

# The anomalous magnetic moment of $\tau$ lepton and its leptonic radiative decays

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work in progress in collaboration with:

S. Eidelman, D. Epifanov, L. Mercolli, C. Ng, M. Passera.

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# Outline

- 1 The SM prediction of the tau  $g - 2$ .
- 2 The tau  $g - 2$ : experimental bounds.
- 3 Tau Radiative Decays and  $g - 2$

## Anomalous Magnetic Moment

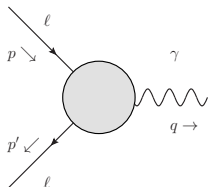
The magnetic moment for a particle of spin  $\vec{S}$  is

$$\vec{\mu} = g \frac{e}{2m} \vec{S}.$$

For a spin-1/2 fermion:

$$(i\partial_\mu - eA_\mu)\gamma^\mu\psi = m\psi \implies g = 2. \quad \text{P. Dirac '28}$$

The  $g$ -factor is modified by quantum correction:  $a = (g - 2)/2$ .



$$\bar{u}(p')\Gamma^\mu u(p) = \bar{u}(p') \left\{ F_1(q^2)\gamma^\mu + F_2(q^2)\frac{i\sigma^{\mu\nu}}{2m_\ell}q_\nu + \dots \right\} u(p)$$

$$a_\ell = F_2|_{\text{on-shell}}$$

The leading contribution  $a_\ell = \alpha/(2\pi)$

J. Schwinger '48

## Leptons Anomalous Magnetic Moment

The present experimental values:

$$a_e = 1\,159\,652\,180.73(28) \cdot 10^{-12}$$

0.24 parts per billion!

Hanneke et al, PRL100 (2008) 120801

$$a_\mu = 1\,165\,920\,89(63) \cdot 10^{-11}$$

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E821 – Final Report: PRD73 (2006) 072003

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$$-0.052 \leq a_\tau \leq 0.013$$

DELPHI - EPJC35 (2004) 159

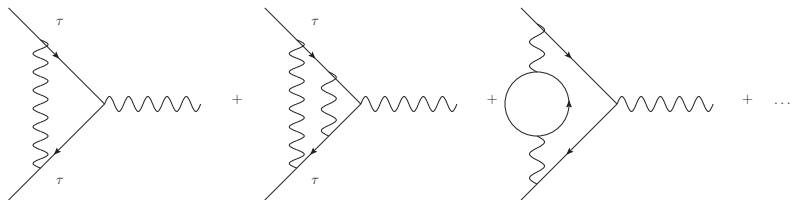
This is not even a test of the leading contribution!

$$a_\tau = \frac{\alpha}{2\pi} + \mathcal{O}(\alpha^2) \approx 0.00116\dots \quad \text{J. Schwinger '48}$$

## The SM Prediction

The Standard Model prediction of the tau  $g - 2$  is:

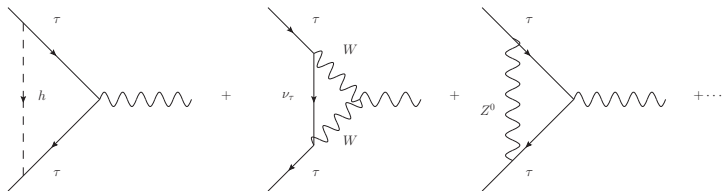
$$\begin{aligned}
 a_{\tau}^{\text{SM}} = & \mathbf{117\,324\,(2)} && \times 10^{-8} && \text{QED} \\
 & + 47.4\,(5) && \times 10^{-8} && \text{EW} \\
 & + 337.5\,(3.7) && \times 10^{-8} && \text{HLO} \\
 & + 7.6\,(2) && \times 10^{-8} && \text{HHO (vac)} \\
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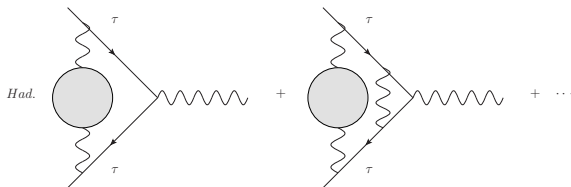
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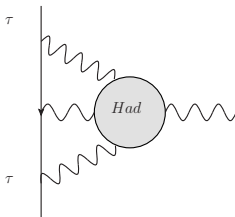




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$$a_{\tau}^{\text{SM}} = 117\,721(5) \times 10^{-8}$$

S. Eidelman, M. Passera, MPLA 22 (2007) 159

## Why is $a_\tau$ important?

- The tau  $g - 2$  is a property of an elementary particle.
- NP associated with a mass scale  $\Lambda$  is expected to modify the  $g - 2$  of a lepton  $\ell$  by a contribution  $a_\ell^{\text{NP}} \approx m_\ell^2 / \Lambda^2$ .
- Given  $m_\tau^2 / m_\mu^2 \sim 283$ ,  $a_\tau$  is much more sensitive than the muon one to EW and NP loop effects which give contributions  $\sim m_\ell^2$ .
- If an NP effect were of the same order of magnitude as that of the EW,  $a_\tau$  would provide a “clean” NP test.

