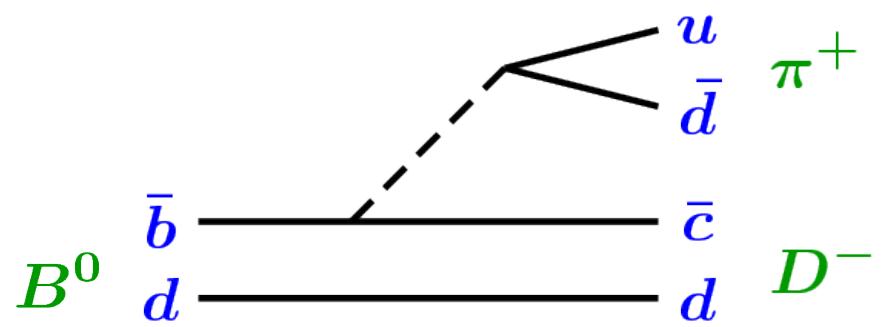
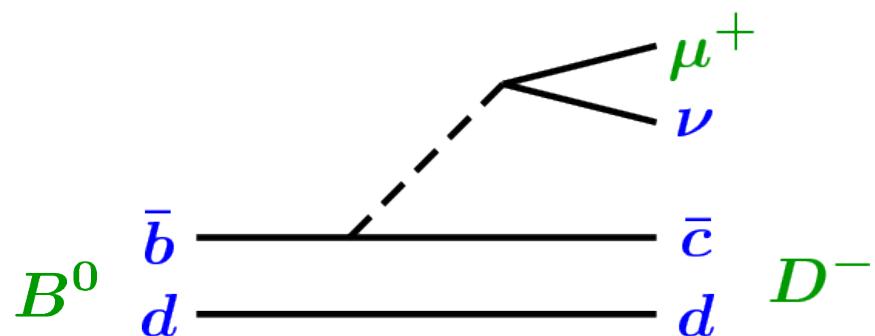


Experimental observables: mixing frequency

- decay time allows to observe time-evolution

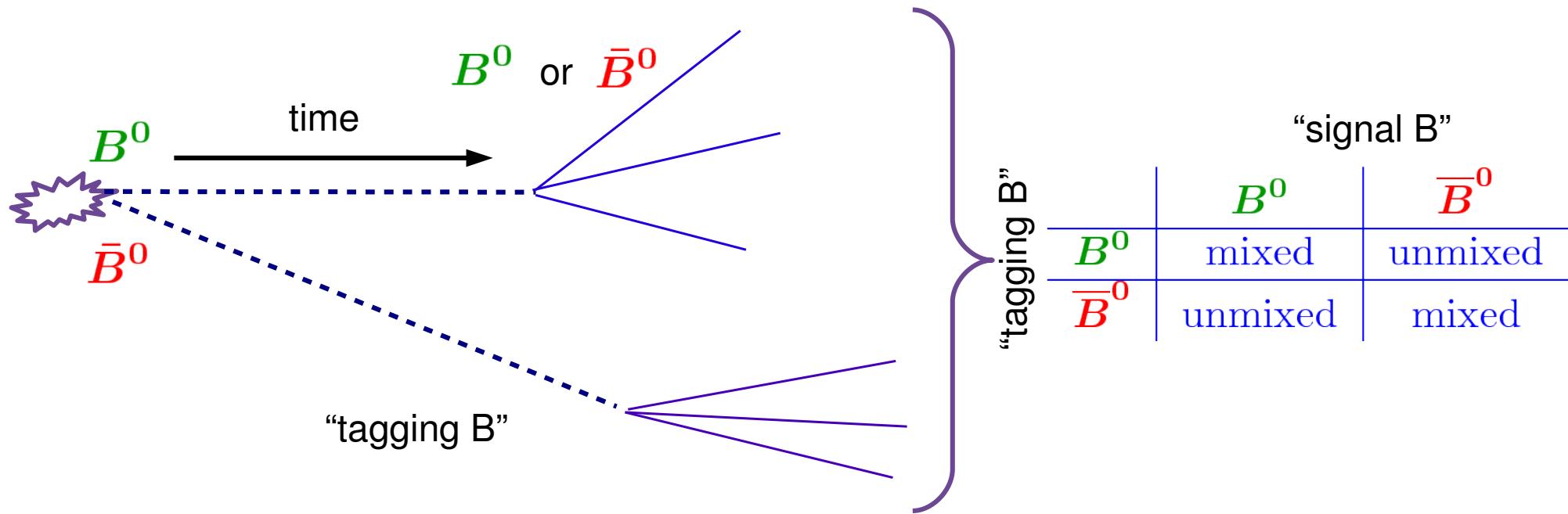


- look at “flavour specific” decays, e.g.



Experimental observables: mixing frequency

- decay time allows to observe time-evolution



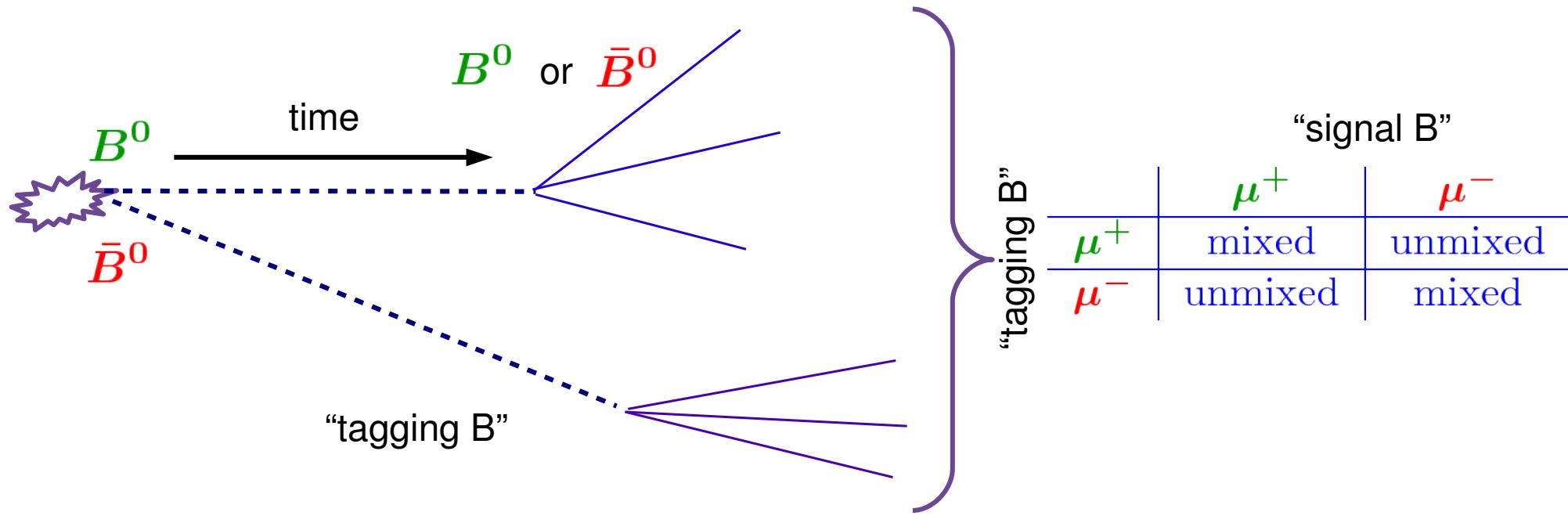
- mixing frequency: count 'mixed' versus 'unmixed'

$$A^{\text{mix}}(t) = \frac{N^{\text{unmixed}} - N^{\text{mixed}}}{N^{\text{unmixed}} + N^{\text{mixed}}} = \cos \Delta M t$$

- measurement requires flavour tagging and decay time measurement

Experimental observables: A_{SL}

- decay time allows to observe time-evolution



- from same observables, create CP asymmetry

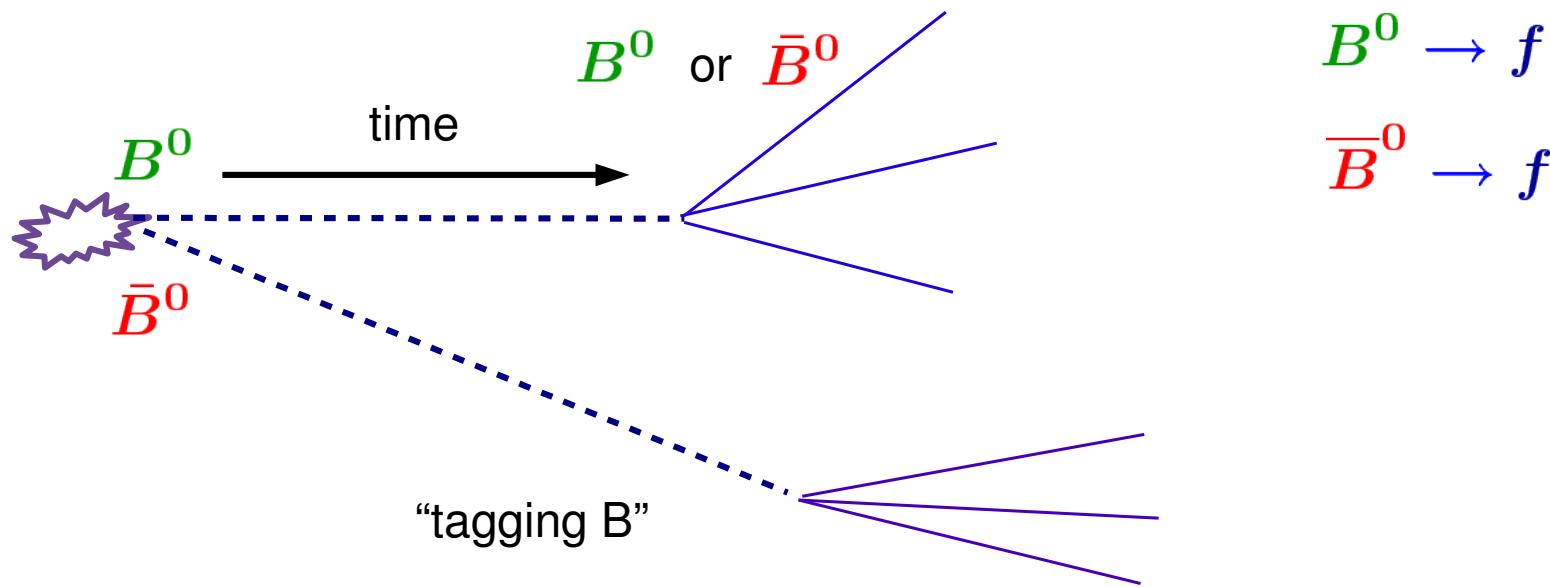
$$a_{fs} \equiv \frac{N^{++} - N^-}{N^{++} + N^-} = \frac{1 - |q/p|^4}{1 + |q/p|^4}$$

decay time
independent

- to probe CPV in mixing, use final state with zero predicted CPV in decay
- sometimes also called "semi-leptonic" asymmetry
- measurement requires very precise knowledge of detection asymmetries

Experimental observables: TD CPV

- now consider “non-flavour specific decays”



- interference between mixing and decay: “time-dependent CP violation”

$$A_{CP}(t) \equiv \frac{N(\bar{B} \rightarrow f) - N(B \rightarrow f)}{N(\bar{B} \rightarrow f) + N(B \rightarrow f)} = S \sin(\Delta m_q t)$$

- gives access to phase of mixing diagram
- measurement requires not only tagging and decay time, but also knowledge of their respective *dilutions*

