

Lepton Flavour Violation

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Overview

- New Physics and Flavour
- Neutral Lepton Flavour Violation
 - Neutrino Oscillation
 - Dirac vs Majorana
- Charged Lepton Flavour Violation
- Effective Operator Approach
- The BSM Flavour Problem
- BSM Models
 - Effective Mass and Seesaw
 - Left–Right Symmetric Models
- Conclusions

New Physics and Flavour



- Distinct pattern of fermion masses and mixing
 - Quarks

$$m_u: m_c: m_t \sim \lambda^8: \lambda^4: 1, \quad m_d: m_s: m_b \sim \lambda^4: \lambda^2: 1, \quad \lambda \sim 0.2,$$

$$\theta_{12}^Q \sim 13^\circ, \quad \theta_{23}^Q \sim 2.4^\circ, \quad \theta_{13}^Q \sim 0.2^\circ.$$

Leptons

New Physics and Flavour



- > What is the origin behind these patterns?
- Most (except neutrino mixing?) seem to be non-random
- Quark and lepton mixing distinctly different
- Is this connected to the light neutrino masses / potential Majorana character?
- Standard Model tells us nothing about the flavour structure (+ neutrinos are massless)
- There is no theory of flavour
- Need to find signs of New Physics to make progress

Neutrino Oscillations



- Neutrino interaction eigenstates different from mass eigenstates
 - Neutrino flavour can change through propagation \rightarrow LFV

$$\nu_{i} = U_{\alpha i} \nu_{\alpha}, \qquad \nu_{i}(t) = e^{-i(E_{i}t - p_{i}x)} \nu_{i}(0)$$
$$\implies P_{\alpha \to \beta} = \sin^{2} 2\theta \sin^{2} \left(1.27 \frac{\Delta m^{2}}{eV^{2}} \frac{L/km}{E/GeV} \right)$$

- Solar Neutrino Oscillations
 - Large Mixing
- Atmospheric Oscillations
 - \approx Maximal Mixing
- Reactor and Accelerator Neutrinos
 - $\sin^2 2\theta_{13} = 0.092 \pm 0.021$
- Experimental Unknowns and Anomalies
 - CP Violation? Sign of Δm_{23} ? Sterile Neutrinos?





Dirac vs. Majorana

 ν_R

Two possibilities to define Fermion mass





 $\nu_L = \overline{\nu}_L$







Dirac mass analogous to other fermions but with ${}^{m_{\nu}}/_{\Lambda_{EW}} \approx 10^{-12}$ couplings to Higgs



Majorana mass, using only a left-handed neutrino → Lepton Number Violation



Dirac vs. Majorana

 $\nu_{L'}$



Two possibilities to define Fermion mass





 $\nu_L = \overline{\nu}_L - \frac{10 \text{ theorists recommend}}{10 \text{ theorists recommend}}$



Dirac mass analogous to other fermions but with ${m_{\nu}}/{\Lambda_{EW}} \approx 10^{-12}$ couplings to Higgs



Majorana mass, using only a left-handed neutrino → Lepton Number Violation



Absolute Neutrino Mass





Neutrinoless Double Beta Decay

Energy Endpoint in Beta Decay

 $m_{\beta\beta} = |\Sigma_i U_{ei}^2 m_{\nu_i}| < 0.2 \cdots 1.0 \text{ eV}$

 $m_{\beta}^2 = \Sigma_i |U_{ei}|^2 m_{\nu_i}^2 < (2.2 \text{ eV})^2$

Impact on Large Scale Structure

 $\Sigma = \Sigma_i m_{\nu_i} < 0.3 \cdots 1.0 \text{ eV}$



Lepton Flavour versus Lepton Number Violation



Neutrinoless double beta decay



 $\Delta L_e = 2, \Delta L_\mu = 0, \Delta L = 2$ Lepton Number Violation





 $\mu^+ \rightarrow e^-$ conversion in nuclei



 $\Delta L_e = 1, \Delta L_{\mu} = -1, \Delta L = 0$ Lepton Flavour Violation

 $\Delta L_e = 1, \Delta L_\mu = 1, \Delta L = 2$ Lepton Flavour Violation + Lepton Number Violation

