

Domenico Bonocore



Next to Eikonal Webs

LHCphenonet Annual Meeting 2012

LHCphenonet

Why Resummation?

Why Resummation?

- *Understanding infrared divergences and large logarithms*

- *Understanding infrared divergences and large logarithms*

example: Drell-Yan $z \rightarrow 1$

$$\frac{d\sigma}{dQ^2} = \sigma_B C_F \left[\delta(1-z) \left(\frac{2\pi^2}{3} - 8 \right) + 4(1+z^2) \left(\frac{\ln(1-z)}{1-z} \right)_+ - 2 \left(\frac{1+z^2}{1-z} \right) \ln z \right]$$

$$\alpha^n \longrightarrow \frac{\log^n(1-z)}{1-z} \equiv L$$

Fixed order calculation near the threshold converges poorly

- *Understanding infrared divergences and large logarithms*

example: Drell-Yan $z \rightarrow 1$

$$\frac{d\sigma}{dQ^2} = \sigma_B C_F \left[\delta(1-z) \left(\frac{2\pi^2}{3} - 8 \right) + 4(1+z^2) \left(\frac{\ln(1-z)}{1-z} \right)_+ - 2 \left(\frac{1+z^2}{1-z} \right) \ln z \right]$$

$$\alpha^n \longrightarrow \frac{\log^n(1-z)}{1-z} \equiv L$$

more in general:

*Fixed order calculation near
the threshold converges poorly*

$$O \sim 1 + \alpha(L+1) + \alpha^2(L^2 + L + 1) + \dots$$

Why Resummation?

- *Understanding infrared divergences and large logarithms*

example: Drell-Yan $z \rightarrow 1$

$$\frac{d\sigma}{dQ^2} = \sigma_B C_F \left[\delta(1-z) \left(\frac{2\pi^2}{3} - 8 \right) + 4(1+z^2) \left(\frac{\ln(1-z)}{1-z} \right)_+ - 2 \left(\frac{1+z^2}{1-z} \right) \ln z \right]$$

$$\alpha^n \longrightarrow \frac{\log^n(1-z)}{1-z} \equiv L$$

more in general:

*Fixed order calculation near
the threshold converges poorly*

$$O \sim 1 + \alpha(L+1) + \alpha^2(L^2 + L + 1) + \dots$$

- *Understanding all order features of gauge theories*

Why Resummation?

- *Understanding infrared divergences and large logarithms*

example: Drell-Yan $z \rightarrow 1$

$$\frac{d\sigma}{dQ^2} = \sigma_B C_F \left[\delta(1-z) \left(\frac{2\pi^2}{3} - 8 \right) + 4(1+z^2) \left(\frac{\ln(1-z)}{1-z} \right)_+ - 2 \left(\frac{1+z^2}{1-z} \right) \ln z \right]$$

$$\alpha^n \longrightarrow \frac{\log^n(1-z)}{1-z} \equiv L$$

more in general:

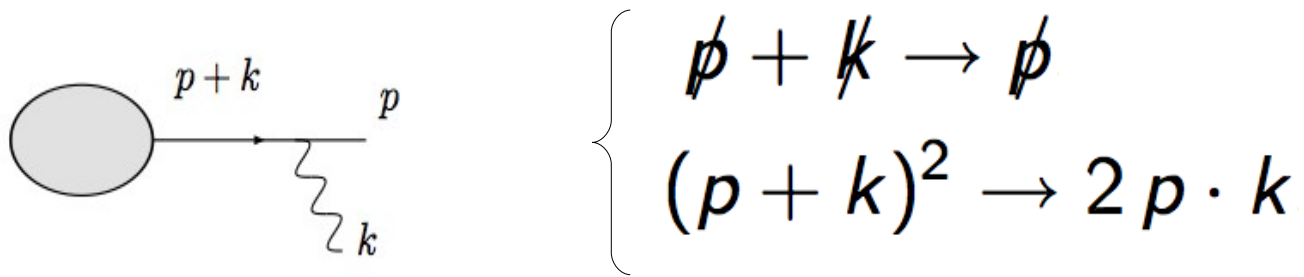
Fixed order calculation near the threshold converges poorly