Di-muon charge asymmetry

- mixing leads to 'wrong charge' combinations in $B \rightarrow X \mu$



- measure CPV in mixing via wrong charge asymmetry

$$A_{\text{fs}}^b \equiv \frac{N_b(\mu^+ \mu^+) - N_b(\mu^- \mu^-)}{N_b(\mu^+ \mu^+) + N_b(\mu^- \mu^-)}$$

- at hadron collider: method does not distinguish between B_d and B_s

--> measure linear combination

$$A_{\text{fs}}^b = C a_{\text{fs}}^d + (1 - C) a_{\text{fs}}^s$$



function of fragmentation fractions,
mixing parameters, lifetimes and
decay time acceptance

CPV parameter

with physics in

$$a_{\text{fs}}^q = \frac{\Delta \Gamma_q}{\Delta M_q} \tan \phi_q$$

- with value of C at Tevatron: $A_{sl}^b(SM) = (-0.023^{+0.005}_{-0.006})\%$

Di-muon charge asymmetry at D0

- experimentally very challenging
 - asymmetries in muon detection efficiency
 - asymmetries in backgrounds (e.g. muons from Kaon decay)
- main player: **D0** at Tevatron
 - low detection asymmetry due frequent field polarity changes
 - control of background asymmetries using single muon asymmetry

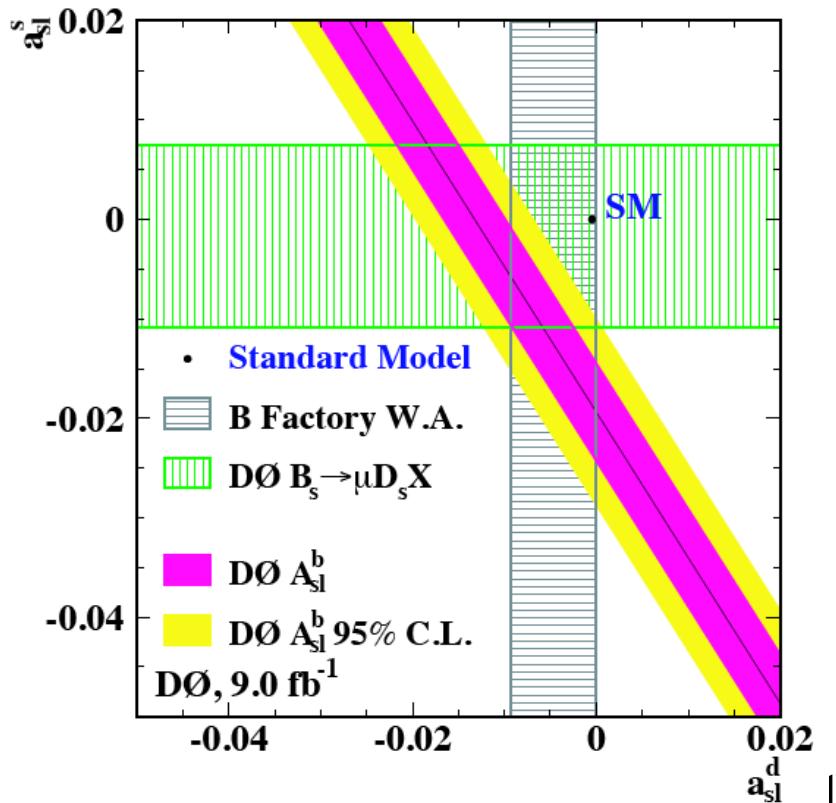
need to be controlled to better than 0.001!

- recent D0 update, arXiv:1106.6308, subm. to PRD:

$$A_{sl}^b = (-0.787 \pm 0.172 \pm 0.093)\%$$

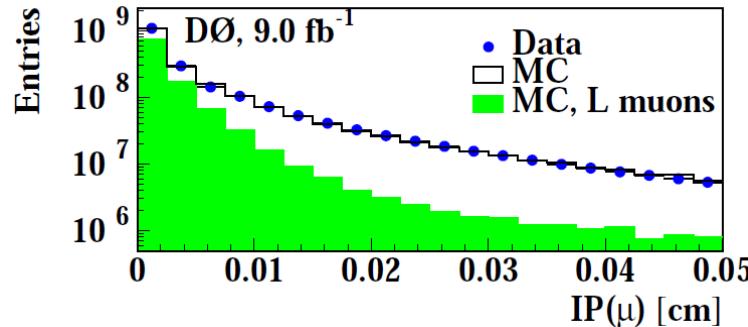
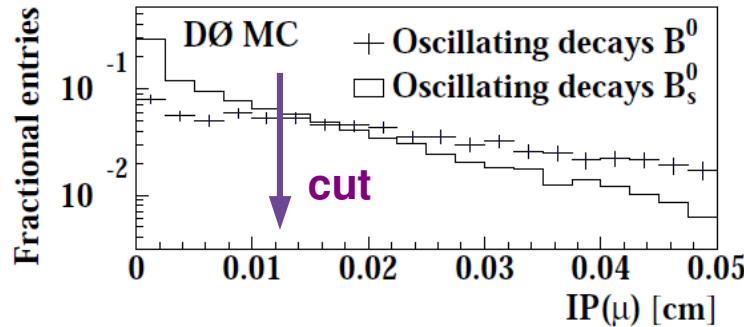
3.9 σ deviation from SM

- consistent with other measurements ... but single one of its kind
- burning question: is it B_d or B_s ?



Di-muon charge asymmetry at D0

- to probe origin, split sample by IP cut to change fraction of Bd



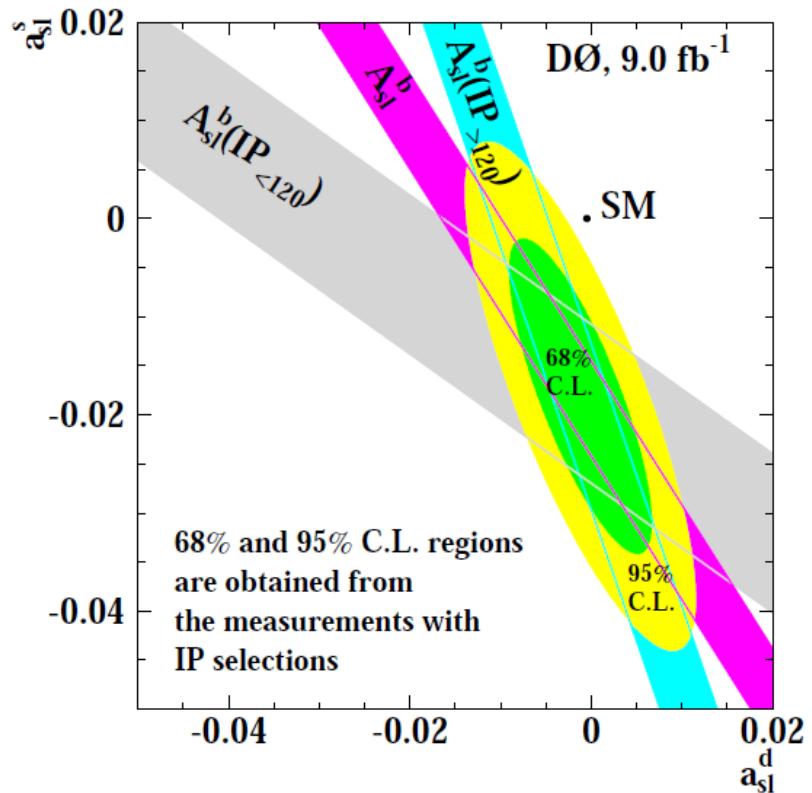
- IP>120 (both muons)
 - less background
 - higher Bd fraction
- result consistent with main result
- extract

$$a_{sl}^d = (-0.12 \pm 0.52)\%$$

$$a_{sl}^s = (-1.81 \pm 1.06)\%$$

with large correlation $\rho_{ds} = -0.799$

- uncertainties too large to be conclusive ...



LHCb plans for a_{fs}

- additional experimental challenge: production asymmetry (no problem for Bs)
- studies concentrate on $B \rightarrow D l \nu$, with fully reconstructed D
 - less background than inclusive muons
 - can distinguish Bs from Bd
 - reasonable decay time resolution --> can do time-dependent measurement
- but ... exclusive reconstruction has a price
 - to regain statistics, perform untagged measurements
 - untagged rate related to flavour-specific asymmetry by

$$a_{fs,unt}(t) = \frac{a_{fs}}{2} \left(1 - \frac{\cos(\Delta M t)}{\cosh(\Delta \Gamma t/2)} \right)$$

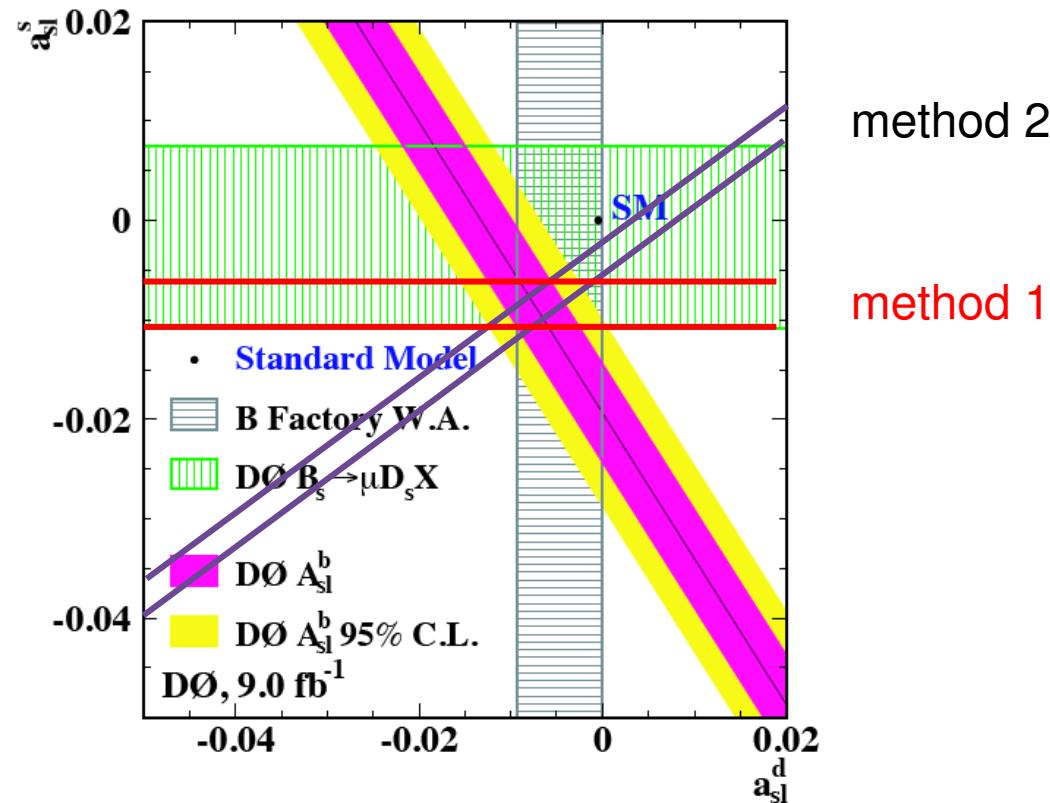
note: in contrast to a_{fs} , not constant in time

LHCb plans for a_{fs}

- two strategies to tackle production/detection asymmetries
 - 1. take only B_s , measure detection asymmetry in control samples
 - 2. measure the (time-dependent) difference

$$\Delta A = A(B_s \rightarrow D_s(KK\pi)\mu X) - A(B_d \rightarrow D(KK\pi)\mu^- X)$$

- a sketch of what this means



- actual attainable uncertainties still very unclear

Oscillations in the Bs system

- oscillation frequency places strong constraints on new physics

- SM prediction: $\Delta m_s^{\text{SM}} = 16.8^{+2.6}_{-1.5} \text{ ps}^{-1}$ PRD.83, 036004 (211)