Charged Lepton Flavour Violation



- Lepton Flavour practically conserved in the vSM
 - LFV is clear sign for BSM physics

$$Br(\mu \to e\gamma) = \frac{3\alpha}{32\pi} \left| \sum_{i} U_{\mu i}^* U_{ei} \frac{\Delta m_{1i}^2}{m_W^2} \right|^2 \approx 10^{-56}$$



- Flavour violation in the quark and neutrino sector
 - Strong case to look for charged LFV
- LFV can shed light on
 - Grand Unification models
 - Flavour symmetries
 - Origin of flavour



Charged Lepton Flavour Violation



- Large number of observables
 - $\mu \rightarrow e\gamma$, $\mu^- \rightarrow e^-e^-e^+$
 - μe conversion in nuclei: $\mu^- + A \rightarrow e^- + A$, $\mu^- + A \rightarrow e^+ + A'$
 - τ decays



Charged Lepton Flavour Violation



- Current experimental situation (future sensitivity)
 - $Br(\mu \to e\gamma) < 5.7 \times 10^{-13}$ @ MEG ($\approx 10^{-13}$ @ MEG)
 - $Br(\mu \to eee) < 1.0 \times 10^{-12}$ @ SINDRUM-I ($\approx 10^{-15}$ @ MuSIC, µ3e)
 - $R(\mu \to e, Au) < 7.0 \times 10^{-13}$ @ SINDRUM-II ($\approx 10^{-16}$ @ COMET, Mu2e)
 - $Br(\tau \rightarrow l\gamma) < 4.0 \times 10^{-8}$ @ Belle, Babar ($\approx 10^{-9}$ @ SuperB factory)



Lepton Flavour Physics



Connection between observables



Effective Operator Approach

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- Describe LFV processes using effective, non-renormalizable interactions,

$$\mathcal{L}_{\rm eff} = \mathcal{L}_{\rm SM} + \frac{1}{M_{\rm NP}^2} \mathcal{L}_5 + \frac{1}{M_{\rm NP}^2} \mathcal{L}_6 + \cdots$$

Dimension-5

$$\mathcal{O}_5 = 1/2 \, (\overline{L}_i \cdot H) (H^+ \cdot L_j)^c$$

Neutrino Majorana masses

 $\mathcal{L}_n = \sum_i C_n^i \mathcal{O}_n^i (\text{SM fields}) + h.c.$

- Dimension-6
 - Two Lepton-Higgs-Photon $\mathcal{O}_6(ll\gamma H) = \overline{L}_i \sigma^{\mu\nu} e_j^c H^+ F_{\mu\nu} \longrightarrow \mu \to e\gamma \text{ etc., } g-2, \text{ EDMs}$
 - Four Lepton $\mathcal{O}_6(llll) = (\overline{L}_i \gamma^{\mu} L_j)(\overline{L}_k \gamma^{\mu} L_l), \text{ etc.}$
 - Two Lepton-Two Quark $\mathcal{O}_6(llqq) = (\bar{L}_i \gamma^{\mu} L_j)(\bar{Q}_k \gamma^{\mu} Q_l), \text{ etc.}$
- $\mu \rightarrow eee$ etc., NSIs in neutrino oscillations

 $\mu \rightarrow e$ conversion in nuclei, Meson decays

Effective Operator Approach

15 / 33

Different processes probe different (combinations of) operators





de Gouvea, Vogel '13

Effective Operator Approach



- Different flavour transitions probe flavour structure
- Compare two example discrete flavour symmetries

 $A_4 \times Z_3 \times U(1)_{FN}$

Dipole operator coefficients

$$C \approx \begin{pmatrix} \lambda^2 \epsilon & \lambda^2 \epsilon^2 & \lambda^2 \epsilon^2 \\ \lambda \epsilon^2 & \lambda \epsilon & \lambda \epsilon^2 \\ \epsilon^2 & \epsilon^2 & \epsilon \end{pmatrix},$$

$$S_4 \times Z_3 \times U(1)_{FN}$$

$$C \approx \begin{pmatrix} \lambda^2 \epsilon^2 & \lambda^2 \epsilon^2 \epsilon' & \lambda^2 \epsilon^2 \epsilon \\ \lambda \epsilon \epsilon' & \lambda \epsilon & \lambda \epsilon \epsilon'^2 \\ \epsilon \epsilon' & \epsilon \epsilon'^2 & \epsilon \end{pmatrix},$$

