## EW and QCD One-Loop Amplitudes with RECOLA

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After the discovery of the Higgs boson:

- Precise investigation of the Standard Model and beyond
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- Precise investigation of the Standard Model and beyond
- Need to have under control potential large corrections for several processes
- QCD corrections are known to be large
- EW corrections can be enhanced:
- in high energy regions (Sudakov log's)
- in Higgs physics
- by photon emission (mass-singular log's)

Let's concentrate on one loop corrections


## Les Houches wishlist 2013 at one loop

- QCD:

$$
p p \rightarrow t \bar{t} H, \quad p p \rightarrow t \bar{t}+j \quad \text { (top decays) }
$$

- EW:

$$
\begin{aligned}
& p p \rightarrow 3 j, \\
& p p \rightarrow t \bar{t}, \quad p p \rightarrow t \bar{t} H, \quad p p \rightarrow t \bar{t}+j \quad \text { (top decays) } \\
& p p \rightarrow V+2 j, \quad p p \rightarrow V V^{\prime}, \quad p p \rightarrow V V+j, \\
& p p \rightarrow V V+2 j \quad p p \rightarrow V V^{\prime} \gamma, \quad p p \rightarrow V V^{\prime} V^{\prime \prime}, \\
& \left(V, V^{\prime}, V^{\prime \prime}=W, Z\right. \text { decay leptonically) }
\end{aligned}
$$

- Many issues at hadronic level:

Multi-channel MCs, Real emission, PDFs, Parton Shower, ...

- At least the partonic processes should be automatized

Many codes have been produced:

| MCFM | Campbell, Ellis |
| :--- | :--- |
| BlackHat | Berger, Bern, Dixon, Febres Cordero, Forde, Ita, Kosower, Maître |
| VBFNLO | Arnold, Bähr, Bozzi, Campanario, Englert, Figy, Greiner, Hackstein, <br> Hankele, Jäger, Klämke, Kubocz, Oleari, Plätzer, Prestel, Worek, <br> Zeppenfeld |
| HELAC-NLO | Bevilacqua, Czakon, Garzelli, van Hameren, Kardos, Papadopoulos, <br> Pittau, Worek |
| GoSam | Cullen, Greiner, Heinrich, Luisoni, Mastrolia, Ossola, Reiter, <br> Tramontano |
| NJet | Badger, Biedermann, Uwer, Yundin |
| AMC@NLO | Hirschi, Frederix, Frixione, Garzelli, Maltoni, Pittau |
| OpenLoops | Cascioli, Maierh"ofer, Pozzorini |

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## REcursive Computation of One Loop Amplitudes

- In the full Standard Model
- Based on recursive relations for off-shell currents


## Off-shell tree currents

Given a process with $L$ external legs:

$$
\underbrace{\mathcal{P}_{1}+\ldots+\mathcal{P}_{L-1}}_{\text {primary }}+\underbrace{\mathcal{P}_{L}}_{\text {last }} \rightarrow 0
$$

Off-shell current of a particle $\mathcal{P}$ with $n$ primary legs:
Def: Amplitude made of $n$ primary on-sheel particles and the off-sheel particle $\mathcal{P}$


List of primary legs

- $w$ is a scalar, spinor or vector
- The off-shell currents for external legs are the wave functions:

$$
\longrightarrow \bullet=u_{\lambda}(p) \quad \longrightarrow=\bar{u}_{\lambda}(p) \quad \rightsquigarrow_{\bullet}=\epsilon_{\lambda}(p) \quad--\bullet=1
$$

- Binary notation for $\left\{l_{1}, \ldots, l_{n}\right\}$ (HELAC):

Binary numbers $1,2,4,8, \ldots, 2^{L-1}$ for the primary legs
$\left\{l_{1}, \ldots, l_{n}\right\}$ can be expressed by $\mathcal{B}_{n}=$ sum of the $n$ binaries
Example: $\quad\{1,2,8\} \quad \rightarrow \quad \mathcal{B}_{3}=1+2+8=11$

Note: The off-shell currents just keep trace of the primary legs used to build them, not the way it has been done.