- method: time-dependence of mixing asymmetry for a flavour-specific final state

$$A^{\text{mix}} = \frac{N^{\text{unmixed}}(t) - N^{\text{mixed}}(t)}{N^{\text{unmixed}}(t) + N^{\text{mixed}}(t)} = \cos(\Delta m t)$$

- experimental requirements

- flavour tag at 'production'
- reconstruction of decay time

$$A^{\text{mix}} = (1 - 2w) \times e^{-\Delta m^2 \sigma_t^2 / 2} \times \cos(\Delta m t)$$

 probability for wrong flavor tag

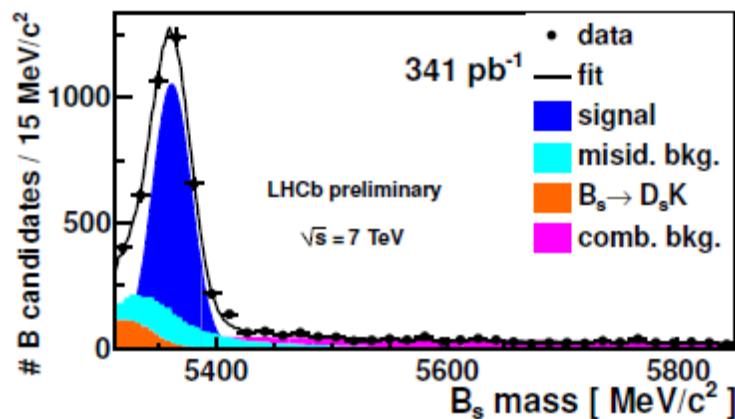
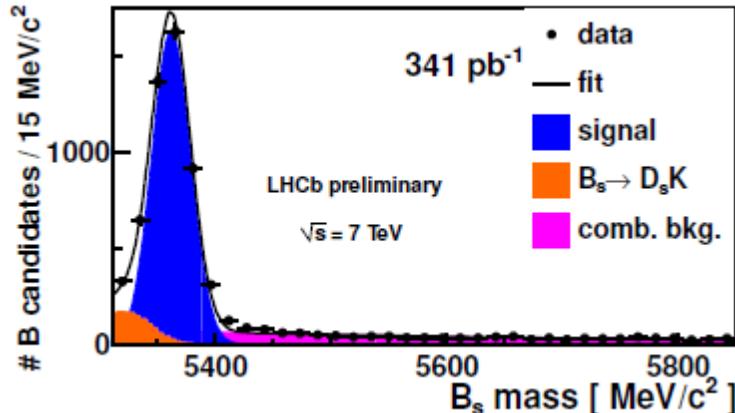
 finite decay time resolution

- optimal time resolution requires fully reconstructed final state

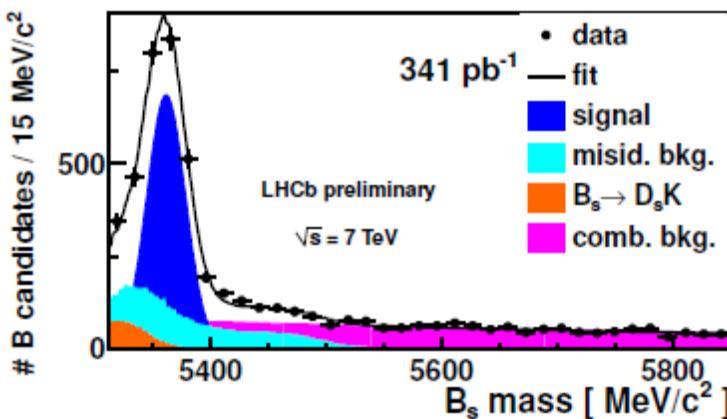
- most easily accessible: $B_s \rightarrow D_s^- \pi^+$

Bs mixing: latest measurement

- LHCb update in 341/pb, presented at Lepton-Photon 2011



decay mode	# signal candidates
$B_s^0 \rightarrow D_s^- (\phi \pi^+) \pi^+$	4371 ± 91
$B_s^0 \rightarrow D_s^- (K^* K^-) \pi^+$	2910 ± 89
$B_s^0 \rightarrow D_s^- \pi^+$ non-resonant	1908 ± 74



- average time resolution: $\sim 44 \text{ fs}$

- flavour tagger performance:

- OST: $\epsilon(1 - 2w)^2 = (3.2 \pm 0.8) \%$
- SST: $\epsilon(1 - 2w)^2 = (1.3 \pm 0.4) \%$

} reduction in signal efficiency due to loss/dilution from tagging

Bs mixing: latest measurement

- LHCb, preliminary result, 341/pb (LHCb-CONF-2011-050)

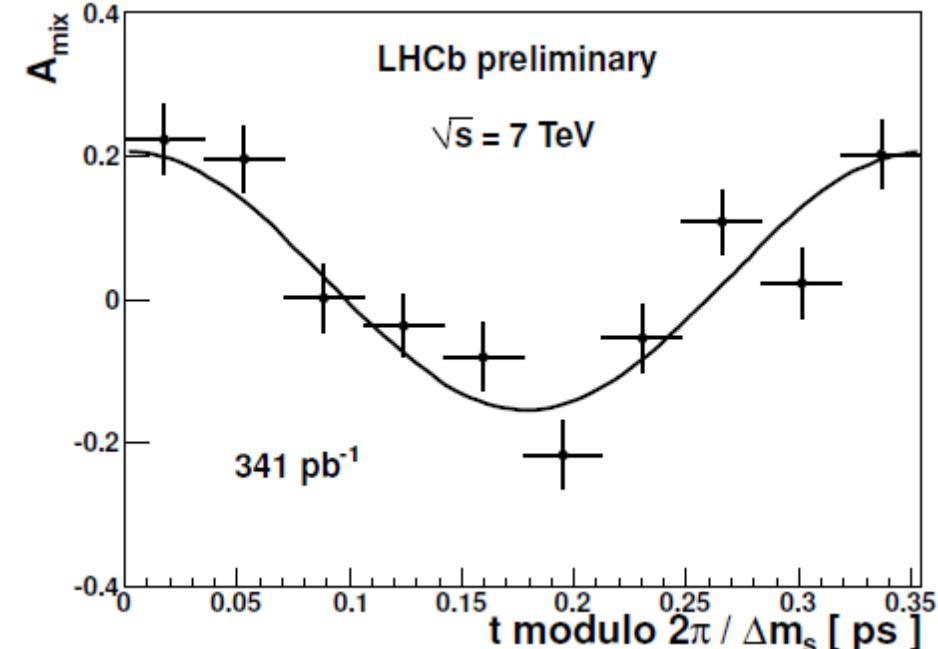
$$\Delta m_s = 17.725 \pm 0.041 \pm 0.026 \text{ ps}^{-1}$$

- dominant systematics
 - z-scale and momentum scale
 - no 'easy' improvements expected

- compare older results

CDF (2006)	$\Delta m_s = 17.77 \pm 0.10 \pm 0.07 \text{ ps}^{-1}$	(PRL97,242003 (2006)))
LHCb, 37/pb:	$\Delta m_s = 17.63 \pm 0.11 \pm 0.03 \text{ ps}^{-1}$	(LHCb-CONF-2011-005)

- new WA: $\Delta m_s^{\text{WA}} = 17.731 \pm 0.045 \text{ ps}^{-1}$



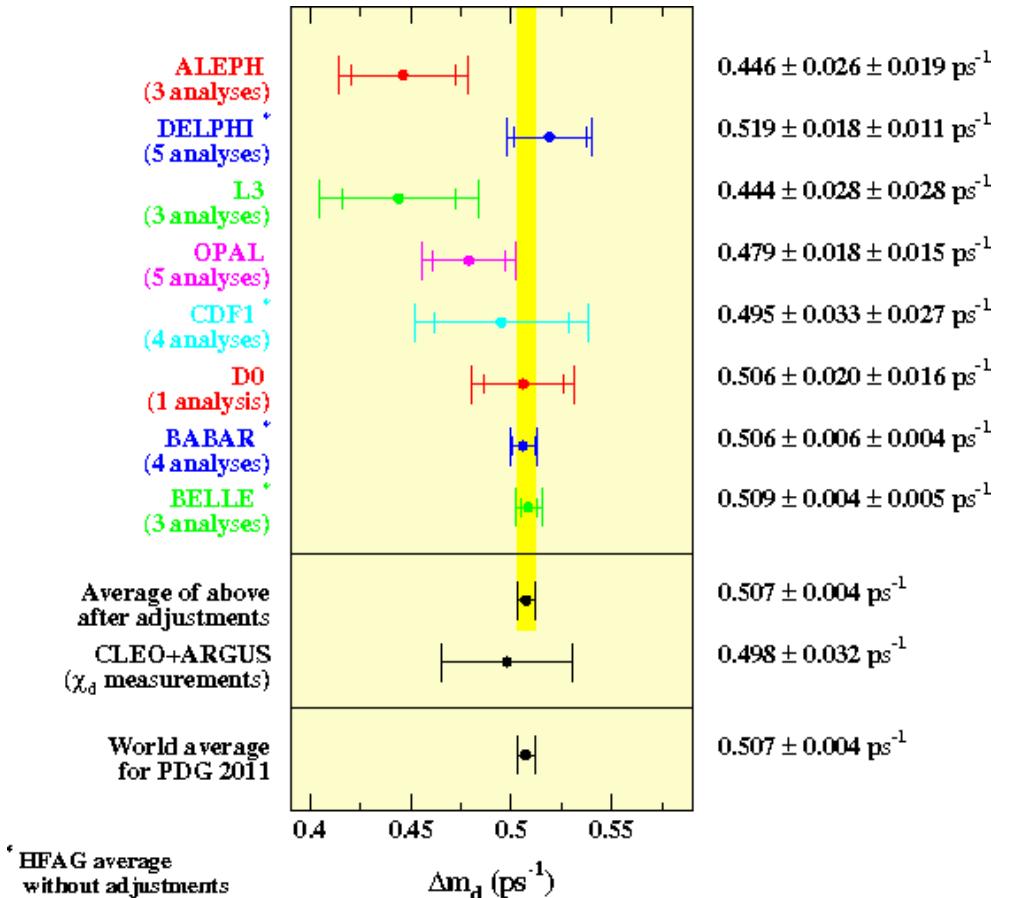
$$\left| \frac{V_{td}}{V_{ts}} \right| = 0.2090 \pm 0.0009 \pm 0.0046$$

exper. lattice

(R. van Kooten, LP2011)

Oscillations in the Bd system

- experimental precision is currently about 0.1%
- LHCb: unlike in B-s system, measurement is sensitive to decay time acceptance



- LHCb preliminary result, using Bd→D⁰ π^{\pm} in 37/pb

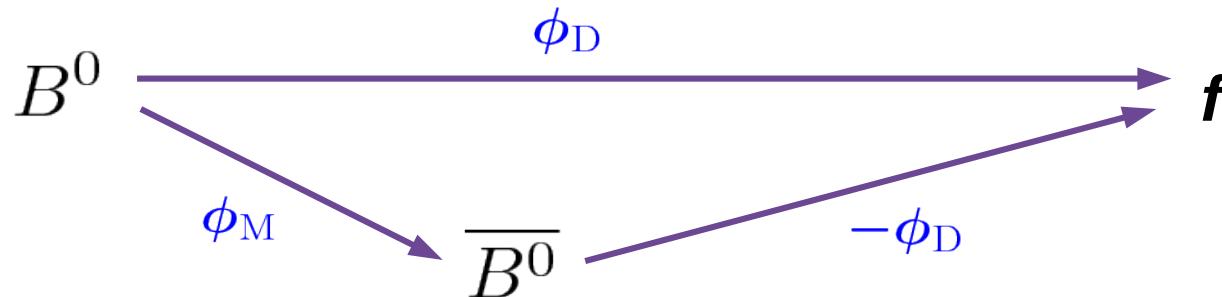
$$\Delta M_d = 0.499 \pm 0.032 \pm 0.003$$

(LHCb-CONF-2011-010, prelim.)