LFV process rates

16 / 33

 $Br(\mu \to e\gamma) \approx Br(\tau \to e\gamma) \approx Br(\tau \to \mu\gamma).$

$$Br(\mu \to e\gamma) \approx Br(\tau \to e\gamma) \gg Br(\tau \to \mu\gamma).$$

The Flavour Problem





Vanille de Madagascar Alextrait de vanille de Madagascar



Caramel



Fraise



Menthe Aux éclats de chocolat noir



Violette



Café Avec du cafe Arabica de Colombie



Framboise



Mandarine



Citron vert



Noix de Coco****



Cassis



Fruit de la Passion



Chocolat



Rhum-Raisins au rhum" des Antilles françaises Avec des rations maceres au rhum"

The BSM Flavour Problem



Stringent limits on NP operators, e.g.

$$Br(l_i \to l_j \gamma) \approx \frac{24\sqrt{2}\pi^3 \alpha}{G_F^3 m_{l_i}^2 M_{NP}^4} \left| C_{ij} \right|^2$$

$$\begin{array}{ll} \circ & Br(\mu \to e\gamma) < 5.7 \times 10^{-13} & \Rightarrow & \left| C_{\mu e} \right| < 5 \times 10^{-9} \left(\frac{M_{NP}}{\text{TeV}} \right)^2 \\ \circ & Br(\tau \to l\gamma) < 4.0 \times 10^{-8} & \Rightarrow & \left| C_{\tau l} \right| < 6 \times 10^{-7} \left(\frac{M_{NP}}{\text{TeV}} \right)^2, \ l = e, \mu \end{array}$$

- > LFV couplings must be suppressed and/or New Physics scale is larger $\approx 10^3$ TeV
- Solutions
 - No New Physics at the TeV scale
 - Specific flavour structure of New Physics
 - Degeneracy
 - Symmetry (e.g. Minimal Flavour Violation)



- Effective operator for Majorana neutrino mass
 - Only dimension-5 operator beyond SM

$$\mathcal{L} \supset \frac{1}{2} \frac{h_{ij}}{\Lambda_{LNV}} (\bar{L}_i^c \cdot H) (H^T \cdot L_j) \xrightarrow[\langle H \rangle]{} \frac{1}{2} (m_v)_{ij} \bar{\nu}_i^c \nu_j$$



Seesaw Mechanism

• Add right-handed neutrinos N_i to SM, $M_N \approx 10^{14} \text{ GeV}$

$$\mathcal{L} \supset Y_{ij}^{\nu} \overline{N}_i L_j \cdot H - \frac{1}{2} M_{ij} \overline{N}_i N_j^c \xrightarrow{\mu \ll M_N} \frac{1}{2} (Y_{ki}^{\nu} M_{kl}^{-1} Y_{lj}^{\nu}) (\overline{L}_i^c \cdot H) (H^T \cdot L_j)$$

Light neutrino mass

$$m_{\nu} \approx 0.1 \text{ eV} \left(\frac{Y_{\nu} \langle H \rangle}{100 \text{ GeV}}\right)^2 \left(\frac{10^{14} \text{ GeV}}{M}\right)$$





Only dimension-5 operator beyond SM

$$\mathcal{L} \supset \frac{1}{2} \frac{h_{ij}}{\Lambda_{LNV}} (\bar{L}_i^c \cdot H) (H^T \cdot L_j) \xrightarrow[\langle H \rangle]{} \frac{1}{2} (m_v)_{ij} \bar{\nu}_i^c \nu_j$$