$$O \sim 1 + \alpha(L+1) + \alpha^2(L^2 + L + 1) + \dots$$

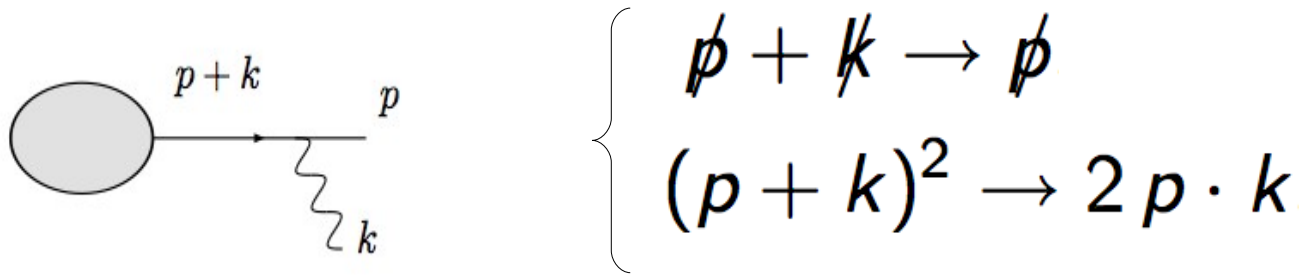
- *Understanding all order features of gauge theories*

\longrightarrow *Let's focus on soft divergences*

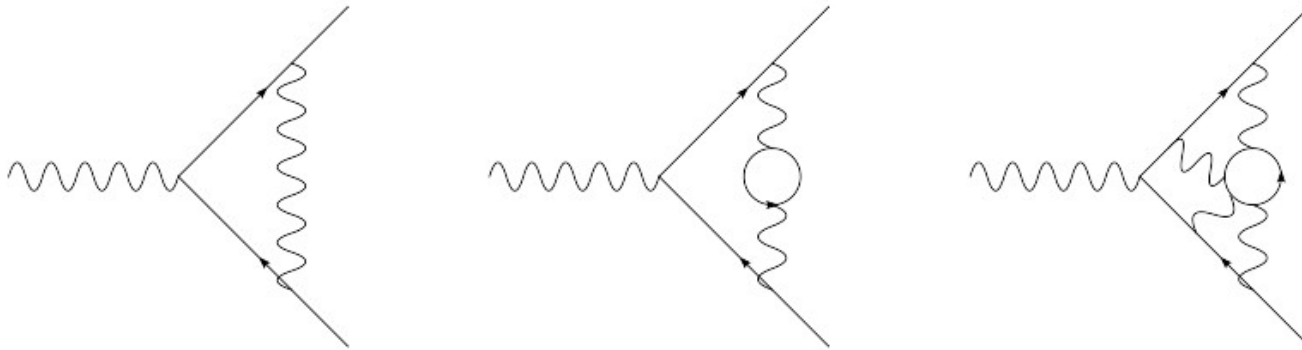
Eikonal Approximation and exponentiation



