

QED corrections in Monte-Carlo event generators

Marek Schönherr

Institute for Particle Physics Phenomenology



24/09/2012

LHCphenonet



Process independent corrections

Will not discuss:

- process specific electroweak corrections
 - SANC Andonov et.al Comput.Phys.Commun.174(2006)481-517
 - WINHAC/ZINHAC Placzek, Jadach EPJC29(2003)325-339
 - HORACE Carloni Calame et.al. Phys.Rev.D69(2004)037301
 - ...

Will discuss:

- process independent/universal approximate corrections
- DGLAP – collinear photon resummation:
 - PHOTOS Barberio, Wąs Comput.Phys.Commun.79(1994)291-308
 - PYTHIA8 Sjöstrand, Mrenna, Skands Comput.Phys.Commun.178(2008)852-867
 - SHERPA Höche, Schumann, Siegert Phys.Rev.D81(2010)034026
- YFS – soft photon resummation
 - HERWIG++ Hamilton, Richardson JHEP07(2006)010
 - SHERPA MS, Krauss JHEP12(2008)018

QED corrections – DGLAP

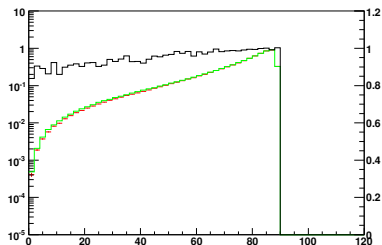
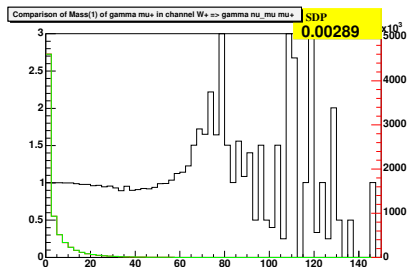
DGLAP

Gribov, Lipatov *Sov.J.Nucl.Phys.*15(1972)438-450, etc.

- resummation of collinear divergences
- strong ordering of emission scales
- soft-photon coherence not trivial
achieved either through reweighting, inclusion of correct soft limit in splitting functions or ordering variable
- QED parton showers:
PHOTOS (unordered); SHERPA/CSSHOWER++, PYTHIA8 (p_{\perp} -ordered)
Seymour Z.Phys.C56(1992)161-170, etc.
→ importance of ordering variable in recovering DGLAP equations,
see *Skands, Weinzierl PRD79(2009)074021*
- dedicated DY implementation in HORACE ($\mathcal{O}(\alpha_{EW})$ matched to QED-DGLAP)

PHOTOS

- implemented for particle decays ($1 \rightarrow n$)
- one-by-one photon emission with DGLAP-like splitting kernels
 → iterated for multiphoton em.
 → no ordering $\Rightarrow \frac{1}{n!}$ sym. factor
- ME corrections for $Z \rightarrow \ell\ell$,
 $W \rightarrow \ell\nu$
- YFS-like soft multiphoton correction
 → approximates multiphoton interferences



QED corrections – DGLAP

Ordered QED parton showers:

- strong ordering of emission scales
→ photons can be identified by their emission scales
- emissions off initial and final state charged particles
- interleaving with QCD evolution crucial to obtain correct emission rates
- need probabilistic formulation
→ in QCD: large- N_C limit
→ in QED: neglect same-sign-charged dipoles
- real emission corrections possible through ME-reweighting
- QED-MEPs (CKKW merging) possible (but not practical)
→ matrix element corrections for multiple hard photon emissions

[Seymour Z.Phys.C56\(1992\)161-170, etc.](#)

[Höche, Schumann, Siegert PRD81\(2010\)034026](#)

QED corrections – YFS

YFS

Yennie, Frautschi, Suura *Ann.Phys.*13(1961)379-452

- resummation of soft-photon logarithms in massive Abelian theories
- construction through sum of dipole eikonals
- no ordering of emission, automatic soft-photon coherence
- universal collinear logarithms can be supplemented order-by-order, but not resummed
→ however, $\exp[-\alpha_{\text{QED}}L^2] \not\approx 1 - \alpha_{\text{QED}}L^2$ in extreme phase space regions
- process dependent fixed order corrections trivial
- used in universal implementation in HERWIG++ and SHERPA
specific processes in e.g. WINHAC, ZINHAC
- heavily used in LEP-time high precision MCs
YFSWW, YFSZZ, KORALW, KORALZ, KKMC, etc.