- first priority is to get much more stats; systematics can probably be reduced too
- Bd→D⁰*μν gives much higher yields, but systematics not clear yet

CPV in mixing via time-dependent CPV

- mixing induced CPV due to interference in decays to common final state



- if f is CP eigenstate, time dependent CP violation with pattern

- b->ccs transitions: decay dominated by TREE amplitude

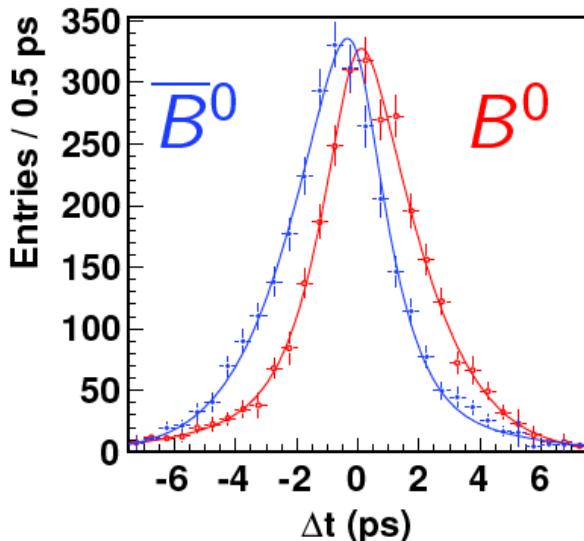
\Rightarrow expect no NP in ϕ_B \Rightarrow CPV probes mixing phase ϕ_M

- 'golden' modes: $B_d \rightarrow \psi K_s$: $S = \sin 2\beta$
 $B_s \rightarrow \psi \phi$: $S = \sin 2\beta_s$

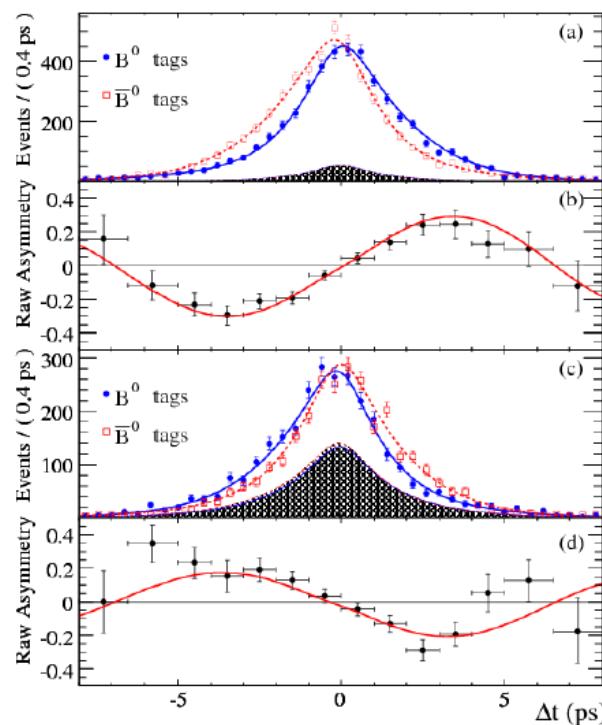
$\sin 2\beta$ from $J/\Psi K_s$, $J/\Psi K_L$ (etc)

- final Belle/Babar datasets

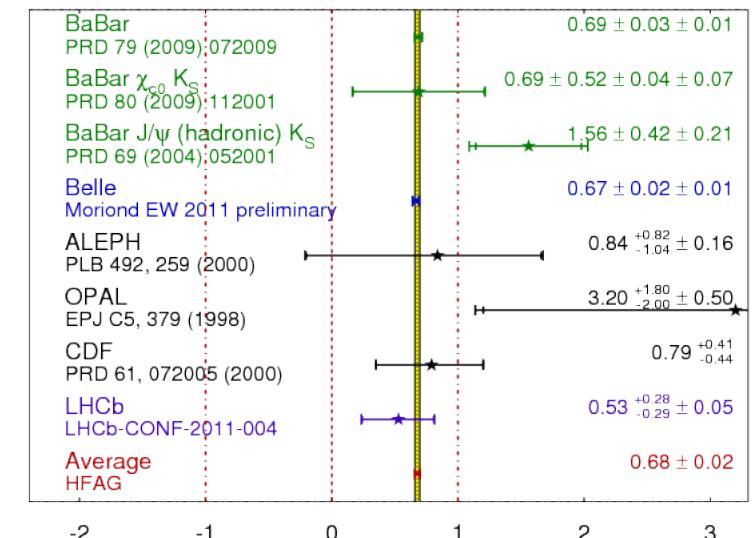
Belle, Moriond 2011



Babar, PRD 79 (2009) 072009



$$\sin(2\beta) \equiv \sin(2\phi_1) \quad \text{HFAG} \quad \text{Beauty 2011 PRELIMINARY}$$



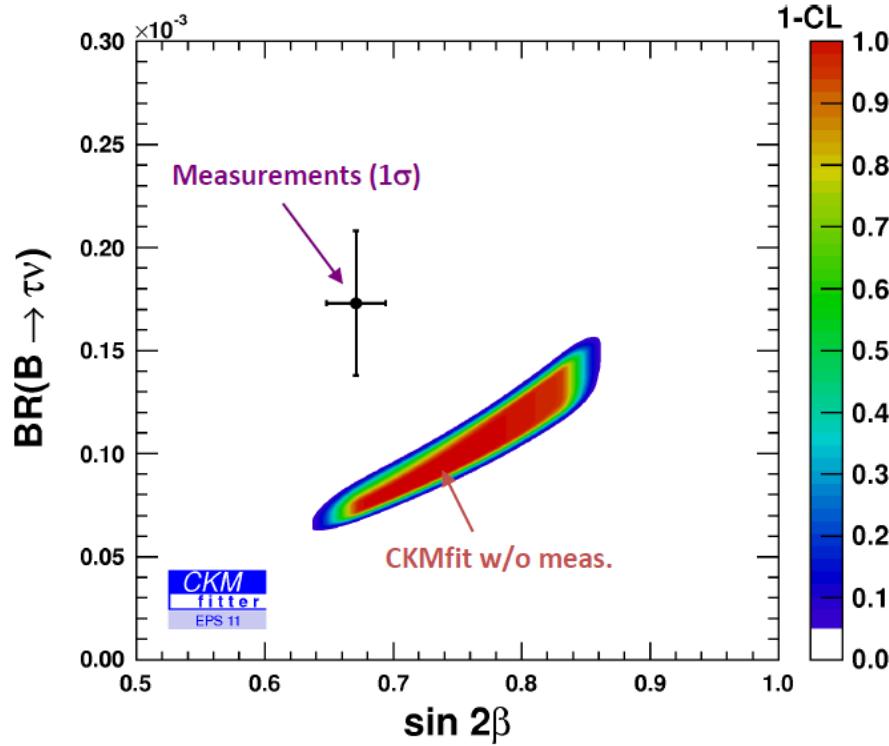
$$\text{HFAG average: } \Phi_{\psi K}^{\text{HFAG}} = (42.8 \pm 1.5)^\circ$$

has some tension with results from global SM fits

$$2\beta^{\text{SM,UTFit}} = (50.4 \pm 3.2)^\circ \quad (\text{summer 2010})$$

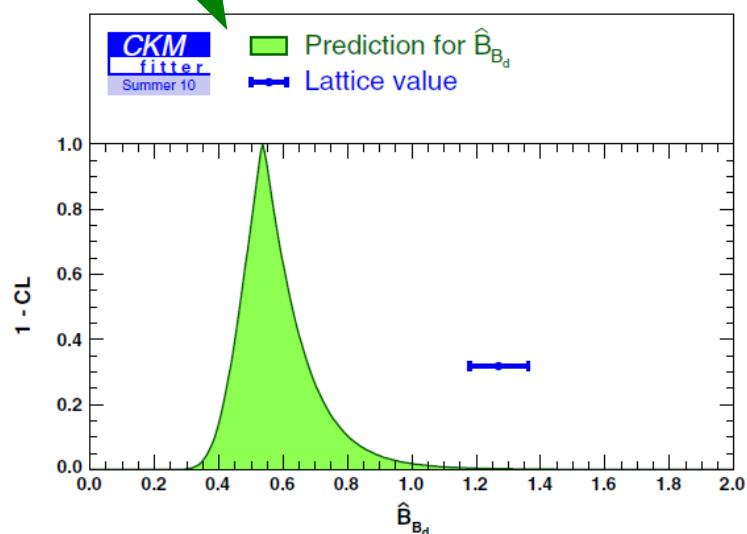
$$2\beta^{\text{SM,CKMFitter}} = (56.3^{+1.4}_{-3.3})^\circ \quad (\text{PRD83,036004})$$

- one part of the problem: correlation with the measurement of $B \rightarrow \tau \nu$ in the fits



Lenz, Nierste, CKMFitter
PRD83, 036004 (2011))

$$\frac{\mathcal{B}(B \rightarrow \tau \nu)}{\Delta m_d} = \frac{3\pi}{4} \frac{m_\tau^2}{m_W^2 S(x_t)} \left(1 - \frac{m_\tau^2}{m_{B^+}^2}\right)^2 \tau_{B^+} \times \frac{1}{\hat{\mathcal{B}}_{B_d} \eta_B} \frac{1}{|V_{ud}|^2} \left(\frac{\sin \beta}{\sin \gamma}\right)^2.$$



penguin pollution

- suppressed contributions may 'pollute' expected CPV

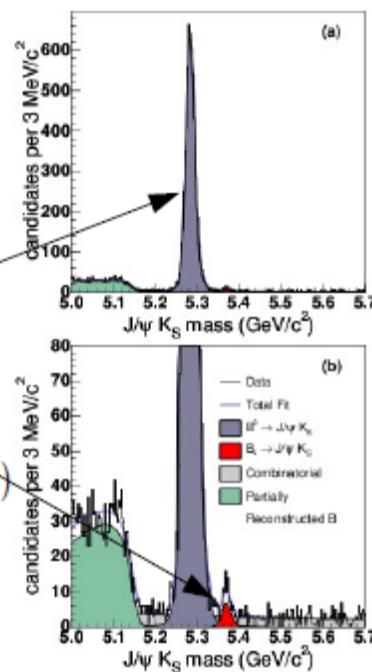
$$S_{\Psi K} = \sin(2\beta + \delta^{\text{SM penguins}} + \phi_d^{\text{NP}})$$

[need picture plus result of U-spin symmetry]

- can estimate contributions with BF of suppressed modes related by symmetry
 - Bd --> J/psi pi0 and Bs --> J/psi Ks (PRD79 (2009) 014030)

CDF experiment
PRD 83 (2011) 052012

$$\frac{B(B_s^0 \rightarrow J/\psi K_S)}{B(B^0 \rightarrow J/\psi K_S)} = 0.041 \pm 0.007 \text{ (stat)} \pm 0.004 \text{ (syst)} \pm 0.005 (f_s/f_d)$$



LHCb experiment
LHCb-CONF-2011-048

NEW

$$\frac{B(B_s^0 \rightarrow J/\psi K_S)}{B(B^0 \rightarrow J/\psi K_S)} = 0.0378 \pm 0.0058 \text{ (stat)} \pm 0.0020 \text{ (syst)} \pm 0.0030 (f_s/f_d)$$

THE UNIVERSITY OF
WARRICK
Tim Gershon
CKM Matrix Overview

(question: based on this, what are estimates of 'penguin' pollution?)