Example: $\quad$ Process $\begin{aligned} & e^{-}+e^{+}+\tau^{+}+\tau^{-} \rightarrow 0 \\ & 1\end{aligned}$


## Recursion relation for tree amplitudes


(incoming currents) $\times($ coupling $) \times($ propagator $)$

## Recursion relation for tree amplitudes


(incoming currents) $\times$ (coupling) $\times($ propagator $)$

- Recursive procedure:

2-leg currents:


## Recursion relation for tree amplitudes


(incoming currents) $\times($ coupling $) \times($ propagator $)$

- Recursive procedure:


3-leg currents:




## Recursion relation for tree amplitudes


(incoming currents) $\times$ (coupling) $\times($ propagator $)$

- Recursive procedure:




4-leg currents:



## Recursion relation for tree amplitudes



- Recursive procedure:

etc. . . .


## Recursion relation for tree amplitudes


(incoming currents) $\times$ (coupling) $\times($ propagator $)$

- Recursive procedure:

etc. . . .
- Amplitude: $\mathcal{A}=w\left(\mathcal{P}_{L}, 2^{L-1}-1\right) \times(\text { propagator })^{-1} \times w\left(\mathcal{P}_{L}, 2^{L-1}\right)$


## Recursion relation for loop amplitudes

General form of the amplitude: Tensor Coefficients (TCs)

$$
\begin{gathered}
\mathcal{A}=\sum_{t} c_{\mu_{1} \ldots \mu_{r_{t}}}^{(t)} T_{(t)}^{\mu_{1} \ldots \mu_{r_{t}}} \rightarrow \text { Tensor Integrals (TIs) } \\
T_{(t)}^{\mu_{1} \cdots \mu_{r_{t}}}=\int \frac{d^{n} q q^{\mu_{1}} \cdots q^{\mu_{r_{t}}}}{D_{0}^{(t)} \cdots D_{k_{t}}^{(t)}} \quad D_{k_{t}}^{(t)}=\left(q+p_{k_{t}}^{(t)}\right)^{2}-\left(m_{k_{t}}^{(t)}\right)^{2}
\end{gathered}
$$

Remark: Indices $\mu_{1}, \ldots, \mu_{r_{t}}$ are computed numerically in $\mathrm{D}=4$ dimensions.
$\rightsquigarrow$ The rational part R2 is computed separatly: Inclusion of effective tree-level Feynman rules (as for the counterterms)
[Draggiotis, Garzelli, Malamos, Papadopoulos, Pittau '09-'10]

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[Draggiotis, Garzelli, Malamos, Papadopoulos, Pittau '09-'10]
Basic idea: Cut the loop line and consider tree diagrams with two more legs.
[A. van Hameren, JHEP 0907 (2009) 088]


Given the loop process

$$
\mathcal{P}_{1}+\ldots+\mathcal{P}_{L} \rightarrow 0
$$

we consider the tree processes

$$
\underbrace{\mathcal{P}_{1}+\ldots+\mathcal{P}_{L}+\mathcal{P}}_{\text {primary }}+\underbrace{\overline{\mathcal{P}}}_{\text {last }} \rightarrow 0 \quad \forall \mathcal{P} \in\{\text { Particle of the SM }\}
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## Problem:

Associated tree diagrams are more than the original loop diagrams:







## Rules to avoid double counting of the associated trees:

Rule 1: $\rightarrow$ Fix starting point of loop flow
The current containing the first external line enters the loop flow first







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Rule 2: $\rightarrow$ Fix direction of loop flow
The 3 currents with the 3 smallest binaries enter the loop flow in fixed order





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- Recursion relation for loop currents

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Remark: Indices $\mu_{1}, \ldots, \mu_{r}$ are symmetrized at each step

- The coefficients $a_{k, r}^{\mu_{1} \cdots \mu_{r}}$ of the last current give the TCs $c_{\mu_{1} \ldots \mu_{r_{t}}}^{(t)}$


## Loop off-shell currents



Sequence of masses in loop propagators


- $i_{k}$ is the tensorial index:

$$
\begin{array}{lll}
i_{k}=0 & \rightarrow & w_{i_{k}}=a_{k, 0} \\
i_{k}=1, \ldots, 4 & \rightarrow & w_{i_{k}}=a_{k, 1}^{\mu_{1}} \\
i_{k}=5, \ldots, 14 & \rightarrow & w_{i_{k}}=a_{k, 2}^{\mu_{2}} \mu_{2}
\end{array}
$$

- Special wave functions for the cutted line:



where the components are $\left(\psi_{i}\right)_{j}=\left(\bar{\psi}_{i}\right)_{j}=\delta_{i j}, \epsilon_{i}^{\mu}=\delta_{i \mu}$.