QED corrections – YFS

- coherent radiation off charged multipole
all interferences due to emissions from different legs present
- unitary implementations for $1 \rightarrow n$
finite virtual corrections affect relative rate of emission and no-emission

- dedicated $\mathcal{O}(\alpha_{\text{QED}})$ corrections
universal collinear emission corrections through CS dipoles (all)
- current limitation: $1 \rightarrow n$ processes
→ applied to hard process by means of narrow-width approximation to production of non-QCD final state
→ applied to all hadronic and τ decays

SHERPA:

matrix element	real $\mathcal{O}(\alpha_{\text{QED}})$	virtual $\mathcal{O}(\alpha_{\text{QED}})$
$V^0 \rightarrow F^+ F^-$	✓	✓
$V^0 \rightarrow S^+ S^-$	✓	✓
$S^0 \rightarrow F^+ F^-$	✓	✓
$S^0 \rightarrow S^+ S^-$	✓	✓
$W^\pm \rightarrow \ell^\pm \nu_\ell$	✓	✓
$\tau^\pm \rightarrow \ell^\pm \nu_\ell \nu_\tau$	✓	✗
$S^0 \rightarrow S^\mp \ell^\pm \nu_\ell$	✓	✓
$S^0 \rightarrow V^\mp \ell^\pm \nu_\ell$	✓	✗

QED corrections – interplay with higher order QCD

QCD+QED DGLAP evolution

- must/can evolve simultaneously
→ compete for phase space
- may have different IR-cutoffs
- photon emissions off quarks
drowned by gluon emissions
→ have little effect besides in
dedicated searches

QCD DGLAP evolution & YFS exponentiation

- YFS does not evolve
→ “simultaneous” emission
- cannot be run interleaved with
QCD DGLAP evolution
- ⇒ need to define mutually distinct
sets on QCD partons and
non-QCD particles
→ sensible only if QCD evolution
leaves non-QCD subset invariant
- ⇒ no QED rad. off quarks
- need some assumption about
internal resonances (possibly
multiple distinct non-QCD
subsets)

⇒ **same applies to PHOTOS + PS**

Comparisons – DGLAP vs. YFS

Definition of objects for pseudo-observables:

- bare electrons (electron object as is generated by MC)
→ due to soft-photon cut-off this includes some amount of soft/collinear photon radiation, not identical to electron of the QED Lagrangian
- dressed electrons (sum of 4-vectors of bare electron and photons within $\Delta R = 0.2$ around bare electron)
- identified photons (isolated by $\Delta R > 0.2$ from electron, $E_\gamma > 1$ GeV)

⇒ look at pseudo-observables characterising the radiation pattern

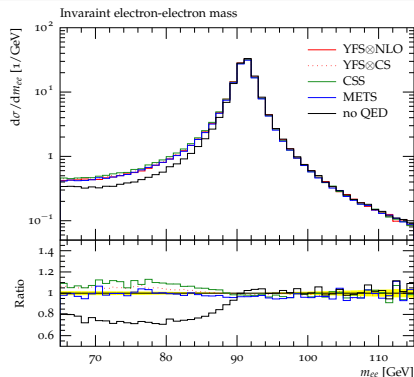
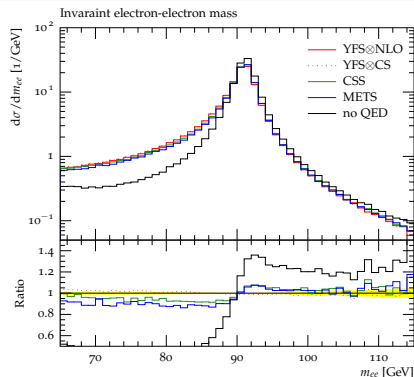
In the following:

left: pseudo-observables with bare electrons

right: pseudo-observables with dressed electrons

→ closer to exp. observables

QED corrections – comparison – DY production



YFS \otimes NLO YFS resummation with exact $\mathcal{O}(\alpha_{\text{QED}})$ correction

