Handling astrophysical uncertainties on direct detection experiments

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- Astrophysical uncertainties
 - i) observations
 - ii) simulations
- Consequences
- Strategies
 - i) integrate out
 - ii) marginalise over
- Parameterising the speed distribution

Introduction

Differential event rate for elastic scattering: (assuming spin-independent coupling and $f_p=f_n$)

$$\frac{dR}{dE} = \frac{\sigma_{\rm p} \rho_0}{\mu_{\rm p,\chi}^2 m_{\chi}} A^2 F^2(E) \int_{v_{\rm min}}^{\infty} \frac{f(v)}{v} dv \qquad v_{\rm min} = \left(\frac{E(m_A + m_{\chi})^2}{2m_A m_{\chi}^2}\right)^{1/2}$$

Particle physics parameters:

WIMP mass and cross-section,

$$m_{\chi}$$
 $\sigma_{
m p}$

Astrophysical input:

local DM density and speed distribution

$$\rho_0$$
 $f(v)$

Experimental constraints on σ - m_χ plane usually calculated using 'standard halo model': isotropic, isothermal sphere, with Maxwell-Boltzmann speed distribution

$$f(\mathbf{v}) \propto \exp\left(-\frac{3|\mathbf{v}|^2}{2\sigma^2}\right)$$
 $\sigma = \sqrt{\frac{3}{2}}v_{\rm c}$

with $v_c=220 \text{ km s}^{-1}$ and local density $\rho_0=0.3 \text{ GeV cm}^{-3}$

Energy spectrum

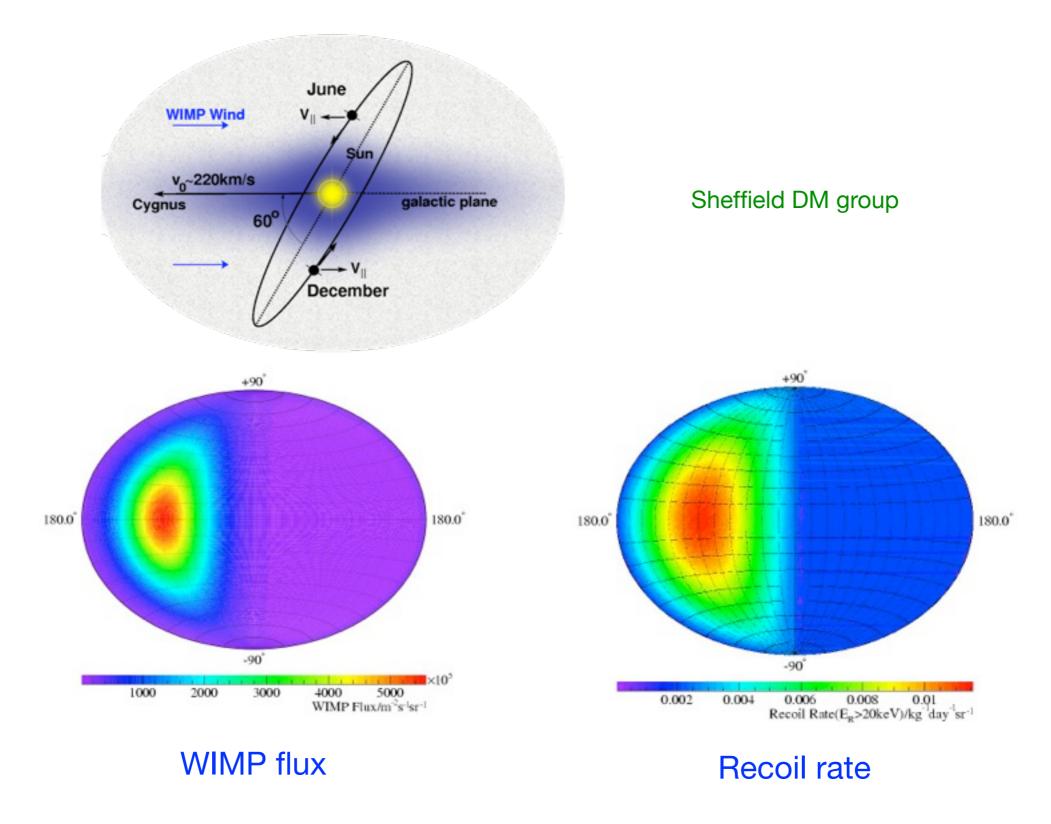
Energy spectrum has characteristic energy which depends on the WIMP mass, target mass and velocity dispersion:

$$E_{\rm R} = \frac{2\mu_{A\chi}^2 v_{\rm c}^2}{m_A} \qquad \qquad \propto m_{\chi}^2 \qquad m_{\chi} \ll m_A$$
$$\sim \text{const} \qquad m_{\chi} \gg m_A$$

$$\log_{10}\left(rac{{
m d}R}{{
m d}E}
ight)$$
 -3 -4 -5 -6 0 20 40 60 80 100 $E/(1\,{
m keV})$