Seesaw Mechanism

- Sterile Neutrino Mass Scale Unknown
 - $\approx 10^{14}$ GeV Naïve Seesaw, GUTs
 - $\gtrsim 10^9$ GeV Thermal Leptogenesis
 - $\approx 10^3$ GeV Resonant Leptogenesis, LFV Contribution Production at LHC
 - $\approx 1 \text{ keV}$

• $\approx 1 \text{ eV}$

- Dark Matter Candidate
- Oscillation, Cosmology, $0\nu\beta\beta$





Only dimension-5 operator beyond SM

$$\mathcal{L} \supset \frac{1}{2} \frac{h_{ij}}{\Lambda_{LNV}} (\bar{L}_i^c \cdot H) (H^T \cdot L_j) \xrightarrow[\langle H \rangle]{} \frac{1}{2} (m_v)_{ij} \bar{\nu}_i^c \nu_j$$



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Seesaw Mechanism

Three possible mediators at tree level









Only dimension-5 operator beyond SM

$$\mathcal{L} \supset \frac{1}{2} \frac{h_{ij}}{\Lambda_{LNV}} (\overline{L}_i^c \cdot H) (H^T \cdot L_j) \xrightarrow{\langle H \rangle} \frac{1}{2} (m_\nu)_{ij} \overline{\nu}_i^c \nu_j$$



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Radiative Generation via Loops

• Alternative to Seesaw, e.g. R-Parity Violating SUSY





Problems of Naïve Seesaw

- Introduces high energy scale
- Right-handed neutrinos are singlets
 - Couple only via small mixture with active neutrinos
- Mechanism not testable with low energy observables
 - Negligible charged LFV
- Possible Solutions
 - SUSY Seesaw
 - Testable LFV effects from sleptons
 - "Bent" Seesaw mechanisms
 - Decouple Λ_{LNV} from heavy neutrino mass
 - Left–Right symmetric models
 - Right-handed neutrinos couple with gauge strength to charged leptons





Small Yukawas vs "Bent" Seesaw

- Seesaw Mechanism with TeV scale heavy neutrinos
 - Standard Seesaw with small Yukawa couplings
 - Charged LFV remains small
 - "Bent" Seesaw mechanisms
- Decouple Λ_{LNV} from Quasi-Degenerate Augorana Neutrinos heavy neutrino mass $M = 10^{3} {
 m GeV}$ • Example 10^{10} 10^{8} $\mathcal{M} = \begin{pmatrix} 0 & Y_{\nu} \langle H \rangle & 0 \\ Y_{\nu} \langle H \rangle & \mu & M \\ 0 & M & \mu \end{pmatrix}$ 0 10^{6} $m_{N,r} [GeV]$ 10⁴ Quasi–Dirac ^{10²} Majorana Neutrinos Standard Seesaw μ 10^{0} f. Inverse Seesaw 10^{-2} Potentially large 10^{-4} 10^{-6} charged LFV 10^{-8} • In the limit $\mu \rightarrow 0$, $m_{\nu} = 0.1 \, \text{eV}$ 10^{-10} no LNV but LFV 10^{-8} 10^{-6} 10^{-4} 10^{-2} 10^{0} 10^{2} 10^{4} 10^{6} 10^{8} 10^{10} 10^{12} 10^{14} μ [GeV]





Left-Right Symmetric Models



Based on $SU(3) \times SU(2)_L \times SU(2)_R \times U(1)_{B-L}$

Pati & Salam '74 Mohapatra & Senjanovic '75

- Higgs Sector
 - Bidoublet (EW Breaking)
 - Left-handed Triplet + Right-handed Triplet (Breaking Lepton Number + Parity + SU(2)_R)
- Generating N_i , W_R , Z_R masses $M_{N_i} \approx M_{W_R} \approx M_{Z_R} \approx \langle H_R \rangle \approx 0.5 5 \text{ TeV}$
- General Seesaw I+II mechanism M_{ν} =