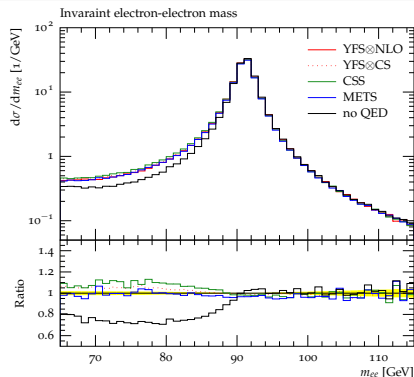
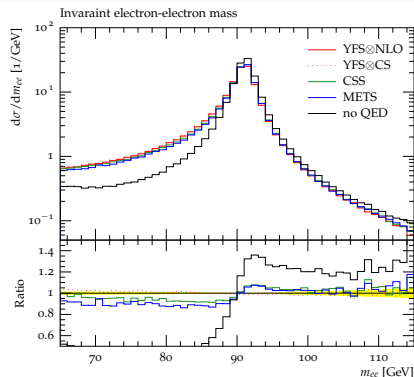
YFS \otimes CS YFS resummation with approximate universal coll. approximation

CSS DGLAP resummation

METS DGLAP res. merged with ME with up to 2 photons ($Q_{\text{cut}} = 1\text{GeV}$)

no QED pure leading order $Z \rightarrow e^+e^-$

QED corrections – comparison – DY production

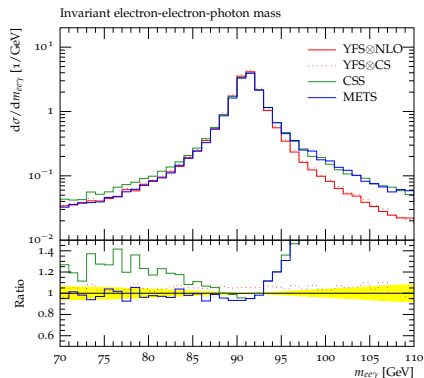
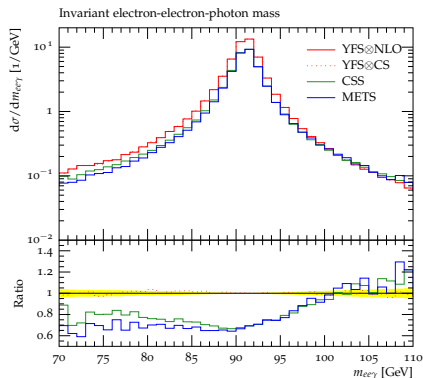


invariant dilepton mass

left: bare electrons, **right:** physical electrons (4-momentare of photons within $\Delta R = 0.2$ recombined bare electron)

bare quantities show rather large differences, but physical quantities show good agreement of all ME-corrected calculations

QED corrections – comparison – DY production

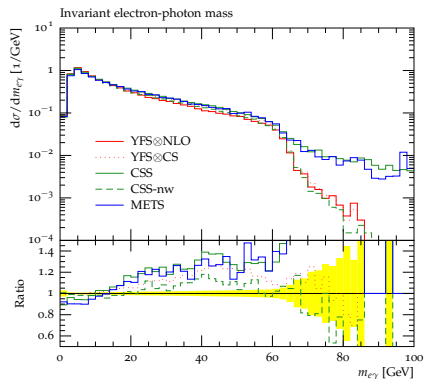
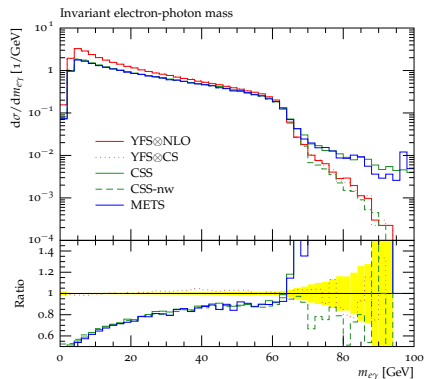


invariant $e^+e^-\gamma$ mass, isolated hard photons ($E_\gamma > 1\text{GeV}$, $\Delta R > 0.2$)

bare quantities show rather large differences, but physical quantities show good agreement of all ME-corrected calculations

difference at large $m_{ee\gamma}$ due to initial state radiation neglected in YFS approach