CPT and $\Delta\Gamma_d/\Gamma_d$

- introduce complex parameter for CPT violation (see e.g. PRD70(2004)012007)

$$z = \frac{\delta m - \frac{i}{2}\delta\Gamma}{\Delta m - \frac{i}{2}\Delta\Gamma} \longrightarrow \begin{aligned} |B_L\rangle &= p\sqrt{1-z}|B^0\rangle + q\sqrt{1+z}|\bar{B}^0\rangle \\ |B_H\rangle &= p\sqrt{1+z}|B^0\rangle - q\sqrt{1-z}|\bar{B}^0\rangle \end{aligned}$$

- extract from simultaneous analysis of CP and flavour specific states
- new result, Belle EPS2011 (prelim.):

$$\text{Re}(z) = (+1.9 \pm 3.7 \pm 3.2) \times 10^{-2}$$

$$\text{Im}(z) = (-5.7 \pm 3.3 \pm 6.0) \times 10^{-3}$$

$$\Delta\Gamma_d/\Gamma_d = (-1.7 \pm 1.8 \pm 1.1) \times 10^{-2}$$

$535 \times 10^6 B\bar{B}$ pairs

compatible with Babar
analysis on 88M/232M BBbar pairs

bonus: strongest single-experiment constraint on $\Delta\Gamma/\Gamma$
compare SM:

$$\frac{\Delta\Gamma_d}{\Gamma_d}^{\text{SM}} = (5.8_{-2.1}^{+1.1}) \cdot 10^{-3}$$

(my comp. using $\Delta\Gamma$ from
SM fit in PRD83,036004 (2011)
and PDG lifetime)

time-dependent CPV in the B_s -system

- time-dependent CPV in $B_s \rightarrow J/\psi \phi$ allows to NP in mixing in Bs system

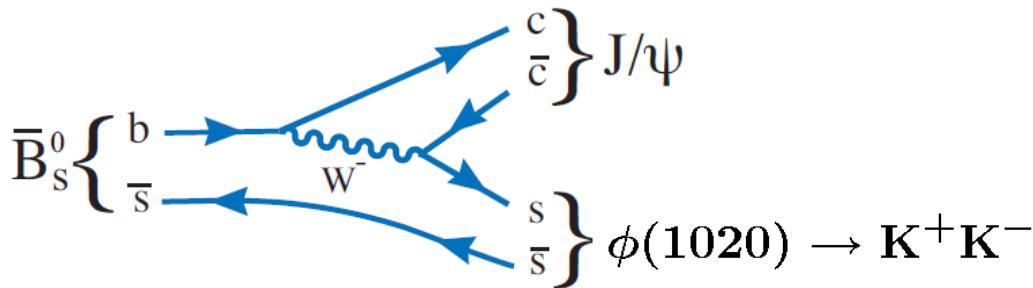
$$S_{\Psi\phi} \equiv " \sin \phi_s " = \sin(-2\beta_s^{\text{SM}} + \phi_s^{\text{NP}})$$

- in contrast to β , CKM-angle β_s is very small

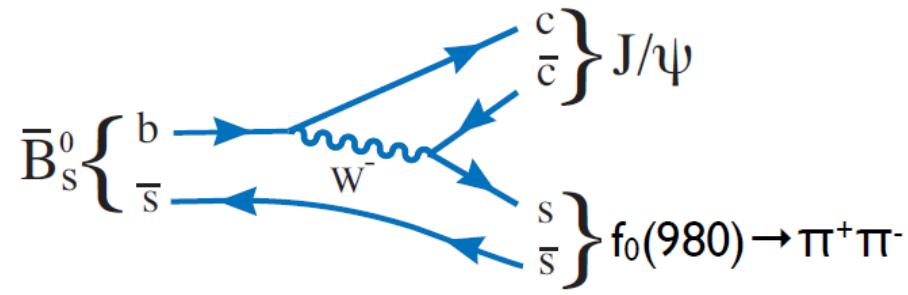
$$-2\beta_s^{\text{SM}} = (-2.08 \pm 0.10)^\circ$$

(PRD83, 036004 (2011))

- two most interesting modes



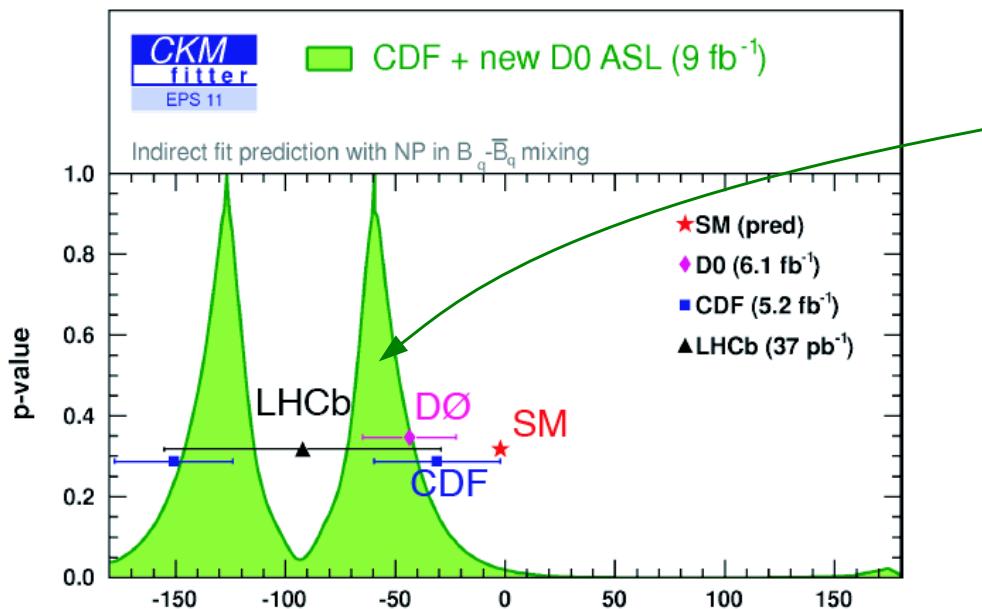
- narrow resonance --> clean
- vector-vector final state (“P-wave”)
 - requires time-dependent angular analysis
- measure also $\Delta\Gamma_s$ w/o external input



- bit lower branching fraction
- vector-pseudoscaler final state (“S-wave”)
 - no angular analysis needed

β_s with $B_s \rightarrow J/\Psi\varphi$

- status quo, before LP2011
 - CDF 5.2/fb, about 6500 events, $\sigma(\varphi_s) \sim 0.5$ rad
 - D0 9/fb, about 5000 events, $\sigma(\varphi_s) \sim 0.35$ rad
 - LHCb 37/pb, about 800 events, , $\sigma(\varphi_s) \sim 0.7$ rad



CKMFitter prediction from fit with
NP phase, using A_{sl} , but not φ_s

another hint of NP in B_s mixing?

- LHCb at LP2011: update with 341/pb
 - about 10x more statistics
 - important improvement: account for S-wave contribution

$B_s \rightarrow J/\psi\phi$ at LHCb

- including S-wave: from 6 to 10 terms in angular/time distributions

LHCb-CONF-2011-049

k	$h_k(t)$	$f_k(\theta, \psi, \varphi)$
1	$ A_0 ^2(t)$	$2\cos^2\psi(1 - \sin^2\theta\cos^2\phi)$
2	$ A_{\parallel}(t) ^2$	$\sin^2\psi(1 - \sin^2\theta\sin^2\phi)$
3	$ A_{\perp}(t) ^2$	$\sin^2\psi\sin^2\theta$
4	$\Re(A_{\parallel}(t)A_{\perp}(t))$	$-\sin^2\psi\sin 2\theta\sin 2\phi$
5	$\Re(A_0(t)A_{\parallel}(t))$	$\frac{1}{2}\sqrt{2}\sin 2\psi\sin^2\theta\sin 2\phi$
6	$\Re(A_0(t)A_{\perp}(t))$	$\frac{1}{2}\sqrt{2}\sin 2\psi\sin 2\theta\cos\phi$
7	$ A_s(t) ^2$	$\frac{2}{3}(1 - \sin^2\theta\cos^2\phi)$
8	$\Re(A_s^*(t)A_{\parallel}(t))$	$\frac{1}{3}\sqrt{6}\sin\psi\sin^2\theta\sin 2\phi$
9	$\Re(A_s^*(t)A_{\perp}(t))$	$\frac{1}{3}\sqrt{6}\sin\psi\sin 2\theta\cos\phi$
10	$\Re(A_s^*(t)A_0(t))$	$\frac{4}{3}\sqrt{3}\cos\psi(1 - \sin^2\theta\cos^2\phi)$

The terms 7–10 are related to the description of the S-wave component, which has been added to this analysis. Expressed in terms of the size $|A_i(0)|$ and phase δ_i of the transversity and S-wave amplitudes at $t = 0$, the time dependent amplitudes are given by

$$|A_0|^2(t) = |A_0|^2 e^{-\Gamma_s t} [\cosh\left(\frac{\Delta\Gamma}{2}t\right) - \cos\phi_s \sinh\left(\frac{\Delta\Gamma}{2}t\right) + \sin\phi_s \sin(\Delta m_s t)], \quad (4)$$

$$|A_{\parallel}(t)|^2 = |A_{\parallel}|^2 e^{-\Gamma_s t} [\cosh\left(\frac{\Delta\Gamma}{2}t\right) - \cos\phi_s \sinh\left(\frac{\Delta\Gamma}{2}t\right) + \sin\phi_s \sin(\Delta m_s t)], \quad (5)$$

$$|A_{\perp}(t)|^2 = |A_{\perp}|^2 e^{-\Gamma_s t} [\cosh\left(\frac{\Delta\Gamma}{2}t\right) + \cos\phi_s \sinh\left(\frac{\Delta\Gamma}{2}t\right) - \sin\phi_s \sin(\Delta m_s t)], \quad (6)$$

$$\Re(A_{\parallel}(t)A_{\perp}(t)) = |A_{\parallel}||A_{\perp}|e^{-\Gamma_s t} [-\cos(\delta_{\perp} - \delta_{\parallel}) \sin\phi_s \sinh\left(\frac{\Delta\Gamma}{2}t\right) - \cos(\delta_{\perp} - \delta_{\parallel}) \cos\phi_s \sin(\Delta m_s t) + \sin(\delta_{\perp} - \delta_{\parallel}) \cos(\Delta m_s t)], \quad (7)$$

$$\Re(A_0(t)A_{\parallel}(t)) = |A_0||A_{\parallel}|e^{-\Gamma_s t} \cos(\delta_{\parallel} - \delta_0) [\cosh\left(\frac{\Delta\Gamma}{2}t\right) - \cos\phi_s \sinh\left(\frac{\Delta\Gamma}{2}t\right) + \sin\phi_s \sin(\Delta m_s t)], \quad (8)$$

$$\Re(A_0(t)A_{\perp}(t)) = |A_0||A_{\perp}|e^{-\Gamma_s t} [-\cos(\delta_{\perp} - \delta_0) \sin\phi_s \sinh\left(\frac{\Delta\Gamma}{2}t\right) - \cos(\delta_{\perp} - \delta_0) \cos\phi_s \sin(\Delta m_s t) + \sin(\delta_{\perp} - \delta_0) \cos(\Delta m_s t)], \quad (9)$$

$$|A_s(t)|^2 = |A_s|^2 e^{-\Gamma_s t} [\cosh\left(\frac{\Delta\Gamma}{2}t\right) + \cos\phi_s \sinh\left(\frac{\Delta\Gamma}{2}t\right) - \sin\phi_s \sin(\Delta m_s t)], \quad (10)$$