## Treatment of the colour

Color-flow representation [Maltoni, Paul, Stelzer, Willenbrock '02]:

"usual" gluon with color index $a=1, \ldots, 8$

$$
i, j=1,2,3
$$

$$
\sum_{i}\left(A_{\mu}\right)_{i}^{i}=0
$$

## Feynman rules:

- Multiply gluon fields $A_{\mu}^{a}$ by $\left(\lambda^{a}\right)_{j}^{i} / \sqrt{2}$ and use properties of $\left(\lambda^{a}\right)_{j}^{i}$
- The color part of the Feynman rules is just product of deltas:

$$
{ }_{i_{1}}^{i_{1}} \infty_{i_{2}}^{j_{2}}={ }_{j_{1}}^{i_{1}} \leftrightarrows{ }_{i} i_{2} \times \frac{-i g_{\mu \nu}}{p^{2}}=\delta_{j_{2}}^{i_{1}} \delta_{j_{1}}^{i_{2}} \times \frac{-i g_{\mu \nu}}{p^{2}}
$$

$$
i \rightarrow \bullet \beta=u_{\lambda}(p) \delta_{\beta}^{i} \quad j \longrightarrow \alpha=\bar{u}_{\lambda}(p) \delta_{j}^{\alpha} \quad{ }_{j}^{i} \nsim \bullet{ }_{\beta}^{\alpha}=\epsilon_{\lambda}(p) \delta_{\beta}^{i} \delta_{j}^{\alpha}
$$

Structure of amplitude:

$$
\mathcal{A}_{\beta_{1} \cdots \beta_{n}}^{\alpha_{1} \cdots \alpha_{n}}=\sum_{P\left(\beta_{1}, \ldots, \beta_{n}\right)} \delta_{\beta_{1}}^{\alpha_{1}} \cdots \delta_{\beta_{n}}^{\alpha_{n}} \mathcal{A}_{P}
$$

- Colour-dressed amplitudes:
$\rightarrow$ Compute $\mathcal{A}_{\beta_{1} \cdots \beta_{n}}^{\alpha_{1} \cdots \alpha_{n}}$ for all possible colours $\left(N_{c}^{2 n}\right)$
Squared amplitude: $\quad \overline{\mathcal{M}^{2}}=\sum_{\alpha_{1} \ldots \alpha_{n}, \beta_{1}, \ldots, \beta_{n}}\left(\mathcal{A}_{\beta_{1} \cdots \beta_{n}}^{\alpha_{1} \cdots \alpha_{n}}\right)^{*} \mathcal{A}_{\beta_{1} \ldots \beta_{n}}^{\alpha_{1} \ldots \alpha_{n}}$
It requires colour-dressed currents
- Structure-dressed (or colour-ordered) amplitudes:
$\rightarrow$ Compute $\mathcal{A}_{P}$ for all possible $P(n!)$
Squared amplitude: $\quad \overline{\mathcal{M}^{2}}=\sum_{P, P^{\prime}} \mathcal{A}_{P}^{*} C_{P P^{\prime}} \mathcal{A}_{P^{\prime}}$
It requires structure-dressed currents


## Structure-dressed off-shell currents

Colour structure of off-shell current:


> with all possible permutations of

$$
\beta_{1}, \ldots \beta_{n}, j
$$

In the recursion procedure:

- Saturated parts of incoming currents multiply
- Open parts of incoming currents are contracted

Optimization: Compute once currents differing just by the colour structure
$\rightsquigarrow$ Overcome lack of colour factorization
Example:


The code RECOLA is structured in two parts:

- Generation of the recursion procedure (to be run once)
- A current-index is given to all currents of the recursion procedure
- A branch-index is given to each step (branch) of the resursion procedure
- Identify currents differing just by the colour structure
- To each branch are associated the relevant indices
- The list of all needed Tls is generated

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- Identify currents differing just by the colour structure
- To each branch are associated the relevant indices
- The list of all needed TIs is generated
- Computation of the currents (to be run at each PS point)
- All needed TIs are computed
- The currents for all polarizations and colors are computed
- The last currents are contracted with the TIs to give the amplidute