QED corrections – comparison – DY production



invariant mass of hardest photon and closest e^\pm , isolated hard photons
 ($E_\gamma > 1\text{GeV}$, $\Delta R > 0.2$)

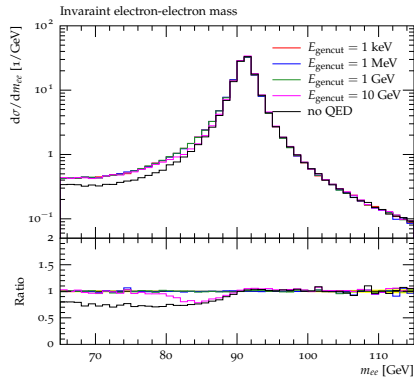
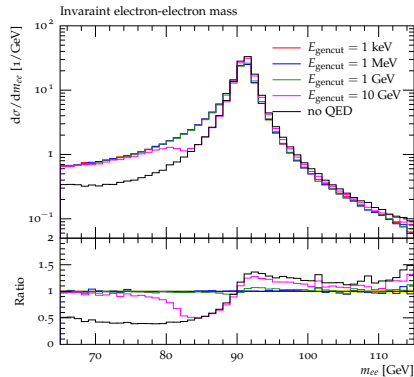
CSS-nw neglects ISR, same large $m_{e\gamma}$ as YFS

in bare spectrum missing collinear resummation of YFS visible

→ could be remedied by inclusion of higher order coll. approximation

QED corrections – comparison – DY production

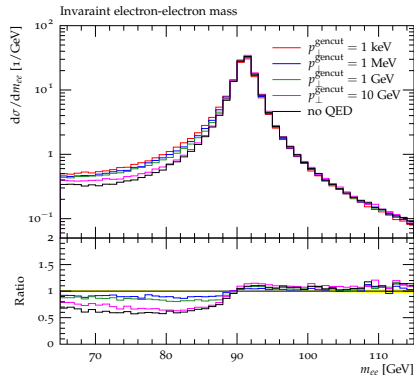
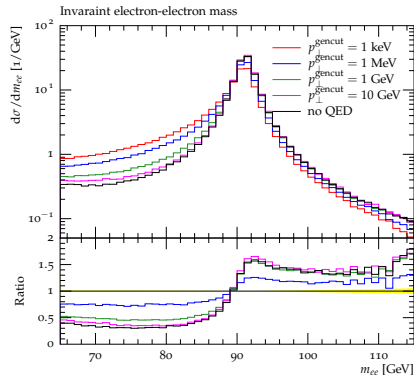
Cut-off dependence in SHERPA/PHOTONS++ (YFS):



cut-off as minimum photon energy in multipole rest frame
 bare quantities stable for $E_{\text{gencut}} < 1\text{GeV}$

QED corrections – comparison – DY production

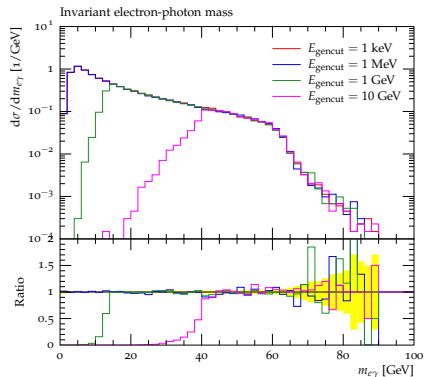
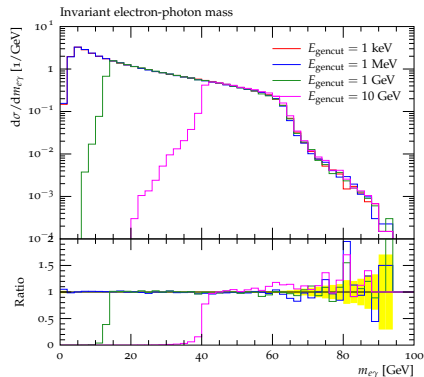
Cut-off dependence in SHERPA/CSSHOWER++ (DGLAP):



cut-off as minimum relative transverse momentum
 physical quantities still show some dependence on $p_{\perp}^{\text{gencut}}$