Eikonal Approximation and exponentiation

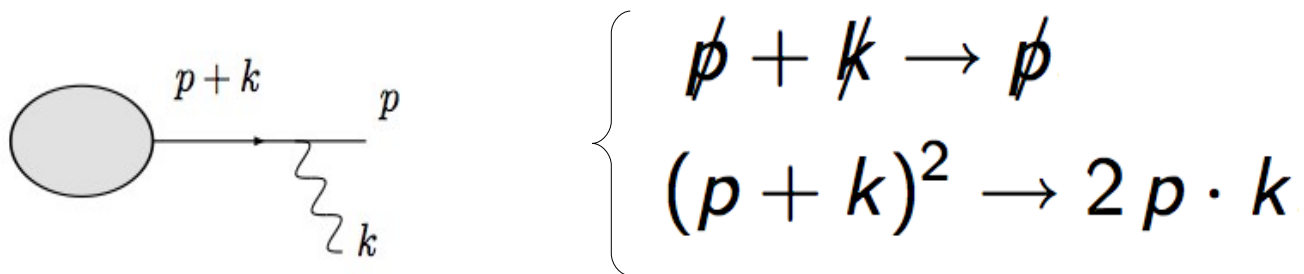


Abelian \longrightarrow **Connected subdiagrams** (~ '60 – Yennie, Frautschi, Suura)

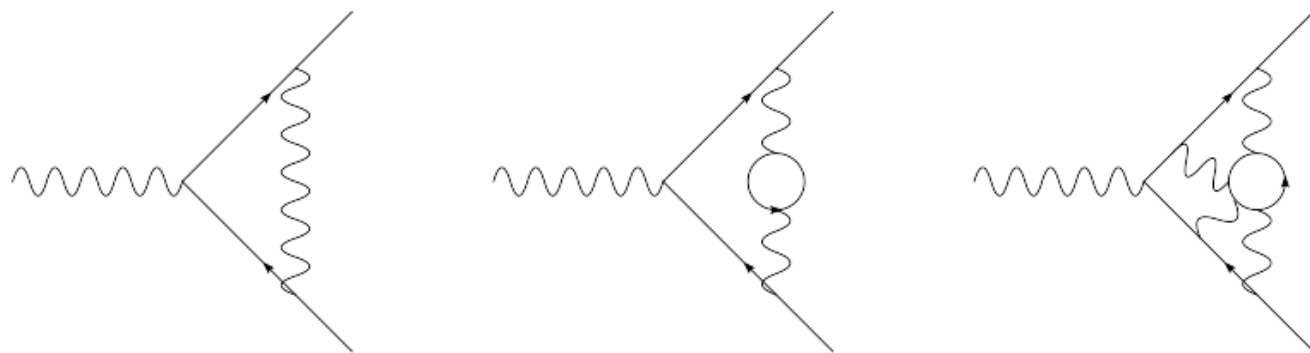


$$\mathcal{A} = \mathcal{A}_0 \exp \left[\sum G_c \right]$$

Eikonal Approximation and exponentiation

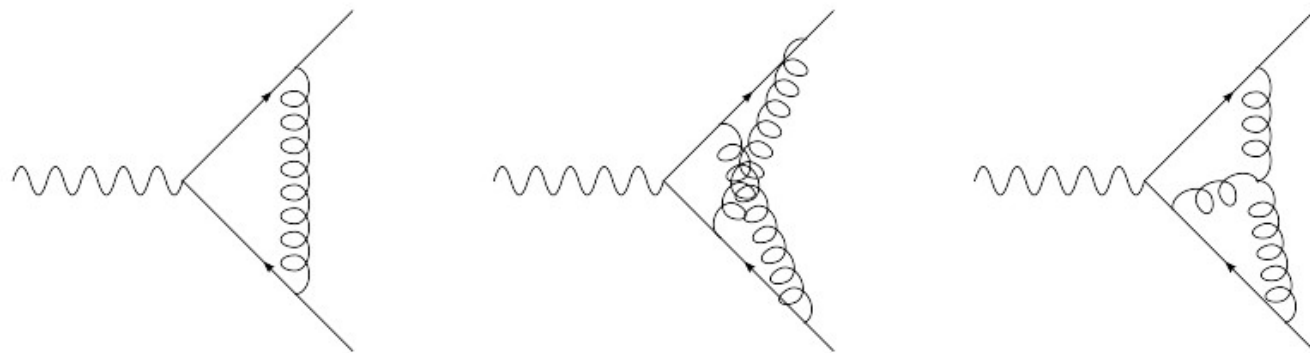


Abelian \longrightarrow **Connected subdiagrams** (~ '60 – Yennie, Frautschi, Suura)