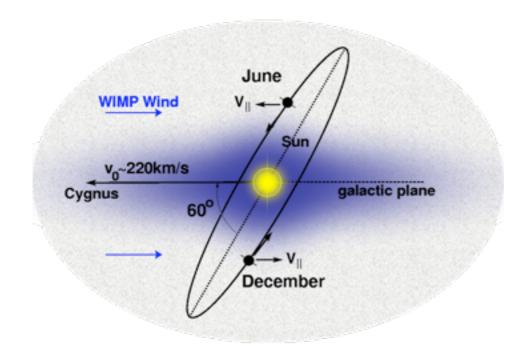
Differential event rate: Ge and Xe $m_X = 50$, 100, 200 GeV

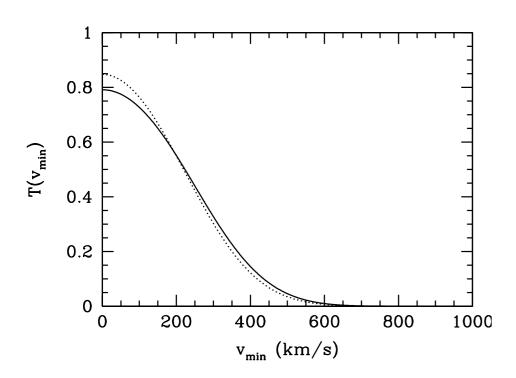
Direction dependence Spergel

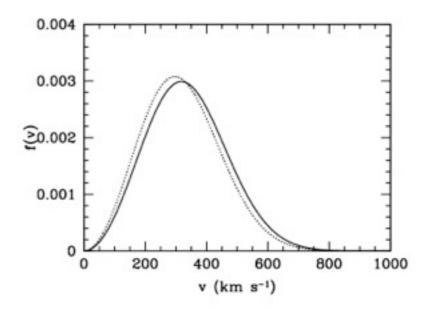


Recoil rate largest in direction opposite to direction of Solar motion. Ratio of rates in rear and forward directions is large.

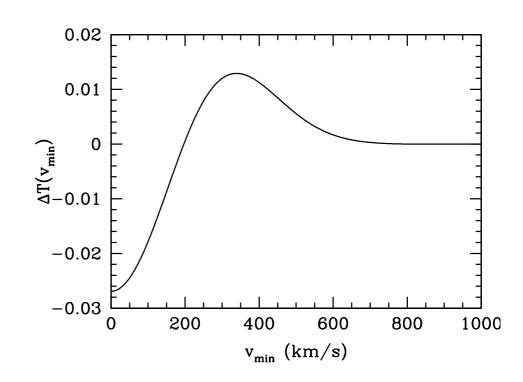
Annual modulation Drukier, Freese & Spergel







Maxwell-Boltzmann speed dist. detector rest frame (summer and winter)



Astrophysical uncertainties i) observations

Local density:

Mass modelling: e.g. Widrow et al., Catena & Ullio, Weber and de Boer, Fornasa & Green in prep model for the MW (luminous components + halo) + multiple data sets (rotation curve, velocity dispersions of halo stars, local surface mass density, total mass...).

~10% statistical errors, central values vary in range $\rho_0 = (0.3-0.4)\,\mathrm{GeV\,cm}^{-3}$

Model independent/minimal assumption methods e.g. Salucci et al. Gabari et al. give consistent values, but with significantly larger errors.

<u>Local circular speed:</u>

Reid & Brunthaler proper motion of Sgr A*:

$$v_{\phi,\odot} \sim (250 \pm 10) \,\mathrm{km \, s}^{-1}$$

Bovy et al. APOGEE data (l.o.s. v of 3000 stars):

$$v_{\phi,\odot} = (242^{+10}_{-3}) \,\mathrm{km} \,\mathrm{s}^{-1}$$

 $v_{\rm c} = (218 \pm 6) \,\mathrm{km} \,\mathrm{s}^{-1}$

implies φ component of Sun's motion wrt Local Standard of Rest (LSR) larger than thought or LSR orbit non-circular.

McMillan & Binney dropping flat rotation curve assumption: $v_c = (200 - 280) \, \mathrm{km \, s}^{-1}$

n.b. Standard halo has one-to-one relationship between circular speed and velocity dispersion & peak speed, but in general this isn't the case.

Local escape speed:

Smith et al, high velocity stars from the RAVE survey assume $f(|\mathbf{v}|) \propto (v_{\rm esc} - |\mathbf{v}|)^k$ with 2.7< k<4.7 (motivated by simulations). $498\,\mathrm{km\,s^{-1}} < v_{\rm esc} < 608\,\mathrm{km\,s^{-1}}$ median likelihood: $v_{\rm esc} = 544\,\mathrm{km\,s^{-1}}$

Summary of observations of MW properties:

Traditional values of circular speed and local density ($v_c=220 \text{ km s}^{-1}$ and $\rho_0=0.3 \text{ GeV cm}^{-3}$), are fairly consistent with recent determinations, which have ~10% statistical errors (but systematic uncertainties from modelling are still significantly larger).