	Electron	Muon	Tau
$a^{\text{EW}} / a^{\text{HAD}}$	1/56	1/45	1/7
$a^{\text{EW}} / \delta a^{\text{HAD}}$	1.6	3	10

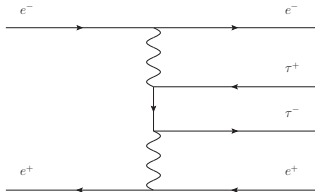
...if only we could measure it!

## Experimental Bounds

- The very short lifetime of the tau ( $2.9 \cdot 10^{-13}$  s) makes it impossible to put  $\tau$  in a storage ring to measure  $a_\tau$  via its spin precession.
- DELPHI's result, from  $e^+ e^- \rightarrow e^+ e^- \tau^+ \tau^-$  total cross-section measurements at LEP 2 (the PDG value):

$$-0.052 \leq a_\tau \leq 0.013 \quad \text{DELPHI - EPJC35 (2004) 159}$$

- González-Sprinberg et al. analysis of LEP1/2 and SLC  $e^+ e^- \rightarrow \tau^+ \tau^-$  data, using effective Lagrangian approach:  
 $-0.007 < a_\tau^{\text{NP}} < 0.005$  (95% CL) González-Sprinberg et al. NPB 582 (2000) 3

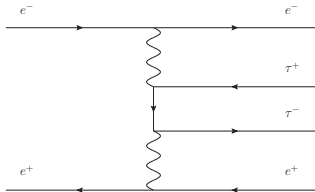


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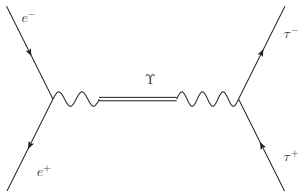
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Bernabéu et al. propose the measurement of  $F_2(q^2 = M_\gamma^2)$  from  $e^+ e^- \rightarrow \tau^+ \tau^-$  production at B factories.



$\implies$  bounds on  $F_2(q^2 = M_\gamma^2)$

Expected sensitivity for  $F_2(q^2 = M_\gamma^2)$  at Super B:  $\sim 10^{-6}$

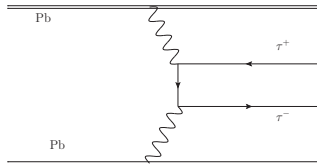
Bernabéu et al. NPB 790 (2008) 160.

Possibility to use heavy ion collision at LHC:  $\text{Pb Pb} \rightarrow \text{Pb Pb} \tau^+ \tau^-$ .

Advantage:  $\text{Pb Pb } \gamma\gamma \approx \text{on-shell}$ .

Expected sensitivity  $\sim 10^{-5}$ .

F. del Aguila et al, PLB 271 (1991) 256



## Radiative Decays of the Tau

What about  $\tau^- \rightarrow \ell^- \bar{\nu}_\ell \nu_\tau \gamma$ ?

It was suggested to search for the  $a_\tau$  in radiative leptonic  $\tau$  decays using the phenomenon of radiation zero: the tree-level amplitude (which corresponds to  $a_\tau = 0$ ) vanishes in the phase space region

$$\cos(\ell, \gamma) = -1, \quad x = \frac{2E_\ell}{m_\tau} = 1 + \frac{m_\ell^2}{m_\tau^2} \quad (\text{in the } \tau \text{ r.f.})$$

M. L. Laursen et al. PRD 29 (1984) 2652

Our aim is to:

- revisit  $\tau$  radiative decays,
- study the dependence of the pol. diff. decay rate on  $a_\tau$ ,
- probe  $a_\tau$  at  $\mathcal{O}(10^{-3})$  with minimal assumption,

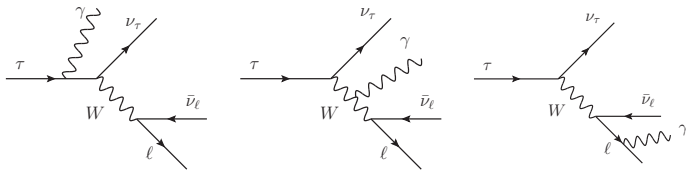
No search for BSM physics!

Discrepancy of  $a_\mu$  scaled as  $\frac{m_\tau^2}{m_\mu^2}$ :  $\sim 10^{-6}$  needed.