QED corrections – comparison – DY production

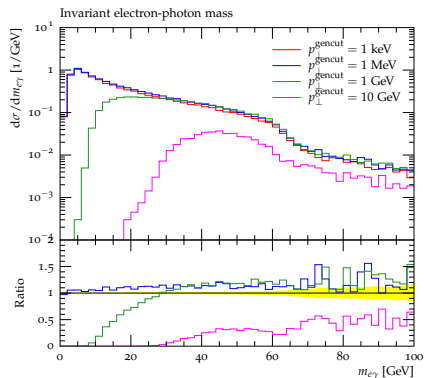
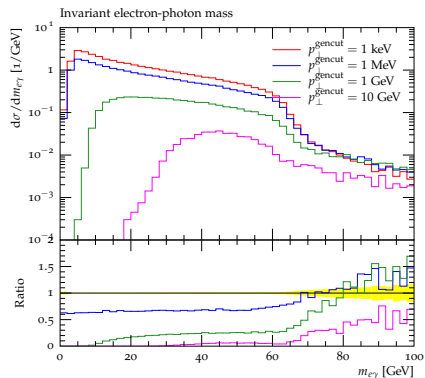
Cut-off dependence in SHERPA/PHOTONS++ (YFS):



cut-off as minimum photon energy in multipole rest frame

QED corrections – comparison – DY production

Cut-off dependence in SHERPA/CSSHOWER++ (DGLAP):

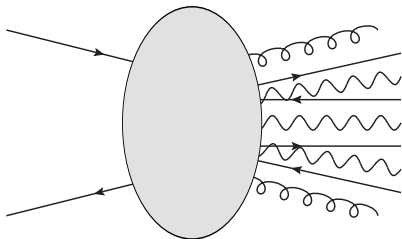


cut-off as minimum relative transverse momentum

QED corrections – effects of resonance assumptions

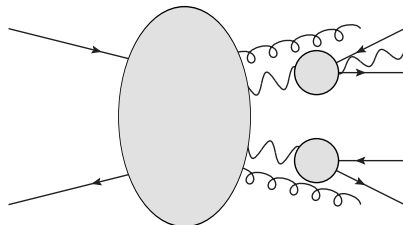
implementation depending on whether resonant decays specified or not, e.g.

$$pp \rightarrow \ell^+ \ell^- \nu_e \bar{\nu}_e + \text{jets}$$



- photons may recoil against full non-QCD system

$$pp \rightarrow Z[\rightarrow \ell^+ \ell^-] Z[\rightarrow \nu_e \bar{\nu}_e] + \text{jets}$$



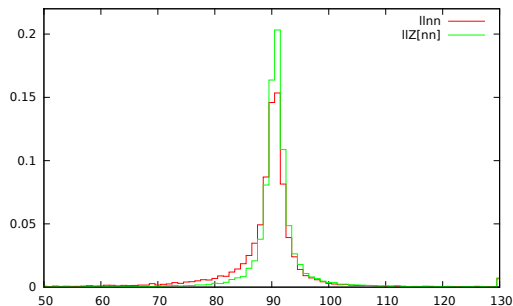
- photons may recoil only within their specified decay subsystem

⇒ different phase space volume for hard wide-angle emissions

⇒ soft and collinear limits the same, differences beyond formal accuracy, must be fixed by exact higher order corrections

QED corrections – effects of resonance assumptions

$m_{\nu\nu}$ in $pp \rightarrow \ell^+ \ell^- \nu_e \bar{\nu}_e + \text{jets}$ vs. $pp \rightarrow Z[\rightarrow \ell^+ \ell^-] Z[\rightarrow \nu_e \bar{\nu}_e] + \text{jets}$



The question is: How much energy is QED-bremsstrahlung allowed to remove from the system? So much that the $Z[\rightarrow \nu_\mu \bar{\nu}_\mu]$ is forced off-shell? Beyond formal accuracy, needs to be answered by exact matrix-element corrections.