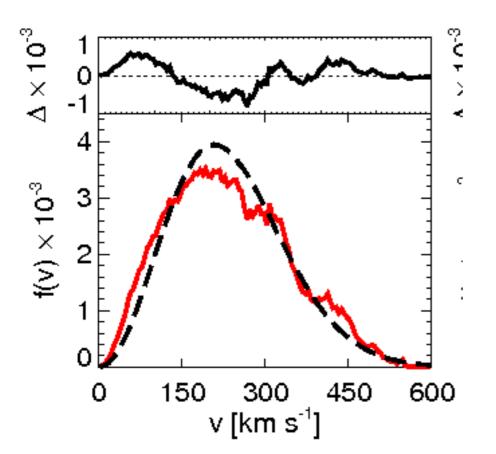
ii) simulations

Systematic deviations from multi-variate gaussian: more low speed particles, peak of distribution lower/flatter.

Features in tail of dist, 'debris flows', incompletely phased mixed material. Lisanti & Spergel; Kuhlen, Lisanti & Spergel

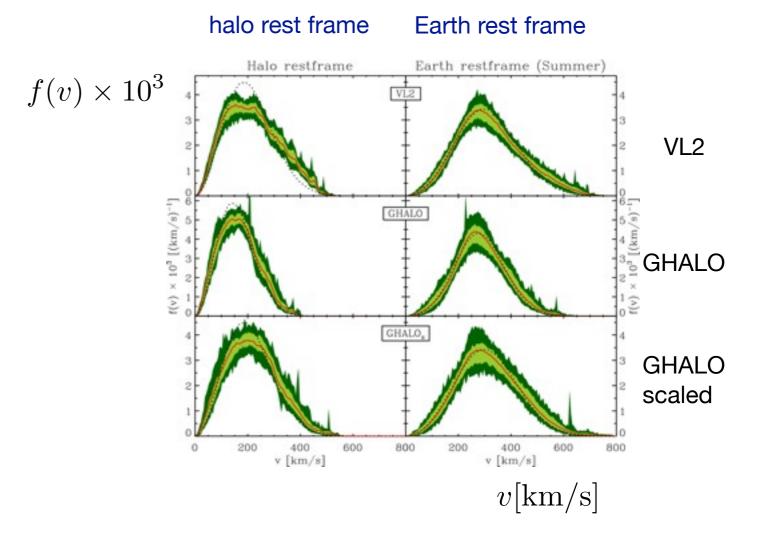
Deviations less pronounced in lab frame than Galactic rest frame.

Vogelsberger et al.



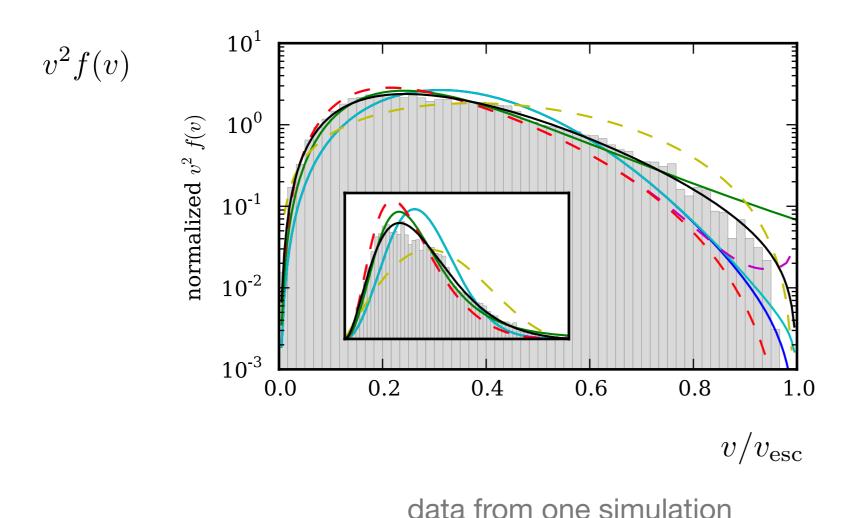
Aquarius simulation data, best fit multi-variate Gaussian

Kuhlen et al.



Various functional forms for f(v) proposed.

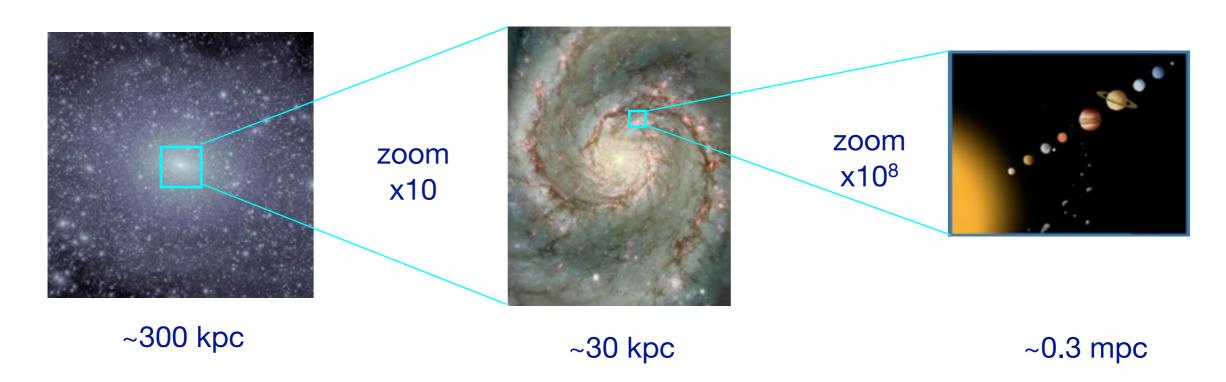
Hard to fit shape of bulk of distribution and tail with a single, simple function:

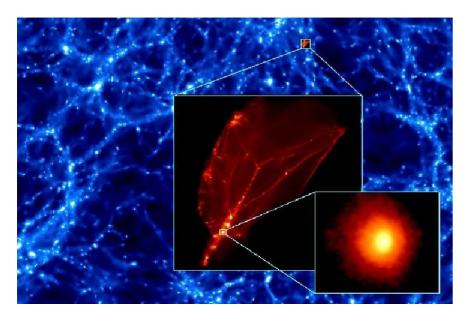


data nom one omidation
 Mao, Strigari & Wechsler
 SHM
 Lisanti et al. double power law
 Tsallis
 Eddington
 Osipkov-Merritt
$\beta = 0.5$

Caveats:

a) scales resolved by simulations are many orders of magnitude larger than those probed by direct detection experiments





microhalo simulation Diemand, Moore & Stadel

Resolution of best Milky Way simulations is many orders of magnitude larger than the mass of the first WIMP microhalos to form

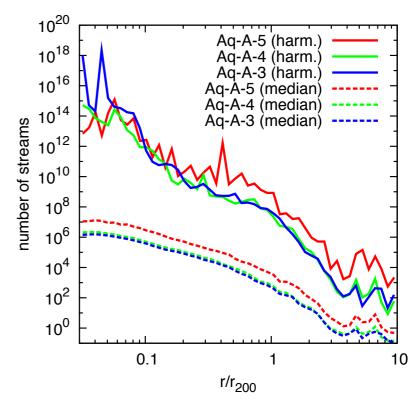
fine structure in ultra-local DM velocity distribution?

Vogelsberger & White:

Follow the fine-grained phase-space distribution, in Aquarius simulations of Milky Way like halos.

From evolution of density deduce ultra-local DM distribution consists of a huge number of streams (but this assumes ultra-local density= local density).

At solar radius <1% of particles are in streams with $\rho > 0.01\rho_0$.



number of streams as a function of radius calculated using harmonic mean/median stream density

Schneider, Krauss & Moore:

Simulate evolution of microhalos. Estimate tidal disruption and heating from encounters with stars, produces 10²-10⁴ streams in solar neighbourhood.

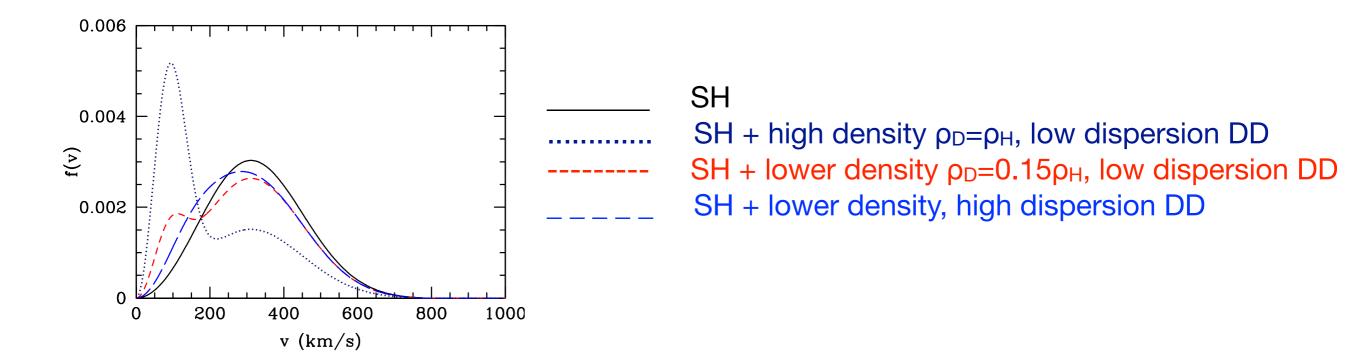
not-so fine structure:

Purcell, Zentner & Wang DM component of Sagittarius leading stream may pass through the solar neighbourhood (as originally suggested by Freese, Gondolo & Newberg).

b) effect of baryons on DM speed distribution?

Sub-halos merging at z<1 preferentially dragged towards disc, where they're destroyed leading to the formation of a co-rotating dark disc. Read et al., Bruch et al., Ling et al.

Could have a significant effect if density is high and velocity dispersion low.



Properties of dark disc are uncertain (simulating baryonic physics and forming Milky Way-like galaxies is hard).