# Effective Lagrangian Approach

Add a minimal  $g - 2$  coupling:

$$\mathcal{L}_{\text{eff}} = e \frac{\tilde{a}}{4m_\tau} \bar{\tau} \sigma^{\mu\nu} \tau F_{\mu\nu}.$$

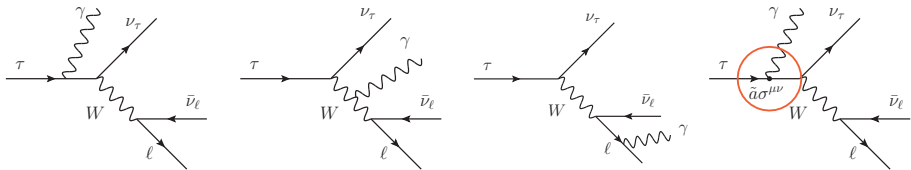




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$$\mathcal{L}_{\text{eff}} = e \frac{\tilde{a}}{4m_\tau} \bar{\tau} \sigma^{\mu\nu} \tau F_{\mu\nu}.$$



The new term contributes to the tree level decay rate:

$$d\Gamma(\tau^- \rightarrow \ell^- \bar{\nu}_\ell \nu_\tau \gamma) = d\Gamma_{\text{SM}} + \tilde{a} d\Gamma_a.$$

Terms  $\mathcal{O}(\tilde{a}^2)$  are neglected since of order  $10^{-6}$ .

## The $\tilde{a}$ Parameter

- Since we want to probe  $a_\tau$  at  $\mathcal{O}(10^{-3})$ , we cannot neglect terms of the same order in the decay rate.
- We need  $d\Gamma$  at order  $\alpha^2$ , i. e. we must include the QED NLO corrections

### What does $\tilde{a}$ parametrize?

The introduction of the new effective coupling shifts the on-shell form factor  $F_2$ :

$$F_2|_{\text{on-shell}} = F_2^{\text{SM}}|_{\text{on-shell}} + \tilde{a}.$$

Therefore the tau  $g - 2$  at leading order is now given by:

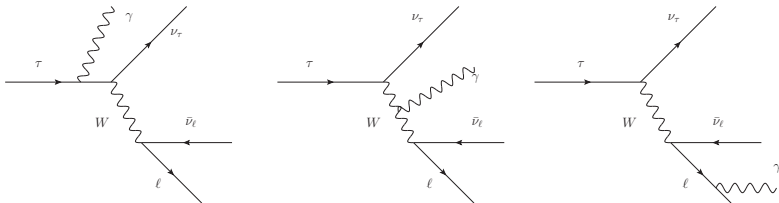
$$a_\tau = F_2^{(1)\text{QED}} + \tilde{a} + \mathcal{O}(\alpha^2) = \frac{\alpha}{2\pi} + \tilde{a} + \mathcal{O}(\alpha^2)$$

After the inclusion of QED one-loop corrections in  $d\Gamma$ ,  $\tilde{a}$  parametrizes all non-QED contributions to the  $g - 2$  of  $\mathcal{O}(\alpha)$ .

## QED one-loop Corrections to Tau Radiative Decay

We computed the polarized decay rate up to order  $\alpha^2$ , including the effective coupling of  $\tilde{a}$ .

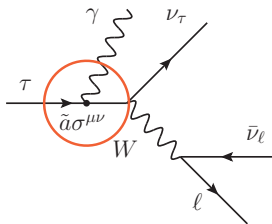
$$d\Gamma = d\Gamma_0 + \frac{m_\tau^2}{m_W^2} d\Gamma_W + \tilde{a} d\Gamma_a + \frac{\alpha}{\pi} d\Gamma_1$$



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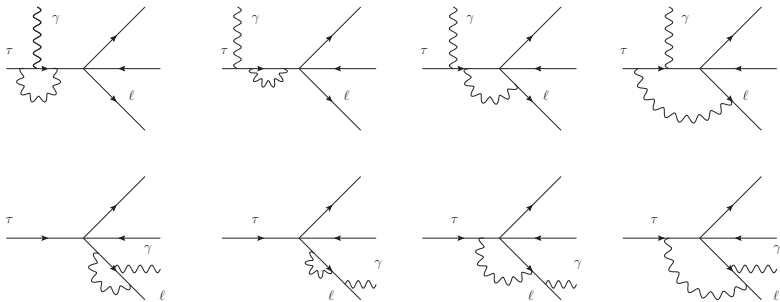
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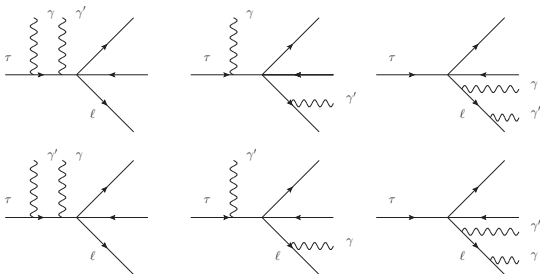
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- Neutrinos momenta integrated over the phase space.
- W propagator effects can be neglected in  $d\Gamma_1$  since  $\mathcal{O}\left(\alpha^2 \frac{m_\tau^2}{m_W^2}\right)$ .
- Fermi four-fermion interaction used in the computation of  $d\Gamma_1$ .
- Agree with previous results: T. Kinoshita, A. Sirlin PRL2(1959)177; Y. Kuno, Y. Okada, RMP73(2001)151; A. Fischer et al. PRD49(1994)3426; A. B. Arbuzov PLB597(2004)285.