Conclusions

- very good description of higher order QED effects not only necessary for precision physics, but to estimate acceptances, isolations, etc.
- DGLAP best describes hard collinear radiation
 - usually gets recombined with charged particle to physical objects
 - hard wide-angle photon emission through fixed-order correction (MEPS)
 - natively incorporates initial state radiation
- YFS best describes comparably soft wide-angle radiation
 - ends up as separate noise depleting energy from its production process
 - hard wide-angle photon emission through fixed-order correction
 - currently limited to $1 \rightarrow n$ type (sub)processes
 - ⇒ good enough for all observables considered so far
- good agreement for physical quantities after (at least) real emission corrections
- good description of rather inclusive quantities needs well understood wide-angle soft emissions

Thank you for your attention!

QED corrections – YFS in SHERPA

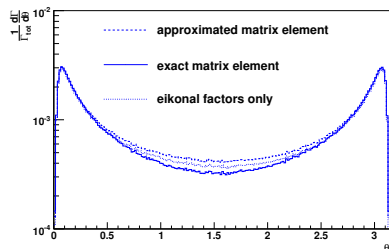
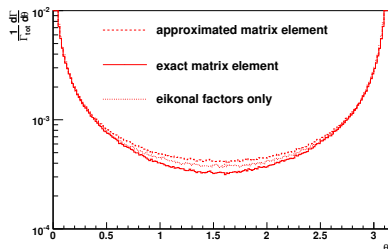
Validation: radiative decay rates

	$\frac{\Gamma(\mu \rightarrow e \nu_e \nu_\mu \gamma)}{\Gamma(\mu \rightarrow e \nu_e \nu_\mu, \text{incl.})}$	$\frac{\Gamma(\tau \rightarrow e \nu_e \nu_\tau \gamma)}{\Gamma(\tau \rightarrow e \nu_e \nu_\tau, \text{incl.})}$	$\frac{\Gamma(\tau \rightarrow \mu \nu_\mu \nu_\tau \gamma)}{\Gamma(\tau \rightarrow \mu \nu_\mu \nu_\tau, \text{incl.})}$
PDG	0.014(4)	0.09(1)	0.021(3)
SHERPA	0.0147(1)	0.0999(3)	0.0233(2)

branching ratios of the radiative leptonic μ and τ decay mode ($E_\gamma > 10\text{MeV}$) in relation to their inclusive leptonic mode calculated by SHERPA/PHOTONS++ and the PDG world average

QED corrections – YFS in SHERPA

Validation: photon emission interferences in $Z \rightarrow \ell\ell$



angle of individual photons in dipole rest frame after radiation
 $Z \rightarrow e^+e^-$ (left), $Z \rightarrow \tau^+\tau^-$ (right)

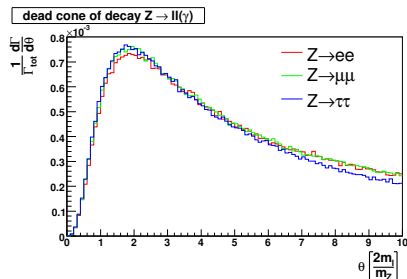
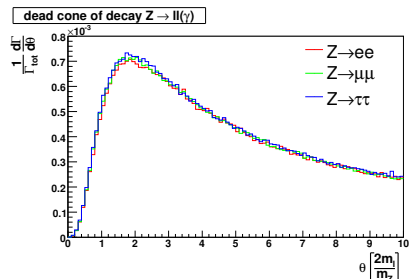
solid exact $\mathcal{O}(\alpha_{\text{QED}})$ correction

dashed universal $\mathcal{O}(\alpha_{\text{QED}})$ collinear approximation

dotted soft eikonals only

QED corrections – YFS in SHERPA

Validation: dead cone of charged massive particle in $Z \rightarrow \ell\ell$



angle between charged lepton and photon in units of m_ℓ/m_Z

left: soft eikonals only, right: exact $\mathcal{O}(\alpha_{\text{QED}})$ ME