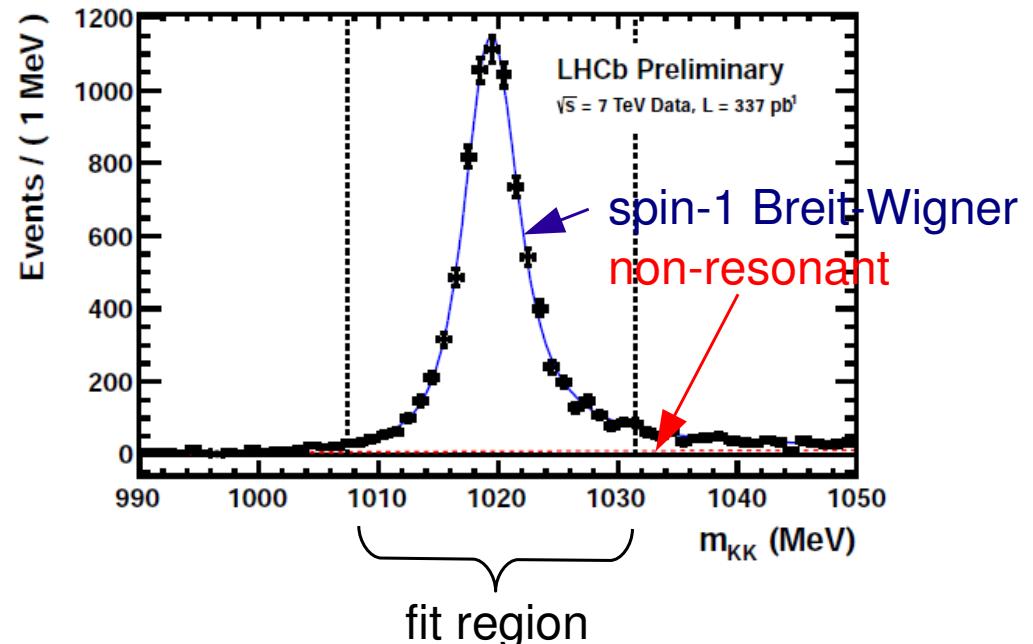
$$\Re(A_s^*(t)A_{\parallel}(t)) = |A_s||A_{\parallel}|e^{-\Gamma_s t} [-\sin(\delta_{\parallel} - \delta_s) \sin\phi_s \sinh\left(\frac{\Delta\Gamma}{2}t\right) - \sin(\delta_{\parallel} - \delta_s) \cos\phi_s \sin(\Delta m_s t) + \cos(\delta_{\parallel} - \delta_s) \cos(\Delta m_s t)], \quad (11)$$

$$\Re(A_s^*(t)A_{\perp}(t)) = |A_s||A_{\perp}|e^{-\Gamma_s t} \sin(\delta_{\perp} - \delta_s) [\cosh\left(\frac{\Delta\Gamma}{2}t\right) + \cos\phi_s \sinh\left(\frac{\Delta\Gamma}{2}t\right) - \sin\phi_s \sin(\Delta m_s t)], \quad (12)$$

$$\Re(A_s^*(t)A_0(t)) = |A_s||A_0|e^{-\Gamma_s t} [-\sin(\delta_0 - \delta_s) \sin\phi_s \sinh\left(\frac{\Delta\Gamma}{2}t\right) - \sin(\delta_0 - \delta_s) \cos\phi_s \sin(\Delta m_s t) + \cos(\delta_0 - \delta_s) \cos(\Delta m_s t)]. \quad (13)$$

where we have chosen a phase convention such that $\delta_0 = 0$. The decay time dependent decay rates for an initial B_s^0 decaying to $J/\psi\phi$ can be obtained from those above by inserting a factor -1 in front of the terms involving mixing ($\sin(\Delta m_s t)$ and $\cos(\Delta m_s t)$).

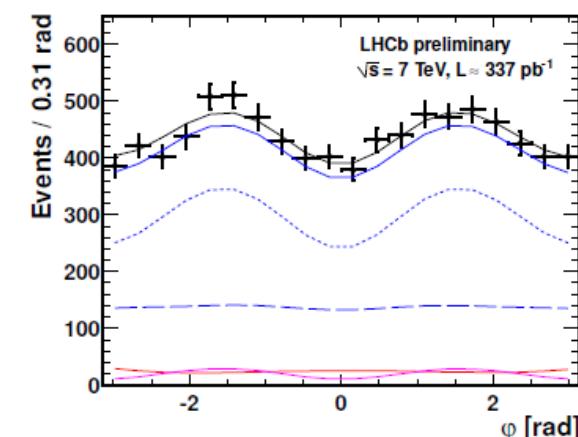
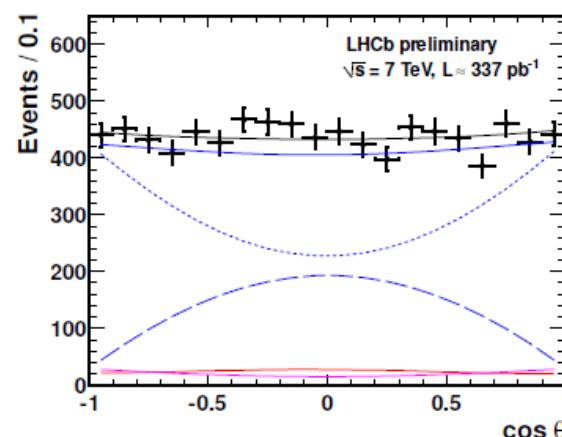
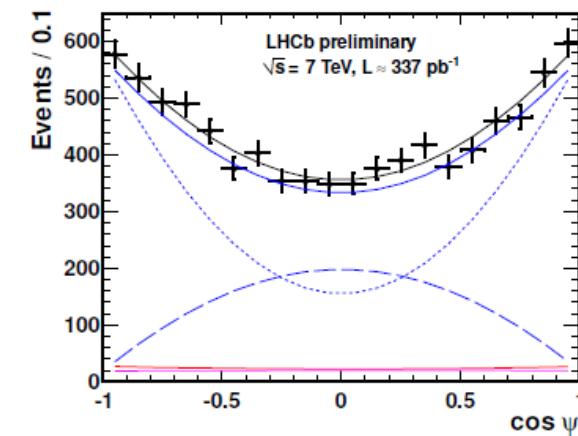
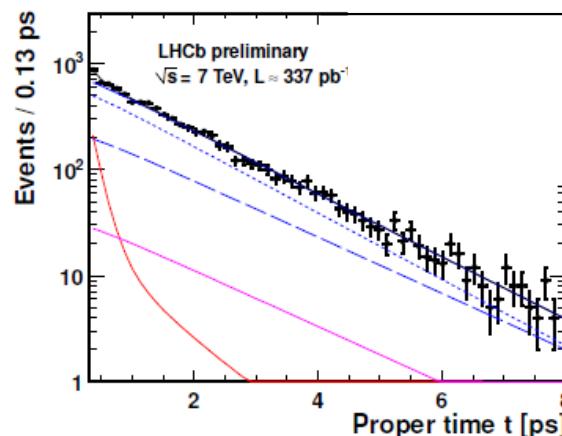
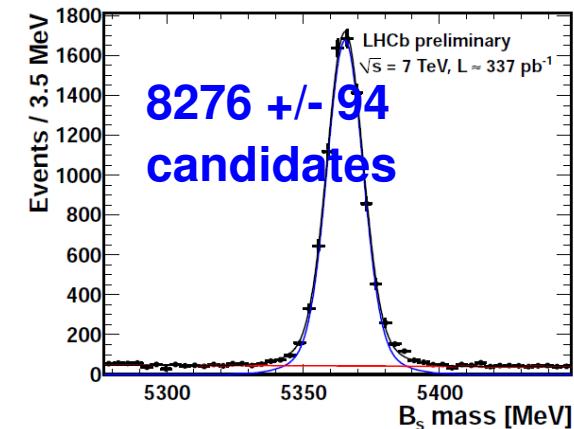
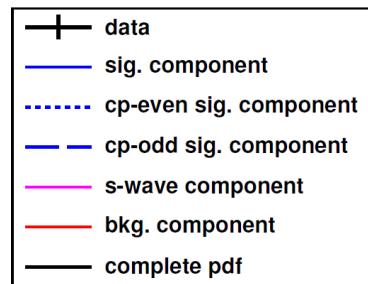
accounts for ~4% “non-resonant” KK in 12 MeV mass window around phi



note: S-wave contribution identified by angular distribution, not by KK mass

ML fit to LHCb the data

- ML fit with 10 physics parameters
 - 7 angular amplitudes and phases
 - Γ_s , $\Delta\Gamma_s$, ϕ_s
- proper-time resolution, calibrated on prompt J/psi gives $\sigma(t) \sim 50$ ps
- only OS flavour tagging used, calibrated on J/psiK+
 $\varepsilon_{\text{tag}} D^2 = (2.08 \pm 0.41)\%$
- goodness of fit, using “point-to-point dissimilarity test” (*) gives P-value of 0.44

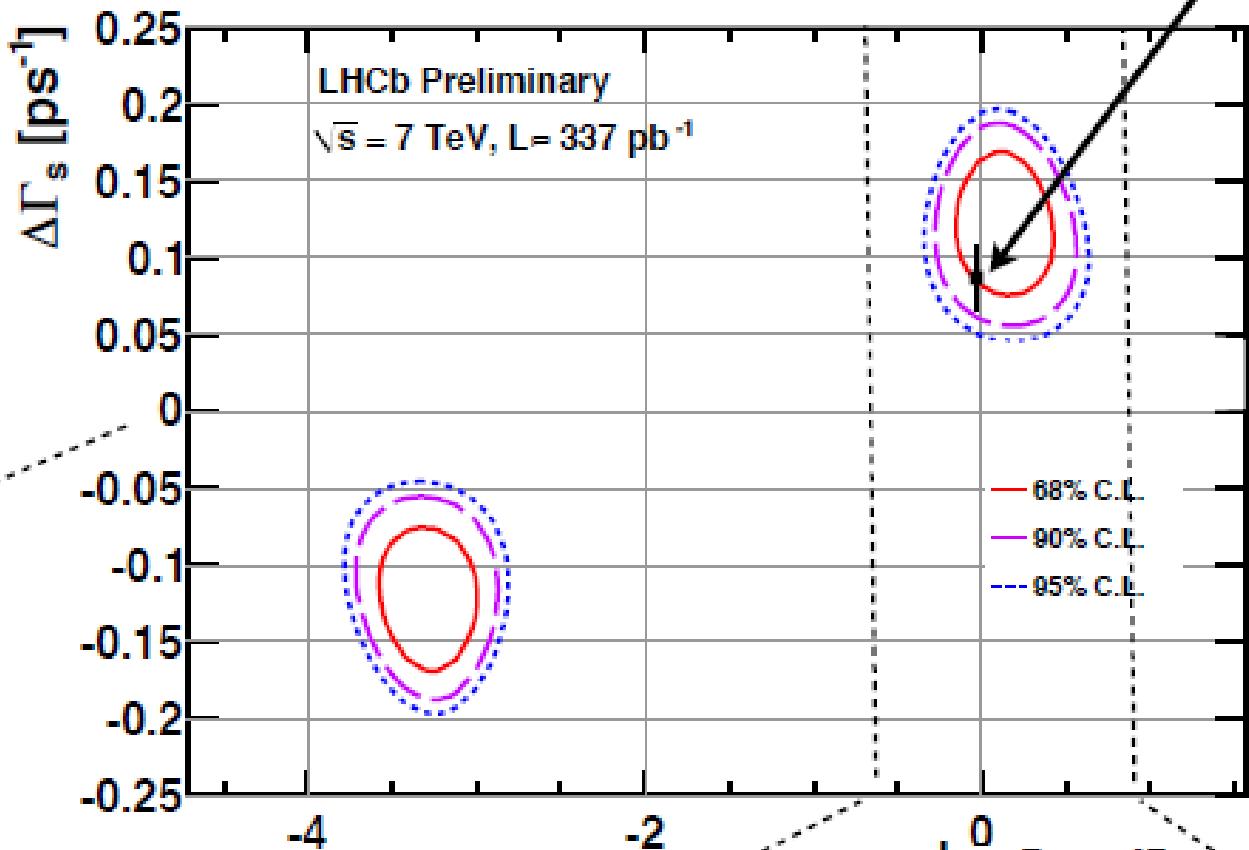
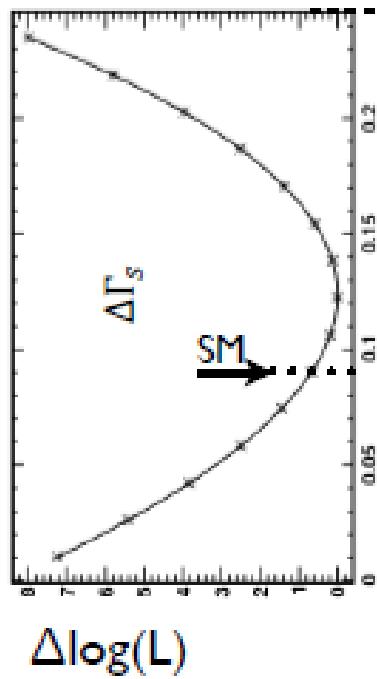