## Opportunities and Challenges

- High statistics is needed.
- Phase space integration kills  $d\Gamma_a$ , since  $d\Gamma_0$  has phase space singularities ( $E_\gamma$  and collinear photon):

$$Br(\tau^- \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau \gamma) \Big|_{E_\gamma > 10\text{MeV}} = 0.361 (38)\% \quad [0.367 + (8.6 \cdot 10^{-4})\tilde{a}] \%$$

CLEO coll. PRL84(2000)830,

our results

$$Br(\tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau \gamma) \Big|_{E_\gamma > 10\text{MeV}} = 1.75 (18)\% \quad [1.837 + (8.2 \cdot 10^{-4})\tilde{a}] \%$$

- A full phase space analysis is needed.
- Feasibility for the search of the tau  $g - 2$  at Belle and future Belle-II is currently under study.
- What can LHC tell us about tau  $g - 2$ ?



# Summary and Outlook

- It is desirable to get a more precise measurement of anomalous magnetic moment of the  $\tau$ .
- There are proposals to improve the PDG value, through the study of the production cross section and decay rate, e.g. radiative decays.
- The possibility to measure  $a_\tau$  in radiative decays at Belle is under investigation.
- A feasibility study for the search of  $a_\tau$  at LHC is needed.

Thanks!

## Decay Rate Formula

The total differential decay for a polarized  $\tau$  lepton in the tau r.f. is

$$\frac{d^3\Gamma}{dx dy d\cos\theta} = \frac{\alpha G_F^2 m_\tau^5}{2(4\pi)^6} y \sqrt{x^2 - 4r^2} \\ \times \left\{ G(x, y, \theta_{\ell\gamma}) + \cos\theta_\ell \sqrt{x^2 - 4r^2} J(x, y, \theta_{\ell\gamma}) + y \cos\theta_\gamma K(x, y, \theta_{\ell\gamma}) \right\},$$

where

$$x = \frac{2E_\ell}{m_\tau}, \quad y = \frac{2E_\gamma}{m_\tau}, \quad r = \frac{m_\ell^2}{m_\tau^2},$$

$\theta_\ell$  ( $\theta_\gamma$ ) is the angle between the lepton (photon) momentum and the polarization vector.

$$G = G_0 + \frac{m_\tau^2}{m_W^2} G_W + \tilde{a} G_a + \frac{\alpha}{\pi} G_{\text{RC}}$$

The functions  $J$  and  $K$  have the same structure.

$$\begin{aligned}
 G_0 = & -\frac{64\pi^2}{3y^2z^2} \left[ r^4 (6xy^2 + 6y^3 - 6y^2z - 8y^2) + r^2 (-4x^2y^2 - 6x^2yz \right. \\
 & - 8xy^3 + 2xy^2z + 6xy^2 + 6xyz^2 + 8xyz + 6xz^2 - 4y^4 + 5y^3z + 6y^3 \\
 & - 2y^2z^2 - 6y^2z - 3yz^3 + 6yz^2 - 6z^3 - 8z^2) + 4x^3yz + 8x^2y^2z \\
 & - 8x^2yz^2 - 6x^2yz - 4x^2z^2 + 6xy^3z - 8xy^2z^2 - 6xy^2z + 6xyz^3 \\
 & - 2xyz^2 + 8xz^3 + 6xz^2 + 2y^4z - 2y^3z^2 - 3y^3z + 2y^2z^3 - 2y^2z^2 \\
 & \left. - 2yz^4 + 5yz^3 + 6yz^2 - 6z^3 \right] \text{Kinoshita \& Sirlin ('59), Kuno \& Okada ('01)}
 \end{aligned}$$

$$\begin{aligned}
 G_a = & -\frac{64\pi^2}{3yz} [xyz + 2xz^2 + y^2z + yz^2 - yz - 2z^3 - 2z^2 \\
 & + r^2 (-y^2 + yz - 3z^2)]
 \end{aligned}$$

$$G_{aa} = -\frac{16\pi^2}{3} [2x^2 + 2xy - 2xz - 3x - yz + 2z + r^2 (-3x - 2y + 4)]$$