Purcell, Bullock & Kaplinghat to be consistent with observed properties of thick disc, MW's merger history must be quiescent compared with typical Λ CDM merger histories, hence DD density must be relatively low, <0.2 ρ_H . Also dispersion larger than stellar thick disk.

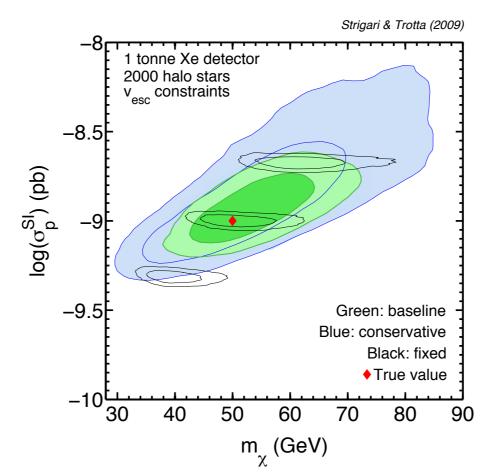
<u>Consequences</u>

Realisation that uncertainties in f(v) will affect signals goes right the way back to the early direct detection papers in the 1980s (e.g. Drukier, Freese & Spergel).

Density:

Event rate proportional to product of σ and ρ , therefore uncertainties in ρ translate directly into uncertainties in σ , same for all DD experiments (but affects comparisons with e.g. collider constraints on σ).

Strigari & Trotta uncertainty leads to bias in determination of WIMP mass:



Circular speed (standard halo):

Shifts exclusion limits, similar, but not identical, effect for all experiments.

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..... $v_c = 195 \text{ km/s}$

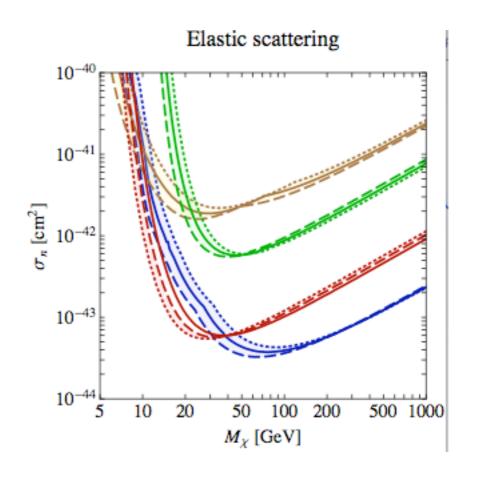
 $_{\rm c}$ v_c=220 km/s

 $- - - v_c = 255 \text{ km/s}$

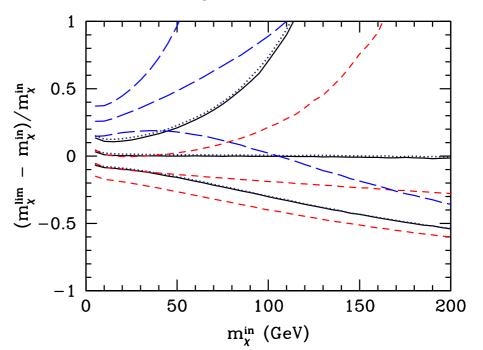
(old)CDMSII Si, CDMSII Ge CRESST, ZENON 10

Bias in future WIMP mass determination:

$$E_{
m R} = rac{2\mu_{A\chi}^2 v_{
m c}^2}{m_A}$$
 $rac{\Delta m_\chi}{m_\chi} = [1+(m_\chi/m_{
m A})] rac{\Delta v_{
m c}}{v_{
m c}}$ $m v_c = 220~km/s$ $m 200~km/s$ $m 280~km/s$



fractional mass limits from a simulated ideal Ge experiment, $\sigma = 10^{-8}$ pb



Shape of velocity distribution

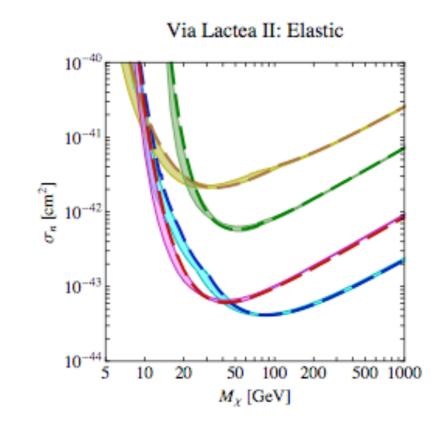
Differential event rate is proportional to integral over speed distribution so exclusion limits are relatively insensitive to exact shape of velocity distribution:

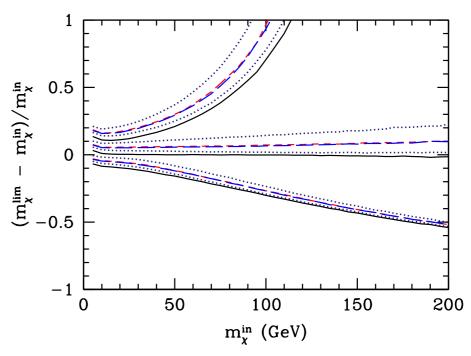
(smallish) change in shape/stochastic uncertainty in exclusion limits.