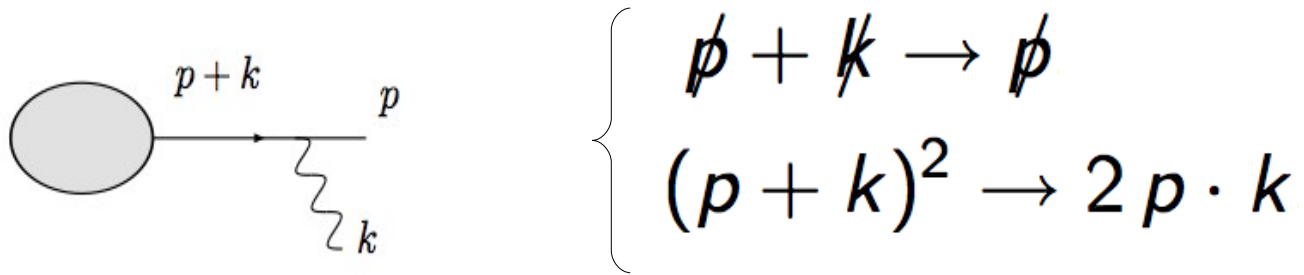
$$\mathcal{A} = \mathcal{A}_0 \exp \left[\sum G_c \right]$$

Non abelian \longrightarrow **Webs** (~ '80 Gatheral, Frenkel, Taylor, Sterman)