$$M_{\nu} = \begin{pmatrix} M_L & M_D \\ M_D^T & M_R \end{pmatrix},$$

Charged current weak interactions

Neglect any Left-Right mixing

$$J_{W}^{\mu-} = \frac{g_L}{\sqrt{2}} \left(\bar{\nu} U_{LL} + \bar{N}^c U_{LR} \right) \gamma^{\mu} e_L + \frac{g_R}{\sqrt{2}} \sin \zeta_W \left(\bar{\nu} U_{RL} + \bar{N} U_{RR} \right) \gamma^{\mu} e_R, \qquad J_{W_L}^{\mu-} \approx \frac{g_L}{\sqrt{2}} U_{\ell i} \bar{\nu}_i \gamma^{\mu} \ell_L, \\ J_{W'}^{\mu-} = -\frac{g_L}{\sqrt{2}} \sin \zeta_W \left(\bar{\nu} U_{LL} + \bar{N} U_{LR} \right) \gamma^{\mu} e_L + \frac{g_R}{\sqrt{2}} \left(\bar{N} U_{RR} + \bar{\nu}^c U_{RL} \right) \gamma^{\mu} e_R, \qquad J_{W_R}^{\mu-} \approx \frac{g_R}{\sqrt{2}} V_{\ell i} \bar{N}_i \gamma^{\mu} \ell_R,$$

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25 / 33

Rare LFV Processes



$$BR(\mu \to e\gamma) \approx 2 \times 10^{-9} \sin^2(2\phi) \left(\frac{\Delta m_{12}^2}{m_{W_R}^2}\right)^2 \left(\frac{2 \text{ TeV}}{m_{W_R}}\right)^4,$$

µ-e conversion in nuclei enhanced via box diagrams

$$R(\mu \to e) \approx Br(\mu \to e\gamma)$$

• $\mu \rightarrow eee$ strongly enhanced due to tree level contribution

 $Br(\mu \rightarrow eee) \approx 300 \times Br(\mu \rightarrow e\gamma)$

- BSM Flavour Problem
 - Small mixing and/or mass differences required





LHC Searches





- Monte Carlo Simulation (PROTOS)
- Main background $t\bar{t}$, Z + jets (Pythia, Alpgen)
- Fast Detector Simulation (AcerDET)
- Selection Criteria
 - Number of jets $N_j > 2$
 - Number of isolated leptons $N_l = 2$
 - Invariant di-lepton mass $m_{ll} > 300 \text{ GeV}$
 - Total invariant mass



LHC Searches







Two heavy neutrinos with maximal mixing and 1% mass splitting

$$\begin{pmatrix} \nu_L & N_e & N_\mu \\ U_{PMNS} & 0 & 0 \\ 0 & \cos\theta & -\sin\theta \\ 0 & \sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} \nu_L \\ N_e \\ N_\mu \end{pmatrix}$$

$$\theta = \pi/4$$

 Correlation with low energy LFV processes





Two heavy neutrinos with maximal mixing and 1% mass splitting

$$\begin{pmatrix} \nu_L & N_e & N_\mu \\ U_{PMNS} & 0 & 0 \\ 0 & \cos\theta & -\sin\theta \\ 0 & \sin\theta & \cos\theta \end{pmatrix} \begin{pmatrix} \nu_L \\ N_e \\ N_\mu \end{pmatrix}$$

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 Correlation with low energy LFV processes





Two heavy neutrinos with maximal mixing and 1% mass splitting

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$$\theta = \pi/4$$

 Correlation with low energy LFV processes









suppressed as

 $\sigma_{LHC} \propto \Delta M_N^2 / (M_N \Gamma_N)$

 Correlation with low energy LFV processes suppressed as

 $\propto \Delta M_N^2/M_N^2$



Conclusion



LFV is crucial probe for BSM physics

- Discovery is smoking gun for BSM physics
- Strong experimental program with high sensitivity $\Lambda \approx 10^{3-4} \text{ TeV}$
- Observation is critical to solve flavour puzzle
- $\circ~$ Experimental problem: Comparatively weak sensitivity to τ

Strong connection to neutrino physics

- Oscillations demonstrate that lepton flavour is violated
- Models of neutrino mass generation predict wildly different LFV rates → LFV can be used to discriminate

Synergy with LHC searches

- Potential to observe charged Lepton Flavour Violation
- Complementarity of Observables