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(old)CDMSII Si, CDMSII Ge CRESST, XENON 10

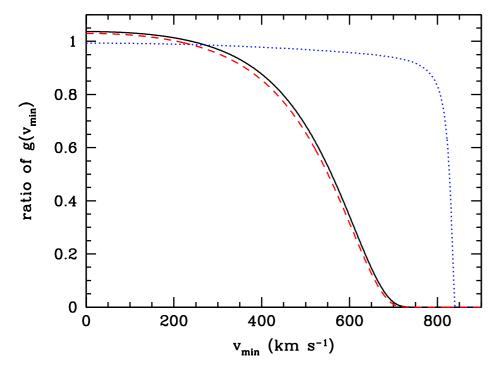
2-5% bias in future WIMP mass determination.





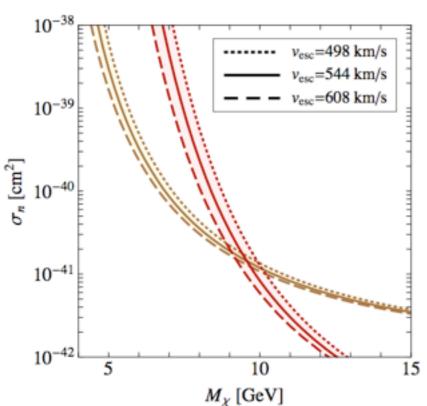
Escape speed & shape of high v tail

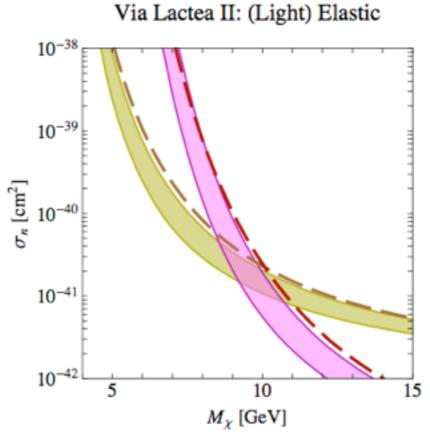
Can have significant effect on event rates/exclusion limits for light WIMPs:



Ratio of speed integral to that of Maxwellian with sharp cut-off at $v_{\rm esc} = 608\,{\rm km\,s}^{-1}$:

same f(v) neglecting Earth's orbit Lisanti et al. k=1.5 $v_{\rm esc} = 498\,{\rm km\,s}^{-1}$ Lisanti et al. neglecting Earth's orbit



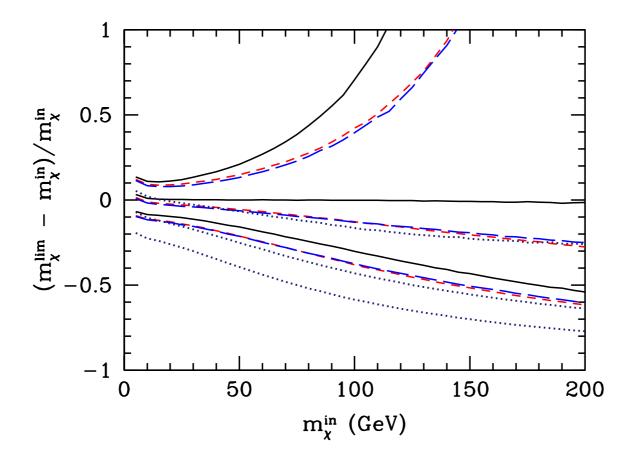


McCabe

(old)CDMSII Si, XENON 10

Dark disc

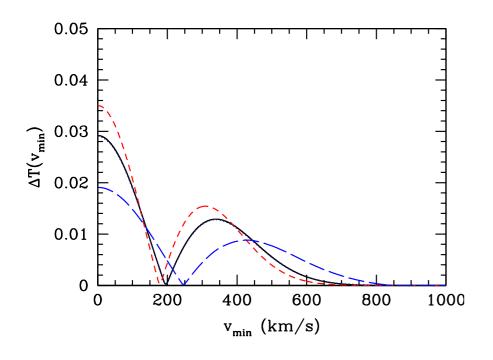
Could significantly bias mass determination, if density sufficiently high and/or velocity dispersion low.



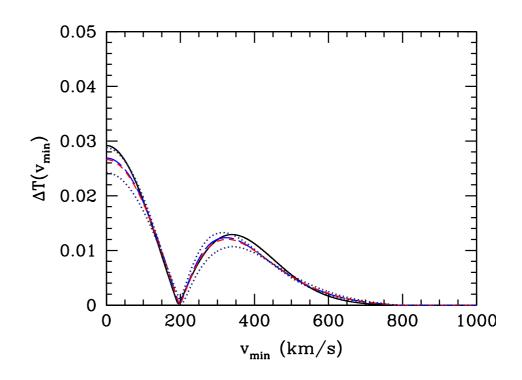
Annual modulation

Arises from small shift in speed distribution due to Earth's orbit.