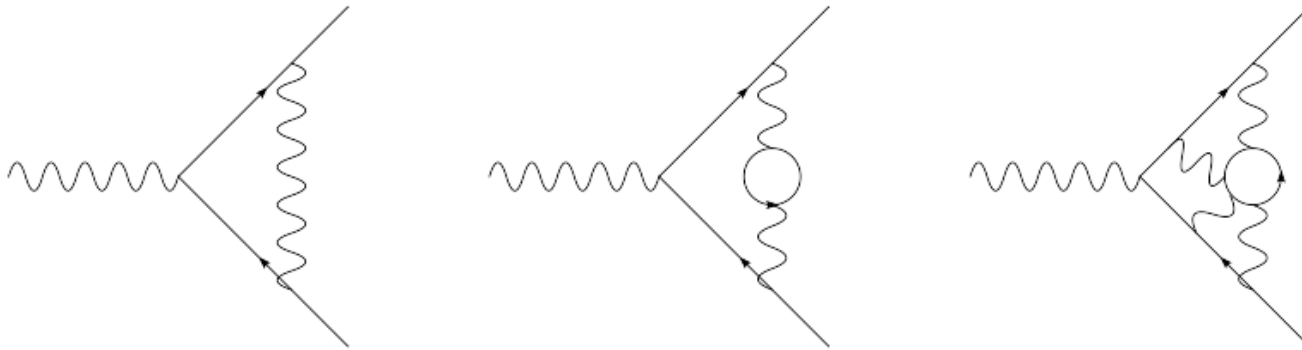


$$\mathcal{A} = \mathcal{A}_0 \exp \left[\sum \bar{C}_W W \right]$$

Eikonal Approximation and exponentiation

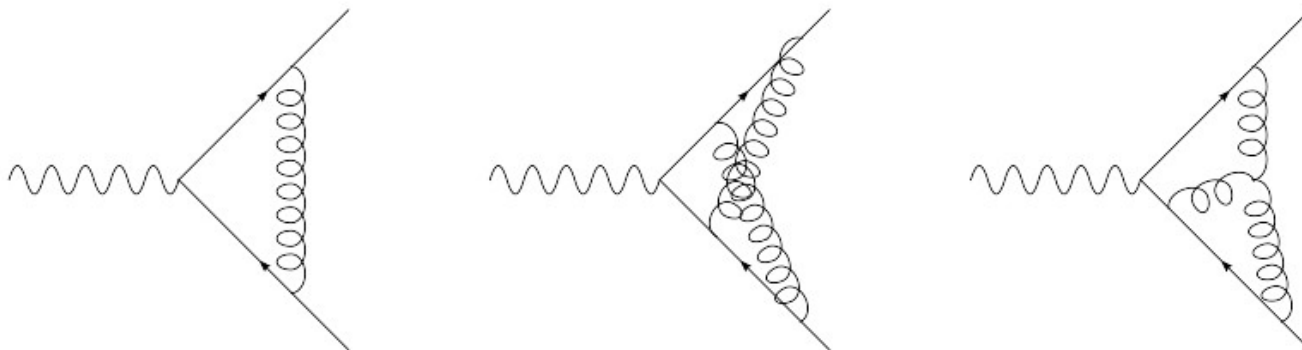


Abelian \longrightarrow **Connected subdiagrams** (~ '60 – Yennie, Frautschi, Suura)



$$\mathcal{A} = \mathcal{A}_0 \exp \left[\sum G_c \right]$$

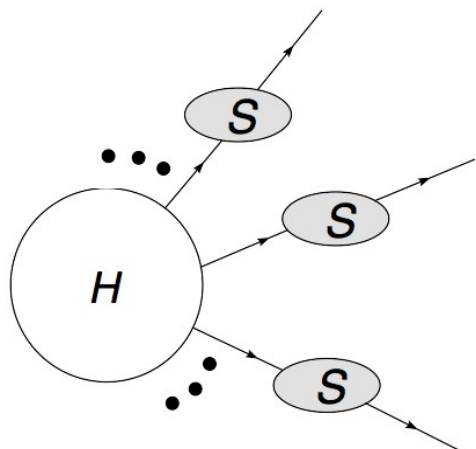
Non abelian \longrightarrow **Webs** (~ '80 Gatheral, Frenkel, Taylor, Sterman)



$$\mathcal{A} = \mathcal{A}_0 \exp \left[\sum \bar{C}_W W \right]$$

(2008: Laenen,Stavenga,White - 0811.2067) **new approach via path integral...**

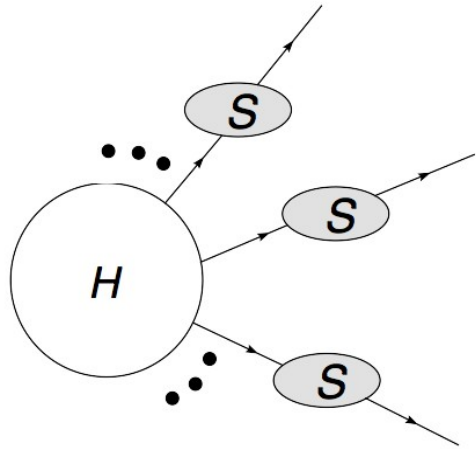
A path integral approach: QED case



$$\int \mathcal{D}A^\mu \equiv \int \mathcal{D}A_s^\mu \mathcal{D}A_h^\mu$$

*Factorization of hard
and soft modes*

$$G(p_1, \dots, p_n) = \int \mathcal{D}A_s^\mu H(x_1, \dots, x_n; A_s^\mu) S(p_1, x_1) \dots S(p_n, x_n) e^{iS[A_s]}$$



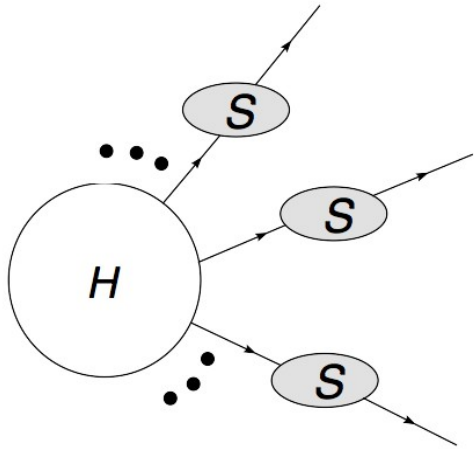
$$\int \mathcal{D}A^\mu \equiv \int \mathcal{D}A_s^\mu \mathcal{D}A_h^\mu$$