(*) see eg. M. Williams, JINST 5 (2010) P09004
[\[arXiv:1006.3019 \[hep-ex\]\]](https://arxiv.org/abs/1006.3019)

$B_s \rightarrow J/\Psi \varphi$: $\Delta\Gamma_s$ vs. ϕ_s

Standard Model
(Lenz, Nierste; arXiv:1102.4274)

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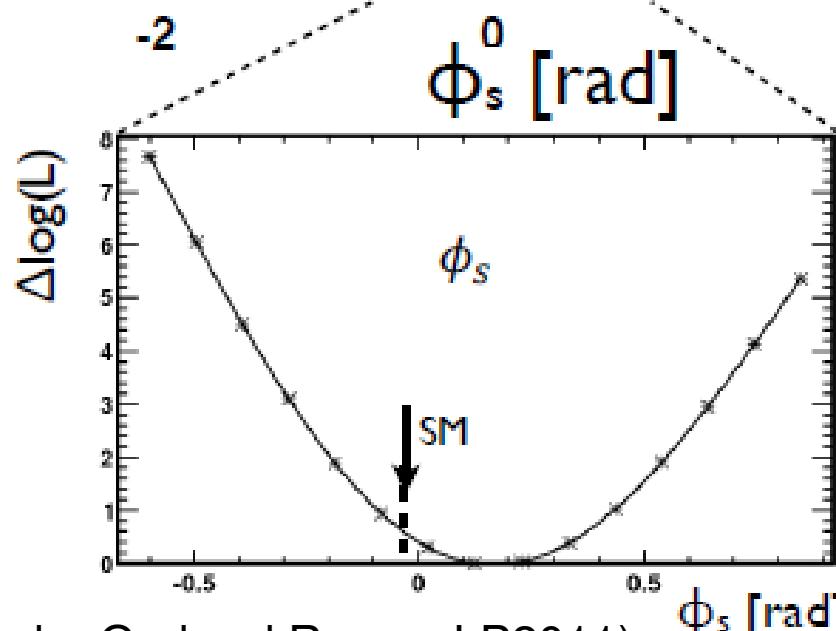


Most precise measurement of ϕ_s

- $\phi_s = 0.13 \pm 0.18$ (stat) ± 0.07 (syst) rad
- Consistent with SM

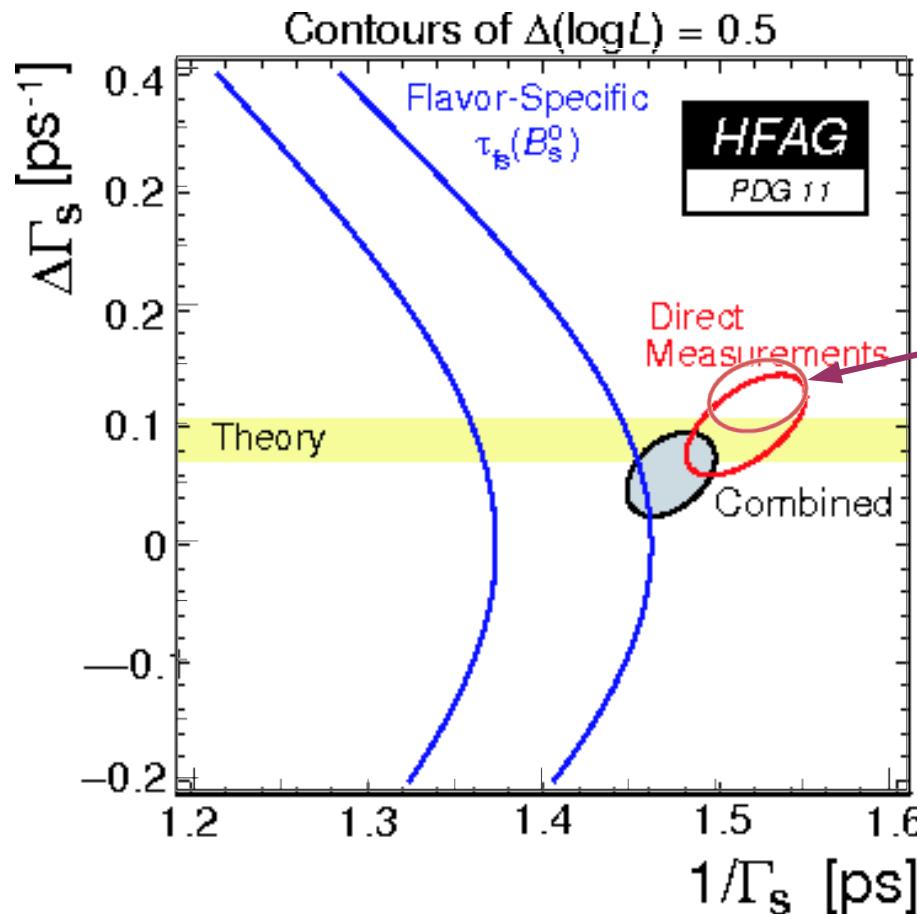
4 σ Evidence for $\Delta\Gamma_s \neq 0$:

- $\Delta\Gamma_s = 0.123 \pm 0.029$ (stat) ± 0.008 (syst) ps⁻¹
- $\Gamma_s = 0.656 \pm 0.009$ (stat) ± 0.008 (syst) ps⁻¹



$\Delta\Gamma_s$

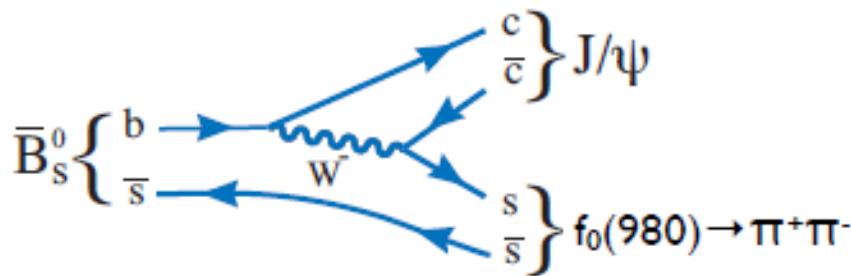
- J/psi phi analysis also gives most precise single measurement of Γ_s and $\Delta\Gamma_s$



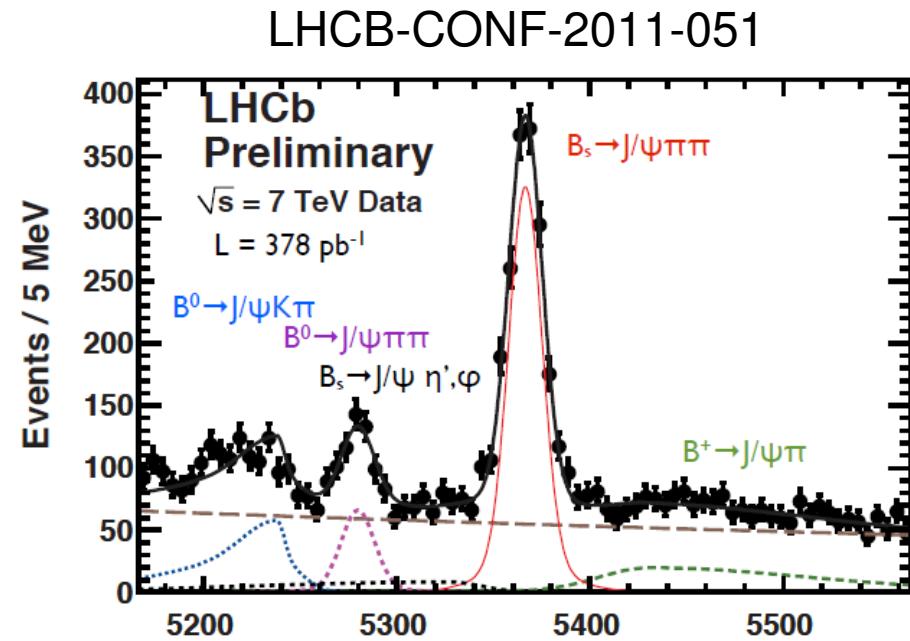
$$\tau(B_s^0)_{\text{fs}} = \frac{1}{\Gamma_s} \frac{1 + \left(\frac{\Delta\Gamma_s}{2\Gamma_s} \right)^2}{1 - \left(\frac{\Delta\Gamma_s}{2\Gamma_s} \right)^2}$$

slight inconsistency in Γ between direct measurement and measurement with flavour specific decays?