Amplitude (and phase) sensitive to detailed shape of speed distribution.



SHM varying v_c



varying shape of f(v)

Direction dependence

Rear-front directional asymmetry is robust, but peak direction of high energy recoils can change. Kuhlen et al.

Strategies i) integrate out

Fox, Liu & Weiner

Compare experiments in g(v_{min}) space:

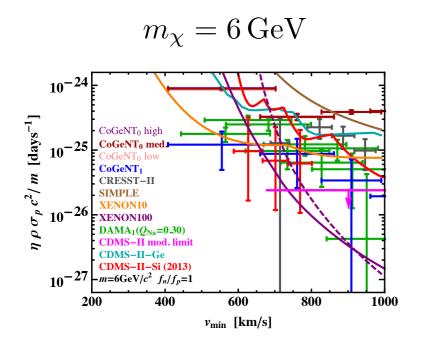
$$g(v_{\min}) = \int_{v_{\min}}^{\infty} \frac{f(v)}{v} dv \qquad v_{\min} = \left(\frac{E(m_A + m_\chi)^2}{2m_A m_\chi^2}\right)^{1/2}$$

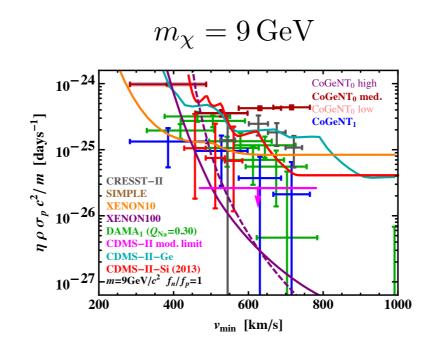
v_{min} values probed by each experiment depend on, unknown, WIMP mass, therefore need to do comparison for each mass of interest.

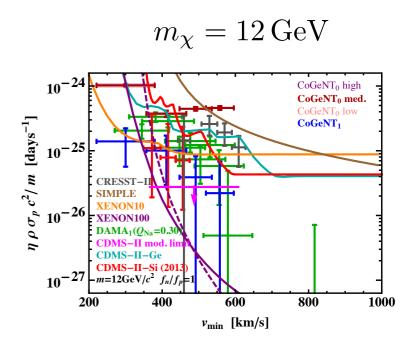
Can incorporate experimental energy resolution and efficiency Gondolo & Gelmini, and also annual modulation signals. Frandsen et al.; Herrero-Garcia, Schwetz & Zupan.

Extremely powerful for checking consistency of signals and exclusion limits. Frandsen et al.; Del Nobile, Gelmini, Gondolo & Huh.

Normalised g(v_{min}) versus v_{min} Del Nobile, Gelmini, Gondolo & Huh





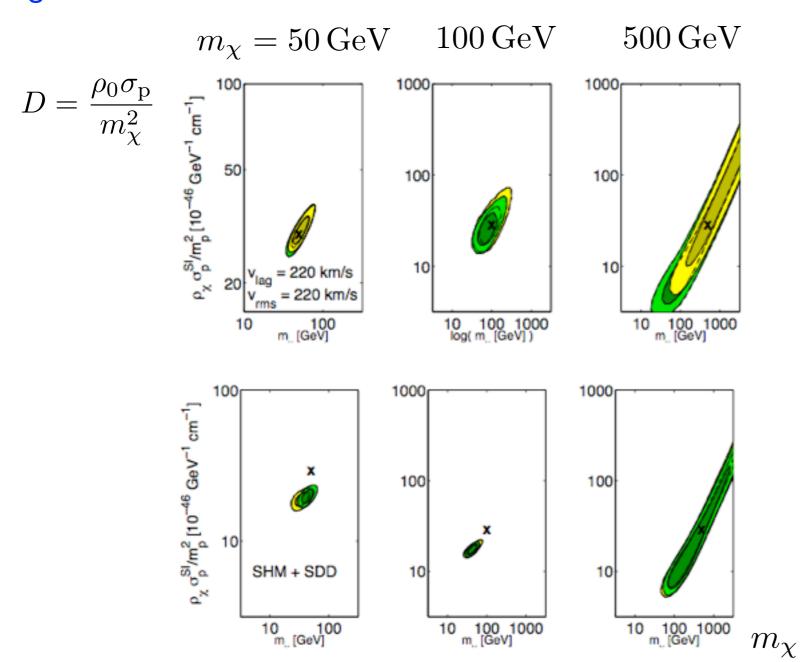


Strategies ii) marginalise over

Parameterize f(v) and/or Milky Way model and marginalise over these parameters, possibly including astrophysical data too e.g. stellar kinematics.