Factorization of hard and soft modes

$$G(p_1, \dots, p_n) = \int \mathcal{D}A_s^\mu H(x_1, \dots, x_n; A_s^\mu) S(p_1, x_1) \dots S(p_n, x_n) e^{iS[A_s]}$$

external propagators $S(p, x)$ as first quantized path integral

$$S(p, x) = \int \mathcal{D}p \mathcal{D}x \exp \left[i \int_0^T dt \left(p \dot{x} - \frac{1}{2} p^2 + (p_f + p) \cdot A_s(x_i + p_f t + x) + \frac{1}{2} \partial \cdot A_s(x_i + p_f t + x) - A_s^2(x_i - p_f t + x) \right) \right]$$



$$\int \mathcal{D}A^\mu \equiv \int \mathcal{D}A_s^\mu \mathcal{D}A_h^\mu$$

Factorization of hard and soft modes

$$G(p_1, \dots, p_n) = \int \mathcal{D}A_s^\mu H(x_1, \dots, x_n; A_s^\mu) S(p_1, x_1) \dots S(p_n, x_n) e^{iS[A_s]}$$

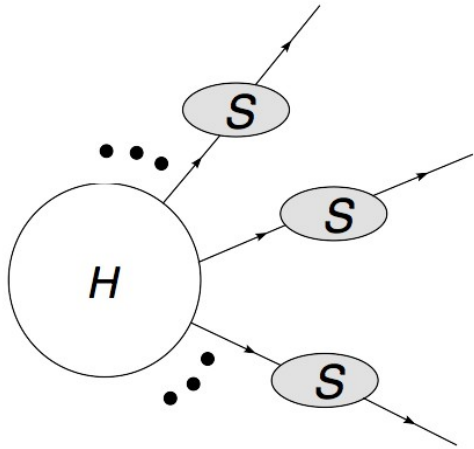
external propagators $S(p, x)$ as first quantized path integral

$$S(p, x) = \int \mathcal{D}p \mathcal{D}x \exp \left[i \int_0^T dt \left(p \dot{x} - \frac{1}{2} p^2 + (p_f + p) \cdot A_s(x_i + p_f t + x) + \frac{1}{2} \partial \cdot A_s(x_i + p_f t + x) - A_s^2(x_i - p_f t + x) \right) \right]$$

going on with computation...

$$S(p_1, \dots, p_n) = \int \mathcal{D}A_s^\mu H(x_1, \dots, x_n; A_s^\mu) e^{-ip_1 x_1} f_1(\infty) \dots e^{-ip_n x_n} f_n(\infty) e^{iS[A_s^\mu]}$$

$$f(\infty) = \int_{x(0)=0} \mathcal{D}x e^{i \int_0^\infty dt \left(\frac{1}{2} \dot{x}^2 + (p_f + \dot{x}) \cdot A(x_i + p_f t + x(t)) + \frac{i}{2} \partial \cdot A(x_i + p_f t + x) \right)}$$



$$\int \mathcal{D}A^\mu \equiv \int \mathcal{D}A_s^\mu \mathcal{D}A_h^\mu$$

Factorization of hard and soft modes

$$G(p_1, \dots, p_n) = \int \mathcal{D}A_s^\mu H(x_1, \dots, x_n; A_s^\mu) S(p_1, x_1) \dots S(p_n, x_n) e^{iS[A_s]}$$

external propagators $S(p, x)$ as first quantized path integral

$$S(p, x) = \int \mathcal{D}p \mathcal{D}x \exp\left[i \int_0^T dt \left(p\dot{x} - \frac{1}{2}p^2 + (p_f + p) \cdot A_s(x_i + p_f t + x) + \frac{1}{2} \partial \cdot A_s(x_i + p_f t + x) - A_s^2(x_i - p_f t + x) \right)\right]$$

going on with computation...