## Features of RECOLA (fortran 95)

- Full Standard Model in the complex mass scheme with:
- Feynman rules for rational parts
- Feynman rules for on-shell Counterterms
- Selection of resonant contributions
- Need external libraries for TIs (link to the COLLIER library implemented)
- Numerical check of cancellation of UV divergences
- Mass and dimensional regularization for collinear and soft singularities
- Select/unselect powers of $\alpha_{s}$ in the amplitude
- Computation of Colour- and Spin-correlations
- Optimizations:
- Helicity sum avoids recalculation of currents
- Use conservation of helicity for massless fermions
- Use partial factorization of colour structure


## Performances

- Memory needed for executables, object files and libraries: negligible
- RAM needed: less than 2 Gbyte also for complicated processes
- CPU time (processor Intel(R) Core(TM) i5-2400 CPU @ 3.10GHz):

| Process | $\begin{gathered} t_{\mathrm{TIs}} \\ (\mathrm{COLLIER}) \end{gathered}$ | $\begin{aligned} & t_{\text {gen }} \quad t_{\mathrm{TCs}} \\ & \text { (single helicity) } \end{aligned}$ | $t_{\text {gen }} \quad t_{\mathrm{TCs}}$ <br> (partial hel. sum) | $t_{\text {gen }} \quad t_{\mathrm{TCs}}$ <br> (helicity sum) |
| :---: | :---: | :---: | :---: | :---: |
| $u \bar{u} \rightarrow \underset{(\text { QCD) }}{\rightarrow W^{+}} W^{-} g$ | 2.8 ms | $\begin{aligned} & 0.3 \mathrm{~s} \quad 0.6 \mathrm{~ms} \\ & \text { (hel: }-+-+- \text { ) } \end{aligned}$ | $\begin{gathered} 0.4 \mathrm{~s} \\ \text { (hel: } \mathrm{S} \mathrm{~S}-+\mathrm{S} \text { ) } \end{gathered}$ | $\begin{aligned} & 1.6 \mathrm{~s} \quad 9.8 \mathrm{~ms} \\ & \text { (hel: S S S S S } \end{aligned}$ |
| $u \bar{d} \underset{(\mathrm{QCD})}{W^{+}} g g g$ | 130 ms | $\begin{aligned} & 14 \mathrm{~s} \quad 14 \mathrm{~ms} \\ & \text { (hel: - + - - - } \end{aligned}$ | $\begin{aligned} & 25 \mathrm{~s} \quad 76 \mathrm{~ms} \\ & \text { (hel: S S - S S S) } \end{aligned}$ | $\begin{array}{lr} 52 \mathrm{~s} & 221 \mathrm{~ms} \\ \text { (hel: S S S S S S) } \end{array}$ |
| $u g \underset{(\mathrm{EW})}{\rightarrow u g} Z$ | 8.2 ms | $\begin{gathered} 0.5 \mathrm{~s} \quad 1.4 \mathrm{~ms} \\ \text { (hel: - -- - - } \end{gathered}$ | $\begin{gathered} 1.0 \mathrm{~s} \\ \text { (hel: S S S S -) } \end{gathered}$ | $\begin{gathered} 2.2 \mathrm{~s} \\ \text { (hel: S S S S S) } \end{gathered}$ |
| $u g \rightarrow \underset{(\mathrm{EW})}{u g} \tau^{-} \tau^{+}$ | 28 ms | $1.3 \mathrm{~s} \quad 2.5 \mathrm{~ms}$ (hel: - - - - + | 2.0 s 14.2 ms (hel: S S S S - +) | $\begin{aligned} & 3.8 \mathrm{~s} \quad 29.0 \mathrm{~ms} \\ & \text { (hel: S S S S S S) } \end{aligned}$ |

$S$ = sum over helicity

## Summary

- Efficient automatization for elementary EW and QCD processes at NLO
- Recursion relations $\rightarrow$ good tool also in the EW sector
- used for EW corrections to $p p \rightarrow Z+2 j \quad \rightarrow$ Talk by Ansgar Denner


## Outlook

- Publication of the code
$\rightarrow$ Robust checks
$\rightarrow$ Implement dynamical running of $\alpha_{s}$
$\rightarrow$ Allow extensions to other Models
$\rightarrow$ Prepare for MC over polarizations and colours
- Let's compute other LHC processes