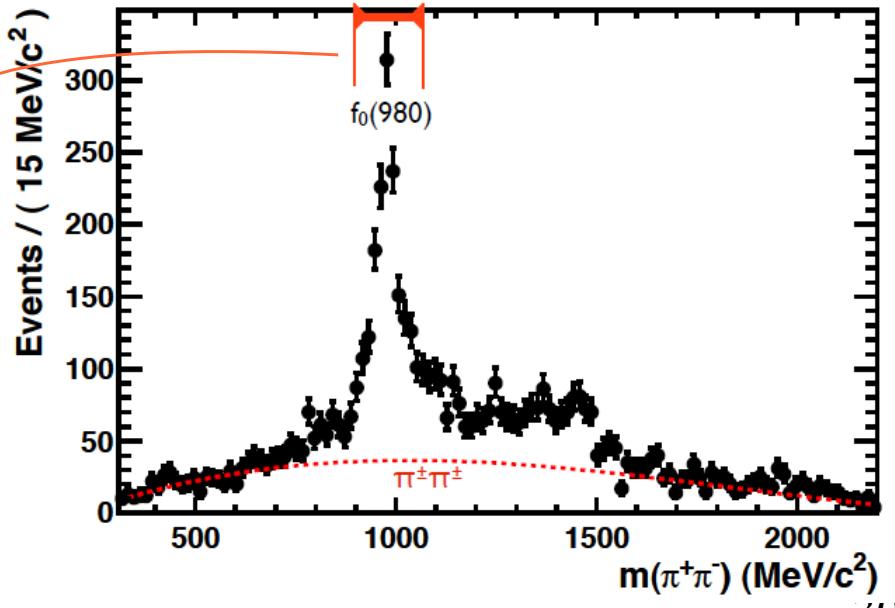
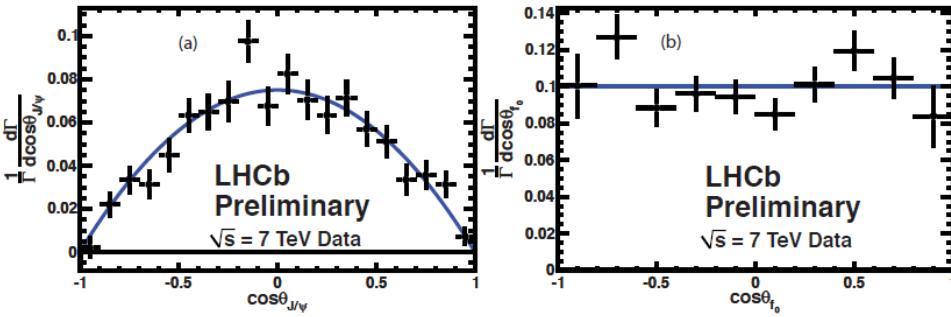
$B_s \rightarrow J/\psi f^0$ at LHCb



- some history
 - predicted by Stone, Zhang (2009)
 - first seen by LHCb (PLB689(2011)115)
 - LP2011: first measurement of CPV

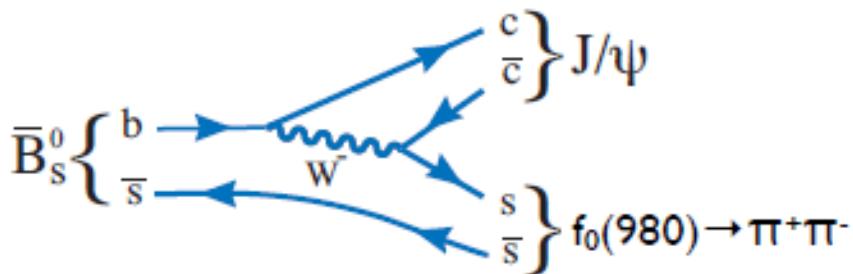


- angular distribution of psi and f0 candidates

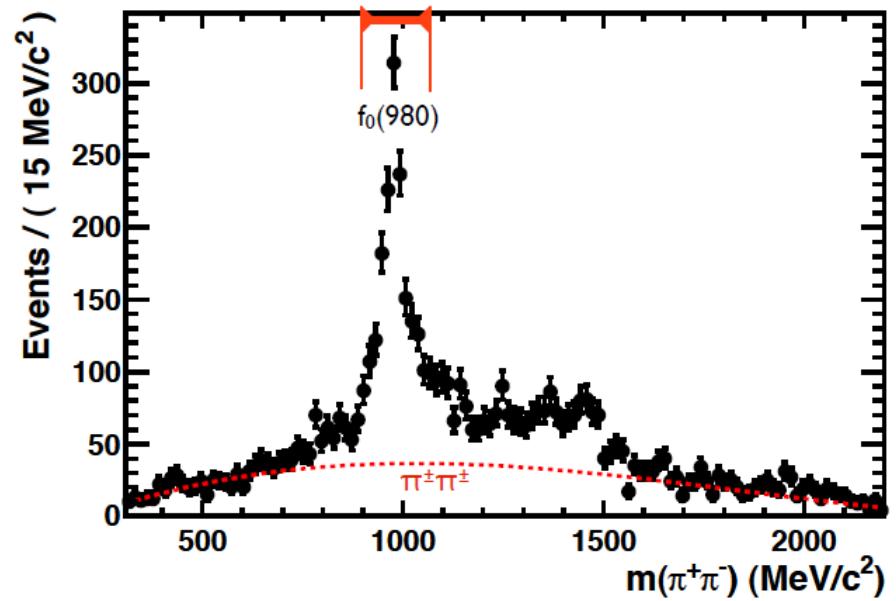
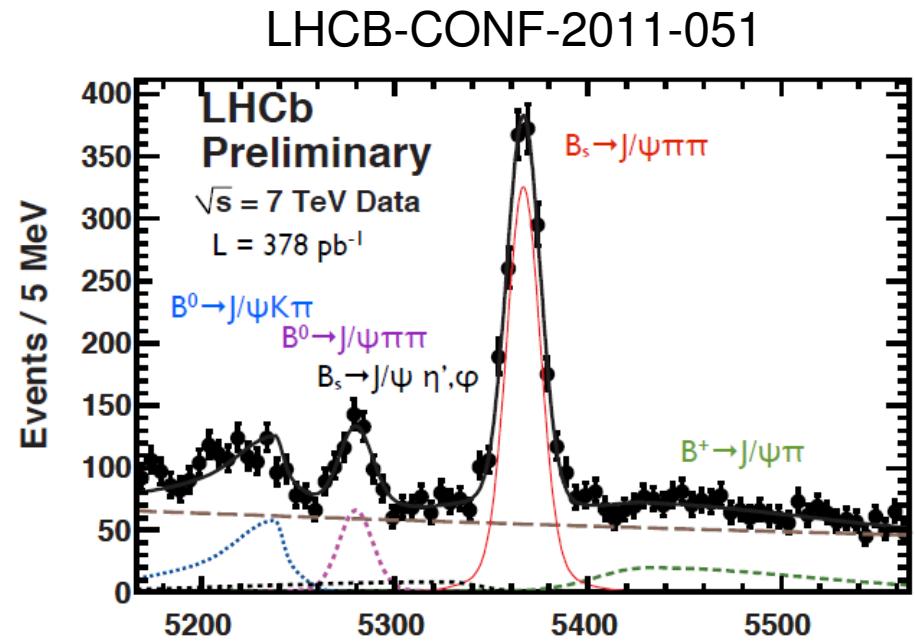
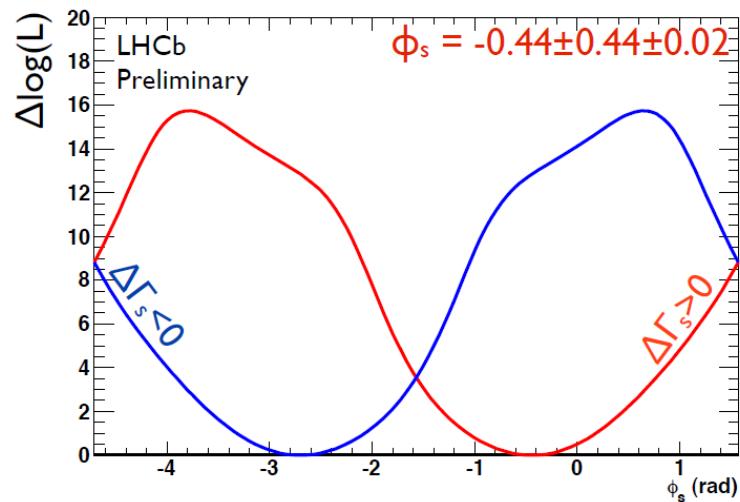


candidates in f(980) region look pure scalar
 --> no angular analysis needed

$B_s \rightarrow J/\psi f^0$ at LHCb



- some history
 - predicted by Stone, Zhang (2009)
 - first seen by LHCb (PLB689(2011)115)
 - LP2011: first measurement of CPV
- result from the LL fit

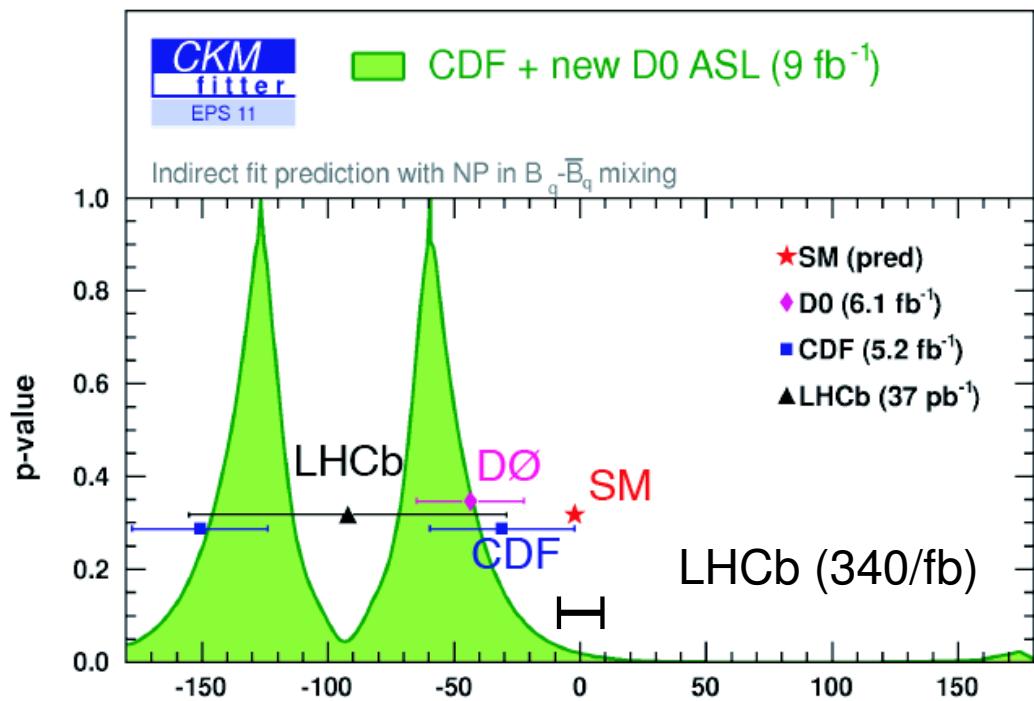


- beware: hadronic uncertainties, e.g. arXiv:1109.1112 [hep-ph]

combination of J/psi phi and J/psi f0

- $B_s \rightarrow J/\psi f_0$ alone cannot constraint both Γ_s , $\Delta\Gamma_s$ and ϕ_s
 - requires combination with $J/\psi \phi$
- simultaneous fit to both samples
$$\phi_s = 0.03 \pm 0.16 \pm 0.07$$

(prelim, LHCb-CONF-2011-056)
- TDCPV gives no evidence (yet) for NP in B_s mixing
- next steps
 - add more data
 - use same-side Kaon tagger
 - resolve two-fold ambiguity by relative S-wave strong phase in bins of $M(KK)$ (Y. Xie et al., JHEP 0909:074, (2009))



conclusions/summary

- precision measurements of mixing observables:
 - mixing frequencies
 - CP asymmetries in final states with no expected CPV from decay amplitude
 - time-dependence CP asymmetries to CP eigenstates dominates by tree diagram
- one tantalizing result: D0 A_SL is 3.9s from SM prediction
- LHCb experiment is working on measurements in all these areas

backup

how about a_{fs}^d ?

- current HFAG average

$$\boxed{\begin{aligned} |\mathbf{q}/\mathbf{p}| &= \mathbf{1.0002 \pm 0.0028} \\ A_{SL} &= \mathbf{-0.0005 \pm 0.0056} \end{aligned}}$$

- many measurements, most constraining:

Babar, 232M BBbar

$$|q/p| - 1 = (-0.8 \pm 2.7(\text{stat.}) \pm 1.9(\text{syst.})) \times 10^{-3},$$

B. Aubert *et al.* (BaBar Collab.), Phys. Rev. Lett. **96**, 251802 (2006).

$$A_{sl} = (-1.1 \pm 7.9(\text{stat}) \pm 8.5(\text{sys})) \times 10^{-3},$$

Belle, 87M BBbar

$$|q/p| = 1.0005 \pm 0.0040(\text{stat}) \pm 0.0043(\text{sys}).$$

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- B factories still have quite a bit more data ...