$$S(p_1, \dots, p_n) = \int \mathcal{D}A_s^\mu H(x_1, \dots, x_n; A_s^\mu) e^{-ip_1 x_1} f_1(\infty) \dots e^{-ip_n x_n} f_n(\infty) e^{iS[A_s^\mu]}$$

$$f(\infty) = \int_{x(0)=0} \mathcal{D}x e^{i \int_0^\infty dt \left(\frac{1}{2} \dot{x}^2 + (p_f + \dot{x}) \cdot A(x_i + p_f t + x(t)) + \frac{i}{2} \partial \cdot A(x_i + p_f t + x) \right)}$$

a theory for the soft field where terms in the exponent work as sources \rightarrow Exponentiation like in usual textbook results!

A path integral approach: QED case

$$S(p_1, \dots, p_n) = \int \mathcal{D}A_s^\mu H(x_1, \dots, x_n; A_s^\mu) e^{-ip_1 x_1} f_1(\infty) \dots e^{-ip_n x_n} f_n(\infty) e^{iS[A_s^\mu]}$$
$$f(\infty) = \int_{x(0)=0} \mathcal{D}x e^{i \int_0^{\infty} dt \left(\frac{1}{2} \dot{x}^2 + (p_f + \dot{x}) \cdot A(x_i + p_f t + x(t)) + \frac{i}{2} \partial \cdot A(x_i + p_f t + x) \right)}$$

A path integral approach: QED case

4

$$S(p_1, \dots, p_n) = \int \mathcal{D}A_s^\mu H(x_1, \dots, x_n; A_s^\mu) e^{-ip_1 x_1} f_1(\infty) \dots e^{-ip_n x_n} f_n(\infty) e^{iS[A_s^\mu]}$$
$$f(\infty) = \int_{x(0)=0} \mathcal{D}x e^{i \int_0^{\infty} dt (\frac{1}{2} \dot{x}^2 + (p_f + \dot{x}) \cdot A(x_i + p_f t + x(t)) + \frac{i}{2} \partial \cdot A(x_i + p_f t + x))}$$

→ *first approx: neglect fluctuation* → *classical trajectory*

$$f(\infty) \propto e^{i \int dx \cdot A(x)} \quad \text{Wilson lines}$$

This is exactly the eikonal limit !

A path integral approach: QED case

4

$$S(p_1, \dots, p_n) = \int \mathcal{D}A_s^\mu H(x_1, \dots, x_n; A_s^\mu) e^{-ip_1 x_1} f_1(\infty) \dots e^{-ip_n x_n} f_n(\infty) e^{iS[A_s^\mu]}$$
$$f(\infty) = \int_{x(0)=0} \mathcal{D}x e^{i \int_0^{\infty} dt (\frac{1}{2} \dot{x}^2 + (p_f + \dot{x}) \cdot A(x_i + p_f t + x(t)) + \frac{i}{2} \partial \cdot A(x_i + p_f t + x))}$$

→ *first approx: neglect fluctuation* → *classical trajectory*

$$f(\infty) \propto e^{i \int dx \cdot A(x)} \quad \text{Wilson lines}$$

This is exactly the eikonal limit !