Strigari & Trotta; Peter x2; Pato et al. x2; Lee & Peter; Billard, Meyet & Santos; Alves, Hedri & Wacker; Kavanagh & Green x2; Friedland & Shoemaker

If actual shape of f(v) is similar to assumed shape this works well, but if not can get significant biases:



Peter simulated data from future tonne scale Xe, Ar & Ge expts, analysed assuming standard halo model (allowing v_{lag} & v_{rms} to vary).

standard halo model in

standard halo model + dark disc in

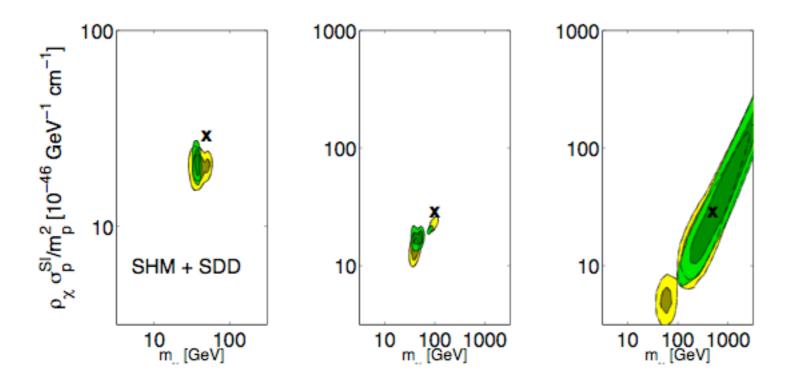
Parameterizing speed distribution

With a single experiment can't say anything about the WIMP mass without making assumptions about f(v) (recoil energies depend on speeds and mass).

But with multiple experiments can break this degeneracy. Drees & Shan; Peter

Peter Use empirical parameterization of f(v), and constrain its parameters along with mass & cross-section.

First approach: piece-wise constant in bins



standard halo model + dark disc in

Better than assuming wrong f(v), but m_{χ} & σ both biased.

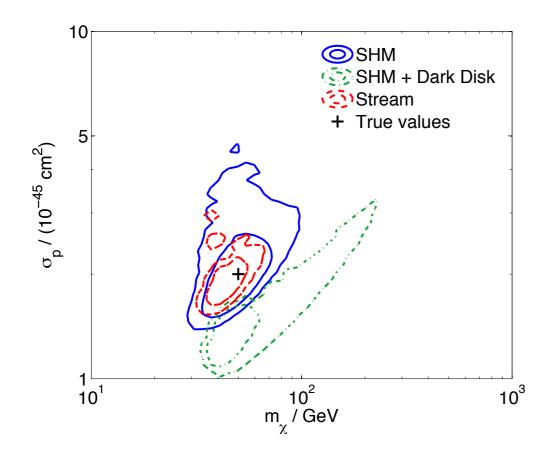
Kavanagh & Green

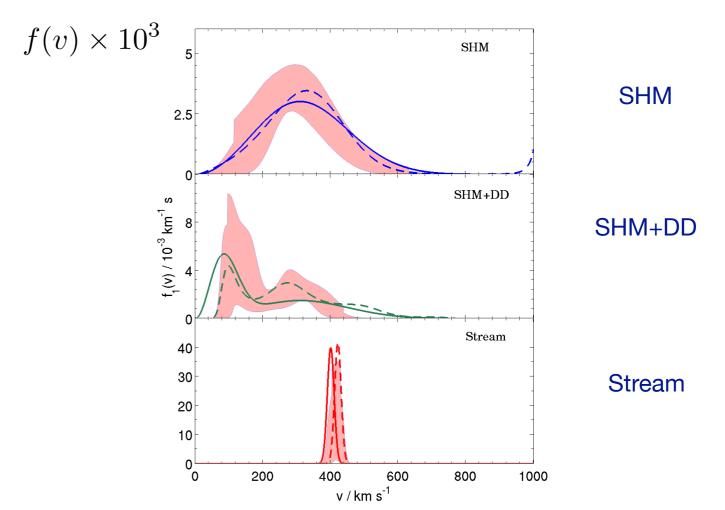
Want parameterisation without fixed scales, and with ability to accommodate features in speed distribution.

Since $f(v) \ge 0$, parameterise log of f(v) in shifted Legendre polynomials:

$$f(v) \propto \exp \left\{ -\sum_{k=0}^{N} a_k \bar{P}_k(v/v_{\text{max}}) \right\}$$

Gives good reconstruction of WIMP mass even for extreme input f(v) (stream or dark disc), and allows f(v) to be reconstructed:





Summary

• Direct detection energy spectrum depends on the local dark matter density, ρ_0 , and velocity distribution, f(v):

Iocal DM density → normalisation of event rate, and hence σ velocity dispersion → characteristic scale of energy spectrum and hence m_χ shape of WIMP velocity distribution → event rate for light WIMPs and amplitude and phase of annual modulation signal

- Determinations of ρ_0 and v_c have ~10% statistical errors, but systematic errors are larger.
- Can assess compatibility of signals/exclusion limits in speed integral, g(v_{min}), space ('integrating out the astrophysics').
- Parameterising f(v)/Milky Way model and marginalising works well **if** actual shape of f(v) is close to assumed shape.
- For unbiased mass measurement use a suitable empirical parameterisation (e.g. shifted Legendre polynomials), and probe f(v) too.