→ *second approx: consider fluctuations over $x(t)$*

$$p_j = \lambda n_j; n_j^2 = 0 \quad \text{expanding in } 1/\lambda:$$

• $\lambda \rightarrow \infty$ → *Eik approx*

• *Subleading term* $\mathcal{O}(\lambda^{-1})$ → *Next to Eikonal approx*

A path integral approach: QED case

4

$$S(p_1, \dots, p_n) = \int \mathcal{D}A_s^\mu H(x_1, \dots, x_n; A_s^\mu) e^{-ip_1 x_1} f_1(\infty) \dots e^{-ip_n x_n} f_n(\infty) e^{iS[A_s^\mu]}$$
$$f(\infty) = \int_{x(0)=0} \mathcal{D}x e^{i \int_0^{\infty} dt (\frac{1}{2} \dot{x}^2 + (p_f + \dot{x}) \cdot A(x_i + p_f t + x(t)) + \frac{i}{2} \partial \cdot A(x_i + p_f t + x))}$$

→ *first approx: neglect fluctuation* → *classical trajectory*

$$f(\infty) \propto e^{i \int dx \cdot A(x)} \quad \text{Wilson lines} \quad \text{This is exactly the eikonal limit !}$$

→ *second approx: consider fluctuations over $x(t)$*

$$p_j = \lambda n_j; n_j^2 = 0 \quad \text{expanding in } 1/\lambda:$$

- $\lambda \rightarrow \infty$ → *Eik approx*
- *Subleading term* $\mathcal{O}(\lambda^{-1})$ → *Next to Eikonal approx*

For both E and NE approx one can find exponentiation and Feynman rules as usual in QFT

A path integral approach: QCD case

Much more difficult: a key point for exponentiation in QED was that everything commute

$$f^{i_1 j_1}(\infty) = \left[\int_{x(0)=x_i} \mathcal{D}x \mathcal{P} e^{i \int_0^\infty dt \left(\frac{1}{2} \dot{x}^2 + (p_f + \dot{x}) \cdot A(x_i + p_f t + x(t)) + \frac{i}{2} \partial \cdot A(x_i + p_f t + x) \right)} \right]^{i_1 j_1}$$

Much more difficult: a key point for exponentiation in QED was that everything commute

$$f^{i_1 j_1}(\infty) = \left[\int_{x(0)=x_i} \mathcal{D}x \mathcal{P} e^{i \int_0^\infty dt \left(\frac{1}{2} \dot{x}^2 + (p_f + \dot{x}) \cdot A(x_i + p_f t + x(t)) + \frac{i}{2} \partial \cdot A(x_i + p_f t + x) \right)} \right]^{i_1 j_1}$$

→ *Replica trick: N copies of the theories, expanding in N, then N=1*

Much more difficult: a key point for exponentiation in QED was that everything commute

$$f^{i_1 j_1}(\infty) = \left[\int_{x(0)=x_i} \mathcal{D}x \mathcal{P} e^{i \int_0^\infty dt \left(\frac{1}{2} \dot{x}^2 + (p_f + \dot{x}) \cdot A(x_i + p_f t + x(t)) + \frac{i}{2} \partial \cdot A(x_i + p_f t + x) \right)} \right]^{i_1 j_1}$$

→ *Replica trick: N copies of the theories, expanding in N, then N=1*

Also in the non abelian case exponentiation become possible at Eikonal and Next to Eikonal level

- *2 parton final legs* (2008: Laenen,Stavenga,White - 0811.2067)
- *multiparton case* (2010: Gardi,Laenen,Stavenga,White - 1008.0098)

State of the art:

Evaluating matrix elements at E and NE level

$$\mathcal{M}^E \quad \mathcal{M}^{NE}$$

State of the art:

Evaluating matrix elements at E and NE level

$$\mathcal{M}^E \quad \mathcal{M}^{NE}$$

*It remains to compute a full cross section
(various attempts by means of other approaches)*

$$\frac{d\sigma}{d\xi} = \int d\text{PS}^{(E)} |\mathcal{M}^{(E)}|^2 + \left[\int d\text{PS}^{(E)} |\mathcal{M}^{(NE)}|^2 + \int d\text{PS}^{(NE)} |\mathcal{M}^{(E)}|^2 \right] + \mathcal{O}(\text{NNE})$$

Thank you!